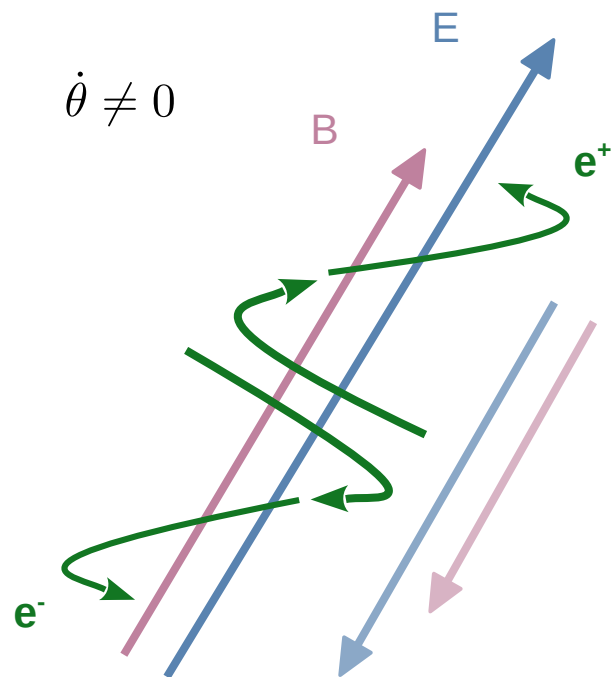


# Axion assisted Schwinger effect



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based on  
[2101.05192](#) and [2204.10842](#)  
w Yohei Ema and Kyohei Mukaida

# Axion assisted Schwinger effect

Schwinger particle production  
in electric field:

$$\sim \exp\left(\frac{-m^2 - p_T^2}{E}\right)$$



Axion assisted Schwinger  
particle production in  
electric field

$$\sim \exp\left(-\frac{m^2}{E}\right) \quad \text{for} \quad \dot{\theta}^2 \gtrsim \frac{m^2 p_T^2}{E}$$



high momentum part of distribution exponentially enhanced  
in presence of large axion velocity

# Outline

- Solving the Dirac equation with axion background field
- Axion assisted Schwinger effect and interpretation
- Basis independence and transient phenomena

# Dirac equation

$$S = \int d^4x \left[ \underbrace{\frac{1}{2} (\partial\phi)^2}_{\text{axion}} - \underbrace{V(\phi)}_{\text{gauge field}} - \underbrace{\frac{1}{4} F_{\mu\nu} F^{\mu\nu}}_{\text{fermion}} + \bar{\psi} \left( i \not{D} - m e^{2i c_m \phi / f_a \gamma_5} \right) \psi + c_A \frac{\alpha}{4\pi f_a} \phi F_{\mu\nu} \tilde{F}^{\mu\nu} + c_5 \frac{\partial_\mu \phi}{f_a} \bar{\psi} \gamma^\mu \gamma_5 \psi \right]$$

axion
gauge field
fermion
shift symmetric axion couplings

→ Dirac equation

$$\left[ i \not{D} - m e^{2i \theta_m \gamma_5} + \partial_\mu \theta_5 \gamma^\mu \gamma_5 \right] \psi = 0, \quad \theta_i = c_i \phi / f_a$$

(note symmetry  $\psi \rightarrow e^{i c \gamma_5 \phi / f_a} : c_5 \rightarrow c_5 - c, c_m \rightarrow c_m + c$ )

→ Solve for constant axion velocity  $\dot{\theta}$  and E-field  $E \hat{e}_z$  (switched on over finite time T)

# Solutions to the Dirac equation

formal solution:

$$\psi = \int \frac{d^3p}{(2\pi)^3} e^{i\vec{p}\cdot\vec{x}} e^{i\theta_5 \gamma_5} \sum_{\lambda=1,2} \left[ B_\lambda u_\lambda e^{-i \int dt \Omega} + D_\lambda^\dagger v_\lambda e^{+i \int dt \Omega} \right]$$

creation/annihilation operators

$$B_\lambda = \sum_{\lambda'=1,2} \left[ \alpha_\lambda^{(\lambda')} b_{\lambda'} - (-)^{\lambda+\lambda'} \beta_\lambda^{(\lambda')*} d_{\lambda'}^\dagger \right]$$

$$D_\lambda^\dagger = \sum_{\lambda'=1,2} \left[ \beta_\lambda^{(\lambda')} b_{\lambda'} + (-)^{\lambda+\lambda'} \alpha_\lambda^{(\lambda')*} d_{\lambda'}^\dagger \right]$$

creation/annihilation operators at infinite past

Bogoliubov coefficients

solutions of Dirac equation for constant background fields,

$$\dot{\theta} = 0, \quad E = -\dot{A}_z = 0$$

positive/negative frequency modes

$$\Omega = \sqrt{(p_z - gQ A_z)^2 + (m^2 + p_T^2)}$$

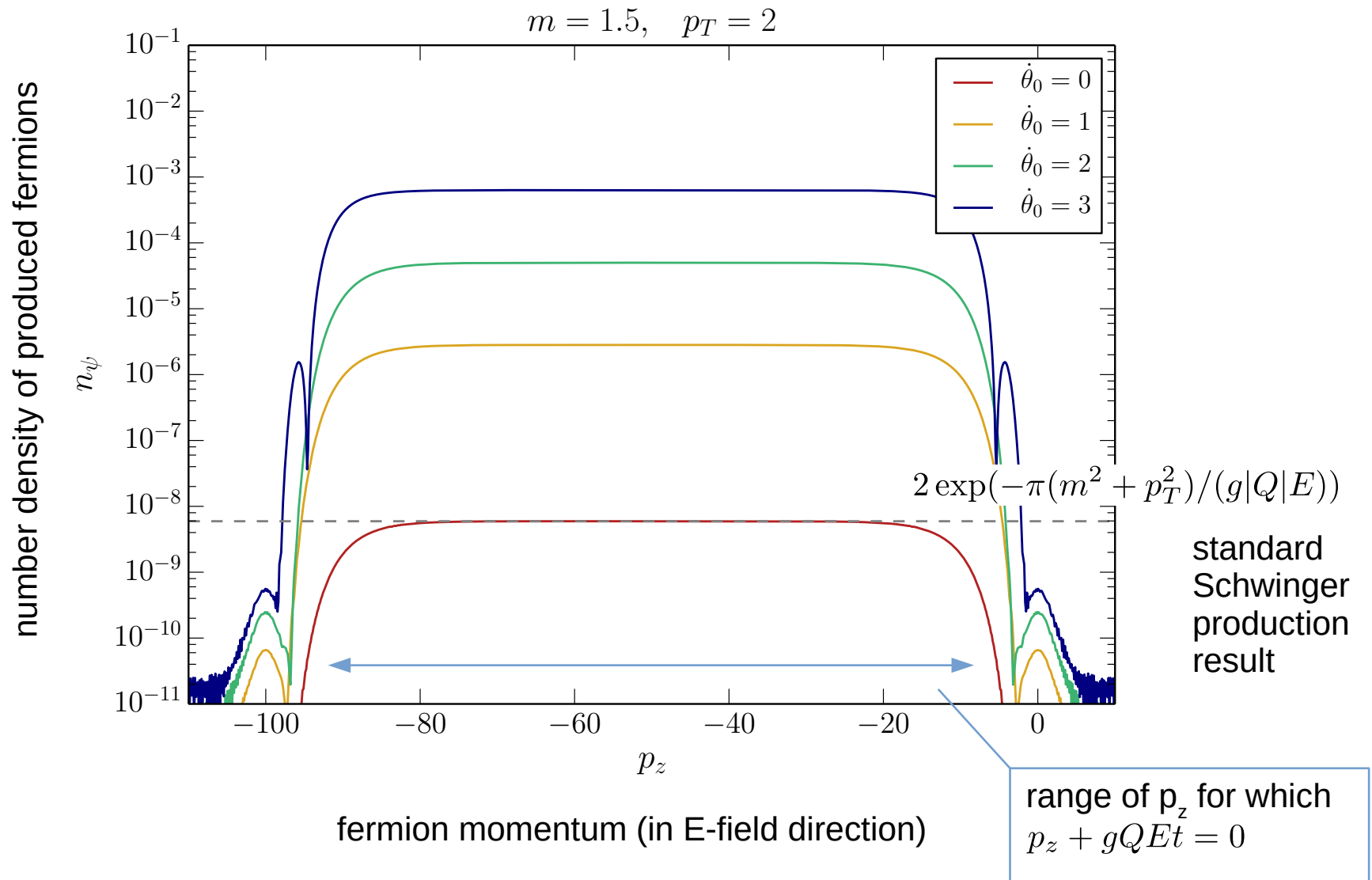
→ analytical solutions for  $u_\lambda, v_\lambda$  (fixes basis)

numerically solve Dirac equation for Bogoliubov coefficients, IC  $\alpha = 1, \beta = 0$

# Outline

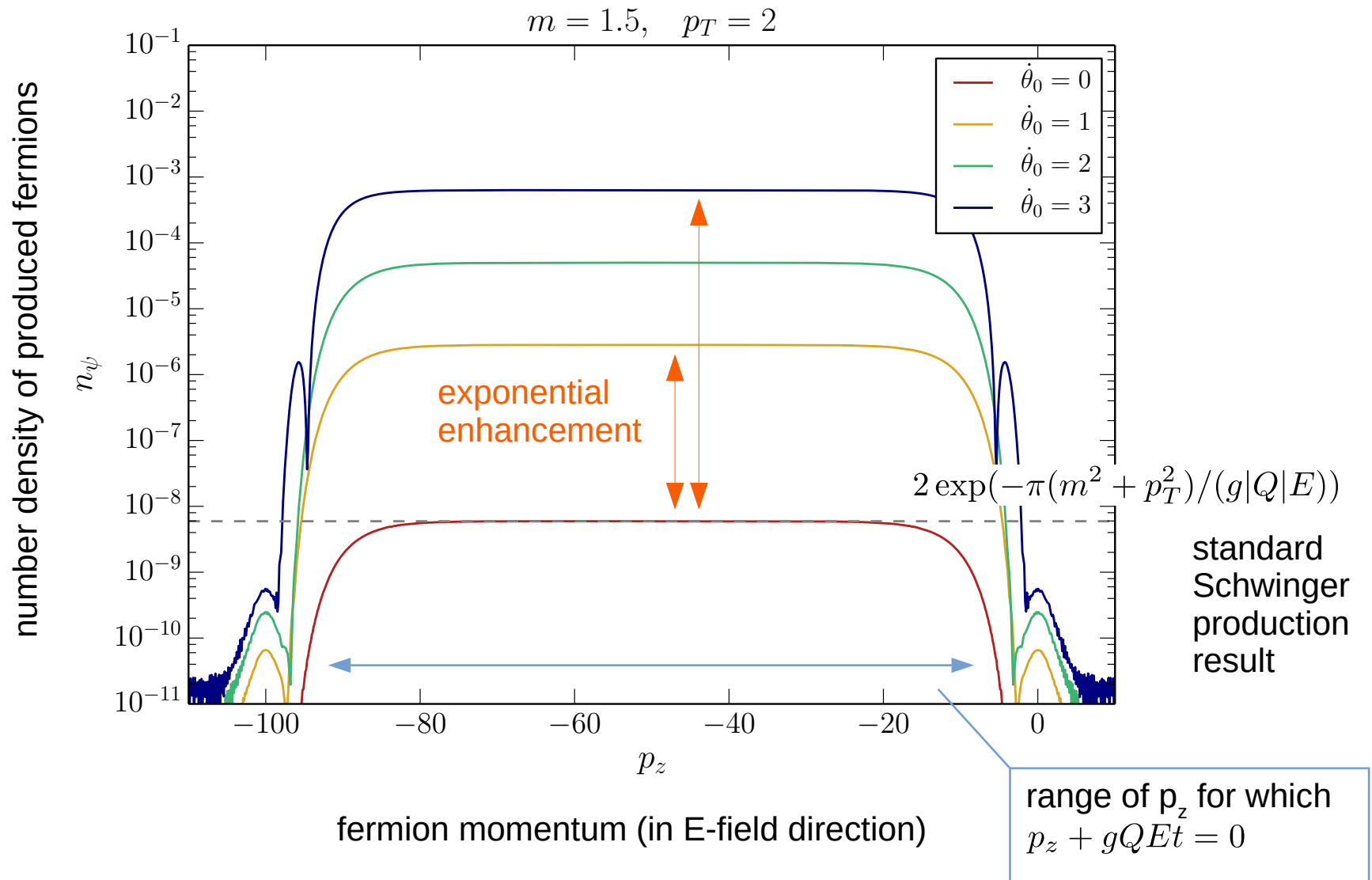
- Solving the Dirac equation with axion background field
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- Basis independence and transient phenomena

# Numerical solutions



units:  
 $g|Q|E = 1$

# Numerical solutions





# Non-relativistic limit

$$\mathcal{L}_\eta = \eta^\dagger i \partial_0 \eta + \frac{1}{2m} \eta^\dagger \left( \Pi^2 - \underbrace{2\dot{\theta}_{5+m} \Pi \cdot \sigma}_{\text{axion induced spin-momentum coupling}} + \dot{\theta}_{5+m}^2 \right) \eta + \mathcal{O}\left(\frac{1}{m^2}\right), \quad \Pi = (p_x, p_y, p_z - gQ A_z)$$

axion induced spin-momentum coupling

$$\longrightarrow \tilde{\Omega}_{\text{NR}}^\pm = \frac{1}{2m} \left( \sqrt{(p_z - gQ A_z)^2 + p_T^2} \pm \dot{\theta}_{5+m} \right)^2$$

minimized for non-zero axion velocity

spin-momentum coupling can compensate  $p_T$  suppression

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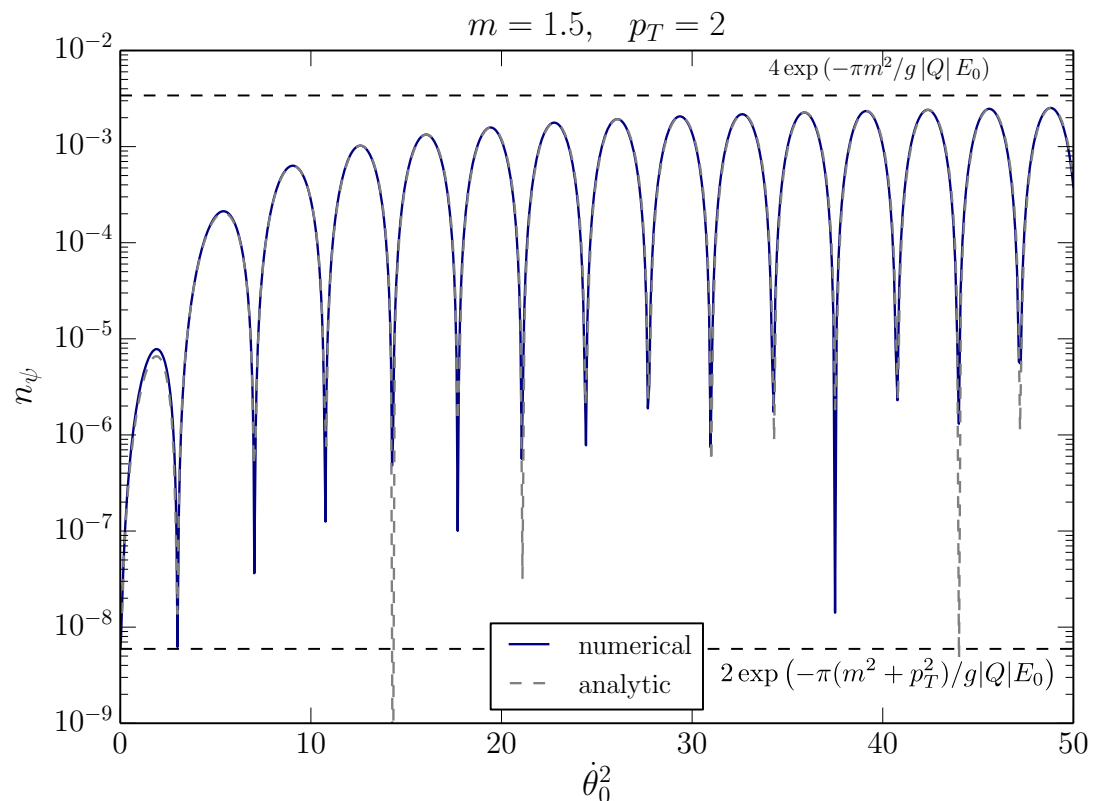
# Analytical solution in FCB

see also [Kitamoto, Yamada '21]

Choose  $c_m = 0$  (fermion current basis, FCB).

relevant energy levels:  $\tilde{\Omega}^{\pm}(t) = \sqrt{\left(\sqrt{(p_z - gQA_z)^2 + p_T^2} \pm \dot{\theta}_{5+m}\right)^2 + m^2}$

solve using  
WKB approximation



# Basis dependence

- for time-independent background, solutions  $u_\lambda, v_\lambda$  are uniquely fixed\*
- for time-dependent background, any complete set of  $u_\lambda, v_\lambda$  is a priori equally good
- for evaluating physical observables in the asymptotic future (= time-independent background) all choices trivially give same result
- when evaluating physical observables at intermediate times (transient phenomena) we need to take care of divergences, regularization and renormalization
- Obviously, if we do everything correctly, physical observables, even at intermediate times, will not depend on the basis choice.

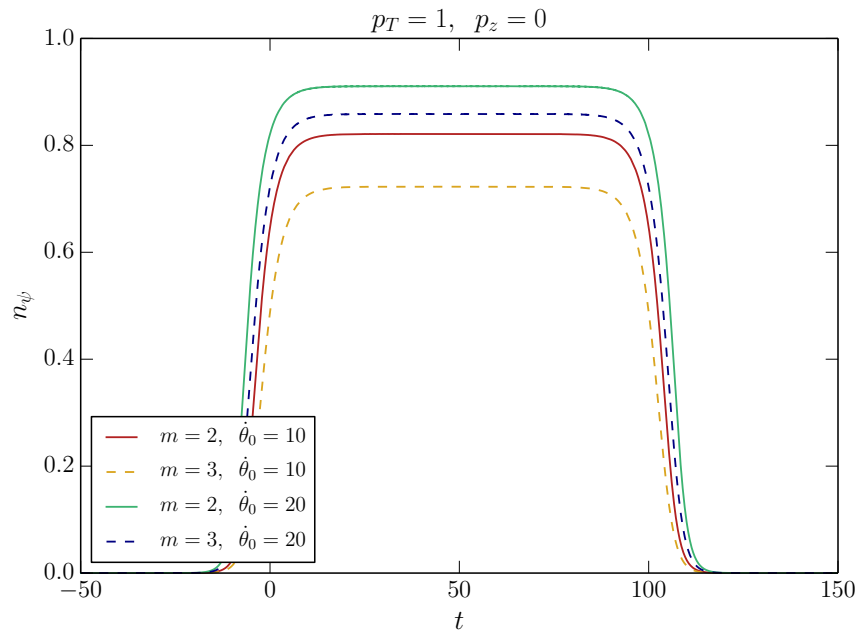
see also [Adshead, Lozanov `21]

\* up to normalization, symmetries, degenerate eigenvalues

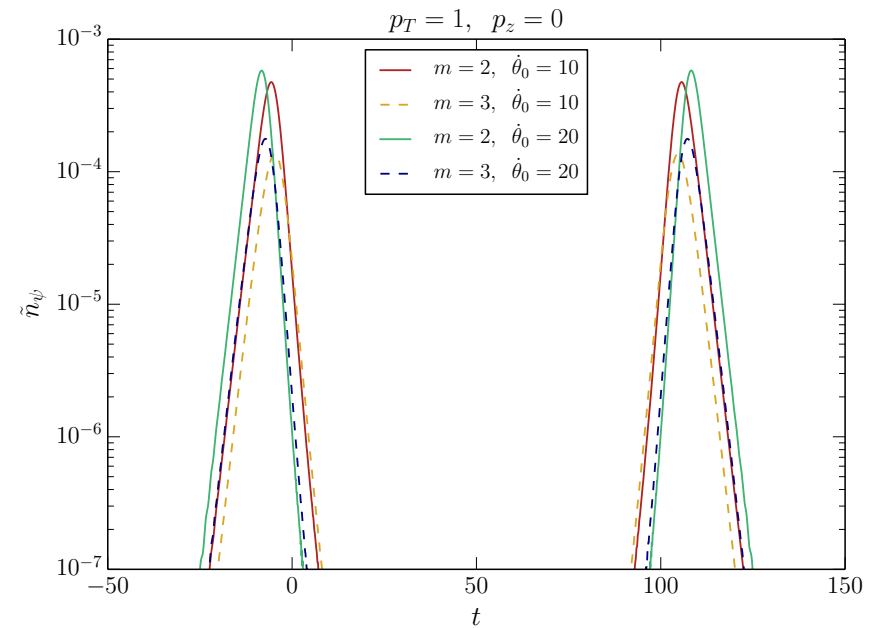
# Pitfalls and resolutions (1)

number density  $n_\psi \sim |\beta|^2$

no E-field,  $\dot{\theta} \neq 0$  for  $0 < t < 100$



Hamiltonian basis  $c_5 = 0$



Fermion current basis  $c_m = 0$

units:

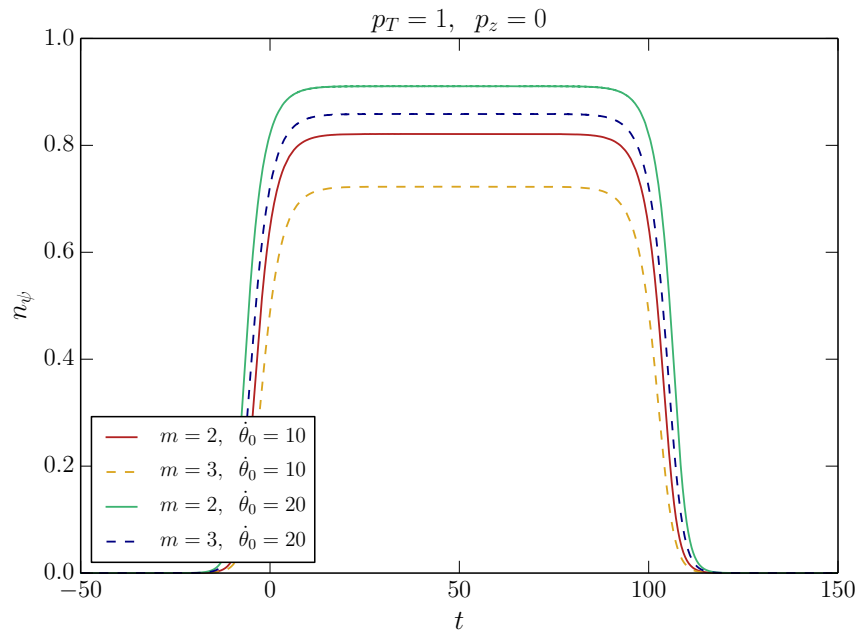
$p_T = 1$

Axion assisted Schwinger effect

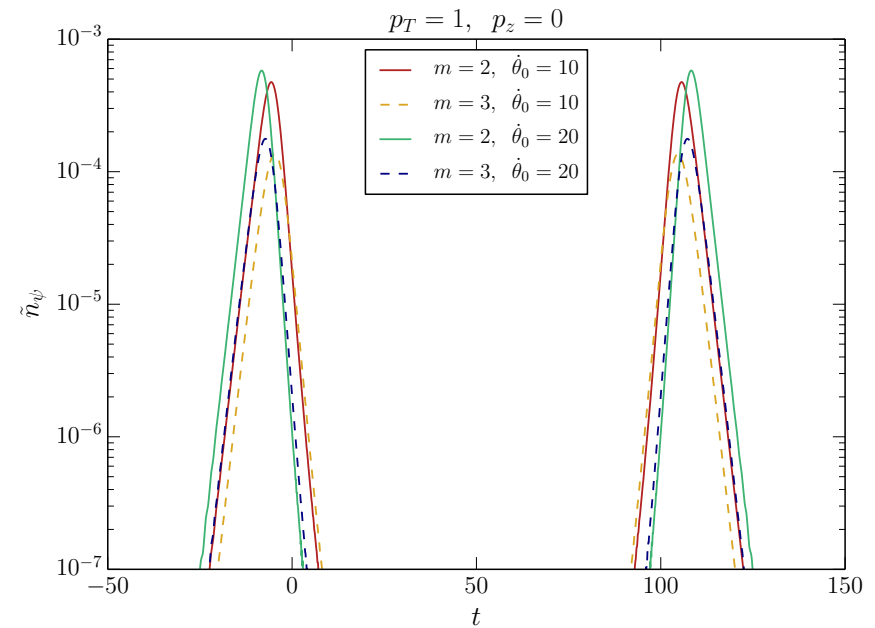
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→  $|\beta|^2$  is not an observable

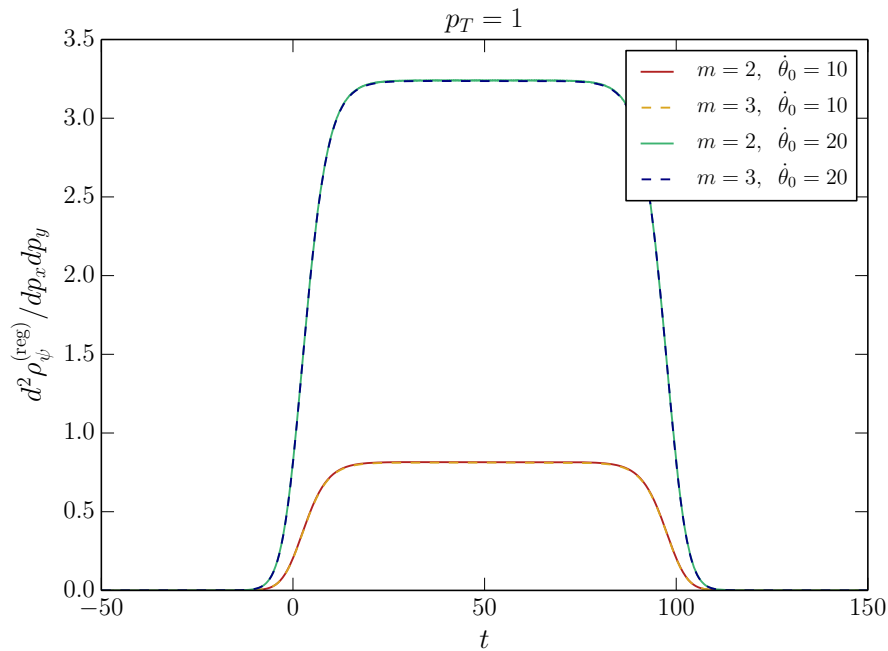
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Axion assisted Schwinger effect

# Pitfalls and resolutions (2)

energy density  $\rho_\psi = \frac{\langle H_\psi \rangle}{\text{Vol}} = \rho_\psi^{(\text{reg})} + \rho_\psi^{(\text{vac})}$

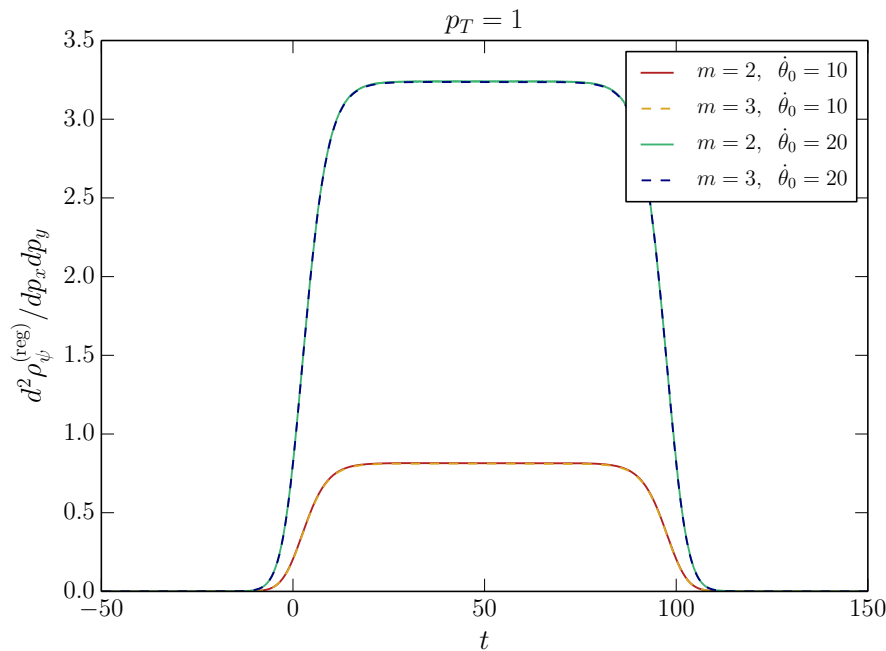


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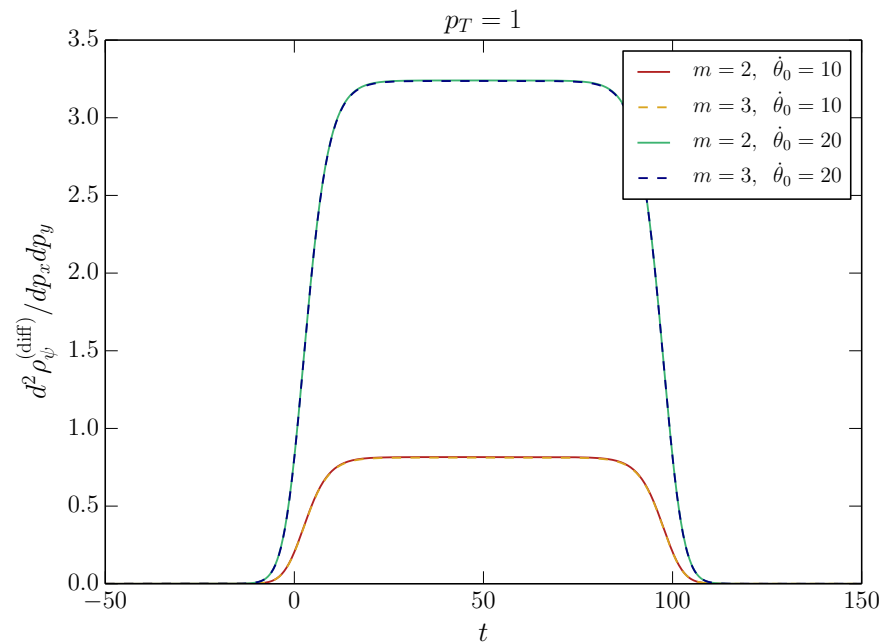
Fermion current basis  $c_m = 0$

# Pitfalls and resolutions (2)

energy density  $\rho_\psi = \frac{\langle H_\psi \rangle}{\text{Vol}} = \rho_\psi^{(\text{reg})} + \rho_\psi^{(\text{vac})}$   $\rho^{(\text{diff})} \equiv: \rho_\psi^{(\text{reg})} : \Big|_{\text{HB}} - : \rho_\psi^{(\text{reg})} : \Big|_{\text{FCB}}$



Hamiltonian basis  $c_5 = 0$



Fermion current basis  $c_m = 0$

- ➔ normal ordering assumes that the argument of the regulator function is an eigenvalue of the energy density operator (btw, dimensional regularization also fails in FCB due to  $\gamma_5$ )



# Final thoughts

- Everything is frame independent, massless limit and anomaly equation are reproduced (as it should!)
- exponential enhancement of large momentum particles in axion assisted Schwinger effect
  - requires large axion velocity
  - pheno consequences ?
- $O(1)$  transient energy density (even no  $\exp(-m^2/E)$  suppression!)
  - requires large axion velocity and specification of UV model
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**Thank you!**

# UV sensitivity

- very large transient effect found in physical observable  $\rho_\psi$  , frame independent
- arises even without any EM background fields
- can be calculated analytically:

$$\rho_\psi^{(\text{diff})} = \frac{m^2 \dot{\theta}_{5+m}^2}{4\pi^2} \log \left( \frac{\Lambda^2}{m^2} \right) + \frac{\dot{\theta}_{5+m}^4}{4\pi^2}$$

(  $\rho_\psi \rightarrow 0$  for  $m \rightarrow 0$  as expected )

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dim-8 operator



(  $\rho_\psi \rightarrow 0$  for  $m \rightarrow 0$  as expected )



low energy axion EFT is non-renormalizable

discussing this dim-8 operator requires specification of UV theory