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Nuclear structure at New Facilities with Rare Isotope Beams

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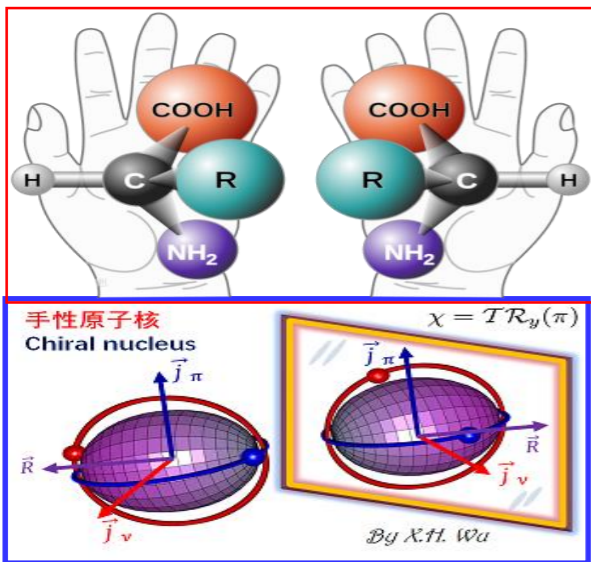
Collaborative Research Centre



国家自然科学基金委员会
National Natural Science Foundation of China

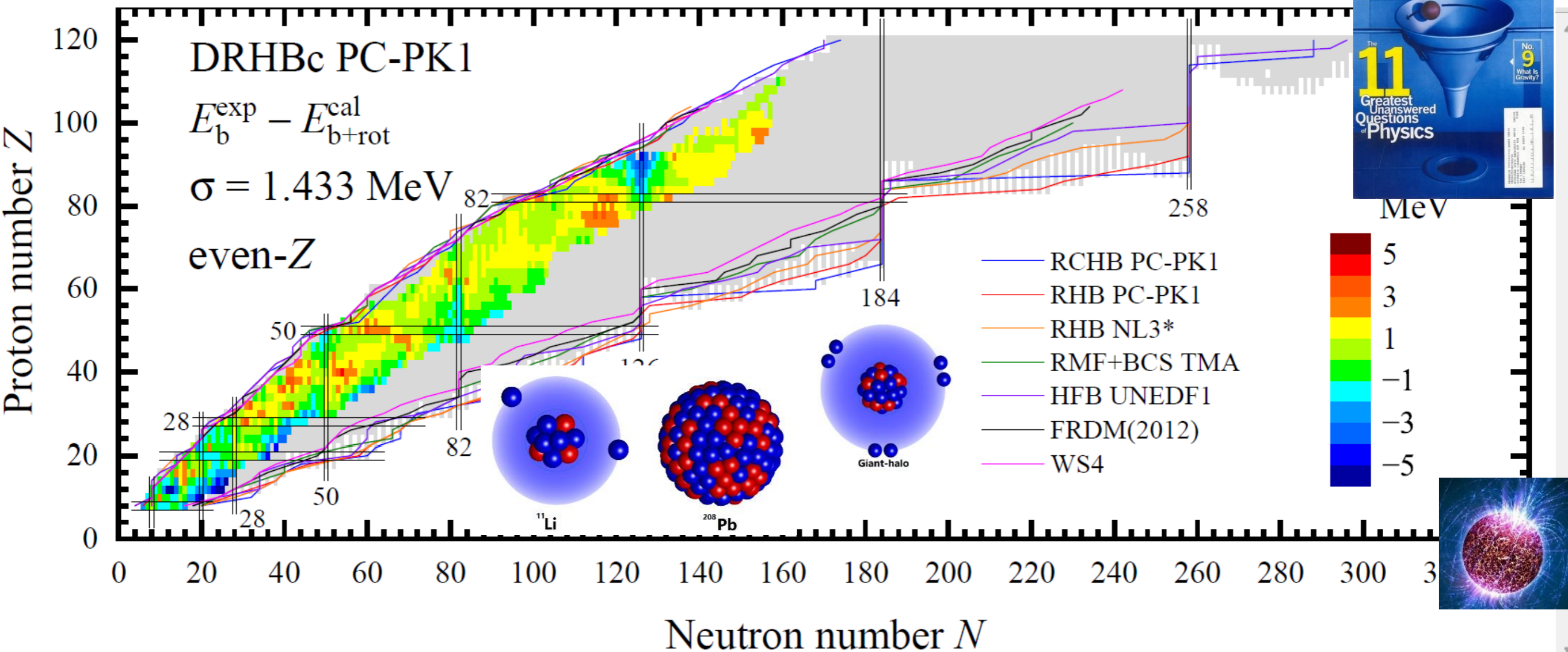
DFG Deutsche
Forschungsgemeinschaft

Existence Limit of nuclei



- stable nuclei
- unstable nuclei observed
- drip-lines (predictions)
- == magic numbers

~ 300 nuclei
~ 2700 nuclei
~ 10000 nuclei

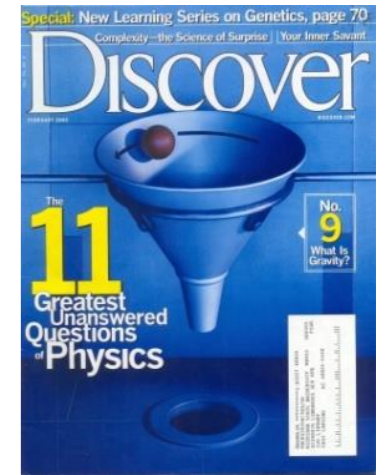


PC-PK1 + DRHBc : 4829 bound even- Z nuclei with $8 \leq Z \leq 120$

DRHBc Collaboration, [At. Data Nucl. Data Tables 144, 101488 \(2022\)](#)

How were the heavy elements from iron to Uranium made ?

Discovery : "11 greatest unanswered questions of Physics"

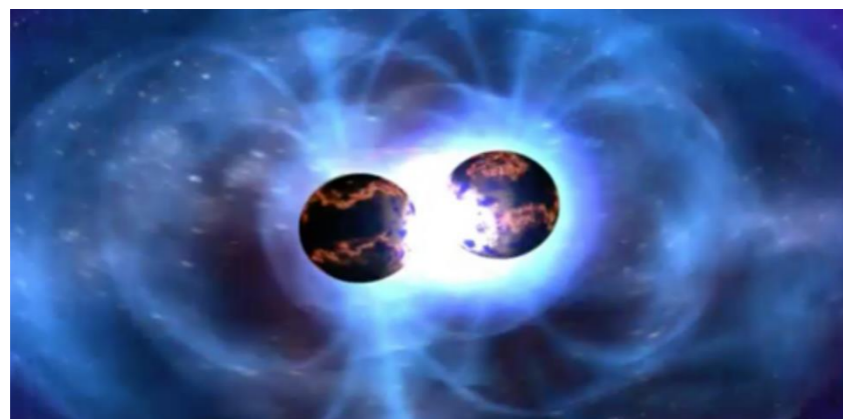


Key mechanism : rapid neutron capture process (r-process)

r-process sites

GW170817 neutron-star-merger event shows that neutron star merger is one of the r-process sites

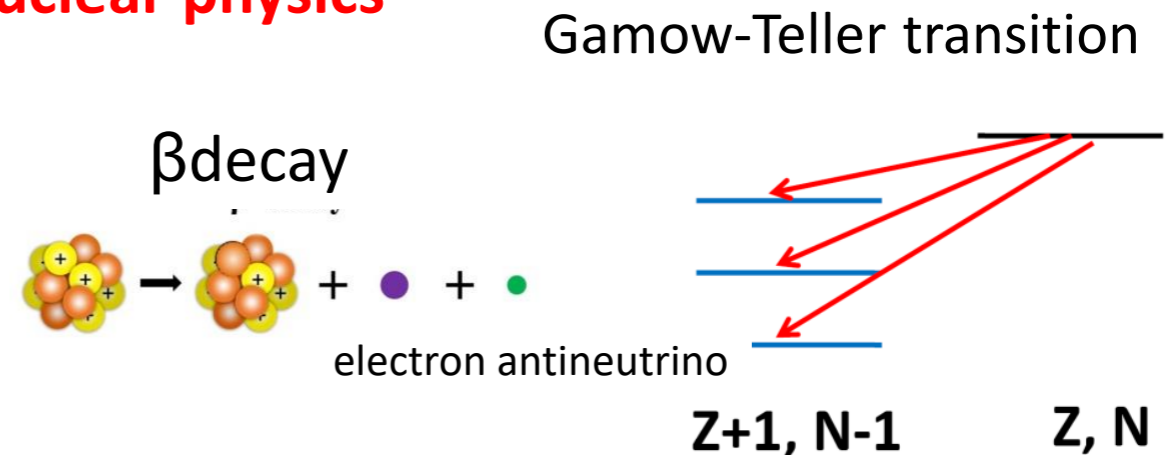
ApJL 848, L12 (2017)



big process in nucleosynthesis study

Systematic nuclear physics knowledge: mass, β -decay half-lives, reaction rates ...

Accurate description: difficult problem in nuclear physics



key : difficulties in experiments

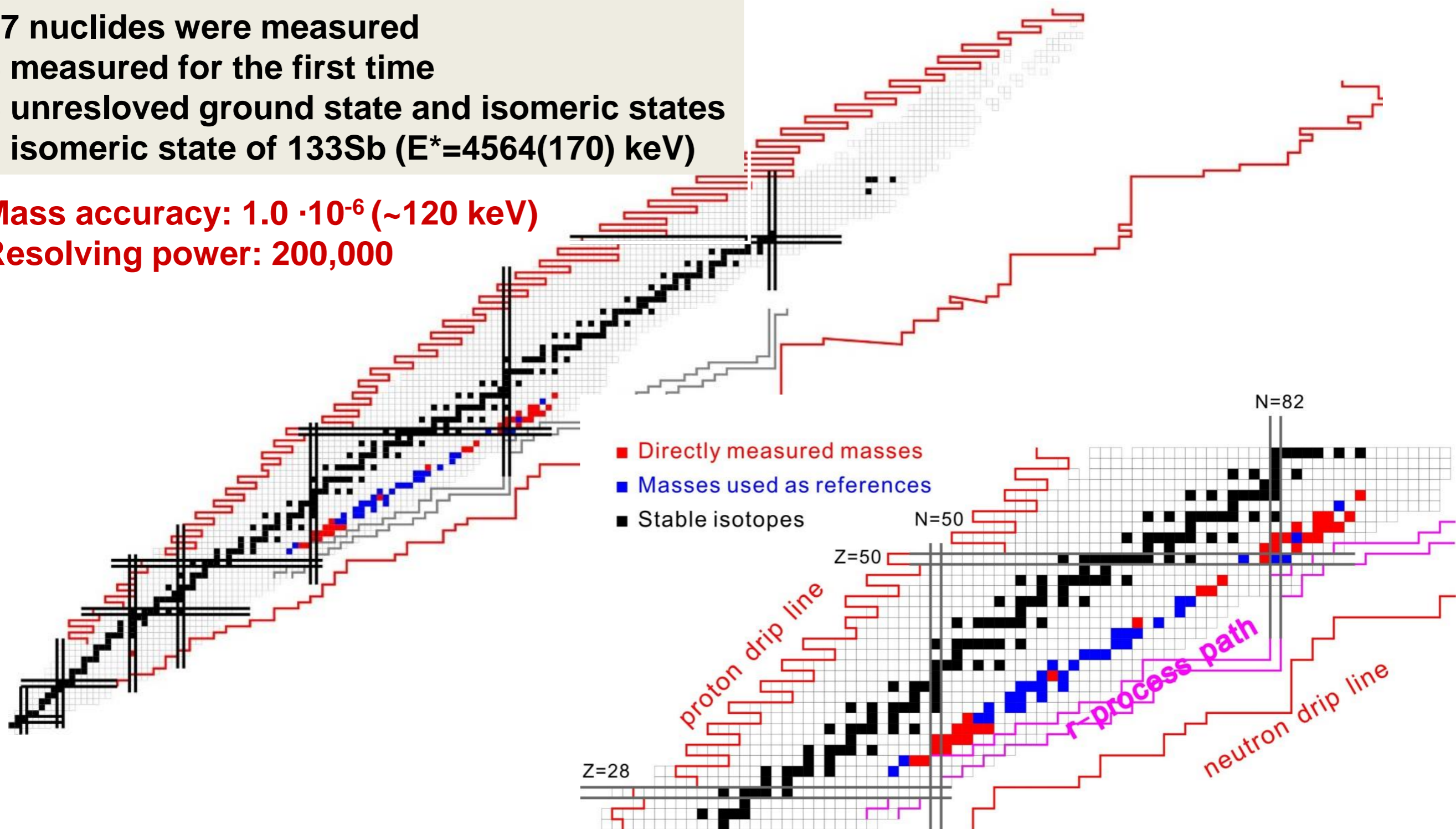
Large-scale accurate theoretical calculations nuclear are demanded

Mass measurement

Short-Lived Neutron-Rich Nuclei with the Novel Large-Scale Isochronous Mass Spectrometry at the FRS-ESR Facility Sun et al. NPA 812 (2008) 1-12

71 nuclides covered
27 nuclides were measured
8 measured for the first time
8 unresolved ground state and isomeric states
1 isomeric state of ^{133}Sb ($E^*=4564(170)$ keV)

Mass accuracy: $1.0 \cdot 10^{-6}$ (~120 keV)
Resolving power: 200,000

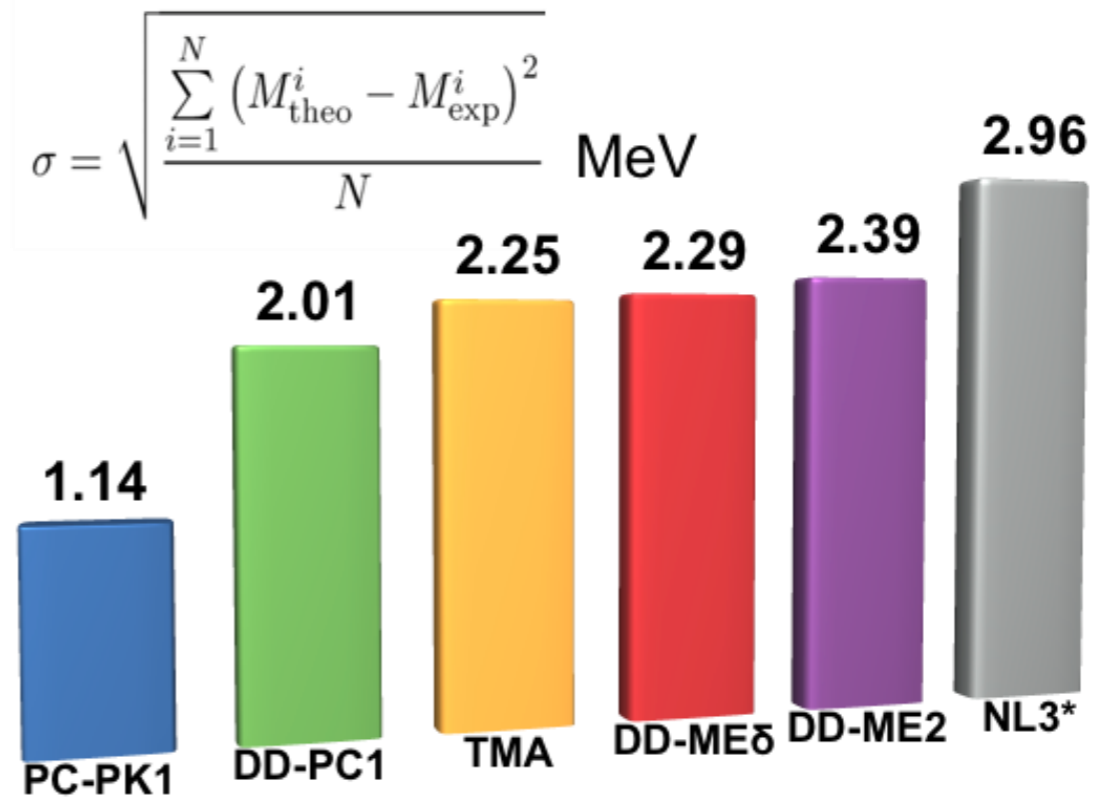
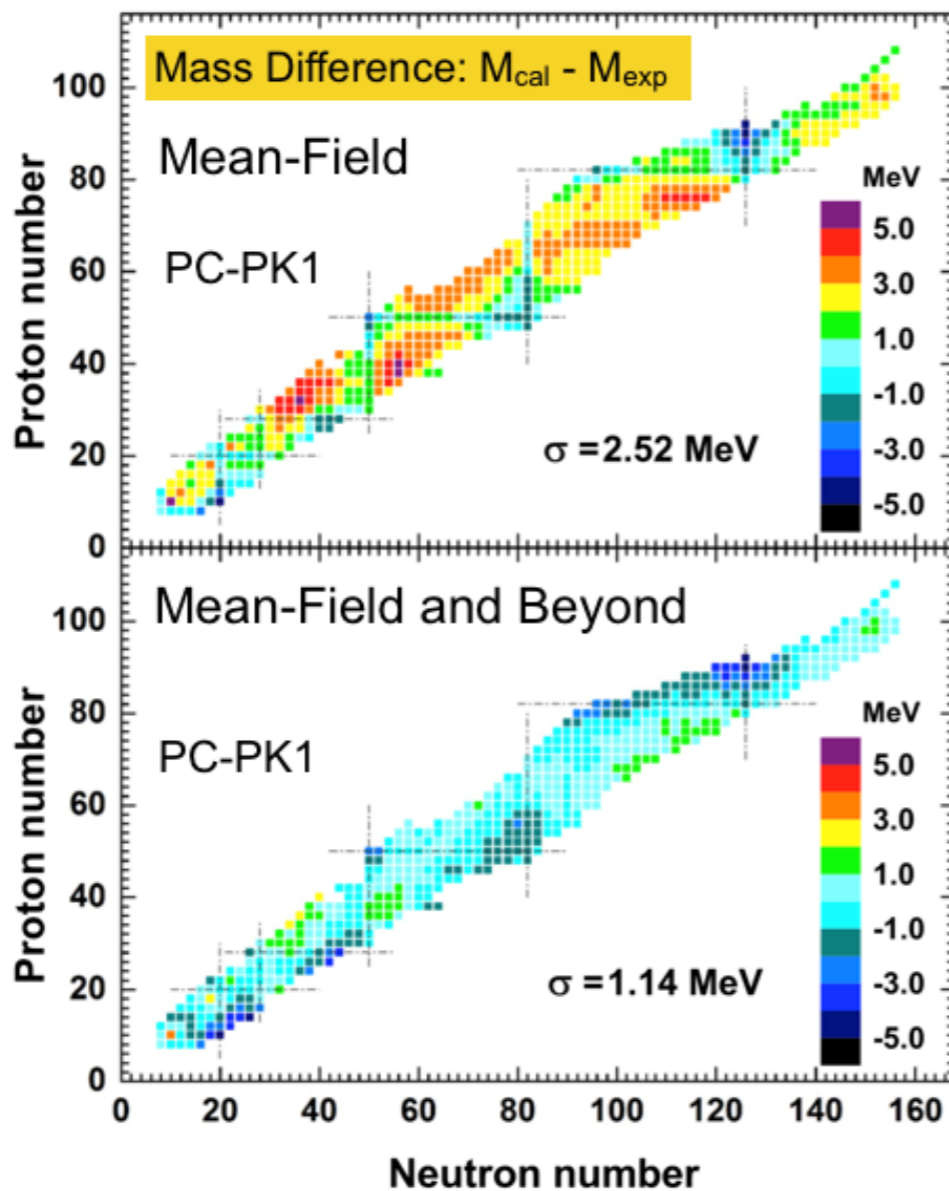


Facility for Antiproton and Heavy Ion Research (FAIR) Darmstadt



- Relativistic Density Functional**
- Physics close to the drip-line**
- DRHBc mass table collaboration**
- Physics along $N=Z$ nuclei**
- Relativistic density functional theory in 3D lattice**
 - Linear alpha-chain**
 - Nuclear fusion**
 - Nuclear fission**
 - Chiral dynamics**
- ReCD theory**
- Toward Relativistic ab initio DFT**

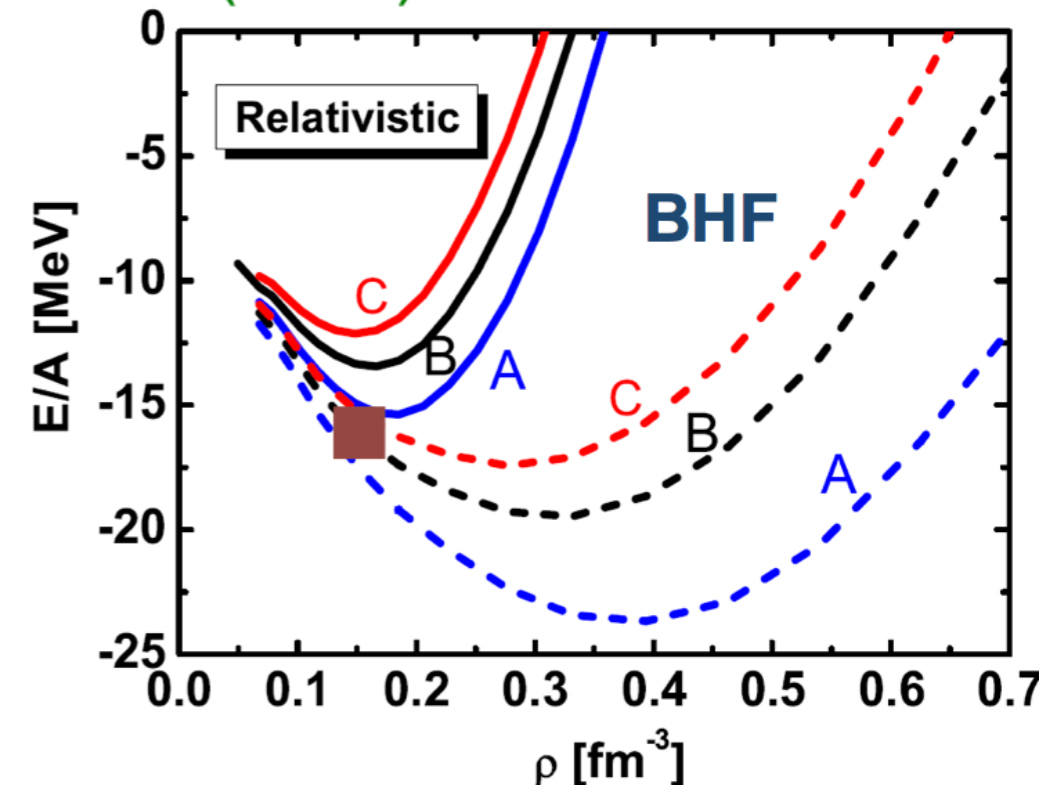
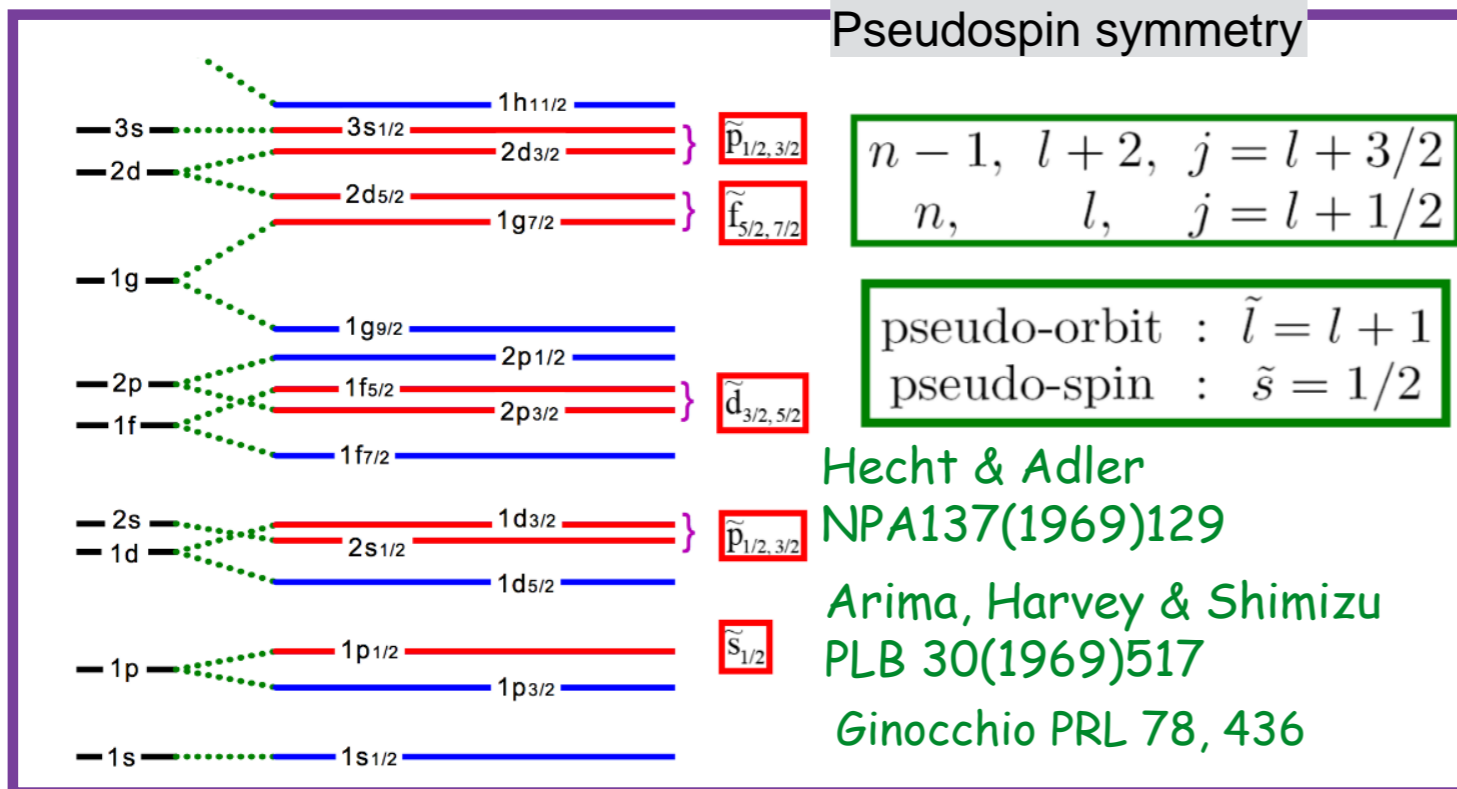
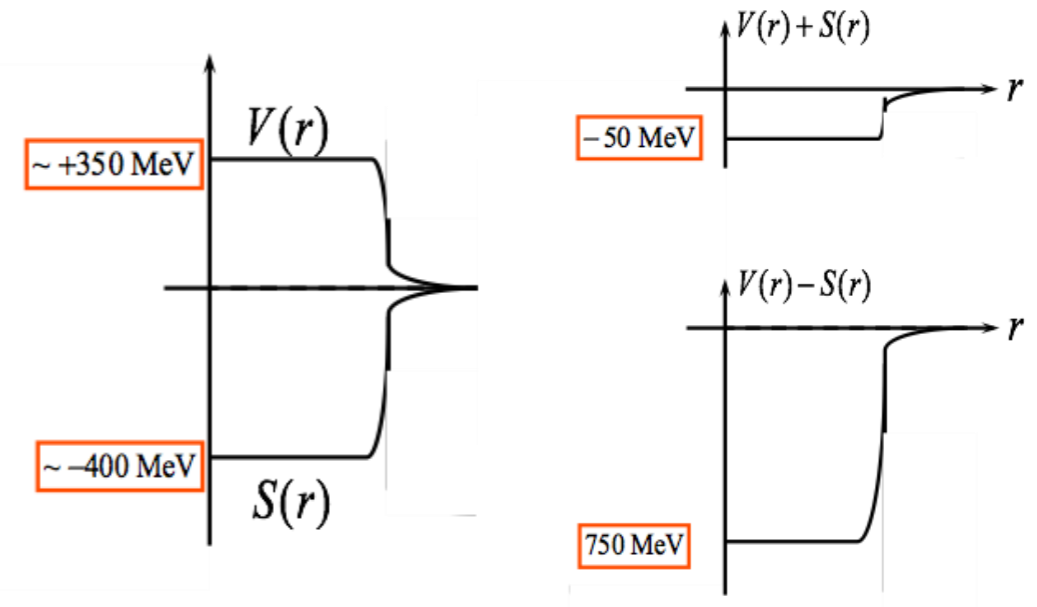
Relativistic functional PC-PK1



Why relativistic?

- ✓ **Spin-orbit** automatically included
- ✓ **Lorentz covariance** restricts parameters
- ✓ **Pseudo-spin Symmetry**
- ✓ Connection to QCD: big $V/S \sim \pm 400$ MeV
- ✓ Consistent treatment of **time-odd fields**
- ✓ Relativistic **saturation mechanism**
- ✓ ... **Liang, Meng, Zhou, Physics Reports 570 : 1-84 (2015).**

Relativistic Density Functional for Nuclear Structure
Meng (Ed.), World Scientific, Singapore (2016)



Brockmann & Machleidt, PRC42, 1965 (1990)

Relativistic Density Functional Theory

Relativistic Density Functional for Nuclear Structure

Meng (Ed.), World Scientific, Singapore (2016)

Elementary building blocks

$$(\bar{\psi} \mathcal{O}_\tau \Gamma \psi) \quad \mathcal{O}_\tau \in \{1, \tau_i\} \quad \Gamma \in \{1, \gamma_\mu, \gamma_5, \gamma_5 \gamma_\mu, \sigma_{\mu\nu}\}$$

Densities and currents

Isoscalar-scalar

$$\rho_S(\mathbf{r}) = \sum_k^{occ} \bar{\psi}_k(\mathbf{r}) \psi_k(\mathbf{r})$$

Isoscalar-vector

$$j_\mu(\mathbf{r}) = \sum_k^{occ} \bar{\psi}_k(\mathbf{r}) \gamma_\mu \psi_k(\mathbf{r})$$

Isovector-scalar

$$\vec{\rho}_S(\mathbf{r}) = \sum_k^{occ} \bar{\psi}_k(\mathbf{r}) \vec{\tau} \psi_k(\mathbf{r})$$

Isovector-vector

$$\vec{j}_\mu(\mathbf{r}) = \sum_k^{occ} \bar{\psi}_k(\mathbf{r}) \vec{\tau} \gamma_\mu \psi_k(\mathbf{r})$$

Energy Density Functional

$$E_{kin} = \sum_k v_k^2 \int \bar{\psi}_k (-\gamma \nabla + m) \psi_k d\mathbf{r}$$

$$E_{2nd} = \frac{1}{2} \int (\alpha_S \rho_S^2 + \alpha_V \rho_V^2 + \alpha_{tV} \rho_{tV}^2) d\mathbf{r}$$

$$E_{hot} = \frac{1}{12} \int (4\beta_S \rho_S^3 + 3\gamma_S \rho_S^4 + 3\gamma_V \rho_V^4) d\mathbf{r}$$

$$E_{der} = \frac{1}{2} \int (\delta_S \rho_S \Delta \rho_S + \delta_V \rho_V \Delta \rho_V + \delta_{tV} \rho_{tV} \Delta \rho_{tV}) d\mathbf{r}$$

$$E_{em} = \frac{e}{2} \int j_\mu^p A^\mu d\mathbf{r}$$

Dirac equation

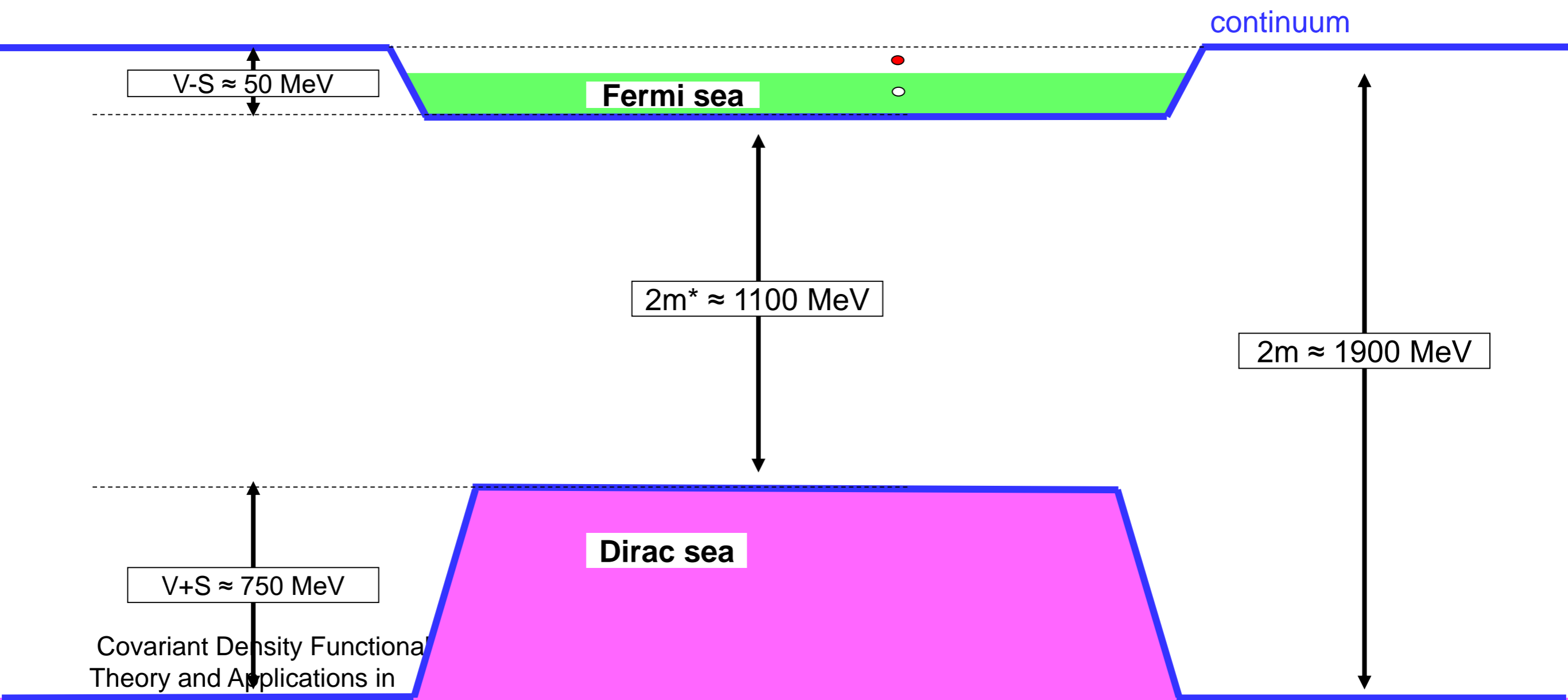
$$\begin{pmatrix} m + S + V & \sigma(p - V) \\ \sigma(p - V) & -m - S + V \end{pmatrix} \begin{pmatrix} f \\ g \end{pmatrix} = \epsilon \begin{pmatrix} f \\ g \end{pmatrix}$$

scalar potential:

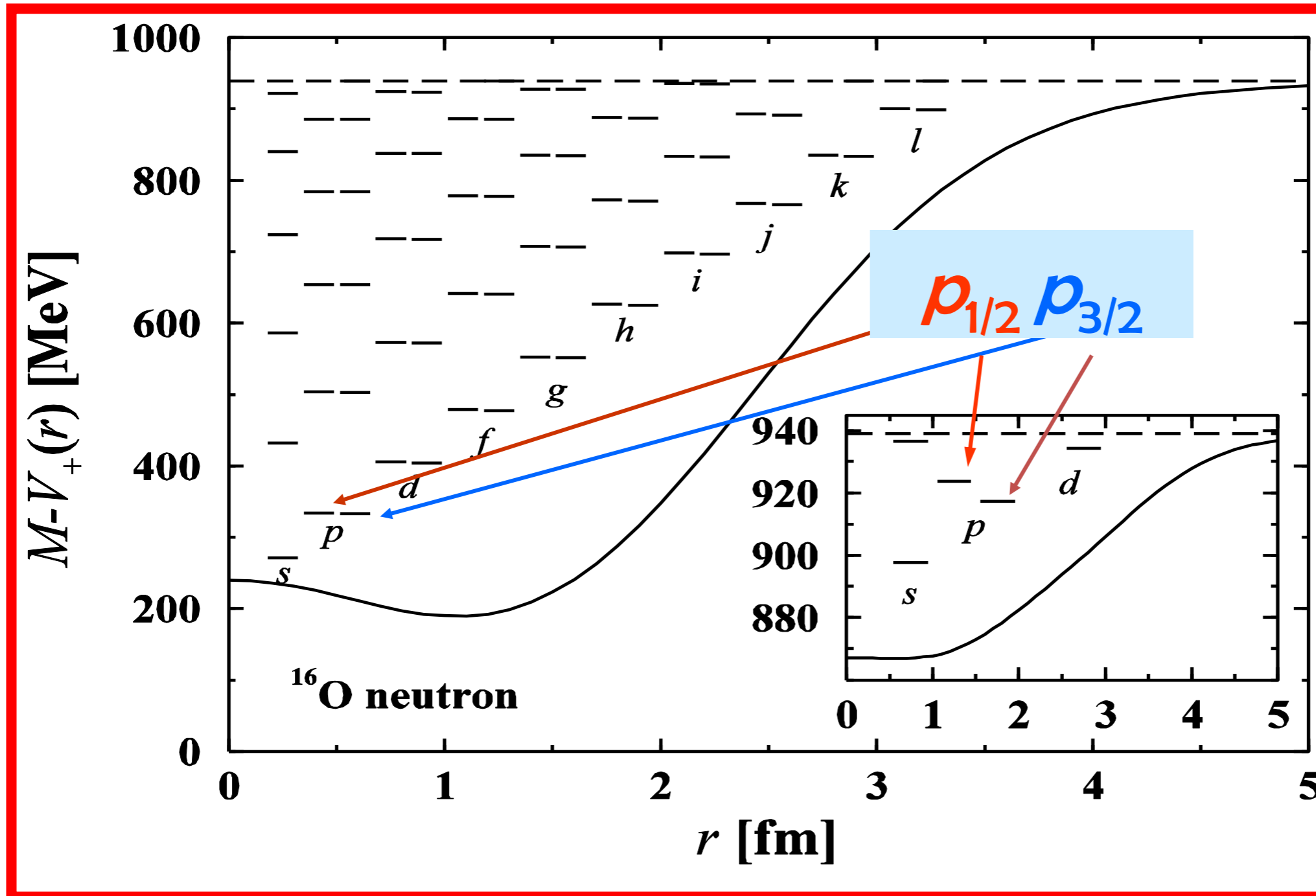
$$S(r) \approx -400 \text{ MeV}$$

vector potential:

$$V(r) \approx 350 \text{ MeV}$$

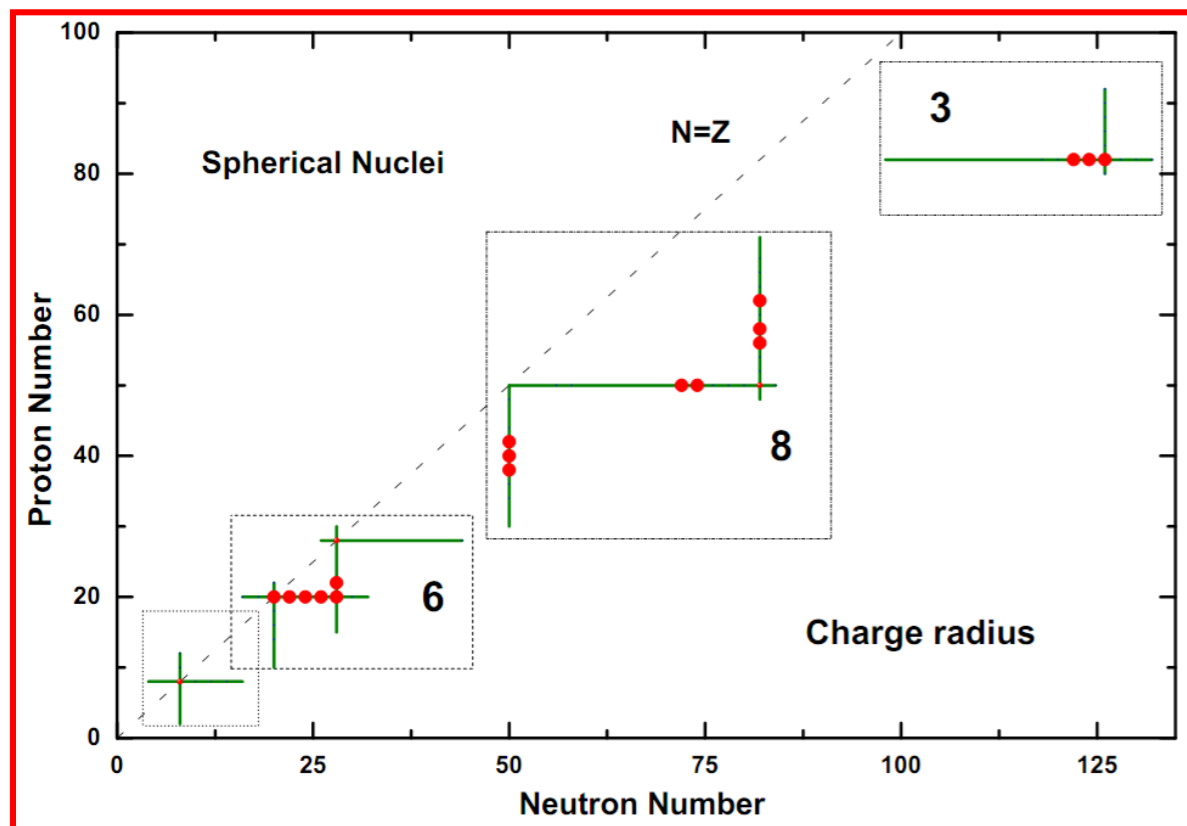
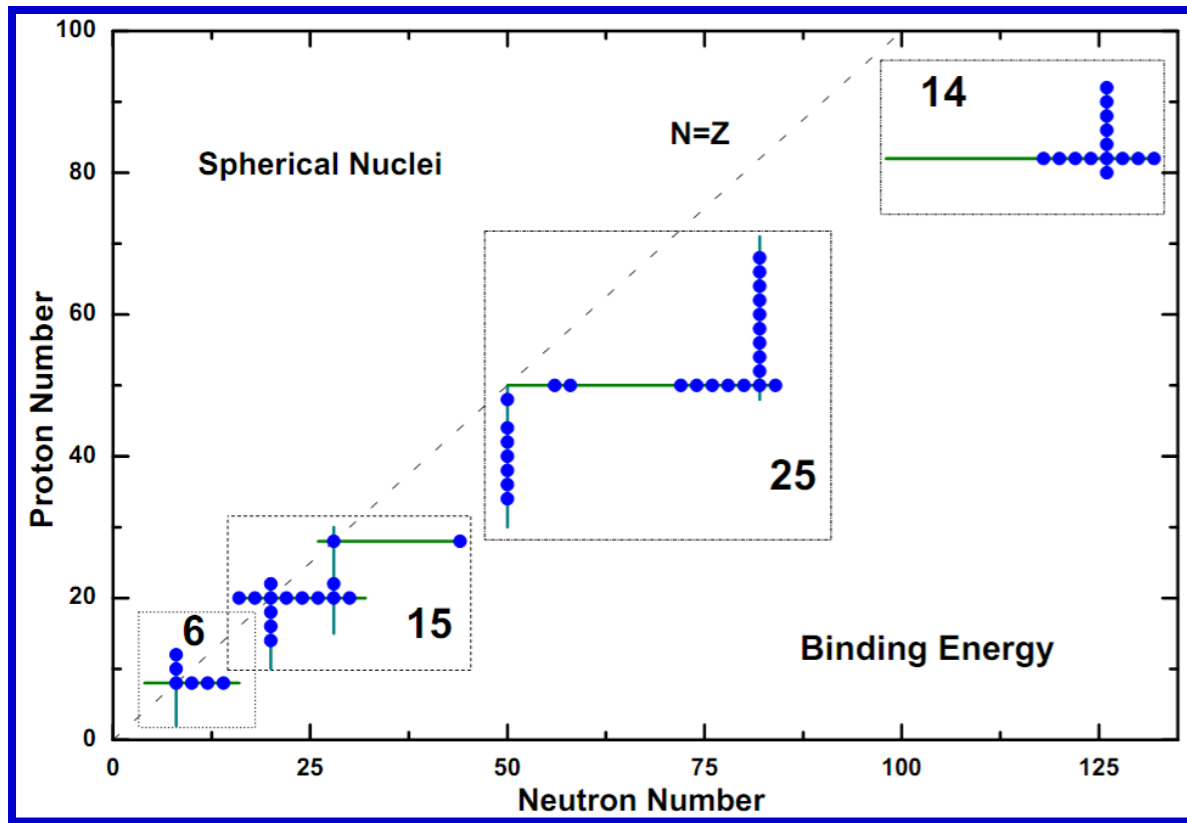


Dirac Sea: Spin symmetry



S. G. Zhou, J. Meng and P. Ring, PRL92(03)262501

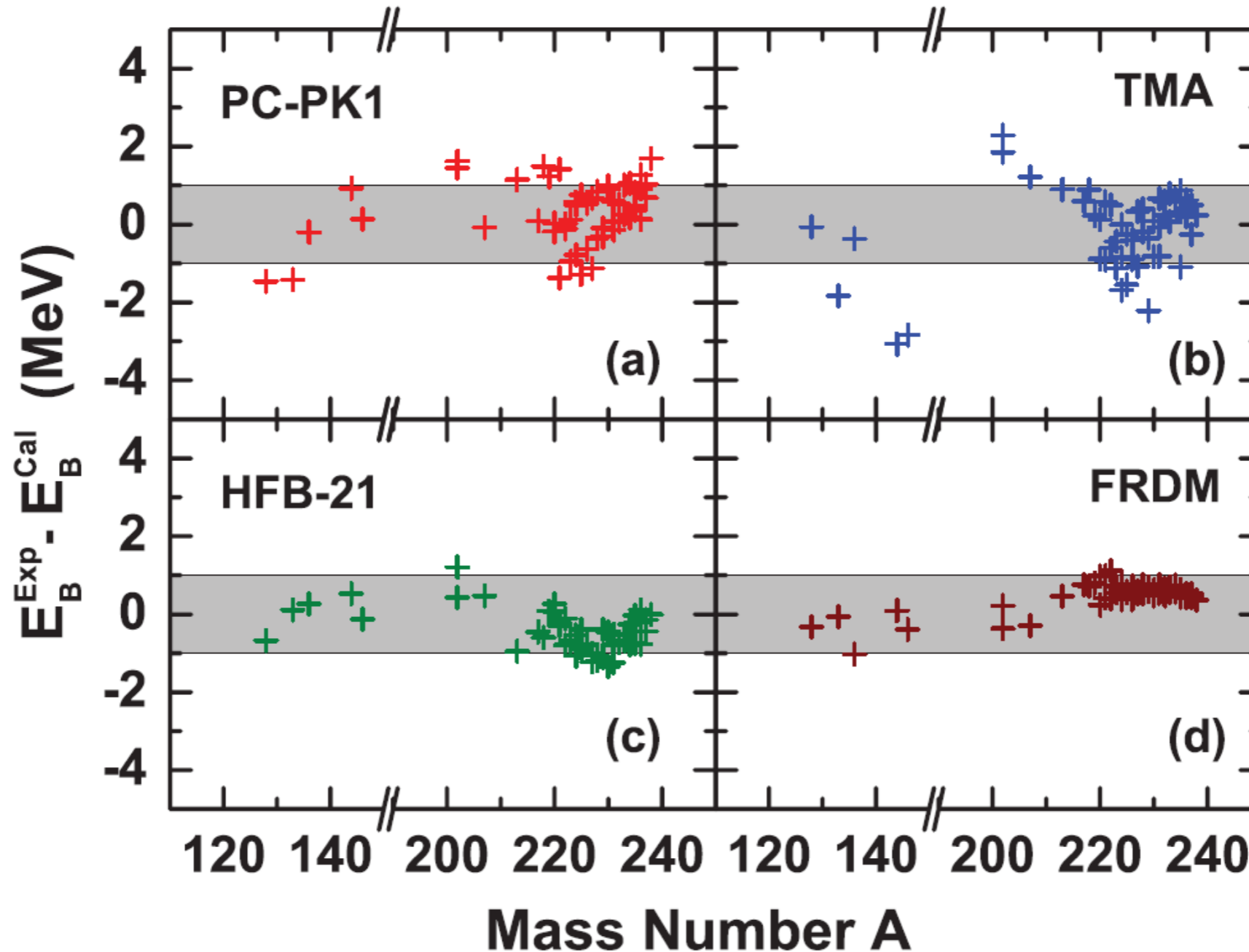
Relativistic functional PC-PK1



Coupl. Cons.	PC-PK1	Dimension
α_S [10 ⁻⁴]	-3.96291	MeV ⁻²
β_S [10 ⁻¹¹]	8.66530	MeV ⁻⁵
γ_S [10 ⁻¹⁷]	-3.80724	MeV ⁻⁸
δ_S [10 ⁻¹⁰]	-1.09108	MeV ⁻⁴
α_V [10 ⁻⁴]	2.69040	MeV ⁻²
γ_V [10 ⁻¹⁸]	-3.64219	MeV ⁻⁸
δ_V [10 ⁻¹⁰]	-4.32619	MeV ⁻⁴
α_{TV} [10 ⁻⁵]	2.95018	MeV ⁻²
δ_{TV} [10 ⁻¹⁰]	-4.11112	MeV ⁻⁴
V_n [10 ⁰]	-349.5	MeV fm ³
V_p [10 ⁰]	-330	MeV fm ³

Zhao, Li, Yao, Meng, PRC 82, 054319 (2010)

Predictive power



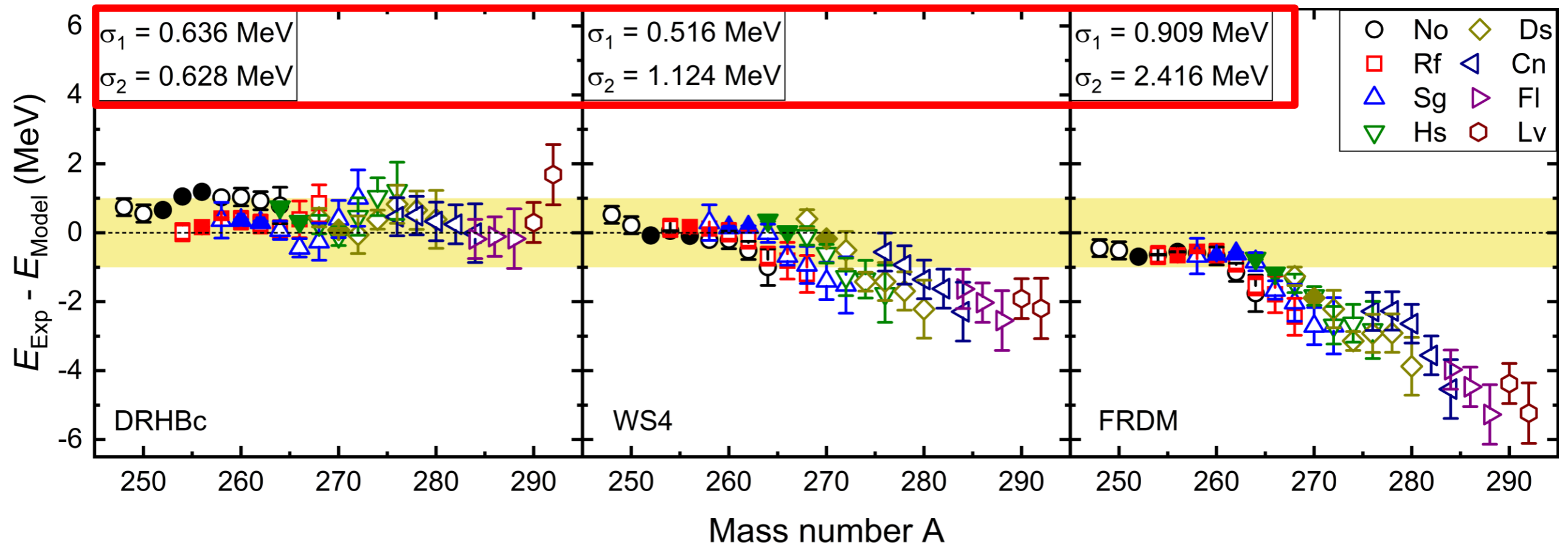
P. W. Zhao, *et al.* Phys. Rev. C, 86 024324 (2012)

Data from L. Chen, *et al.* Nucl. Phys. A 882 71 (2012)

- ✓ 53 new mass measured at GSI are reproduced well by PC-PK1 (only 11 parameters) with a rms deviation of 0.859 MeV.

Predictive power

Z = 102-116

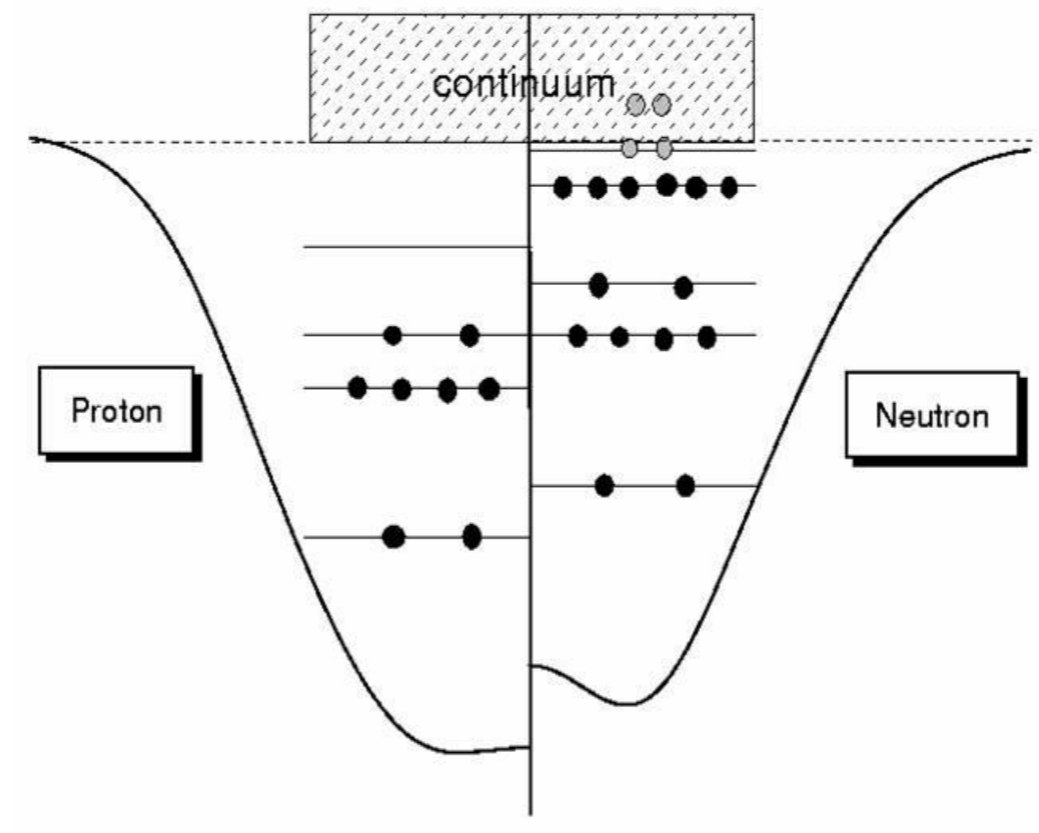


Kaiyuan Zhang, et al

Phys. Rev. C104 (2021) L021301

Predictive power for superheavy nuclear mass and possible stability beyond the neutron drip line in deformed relativistic Hartree-Bogoliubov theory in continuum

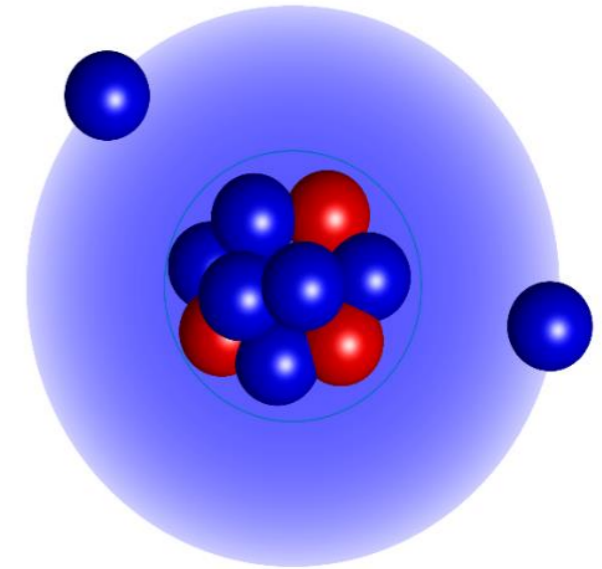
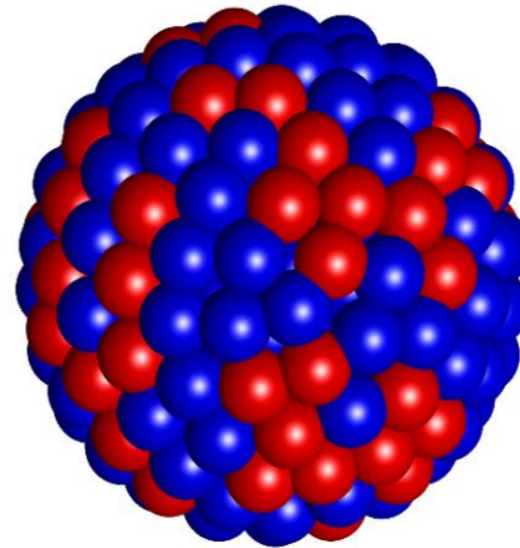
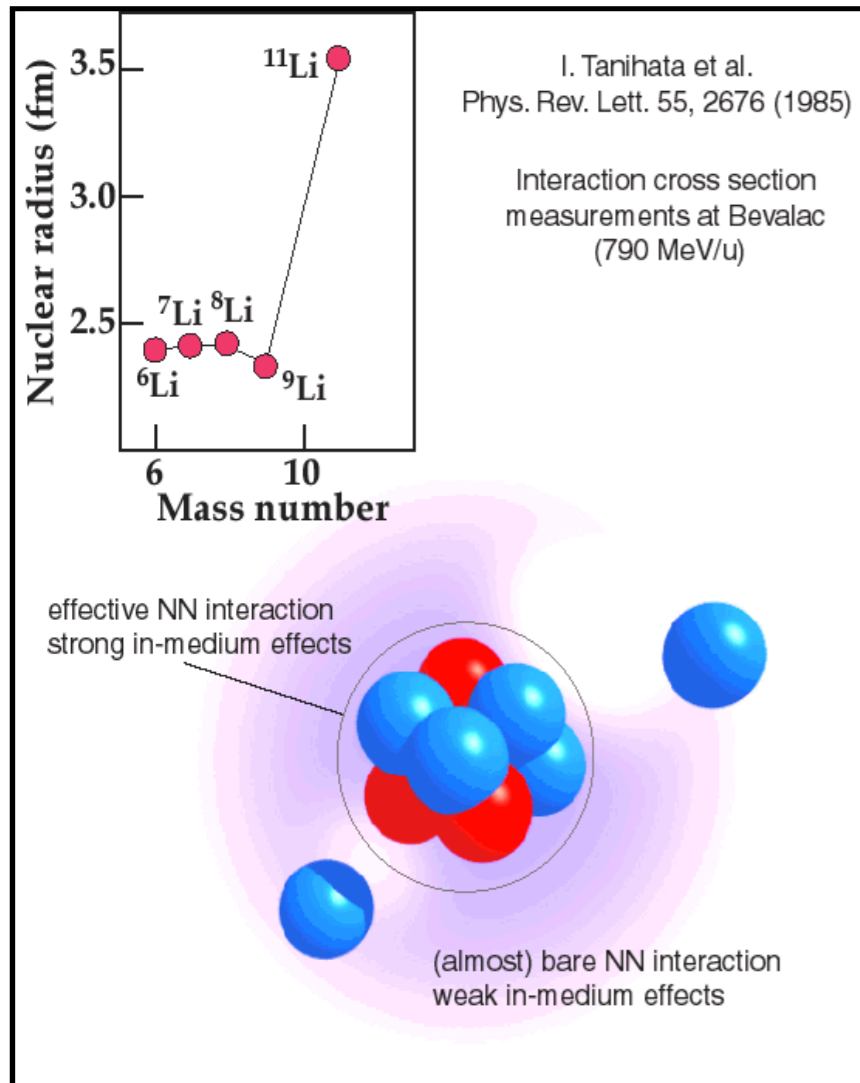
Physics close to the drip-line



Continuum, deformation, clustering, ...

New phenomena in exotic nucleus

I. Tanihata, et al Phys. Rev. Lett. 55 (1985) 2676



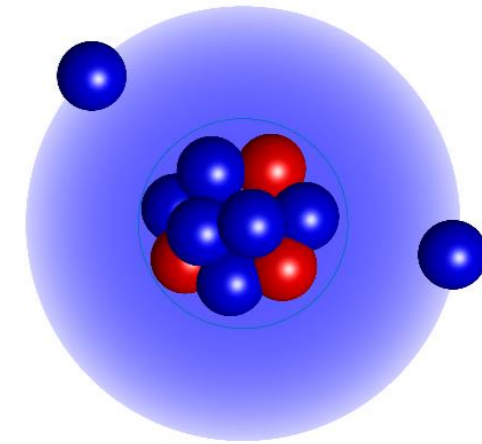
From Xin-Hui Wu

Meng and Ring, Phys. Rev. Lett. 77 (1996) 3963
Meng and Ring, Phys. Rev. Lett. 80 (1998) 460

Shell structure, low density, continuum, bound state, spatial distribution, pairing correlation, coupling between bound state and continuum...

- Meng, Toki, Zhou, Zhang, Long & Geng, Prog. Part. Nucl. Phys. 57 (2006) 470
- Meng and Zhou, J. Phys. G: 42 (2015) 093101

Halo nuclei

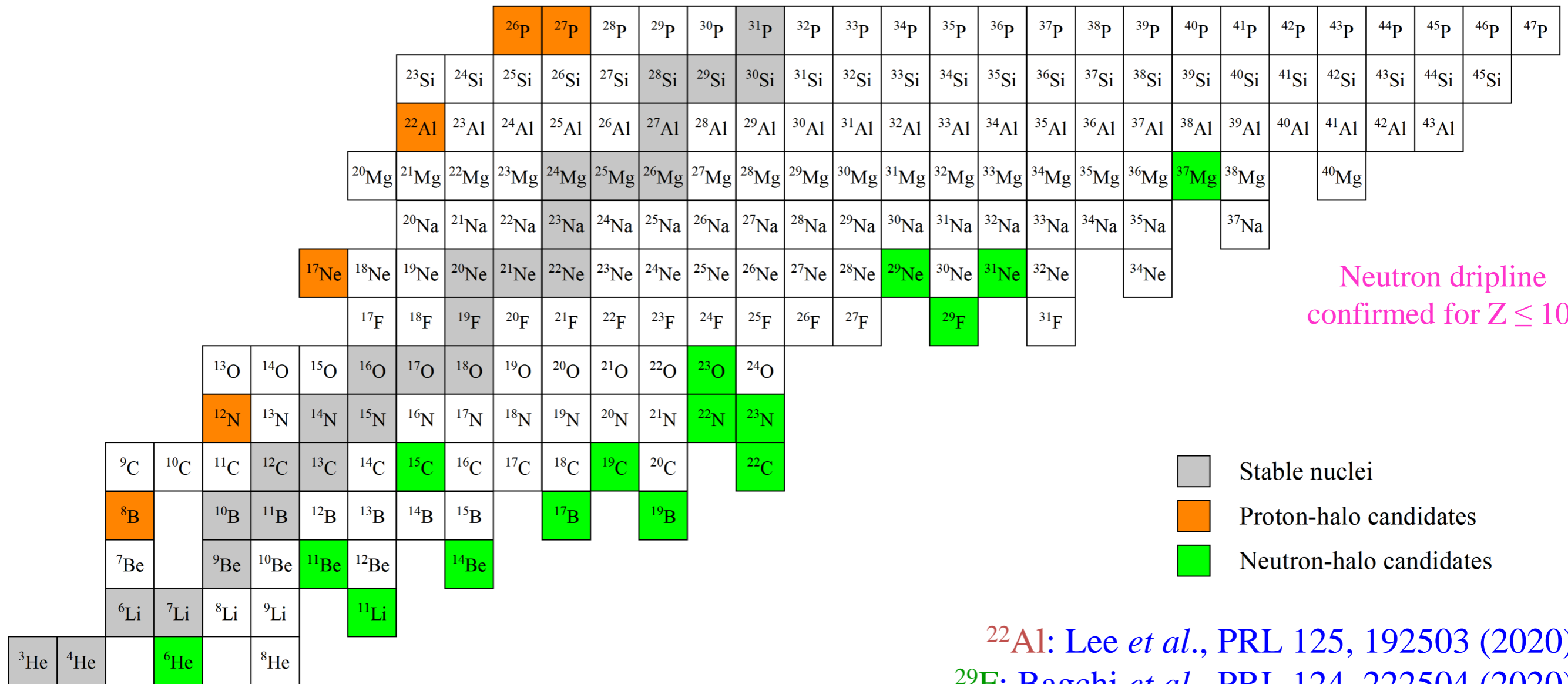


¹¹Li

➤ Halo nuclei have attracted lots of attention since the discovery of the halo phenomenon in ¹¹Li.

Tanihata *et al.*, PRL 55, 2676 (1985)

Tanihata *et al.*, PPNP 68, 215 (2013)



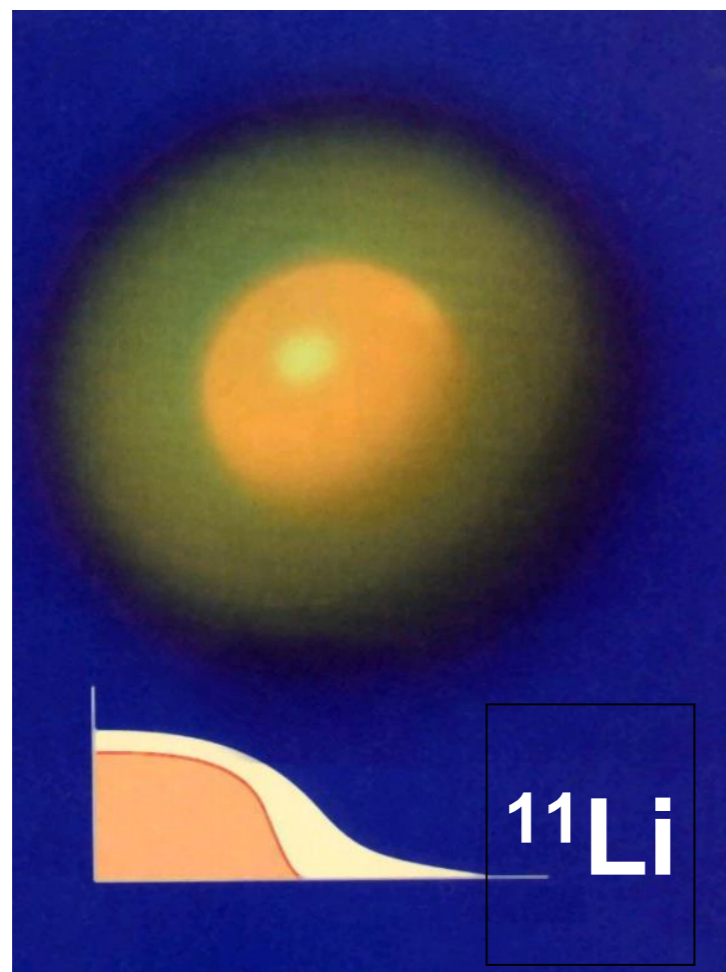
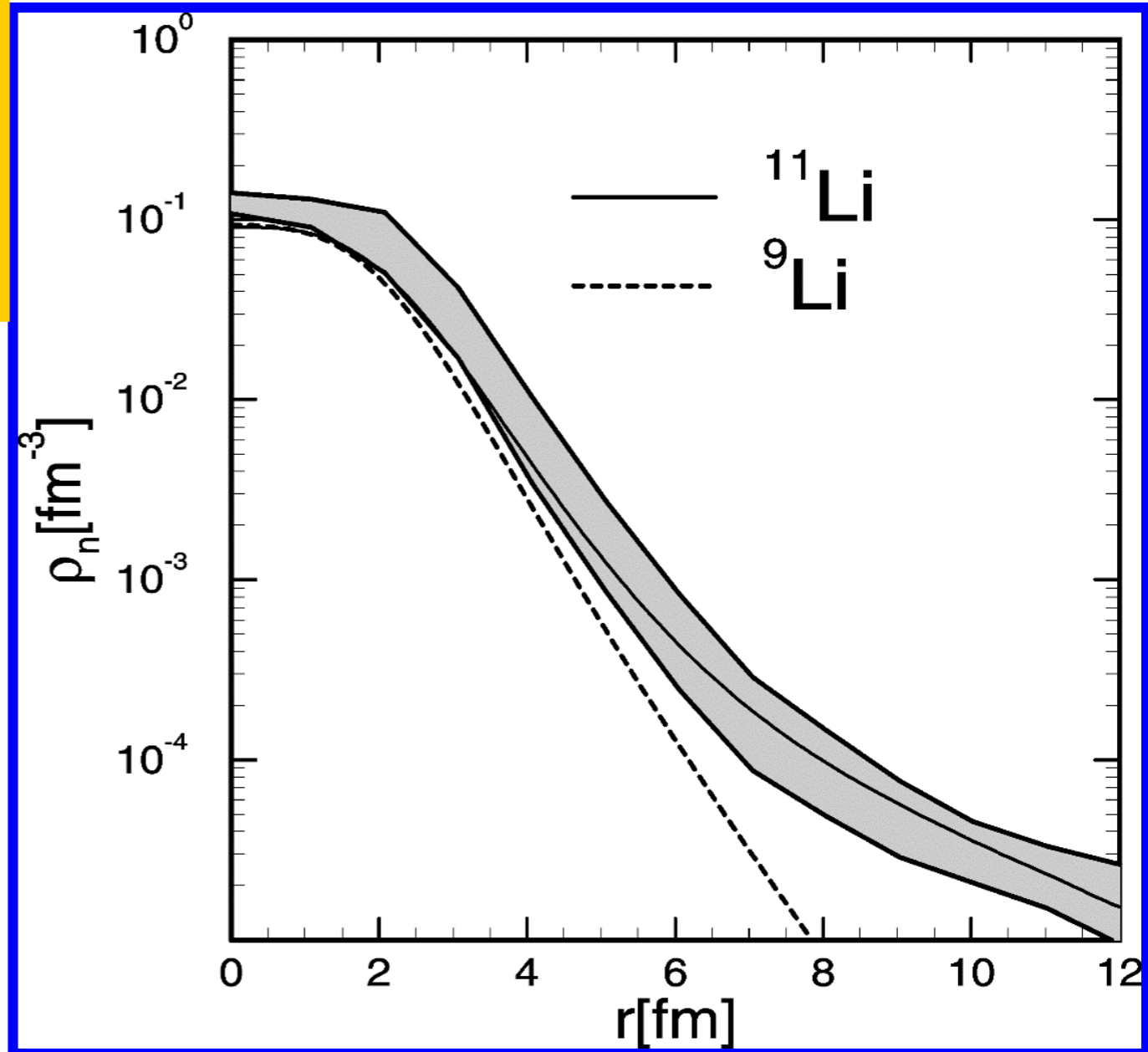
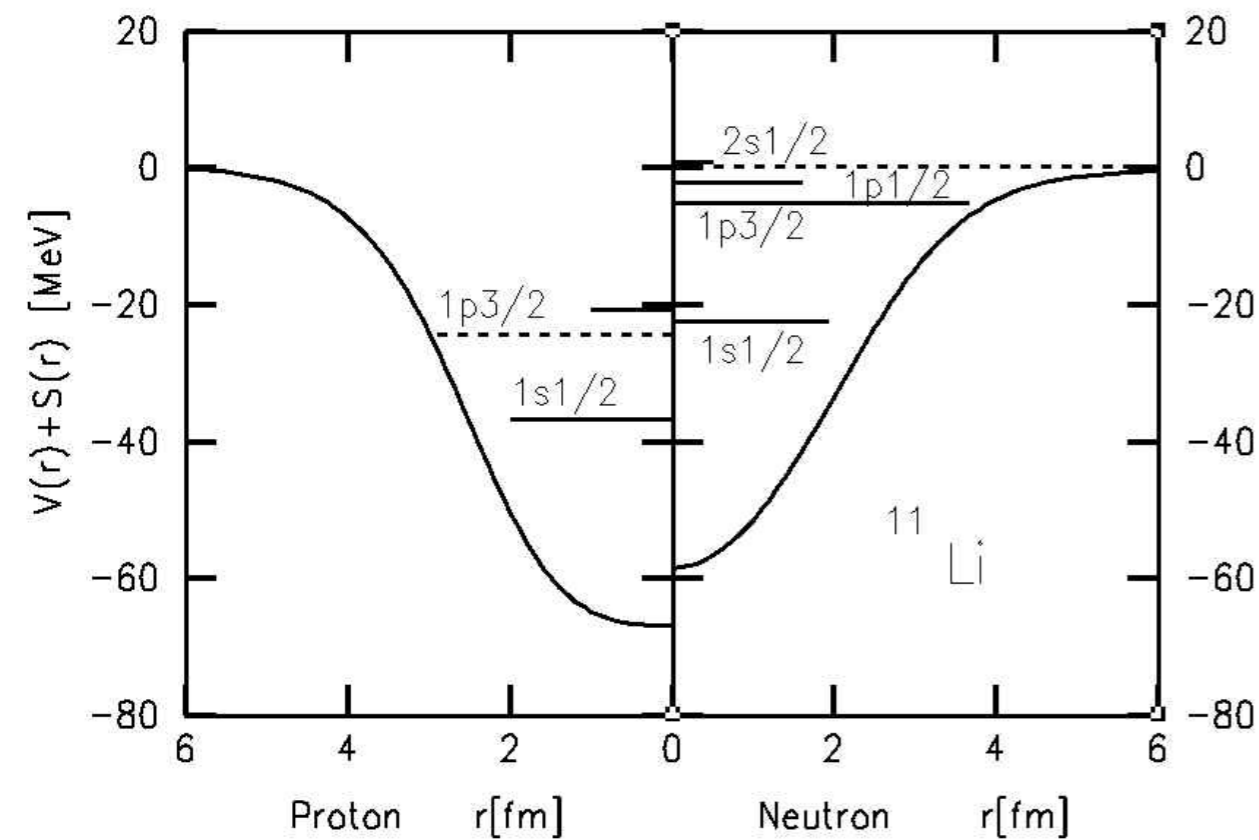
²²Al: Lee *et al.*, PRL 125, 192503 (2020)

²⁹F: Bagchi *et al.*, PRL 124, 222504 (2020)

³⁷Mg: Kobayashi *et al.*, PRL 112, 242501 (2014)

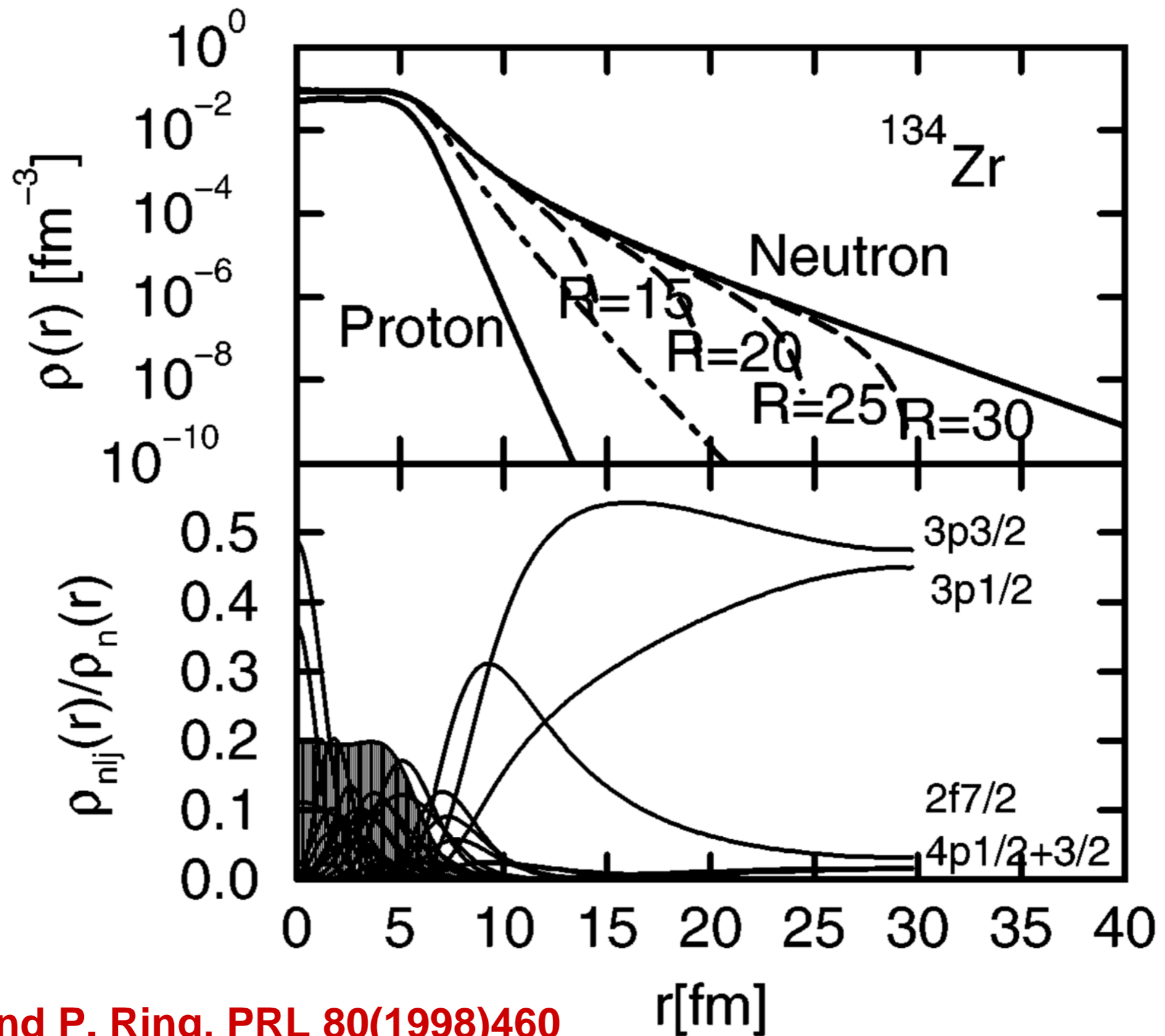
Density in ^{11}Li

RCHB: Relativist. Continuum Hartree-Bogoliubov theory with density dependent pairing force



J. Meng and P. Ring, PRL 77, 3963 (1996),

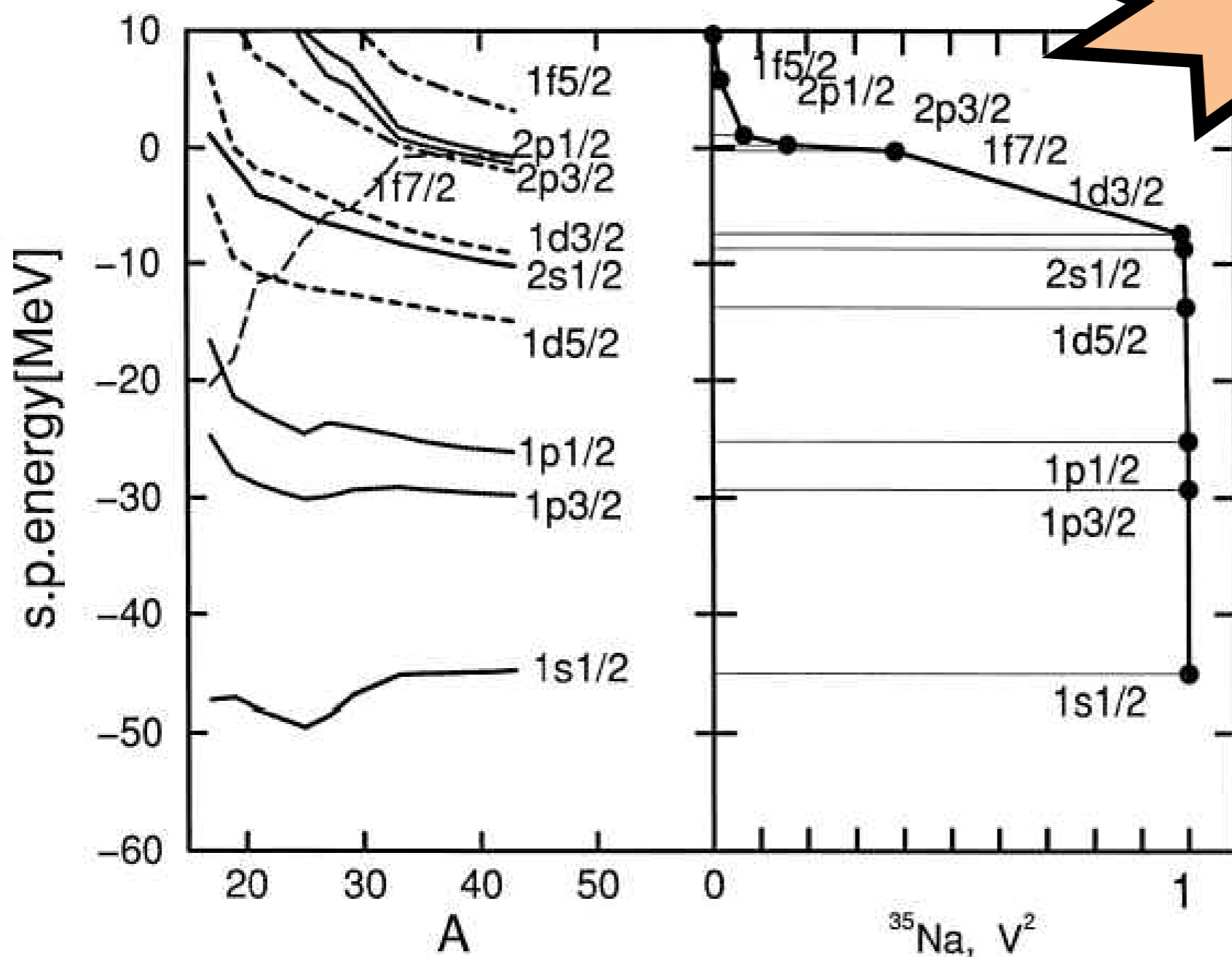
Giant halos



Halo in ^{35}Na

Meng, Tanihata & Yamaji, PLB419, 1(1998)

Prediction of giant halos
in Na isotopes



Deformed halos

✓ There has been controversy over the existence of deformed halo nuclei.

Otsuka, Muta, Yokoyama, Fukunishi, and Suzuki, NPA 588, 113c (1995)

Misu, Nazarewicz, and Aberg, NPA 614, 44 (1997)

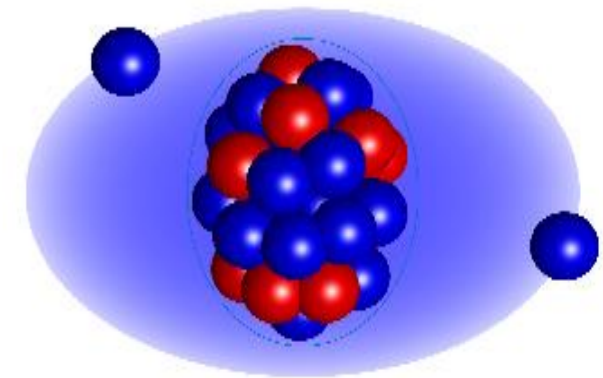
Tanihata, Hirata, and Toki, NPA 583, 769 (1995)

Nunes, NPA 757, 349 (2005)

✓ Considering deformation, pairing, and continuum effects, **the deformed relativistic Hartree-Bogoliubov theory in continuum (DRHBc)** predicts deformed halo nuclei.

Zhou, Meng, Ring, and Zhao, PRC 82, 011301(R) (2010)

Li, Meng, Ring, Zhao, and Zhou, PRC 85, 024312 (2012)



✓ Candidates of deformed halo nuclei have been suggested in experiment, ^{31}Ne and ^{37}Mg .

Nakamura *et al.*, PRL 112, 142501 (2014)

Kobayashi *et al.*, PRL 112, 242501 (2014)

✓ DRHBc theory has been applied for halo and other exotic phenomena.

Sun, Zhao, and Zhou, PLB 785, 530 (2018)

Zhang, Wang, and Zhang PRC 100, 034312 (2019)

Sun, Zhao, and Zhou, NPA 1003, 122011 (2020)

Yang *et al.*, PRL 126, 082501 (2021)

Sun, PRC 103, 054315 (2021)

Zhang *et al.*, PRC 104, L021301 (2021)

Pan *et al.*, PRC 104, 024331 (2021)

He *et al.*, CPC 45, 101001 (2021)

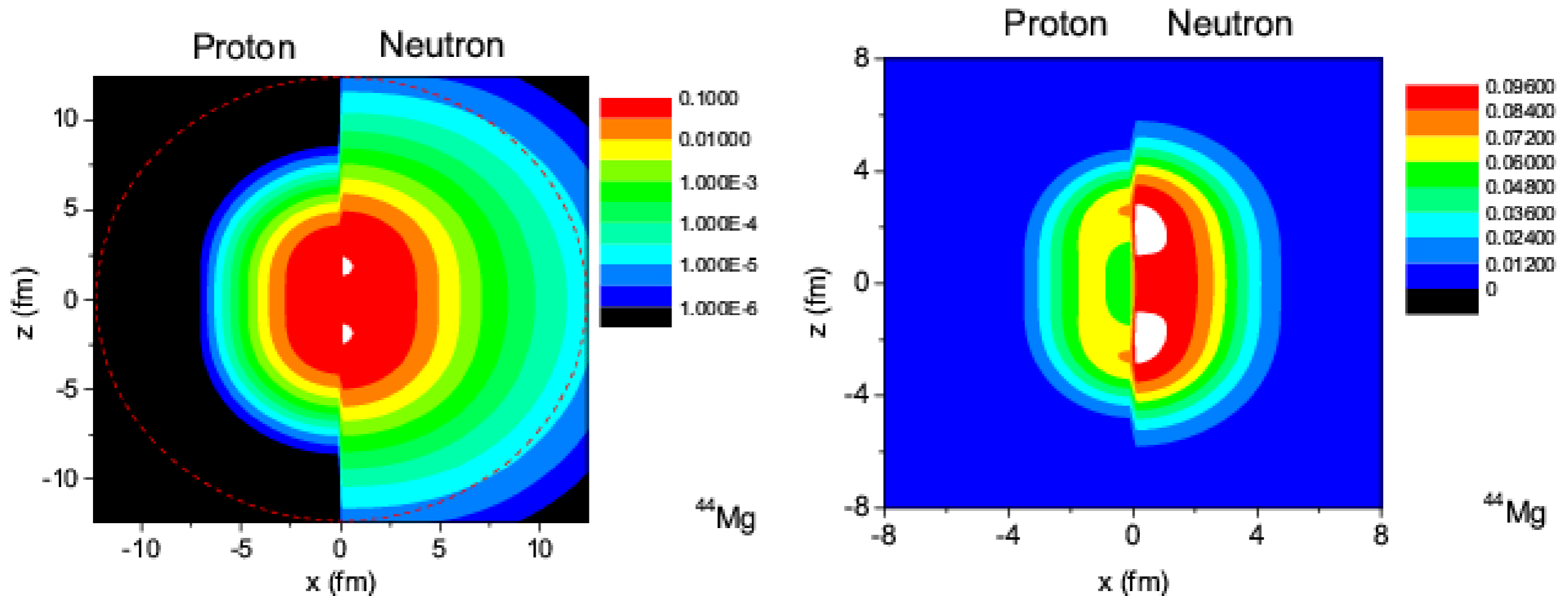
^{44}Mg : Density distributions

Zhou_Meng_Ring_Zhao2010_PRC82-011301R

Zhou_Meng_Ring_Zhao2011_JPConfProc312-092067

Li_Meng_Ring_Zhao_Zhou2012_PRC85-024312

- Prolate deformation
- Large spatial extension in neutron density distribution

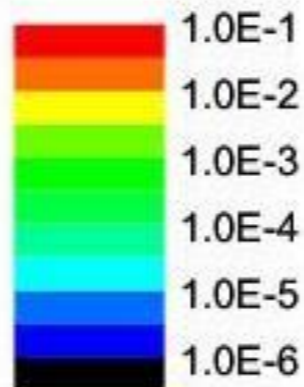
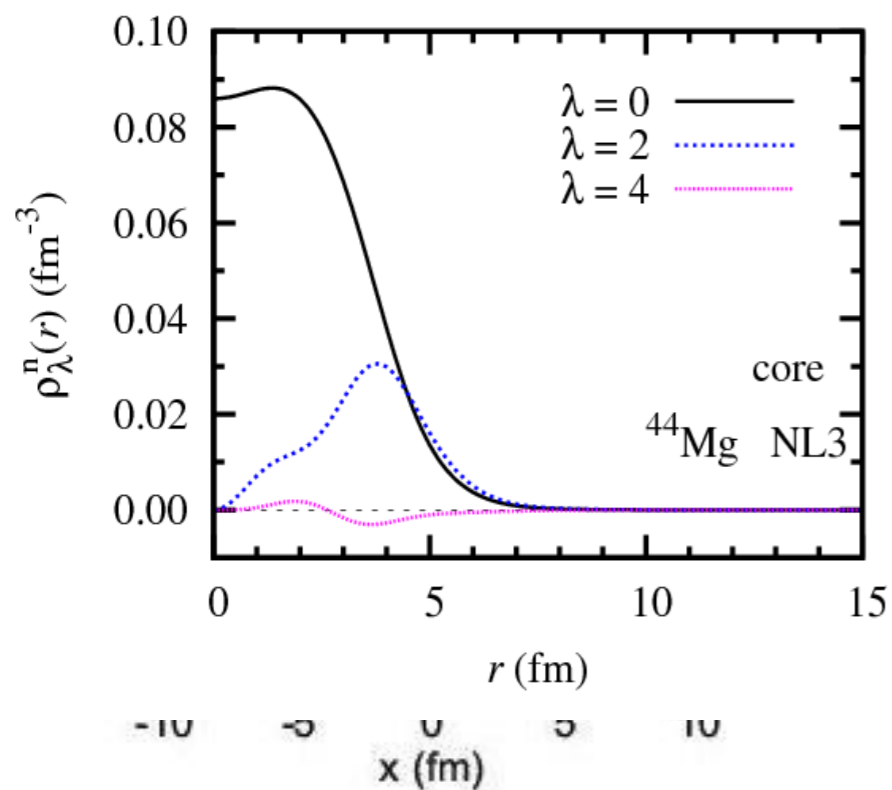


Viewpoint: A Walk Along the Dripline by Paul Cottle and Kirby Kemper

<http://link.aps.org/doi/10.1103/Physics.5.49>

Prolate core & oblate halo

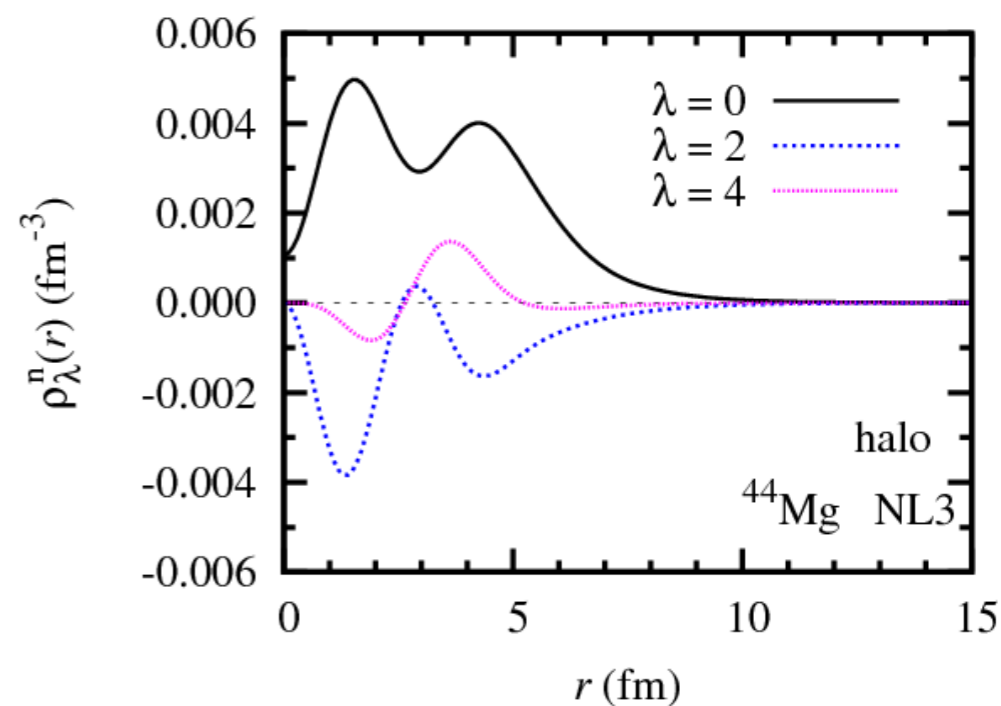
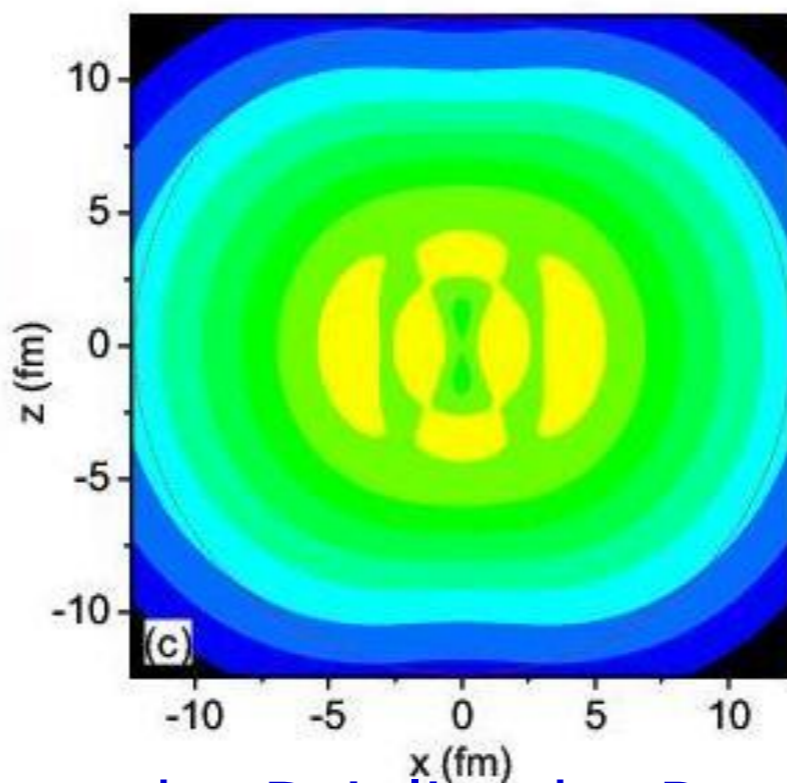
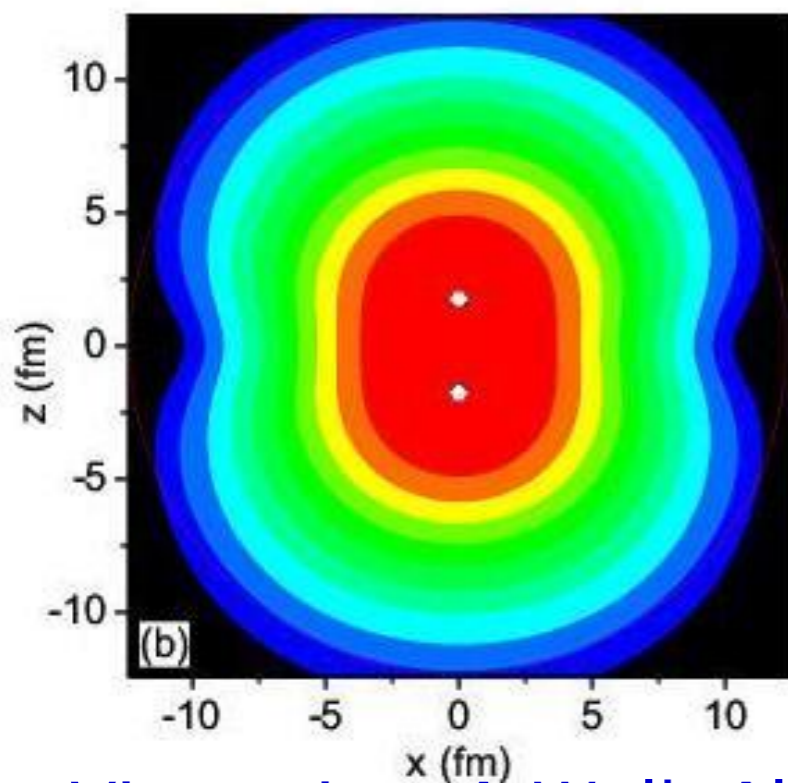
Zhou, Meng, Ring & Zhao, Phys. Rev. C 82, 011301 (2010)



^{44}Mg

❖ Prolate core, but slightly oblate halo with sizable hexadecapole component !

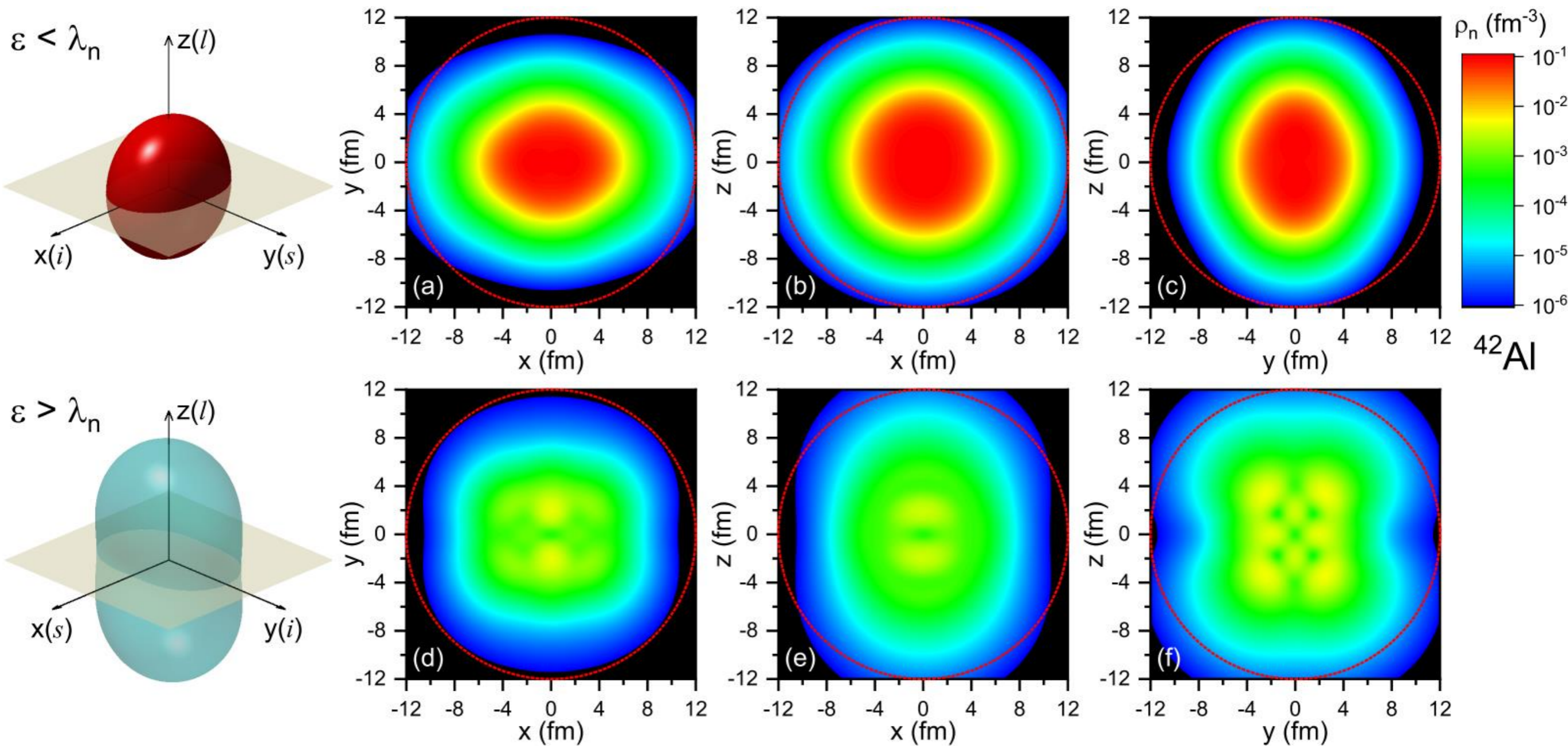
❖ Decoupling of deformation between core & halo



Viewpoint: A Walk Along the Dripline by Paul Cottle and Kirby Kemper

<http://link.aps.org/doi/10.1103/Physics.5.49>

Halo in triaxial nucleus ^{42}Al



✓ Core: $r = 3.85$ fm, $\beta = 0.38$, $\gamma = 50^\circ$, $z > x > y$

✓ Halo: $r = 5.26$ fm, $\beta = 0.79$, $\gamma = -23^\circ$, $z > y > x$

DRHBc mass table collaboration

Since December 5, 2018, PC-PK1 + DRHBc



RCHB mass table

First nuclear mass table including continuum effects

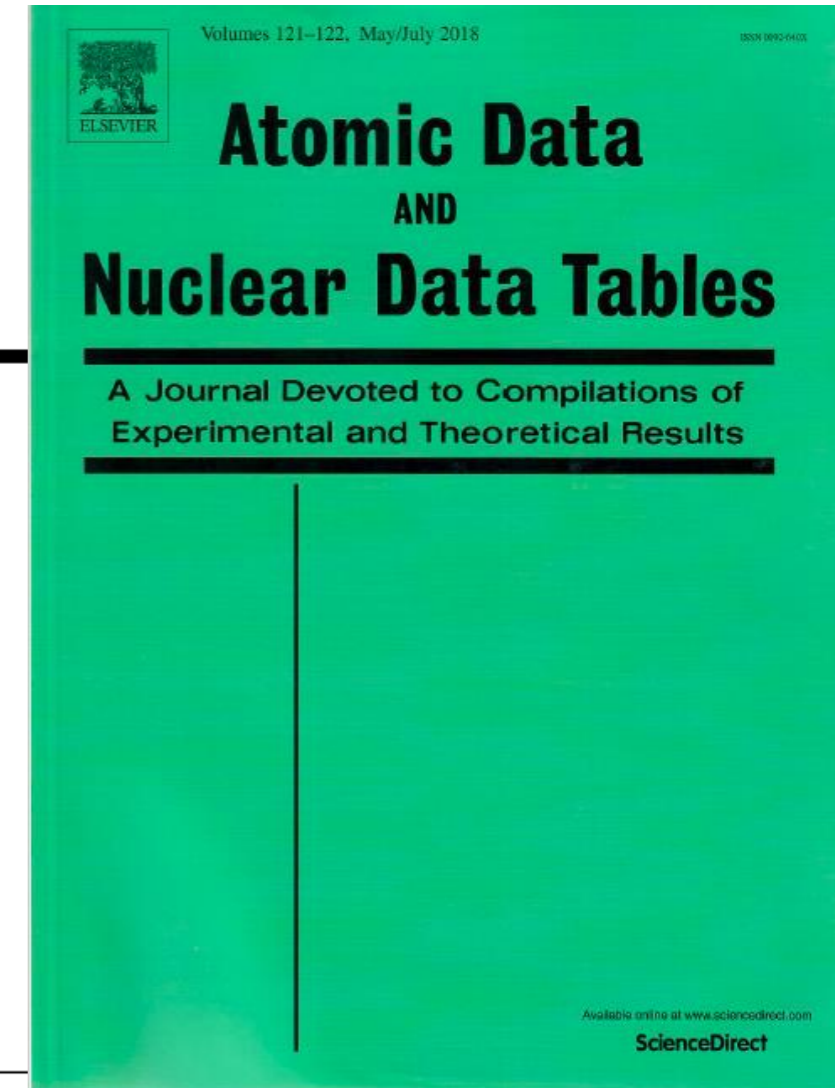
Atomic Data and Nuclear Data Tables 121–122 (2018) 1–215



Contents lists available at ScienceDirect

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journal homepage: www.elsevier.com/locate/adt



The limits of the nuclear landscape explored by the relativistic continuum Hartree–Bogoliubov theory

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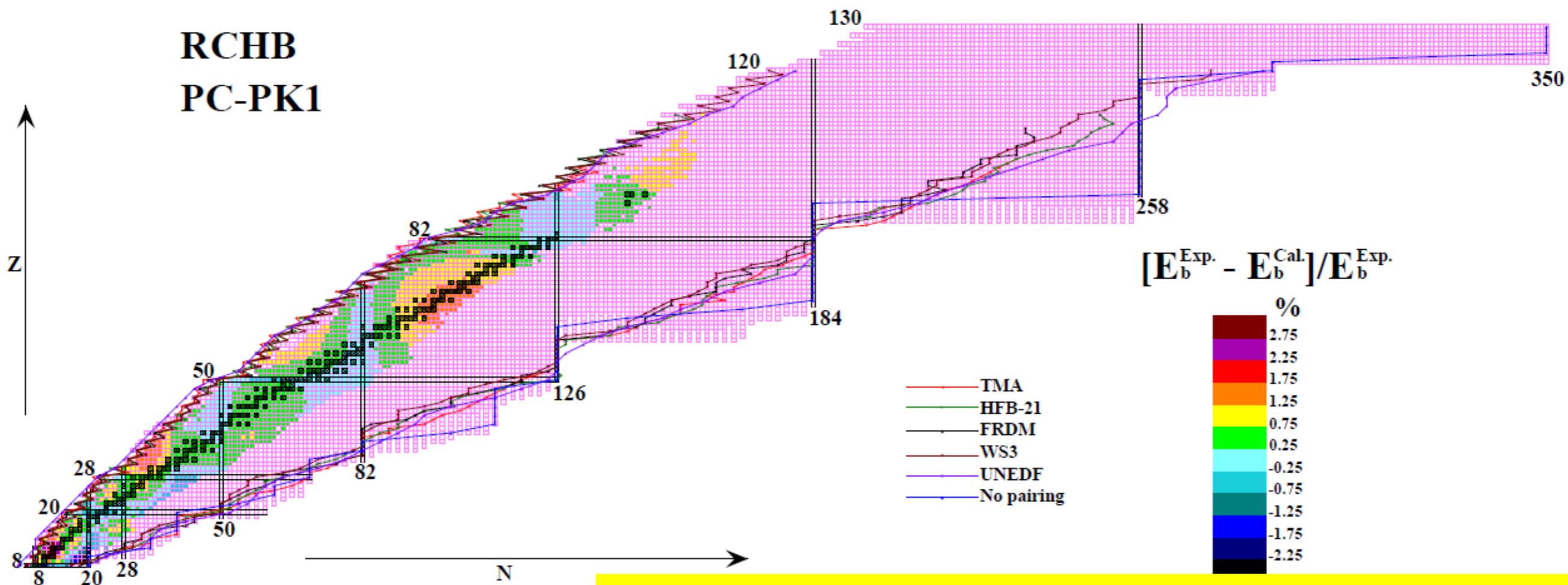
Available online 1 November 2017

ABSTRACT

The ground-state properties of nuclei with $8 \leq Z \leq 120$ from the proton drip line to the neutron drip line have been investigated using the spherical relativistic continuum Hartree–Bogoliubov (RCHB) theory with the relativistic density functional PC-PK1. With the effects of the continuum included, there are totally 9035 nuclei predicted to be bound, which largely extends the existing nuclear landscapes predicted with other methods. The calculated binding energies, separation energies, neutron and proton Fermi surfaces,

Drip-lines in variant models

The number of bound nuclides with between 2 and 120 protons is around 7,000 28 JUNE 2012 | VOL 486 | NATURE | 509



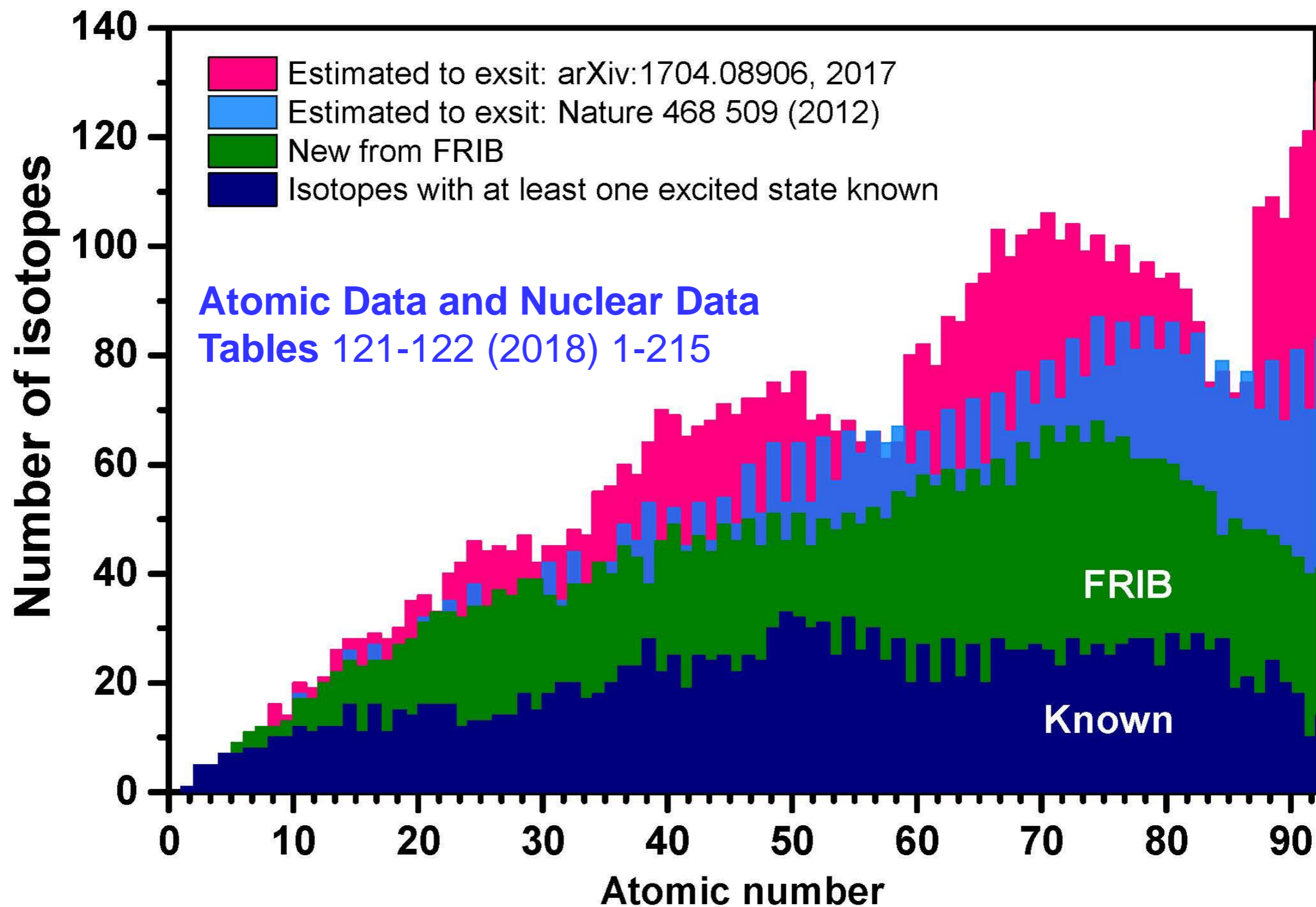
$8 \leq Z \leq 120$: 9035 nuclei predicted to be bound
[Atomic Data and Nuclear Data Tables 121-122 \(2018\) 1-215](#)

10532 bound nuclei from $Z=8$ to $Z=130$ predicted by RCHB theory with PC-PK1. For 2227 nuclei with data, binding energy differences between data and calculated results are shown in different color. The nucleon drip-lines predicted TMA, HFB-21, WS3, FRDM, UNEDF and without pairing correlation are plotted for comparison.

See also: Afanasjev, Agbemava, Ray, Ring, PLB726(2013)680

Possible existing isotopes

Atomic Data and Nuclear Data Tables 121-122(2018)1-215



DRHBc mass table collaboration

PC-PK1 + DRHBc

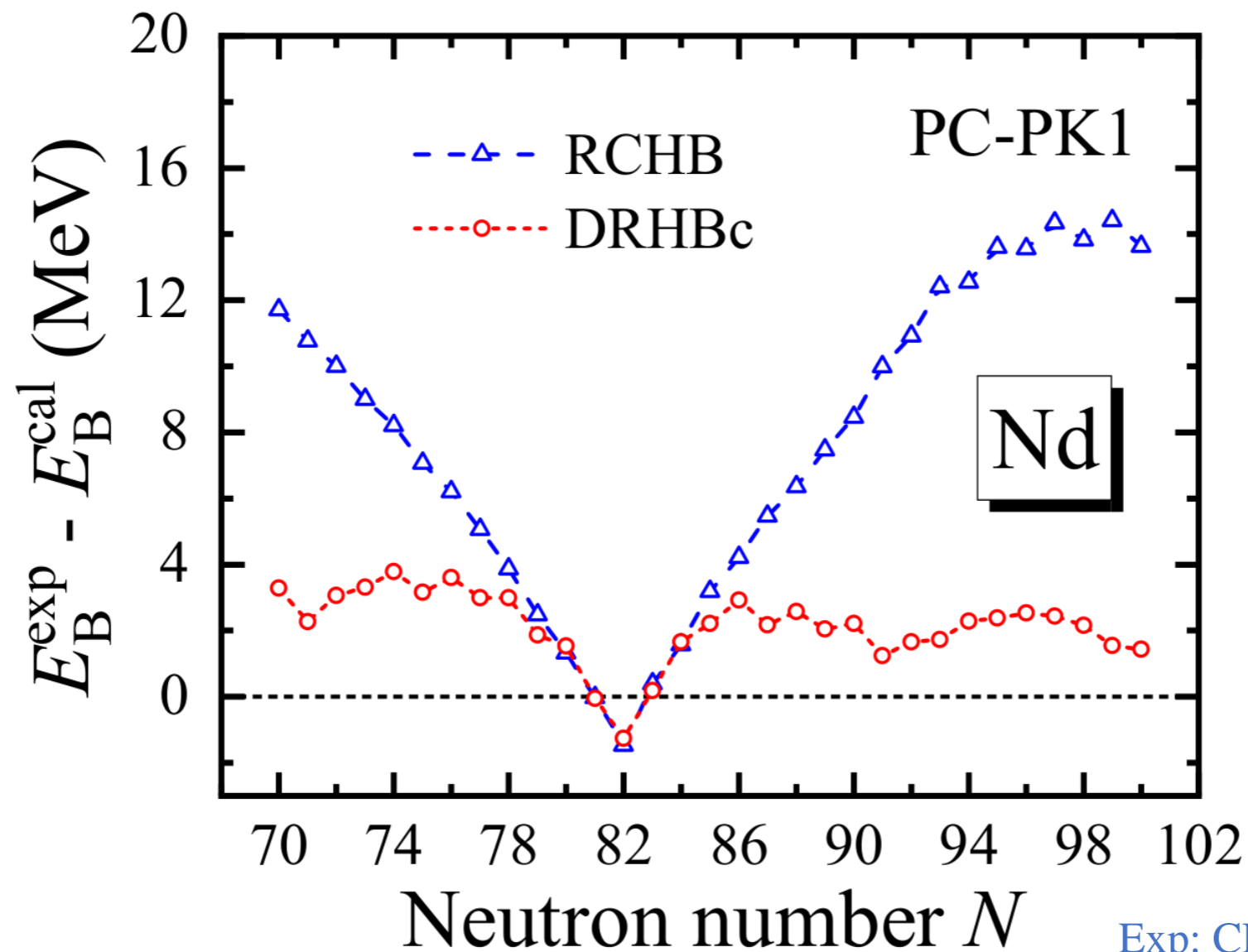
I. Even-even nuclei

II. Even Z-Odd N nuclei

III. Odd-Z nuclei



Deformation improve the accuracy



Exp: CPC 45, 030003 (2021)
Even-even: PRC 102, 024314 (2020)

- ✓ With deformation included, the data can be better reproduced
DRHBc 2.38 MeV **RCHB 9.08 MeV.**
- ✓ The rotational correction energy is expected to further improve the results for odd nuclei.

Deformed relativistic Hartree-Bogoliubov theory in continuum with a point-coupling functional: Examples of even-even Nd isotopes

Kaiyuan Zhang (张开元) *et al.* (DRHBc Mass Table Collaboration)

Phys. Rev. C **102**, 024314 – Published 12 August 2020

PHYSICAL REVIEW C **102**, 024314 (2020)

Article

Deformed relativistic Hartree-Bogoliubov theory in continuum with a point-coupling functional: Examples of even-even Nd isotopes

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(DRHBc Mass Table Collaboration)

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⁶School of Physics, University of Chinese Academy of Sciences, Beijing 100049, China



DRHBc Mass Table: even-even Nuclei

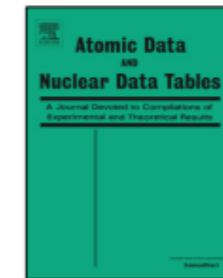
Atomic Data and Nuclear Data Tables 144 (2022) 101488



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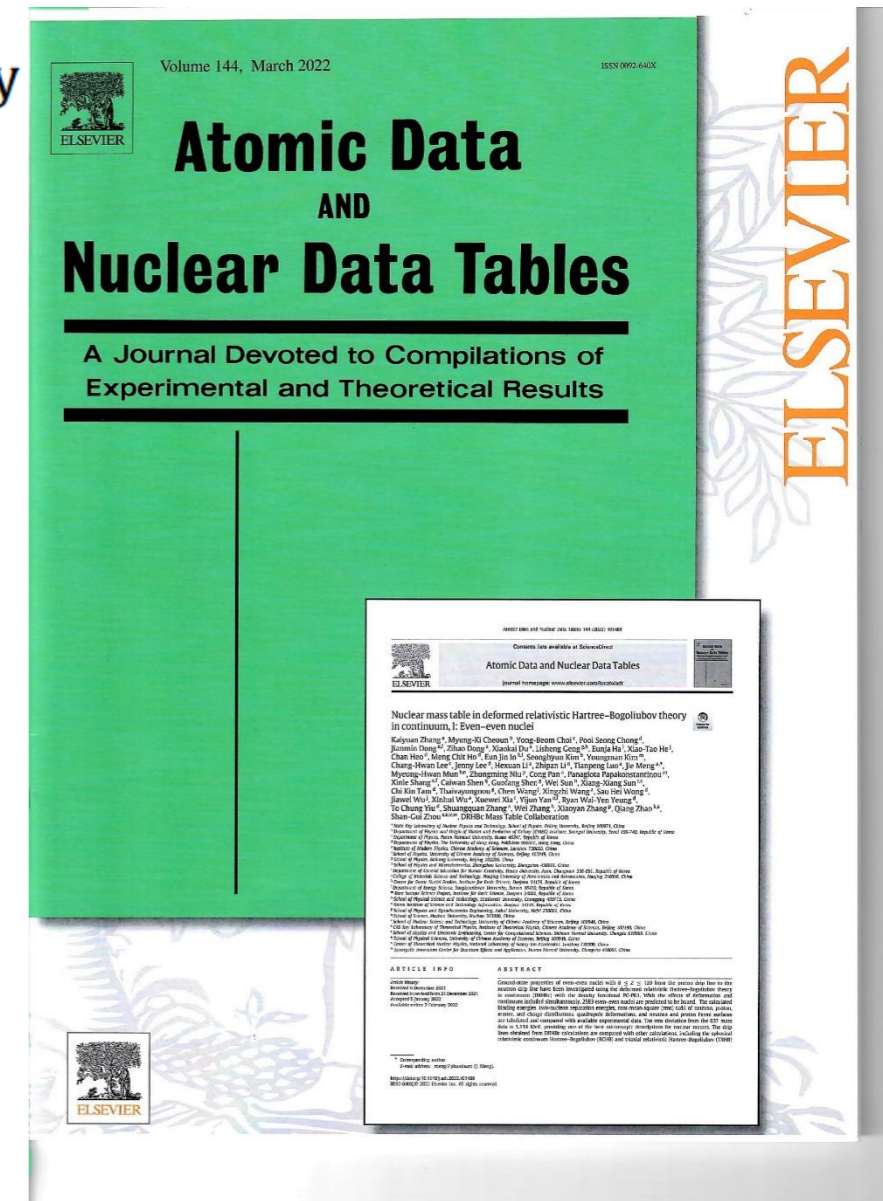
Atomic Data and Nuclear Data Tables

journal homepage: www.elsevier.com/locate/adt



Nuclear mass table in deformed relativistic Hartree–Bogoliubov theory in continuum, I: Even–even nuclei

Kaiyuan Zhang^a, Myung-Ki Cheoun^b, Yong-Beom Choi^c, Pooi Seong Chong^d, Jianmin Dong^{e,f}, Zihao Dong^a, Xiaokai Du^a, Lisheng Geng^{g,h}, Eunja Haⁱ, Xiao-Tao He^j, Chan Heo^d, Meng Chit Ho^d, Eun Jin In^{k,l}, Seonghyun Kim^b, Youngman Kim^m, Chang-Hwan Lee^c, Jenny Lee^d, Hexuan Li^a, Zhipan Liⁿ, Tianpeng Luo^a, Jie Meng^{a,*}, Myeong-Hwan Mun^{b,o}, Zhongming Niu^p, Cong Pan^a, Panagiota Papakonstantinou^m, Xinle Shang^{e,f}, Caiwan Shen^q, Guofang Shen^g, Wei Sunⁿ, Xiang-Xiang Sun^{r,s}, Chi Kin Tam^d, Thaivayongnou^g, Chen Wang^j, Xingzhi Wang^a, Sau Hei Wong^d, Jiawei Wu^j, Xinhui Wu^a, Xuewei Xia^t, Yijun Yan^{e,f}, Ryan Wai-Yen Yeung^d, To Chung Yiu^d, Shuangquan Zhang^a, Wei Zhang^h, Xiaoyan Zhang^p, Qiang Zhao^{k,a}, Shan-Gui Zhou^{s,u,v,w}, DRHBc Mass Table Collaboration



DRHBc for Odd-A and O-O Nuclei

PHYSICAL REVIEW C

covering nuclear physics

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Deformed relativistic with a point-coupling

Cong Pan (潘琮) *et al.* (DRHBc Ma
Phys. Rev. C **106**, 014316 – Publist

Article

PHYSICAL REVIEW C **106**, 014316 (2022)

Deformed relativistic Hartree-Bogoliubov theory in continuum with a point-coupling functional. II. Examples of odd Nd isotopes

Cong Pan (潘琮),¹ Myung-Ki Cheoun,² Yong-Beom Choi,³ Jianmin Dong (董建敏),^{4,5} Xiaokai Du (杜晓凯),¹ Xiao-Hua Fan (范小华),⁶ Wei Gao (高威),⁷ Lisheng Geng (耿立升),^{8,7} Eunja Ha,⁹ Xiao-Tao He (贺晓涛),¹⁰ Jinke Huang (黄靳苛),⁷ Kun Huang (黄坤),¹⁰ Seonghyun Kim,² Youngman Kim,¹¹ Chang-Hwan Lee,³ Jenny Lee,¹² Zhipan Li (李志攀),⁶ Zhi-Rui Liu (刘治瑞),¹⁰ Yiming Ma (马艺铭),¹³ Jie Meng (孟杰),^{1,*} Myeong-Hwan Mun,^{2,14} Zhongming Niu (牛中明),¹⁵ Panagiota Papakonstantinou,¹¹ Xinle Shang (尚新乐),^{4,5} Caiwan Shen (沈彩万),¹⁶ Guofang Shen (申国防),⁸ Wei Sun (孙玮),⁶ Xiang-Xiang Sun (孙向向),^{17,18} Jiawei Wu (吴佳威),¹⁰ Xinhui Wu (吴鑫辉),¹ Xuwei Xia (夏学伟),¹⁹ Yijun Yan (晏一珺),^{4,5} To Chung Yiu,¹² Kaiyuan Zhang (张开元),^{1,20} Shuangquan Zhang (张双全),¹ Wei Zhang (张炜),⁷ Xiaoyan Zhang (张晓燕),¹⁵ Qiang Zhao (赵强),^{21,1} Ruyou Zheng (郑茹尤),⁸ and Shan-Gui Zhou (周善贵)^{18,22,23,24}
(DRHBc Mass Table Collaboration)

¹State Key Laboratory of Nuclear Physics and Technology, School of Physics, Peking University, Beijing 100871, China

²Department of Physics and Origin of Matter and Evolution of Galaxy Institute, Soongsil University, Seoul 156-743, Korea

³Department of Physics, Pusan National University, Busan 46241, Korea

⁴Institute of Modern Physics, Chinese Academy of Sciences, Lanzhou 730000, China

⁵School of Physics, University of Chinese Academy of Sciences, Beijing 100049, China

⁶School of Physical Science and Technology, Southwest University, Chongqing 400715, China

⁷School of Physics and Microelectronics, Zhengzhou University, Zhengzhou 450001, China

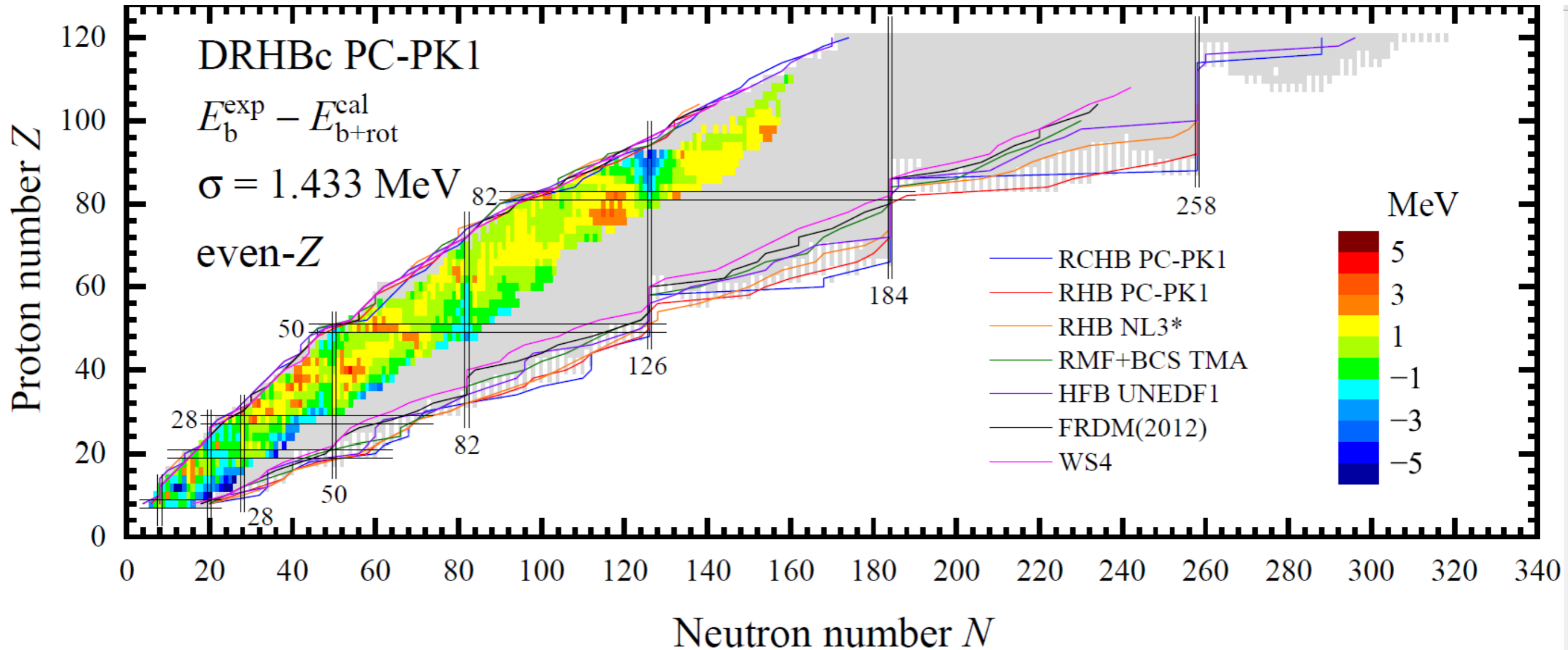
DRHBc Mass Table: even Z Nuclei



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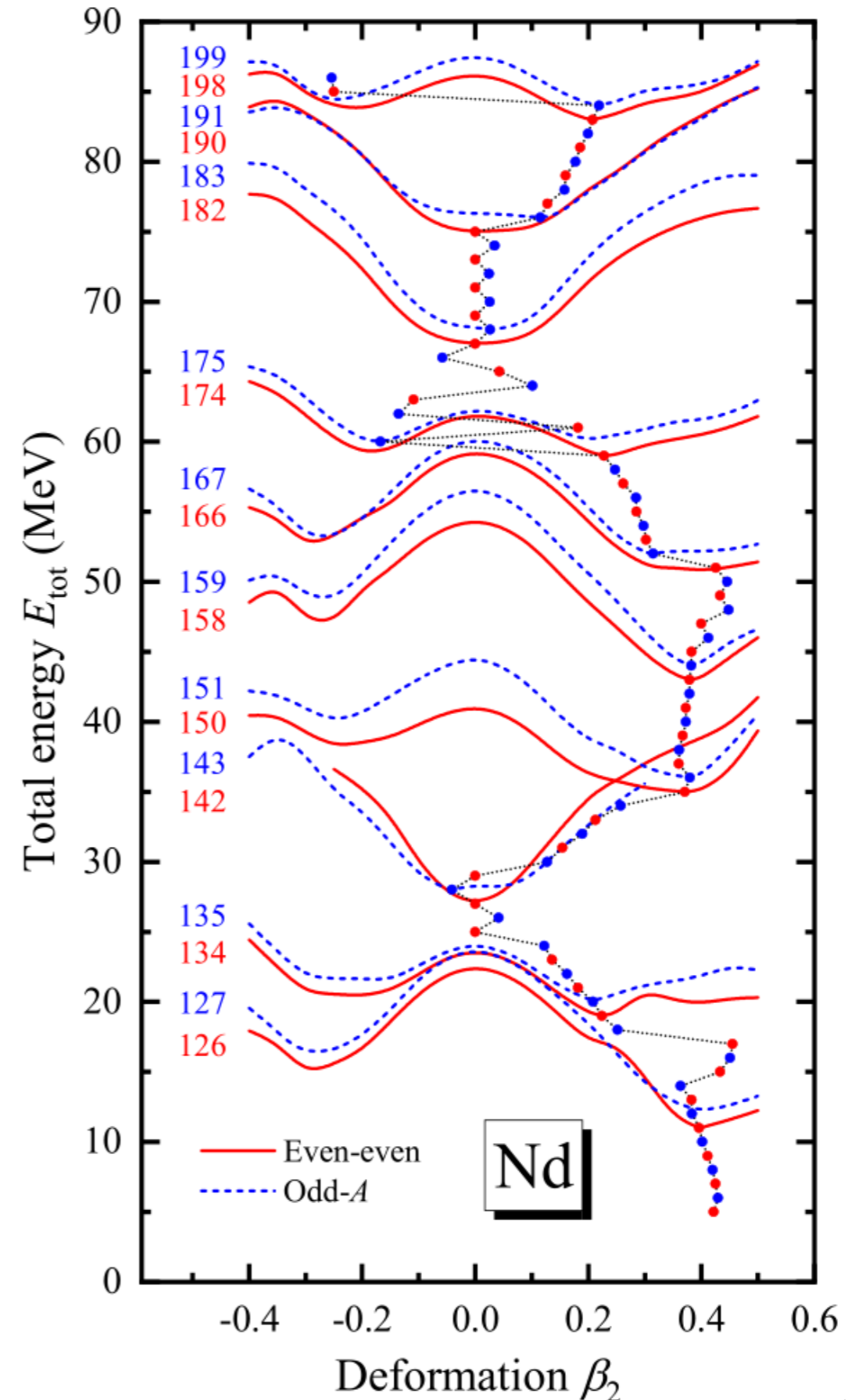


PC-PK1 + DRHBc : 4829 bound even-Z nuclei with $8 \leq Z \leq 120$

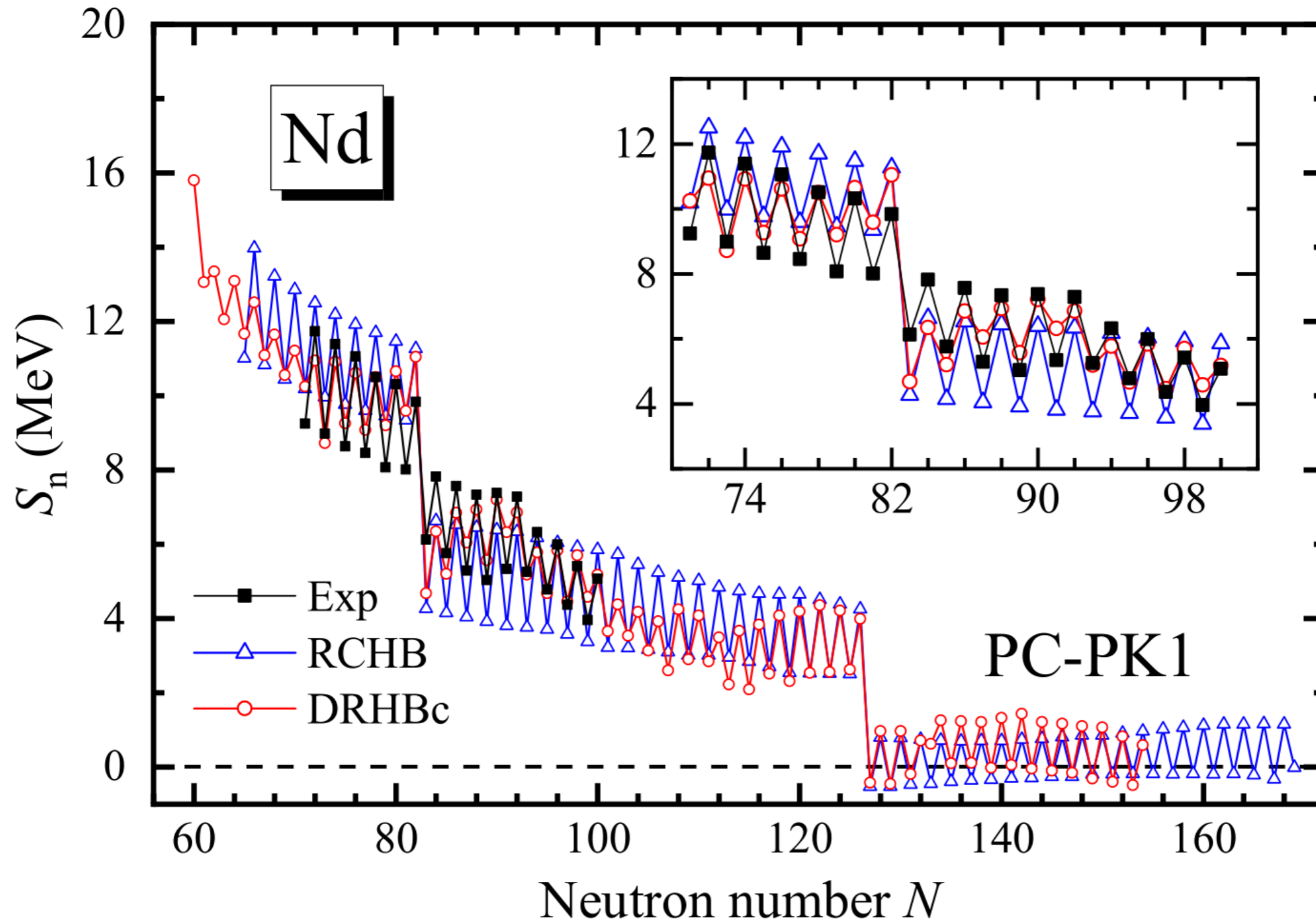
DRHBc Collaboration, [At. Data Nucl. Data Tables 144, 101488 \(2022\)](#)

Potential energy curve (PEC)

- ✓ Ground state for odd-A nuclei are double checked by different blocking.
- ✓ The PEC of odd-A nuclei are similar to their even-even neighbors.
- ✓ Sudden change of β_2 corresponds to possible shape coexistence.

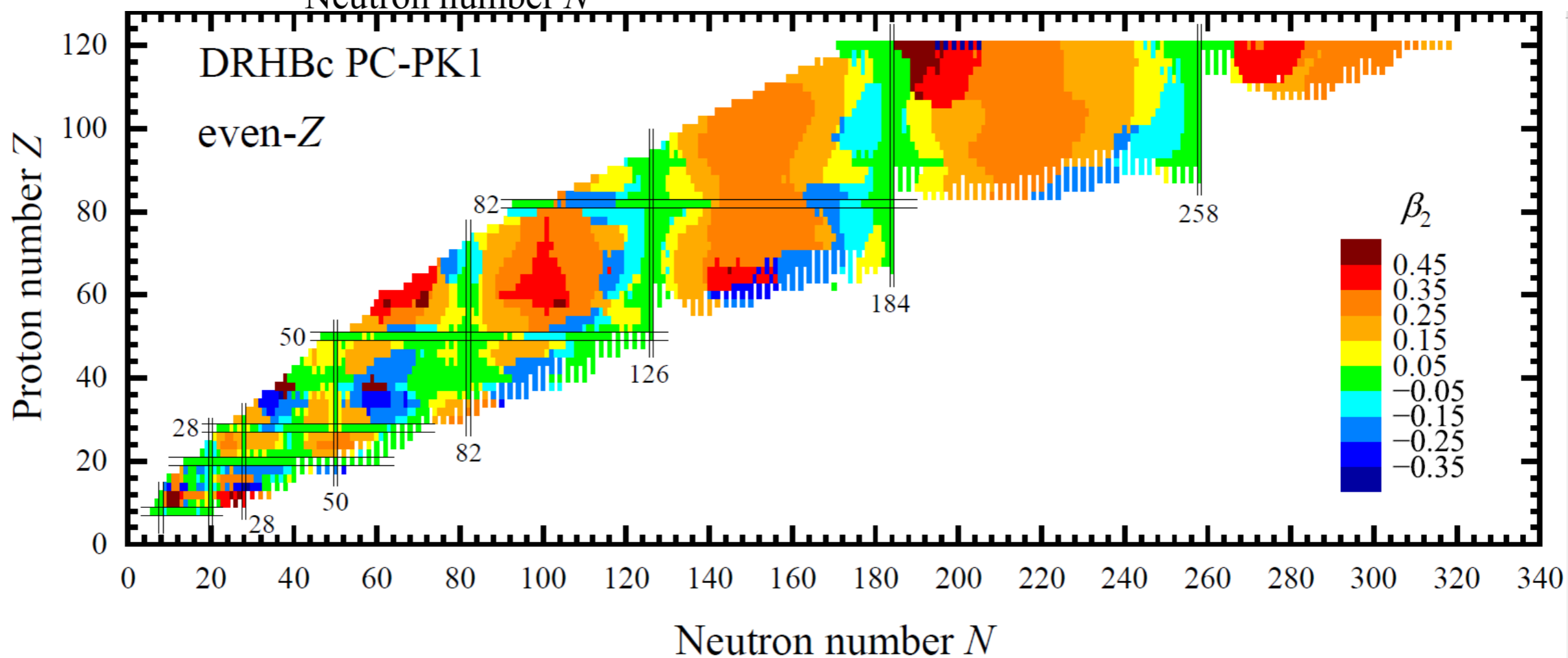
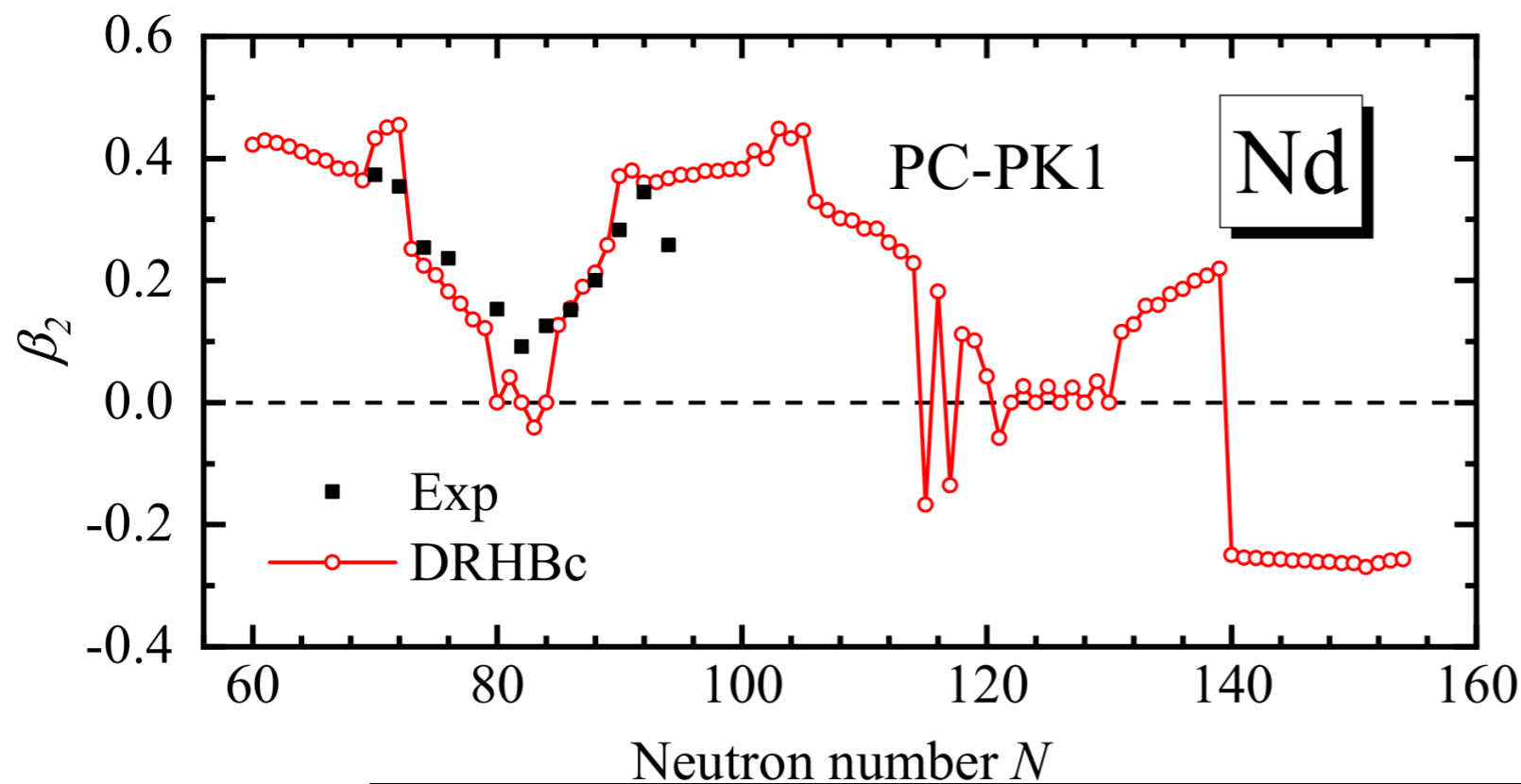


One-neutron separation energy

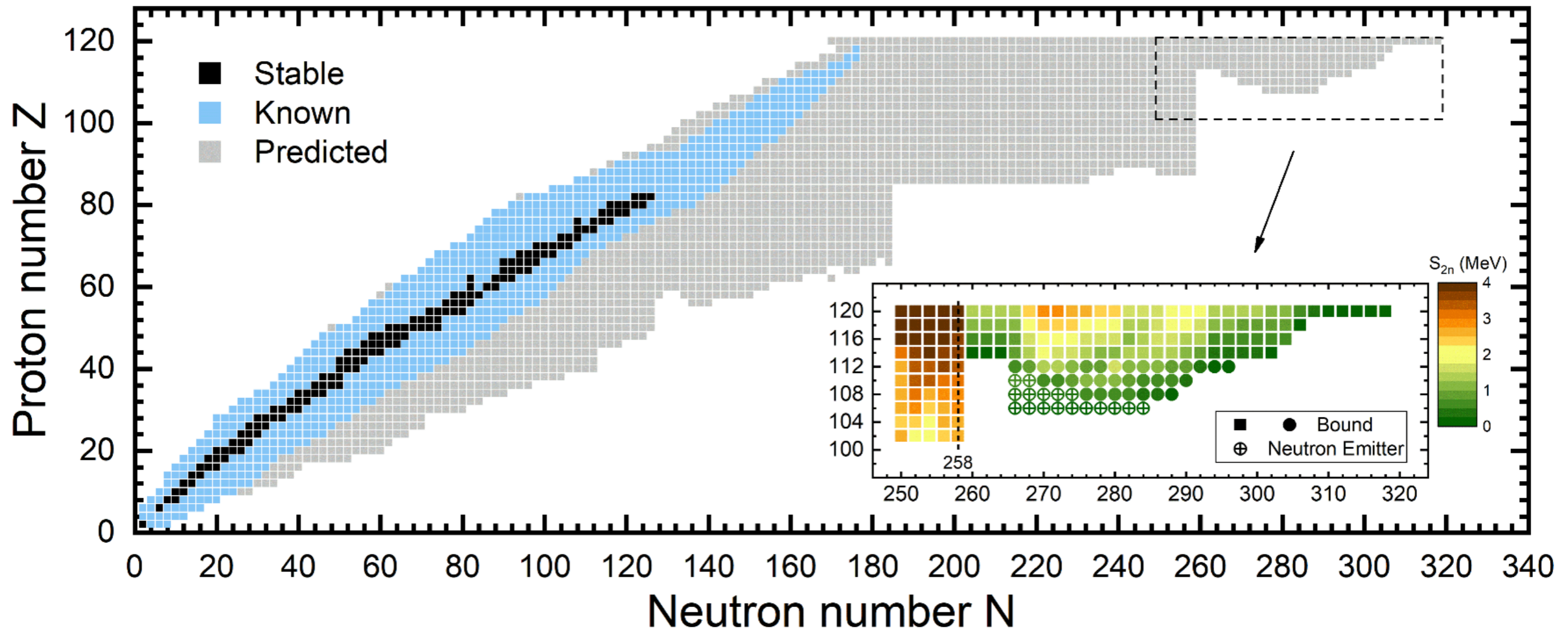


- ✓ Rms deviations: **DRHBc: 0.74 MeV** **RCHB: 1.10 MeV**
- ✓ One-neutron drip line: **DRHBc: $N = 126$** **RCHB: $N = 126$**

Quadrupole deformation



Possible stability beyond the neutron drip line in DRHBc

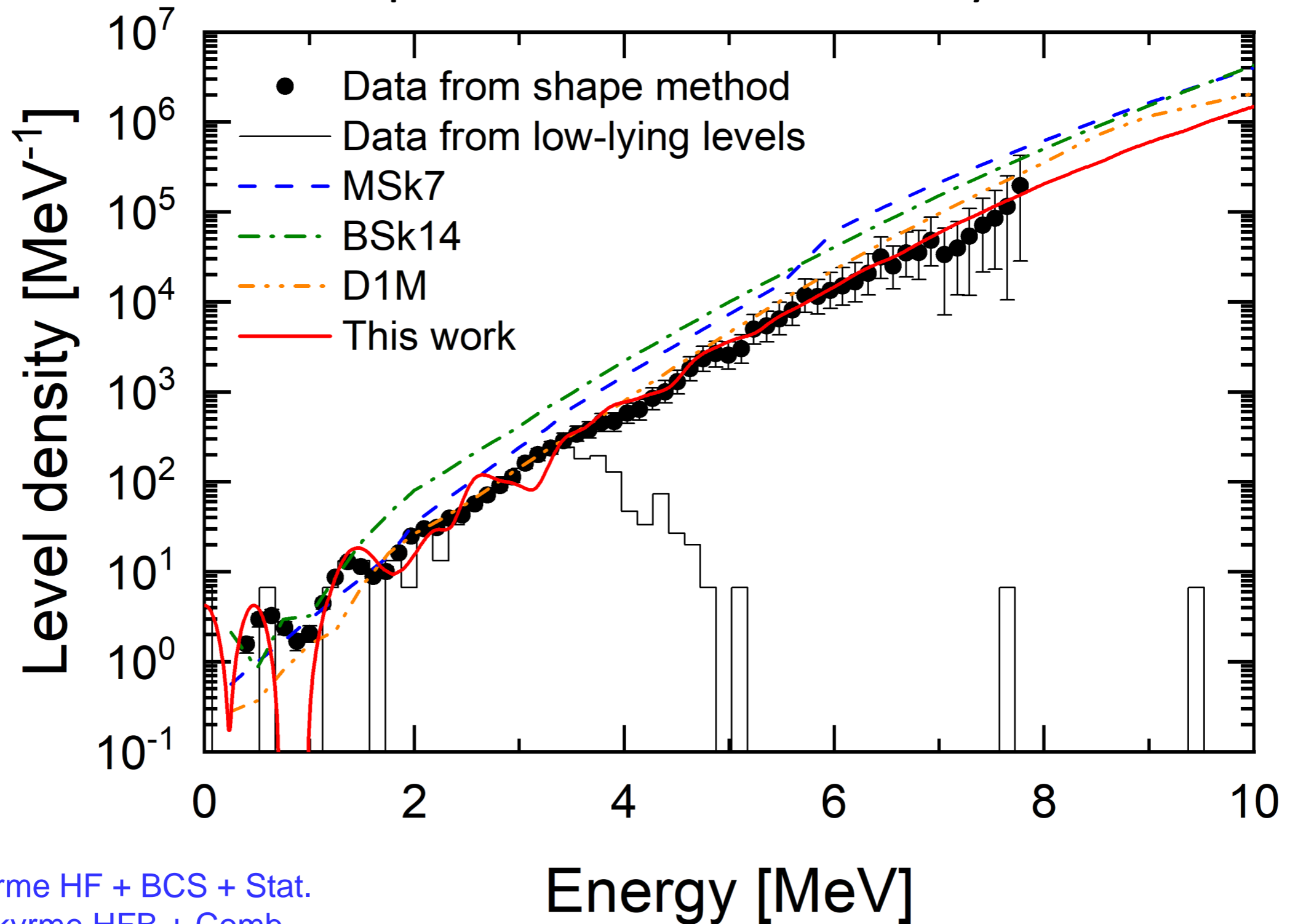


aiyuan Zhang, Xiaotao He, Jie Meng, Cong Pan, Caiwan Shen, Chen Wang, and Shuangquan Zhang, **Phys. Rev. C** **104**, L021301 – Published 5 August 2021

K. Y. Zhang *et al.* (DRHBc Collaboration), Nuclear mass table in deformed relativistic Hartree-Bogoliubov theory in continuum, I: Even-even nuclei, **At. Data Nucl. Data Tables** **144**, 101488 (2022)

Nuclear level density

CDFT + combination + Strutinski well
reproduce the level density in ^{112}Cd

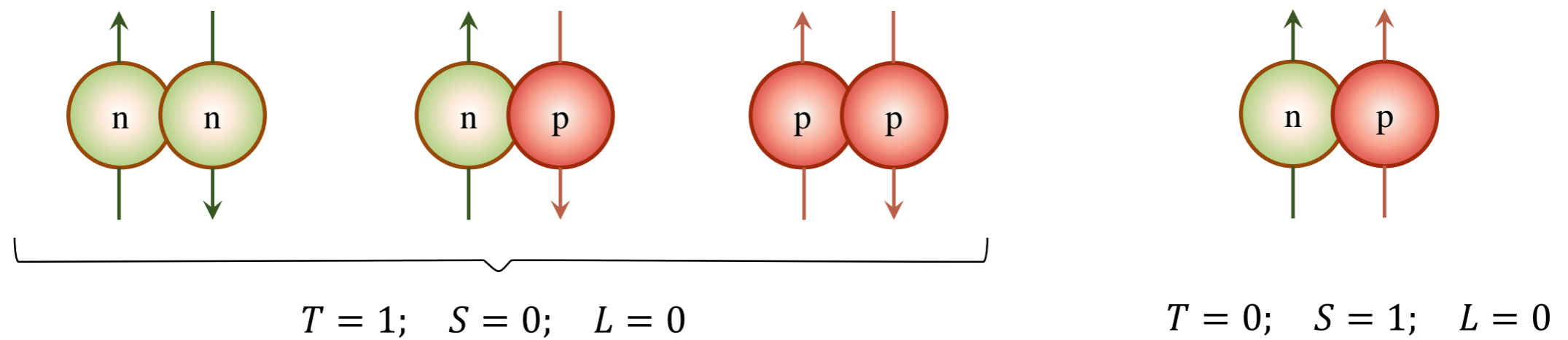


MSk7: Skyrme HF + BCS + Stat.

BSk14: Skyrme HFB + Comb.

D1M: Gogny THFB + Comb.

Physics along $N=Z$ nuclei

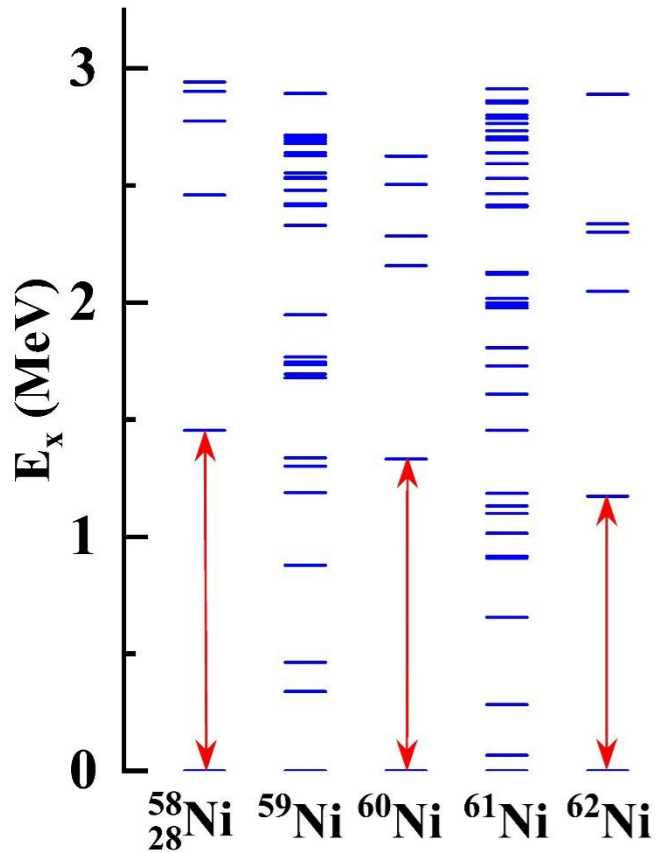


pp + nn + pn pairing

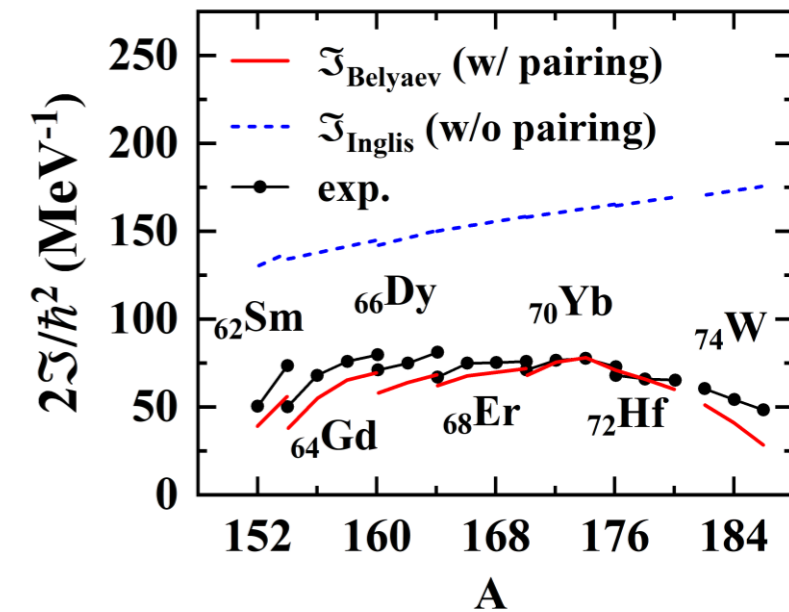
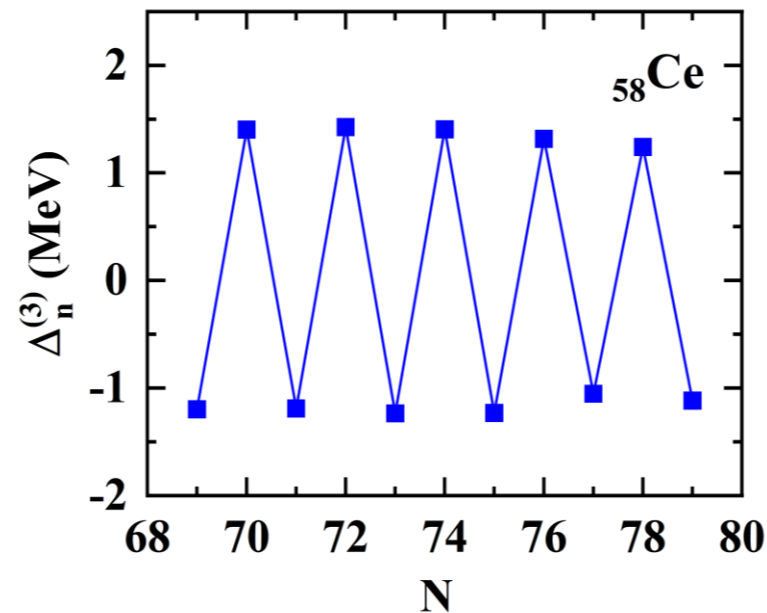
Evidence for pairing correlations

① Low-lying spectra ② 0-E mass differences

③ MoI



$$\Delta_n^{(3)}(N, Z) = B(N, Z) - \frac{B(N+1, Z) + B(N-1, Z)}{2}$$

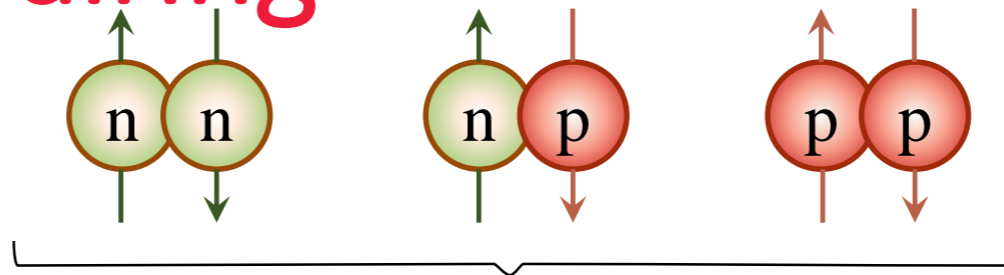


Data from National Nuclear Data Center

Bohr, Mottelson, Pines, PR 110, 936 (1958)

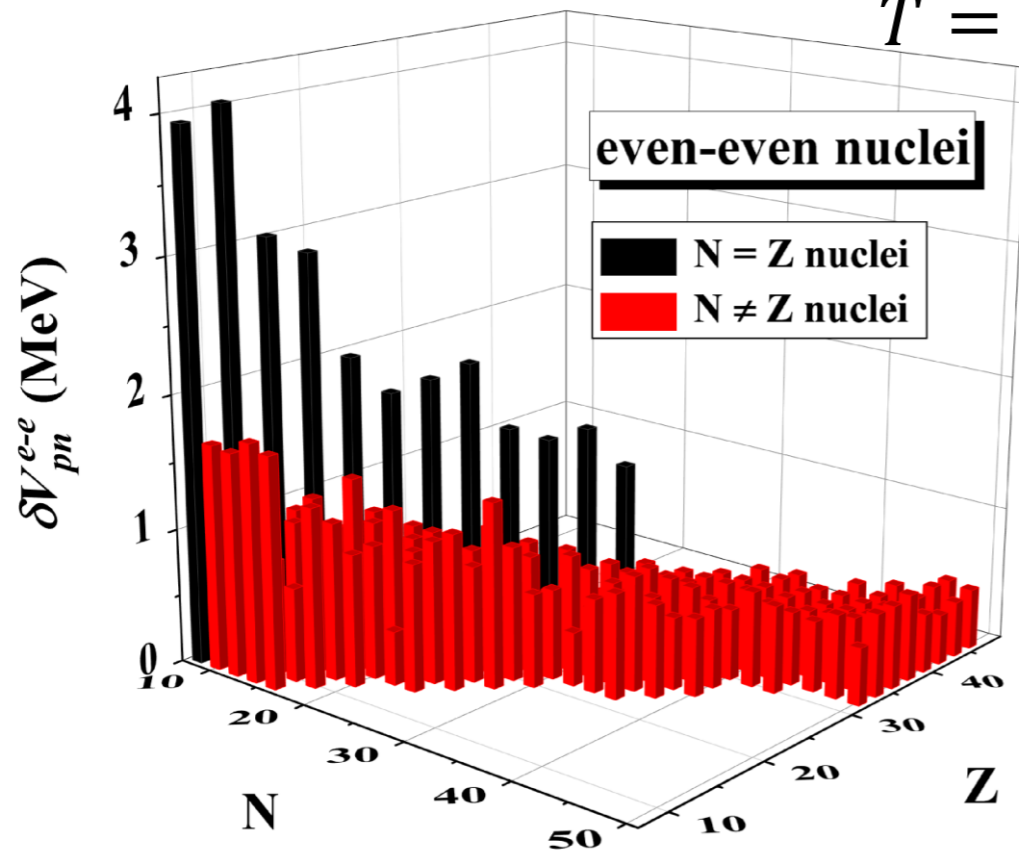
Evidence for pn pairing

□ Wigner terms



$T = 1; S = 0; L = 0$

$T = 0; S = 1; L = 0$



Z	${}^{62}\text{Ge}$	${}^{64}\text{Ge}$
$Z - 2$	${}^{60}\text{Zn}$	${}^{62}\text{Zn}$
	$N - 2$	N

$$\delta V_{pn}(N, Z) = \frac{1}{4} [B(N, Z) + B(N - 2, Z - 2) - B(N - 2, Z) - B(N, Z - 2)] \sim \frac{\partial^2 B}{\partial N \partial Z}$$

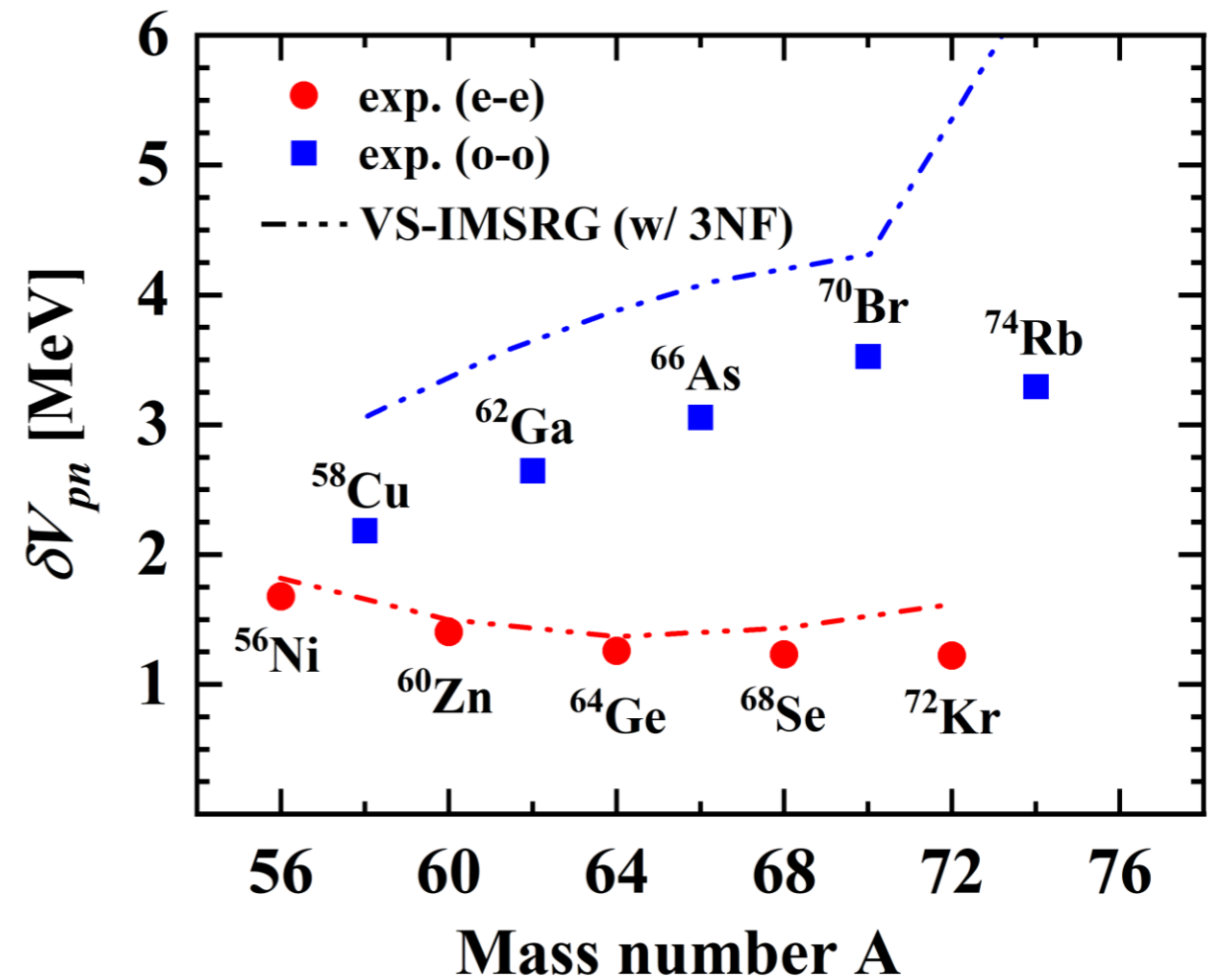
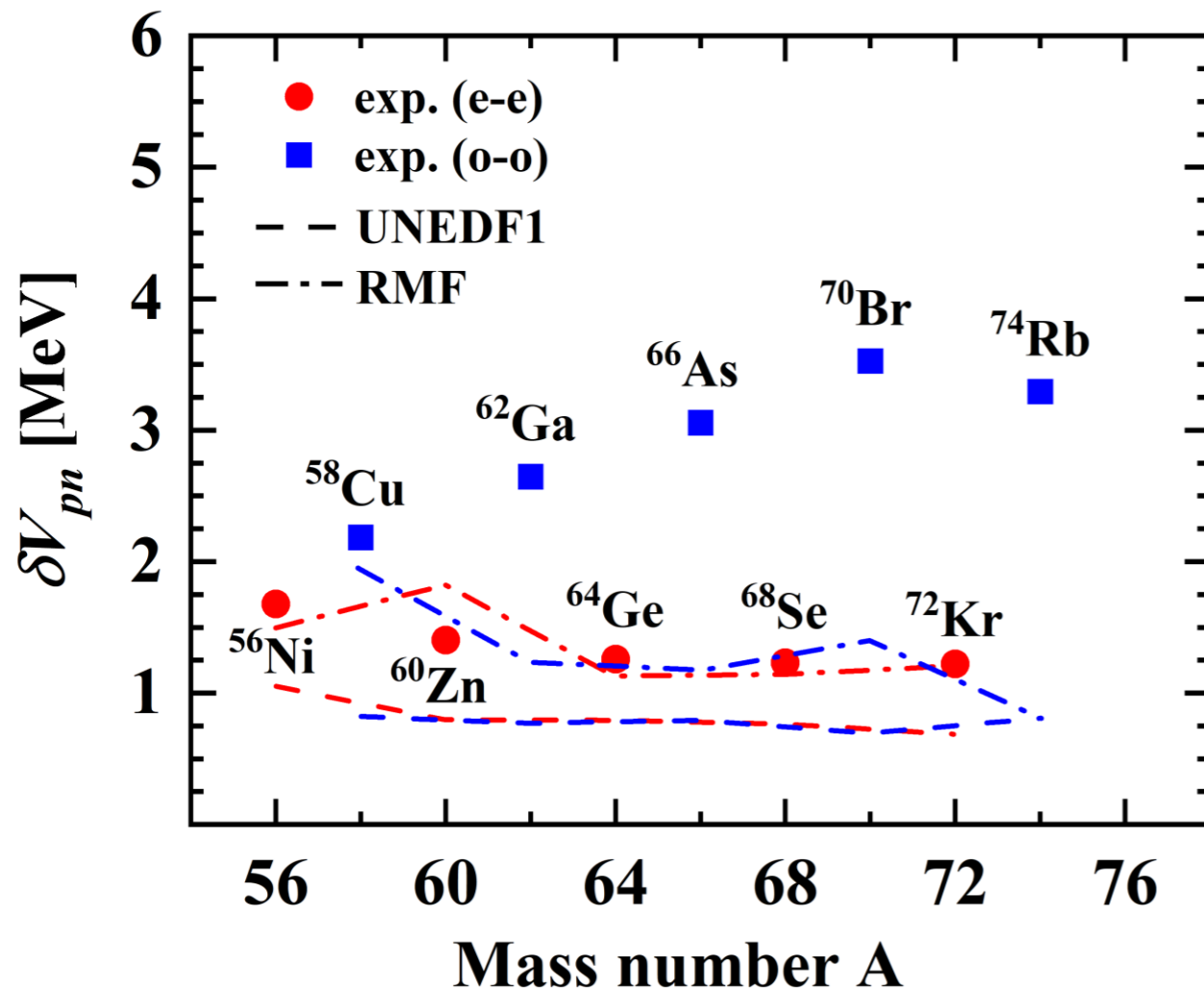
Zhang, Casten, Brenner, PLB 227, 1 (1989); Data from National Nuclear Data Center

Larger δV_{pn} reflects larger binding energy with $N = Z$

Warner, Bentley, Van Isacker, Nature 2, 311 (2006)

Abnormal bifurcation

The double binding energy differences δV_{pn} between the odd-odd and even-even nuclei along the $N=Z$ line



RMF: Geng *et al.*, PTP 113, 785 (2005)

Wang *et al.*, PRL 130, 192501 (2023)

UNEDF1: Kortelainen *et al.*, PRC 85, 024304 (2012)

3-body force doesn't work!

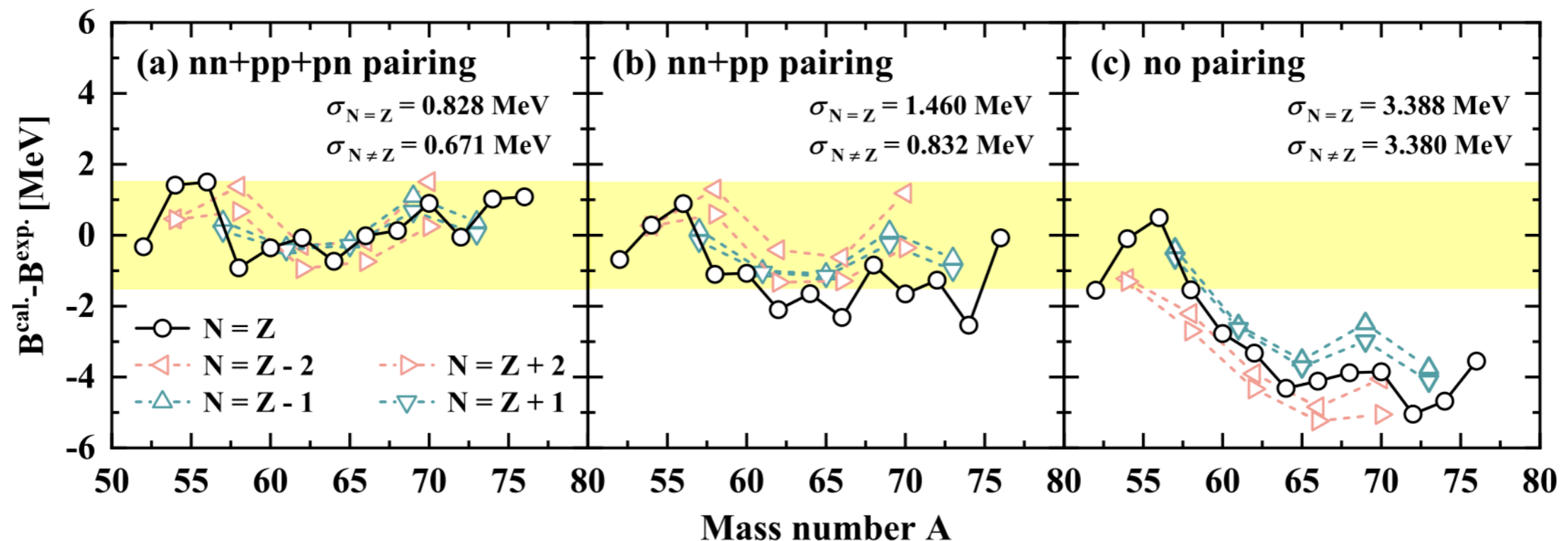
Relativistic density functional theory + SLAP

Shell model Like Approach

Yang & Zeng, Acta Physica Sinica 20, 846 (1964)

Zeng & Cheng, NPA 405, 1 (1983)

- **Treating the blocking effects exactly**
- **good particle number**
- **treating simultaneously nn, pp, pn pairing**



Y. P. Wang, Y. K. Wang, F. F. Xu, P. W. Zhao, and J. Meng
Phys. Rev. Lett. **132**, 232501 – Published 4 June 2024

Relativistic density functional theory + SLAP

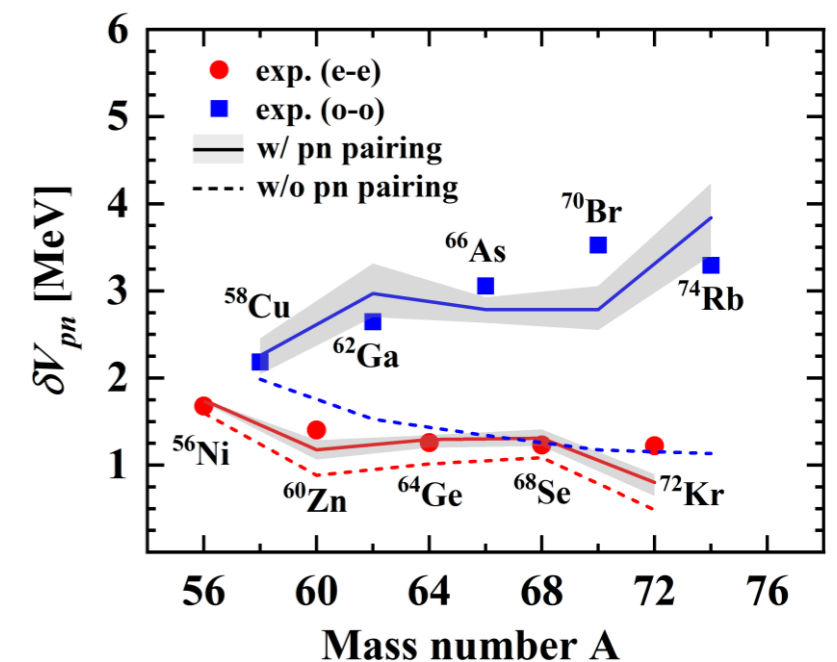
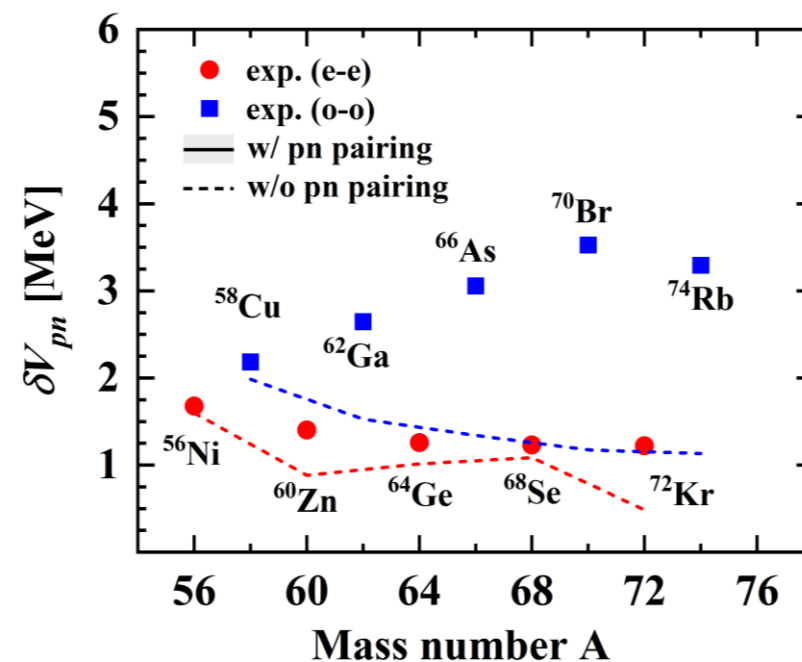
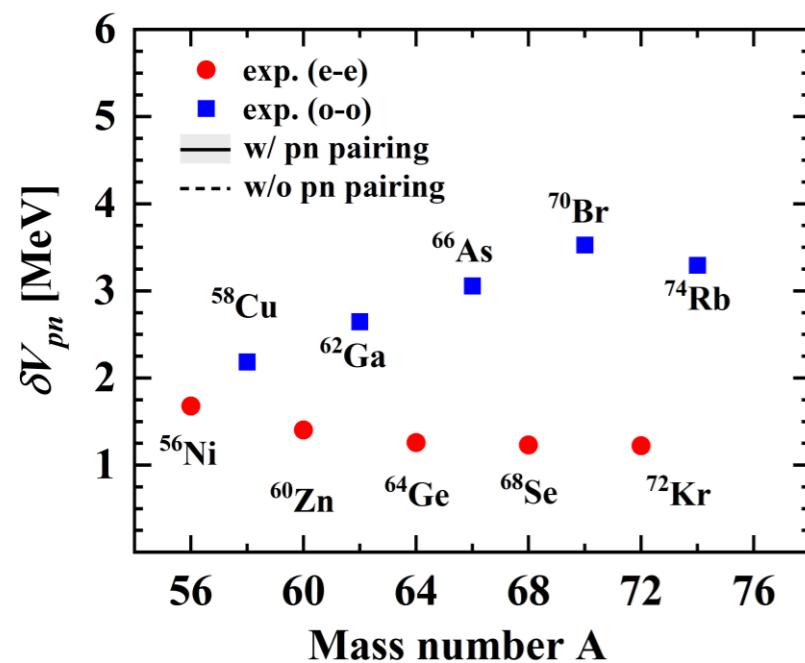
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provide excellent interpretation for the abnormal δV_{pn} bifurcation, and signal for the pn pairing correlations for nuclei close to the $N = Z$ line.

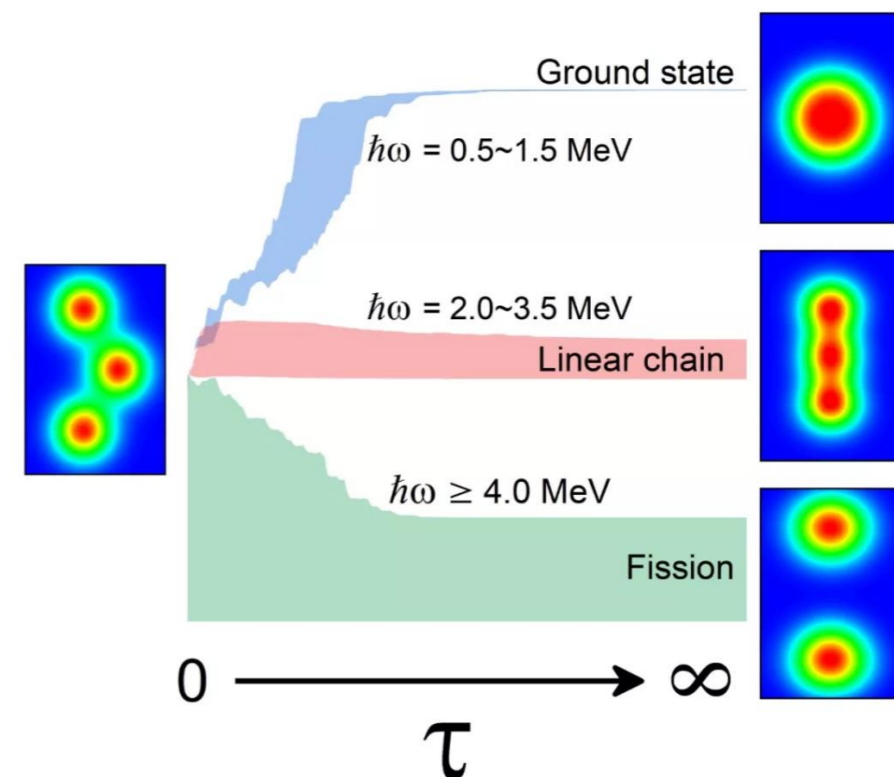


Y. P. Wang, Y. K. Wang, F. F. Xu, P. W. Zhao, and J. Meng

Phys. Rev. Lett. 132, 232501 – Published 4 June 2024

Relativistic density functional theory in 3D lattice

- **Linear alpha-chain**
- **Nuclear fusion**
- **Nuclear fission**
- **Chiral dynamics**



Time-dependent CDFT

The many-body problem is mapped onto a one-body problem!

Runge-Gross Theorem

There is a **unique mapping** between the **time dependent external potential** and the **density**, for many body systems evolving from a **given initial state**.

Runge and Gross, PRL 52, 997 (1984)

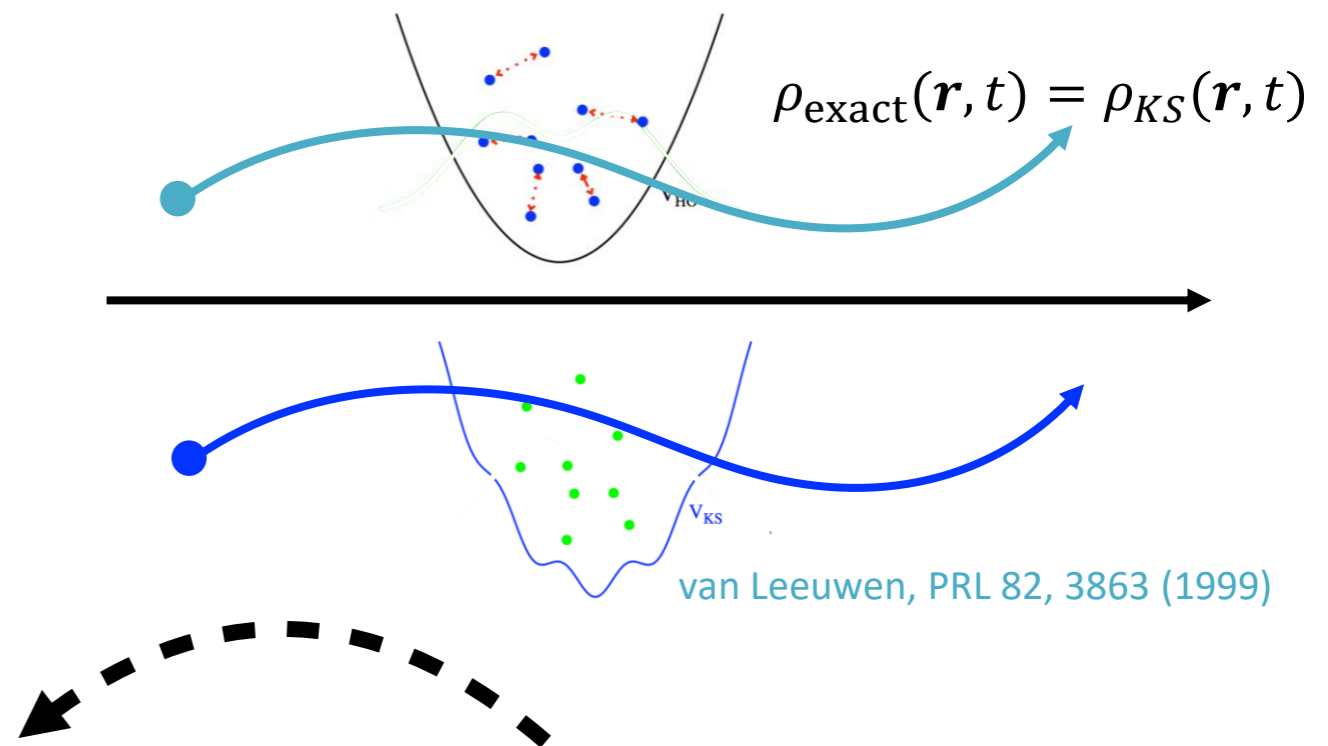
Ren, Zhao, Meng, PLB 801,135194 (2020)

$$i\partial_t \begin{pmatrix} f \\ g \end{pmatrix} = \begin{pmatrix} m + V + S & \boldsymbol{\sigma} \cdot \mathbf{p} - \boldsymbol{\sigma} \cdot \mathbf{V} \\ \boldsymbol{\sigma} \cdot \mathbf{p} - \boldsymbol{\sigma} \cdot \mathbf{V} & -m + V - S \end{pmatrix} \begin{pmatrix} f \\ g \end{pmatrix}$$

$$\rho(\mathbf{r}, t) = \sum_i^N f_i^2 + g_i^2$$

$V[\rho](\mathbf{r}, t)$ No memory effects!

Time-dependent Kohn-Sham DFT



van Leeuwen, PRL 82, 3863 (1999)

Applications of the TD-CDFT

3D Lattice: no spatial symmetry restriction

✓ Applications include:

Linear alpha-chain

PRL 115, 022501 (2015)

PLB 801,135194 (2020)

Nuclear fission

PRL 128, 172501 (2022)

PRC 105, 044313 (2022)

PRC 107, 014303 (2023)

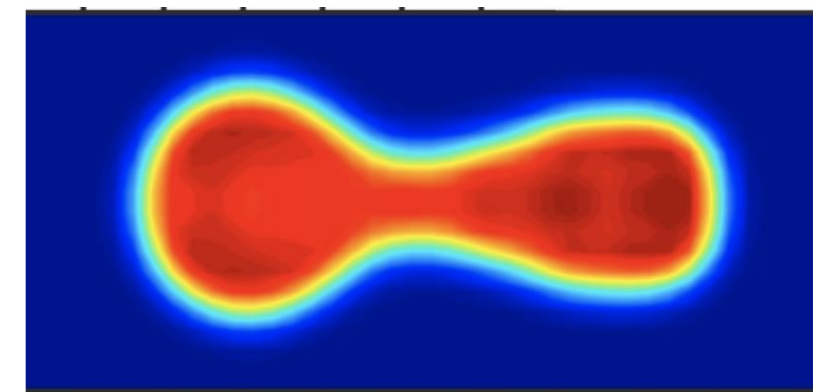
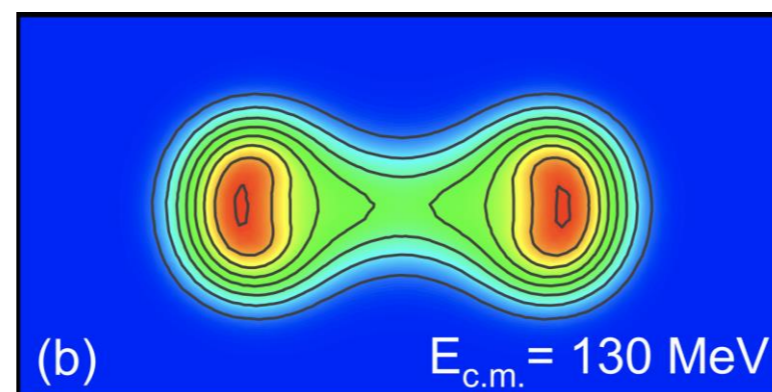
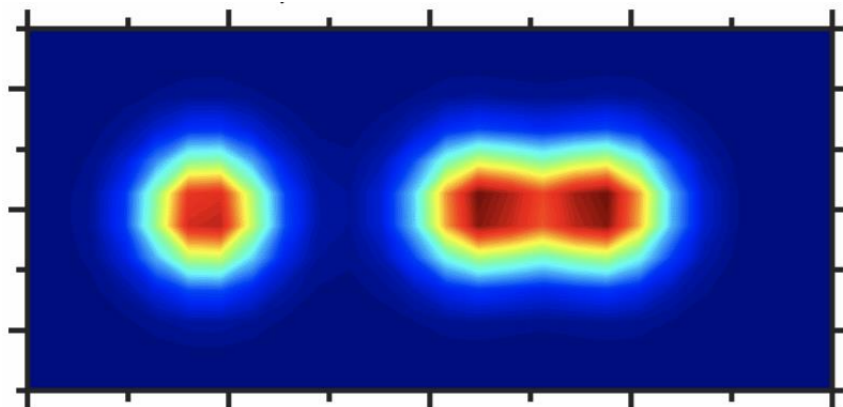
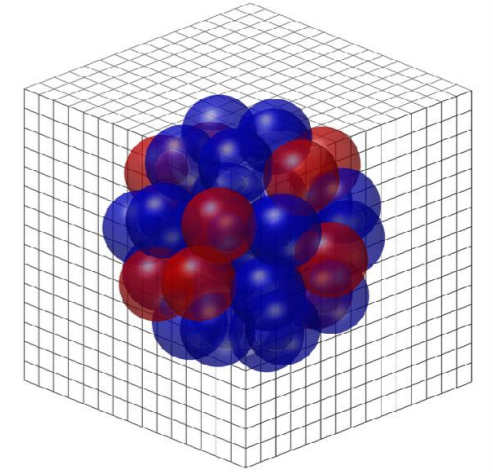
Chiral dynamics

PRC, 105, L011301 (2022)

$^{16}\text{O} + ^{16}\text{O}$ reaction

PRC 102, 044603 (2020)

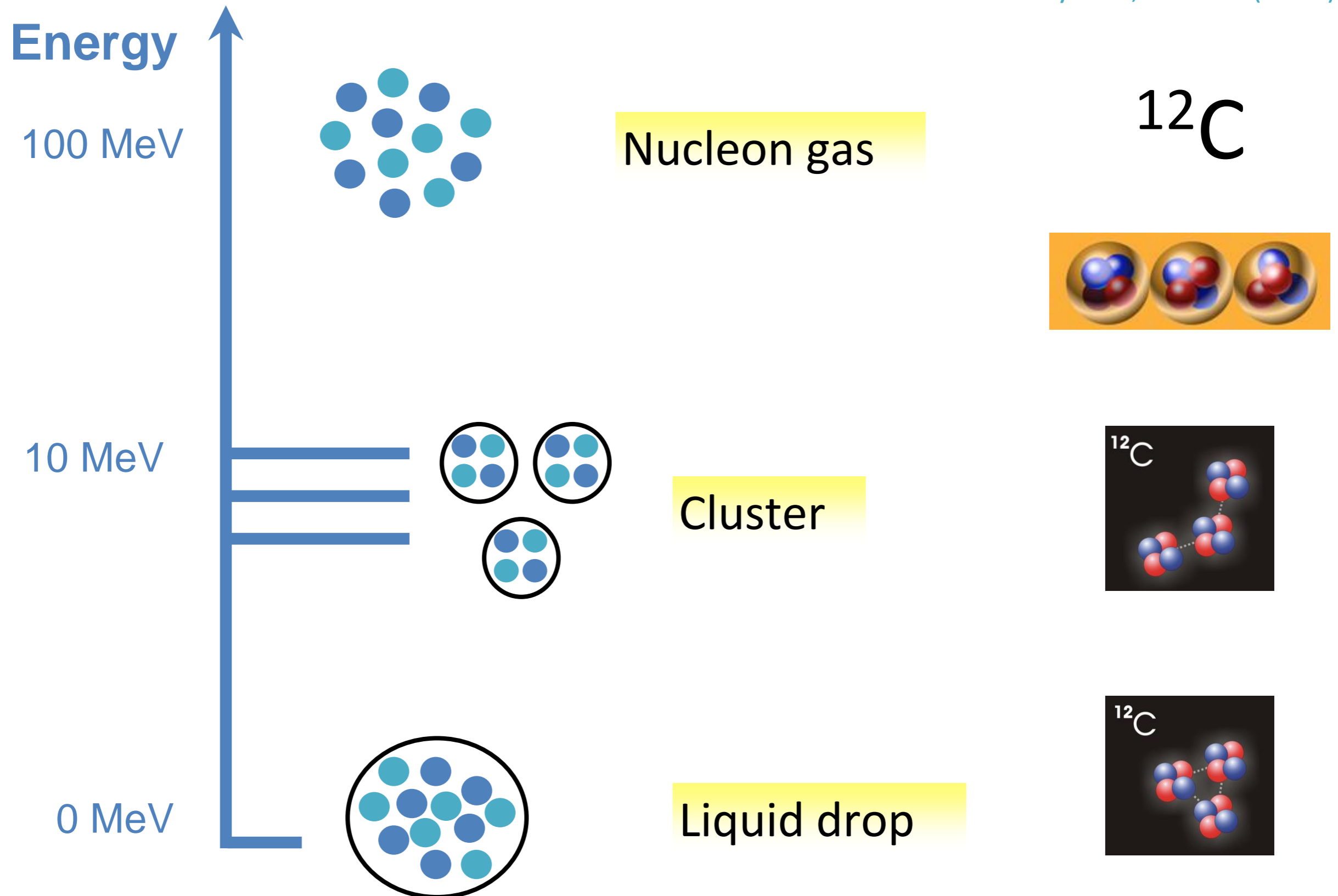
...



Linear alpha-chain

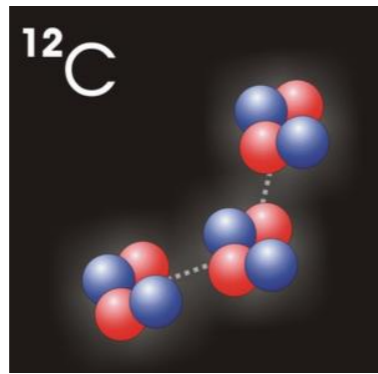
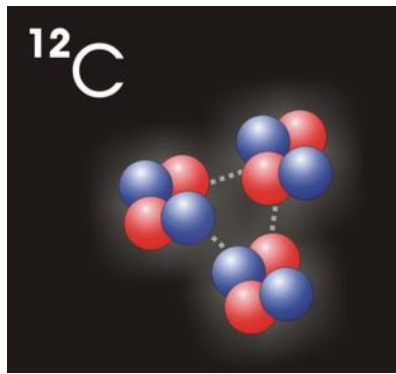
Nuclear cluster

Rev. Mod. Phys. 90, 035004 (2018)



Linear-chain clustering

Strongly deformed states [towards a hyper-deformation](#) may exist in light $N = Z$ nuclei due to a cluster structure.



- ◆ the linear alpha cluster chain has been searched more than 60 years.
- ◆ new radioactive beams provide new opportunities in realizing the linear chain state.

No firm evidence



Two difficulties

- ✓ antisymmetrization effects
- ✓ weak-coupling nature

Two mechanisms

- ✓ adding neutrons (isospin)
- ✓ rotating the system (spin)

Itagaki, PRC2001; Maruhn, NPA2010;
Ichikawa, PRL2011

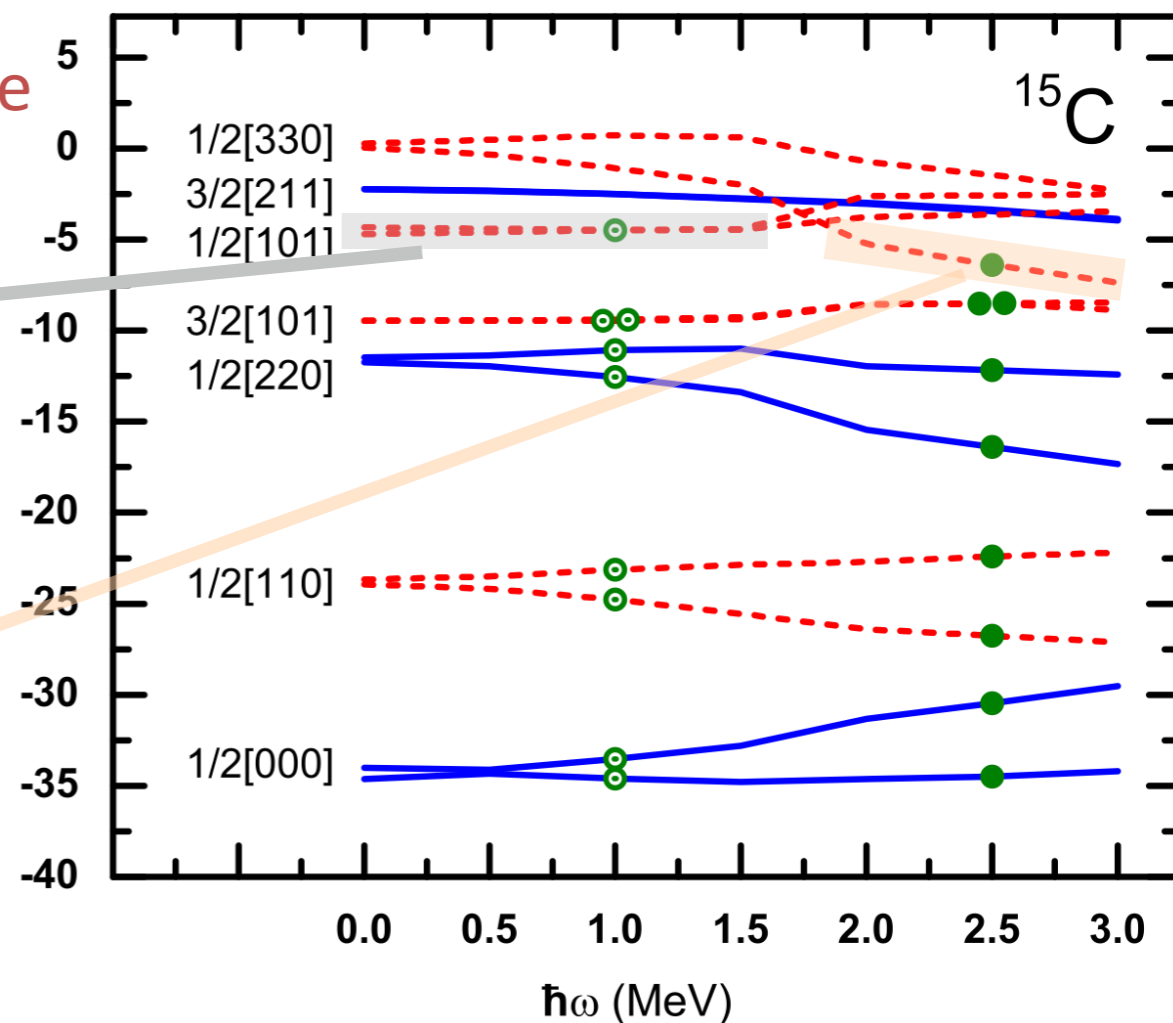
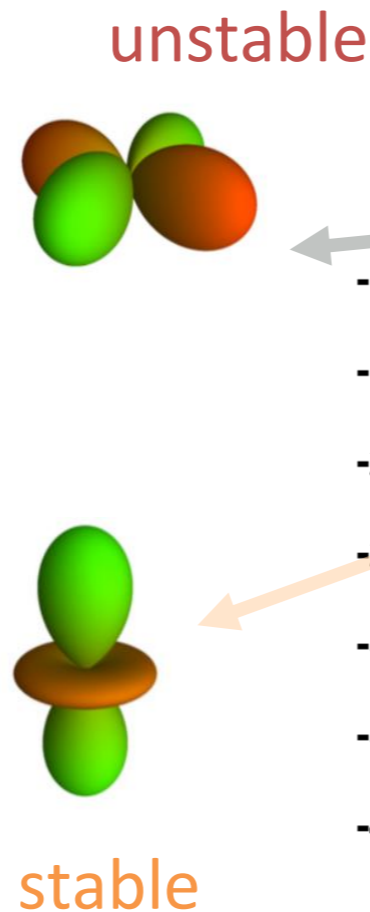
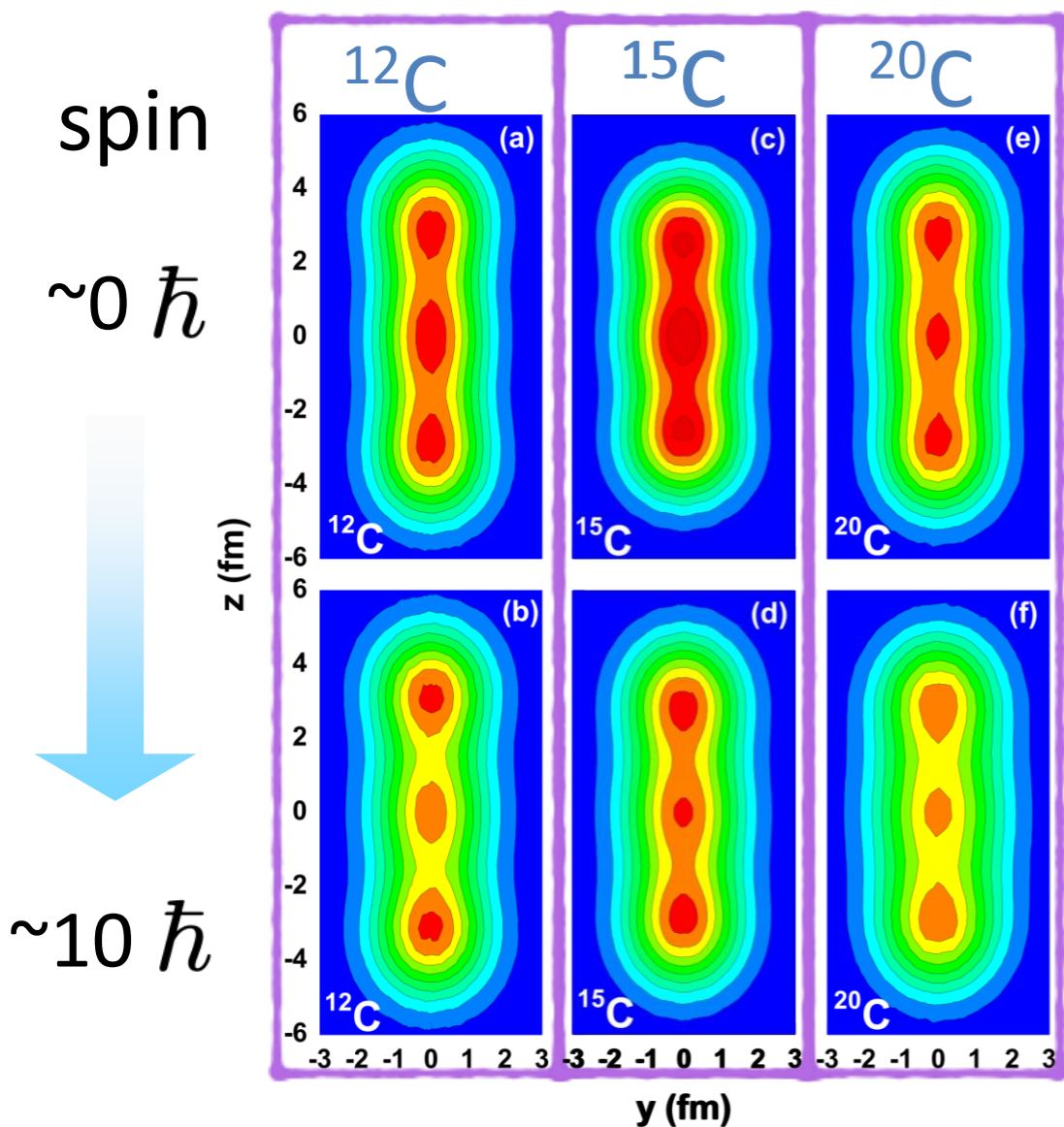
CDFT is employed without assuming clustering a priori.

Spin and Isospin Coherent Effects

Static calculations with reflection-symmetry

Proton density distribution

neutron orbitals



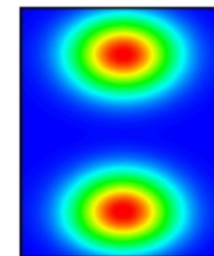
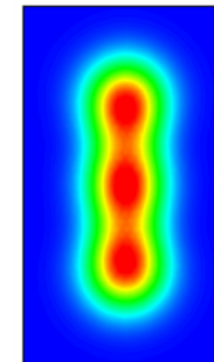
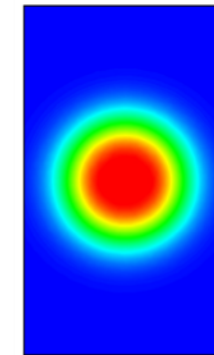
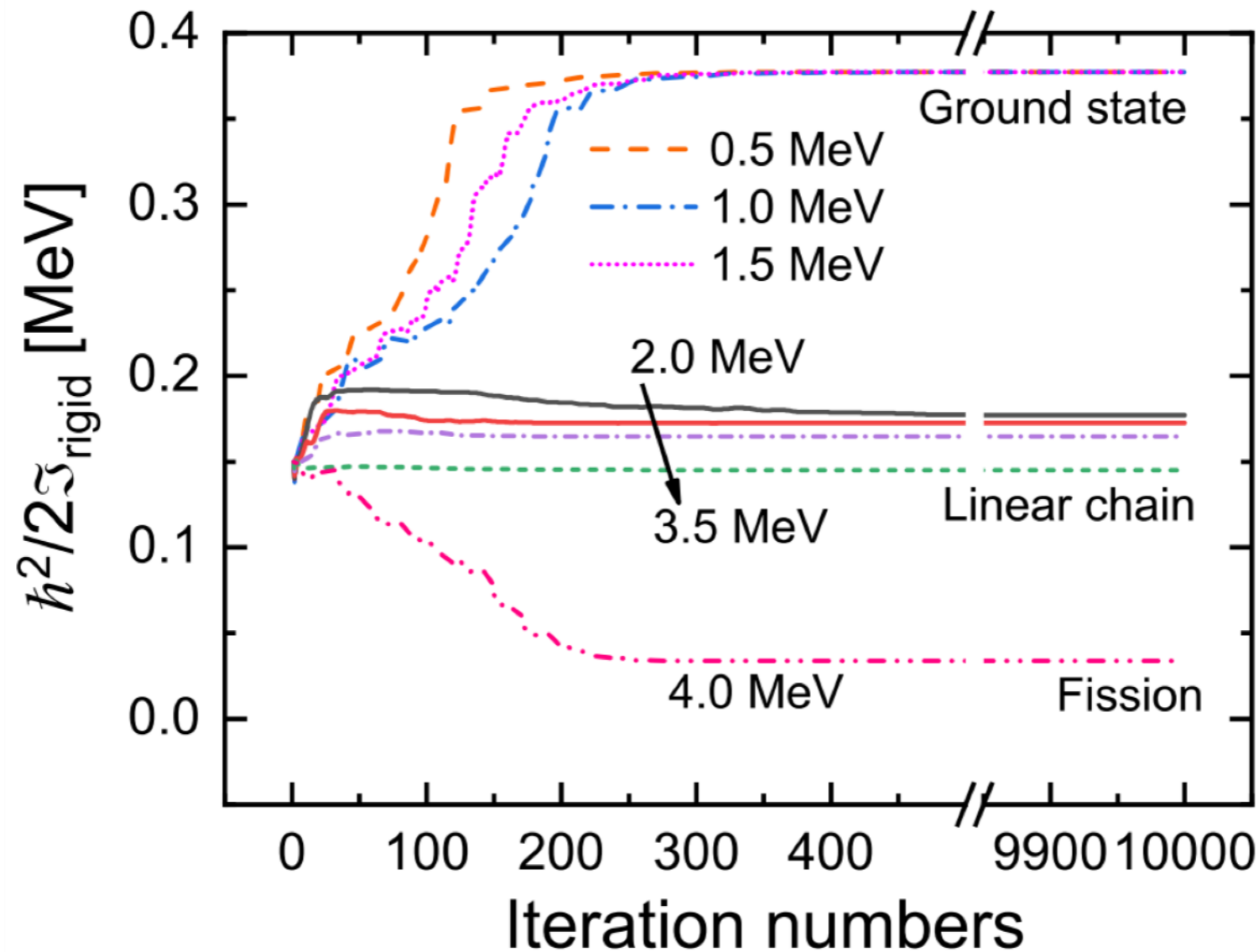
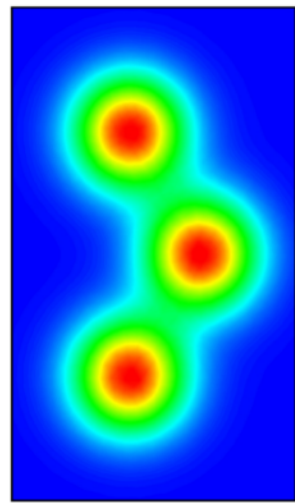
Zhao, Itagaki, Meng, PRL 115, 022501 (2015)

Rod shapes could be realized towards extreme spin and isospin!

Rod shape against bending and fission

Static calculations in 3D lattice

^{12}C

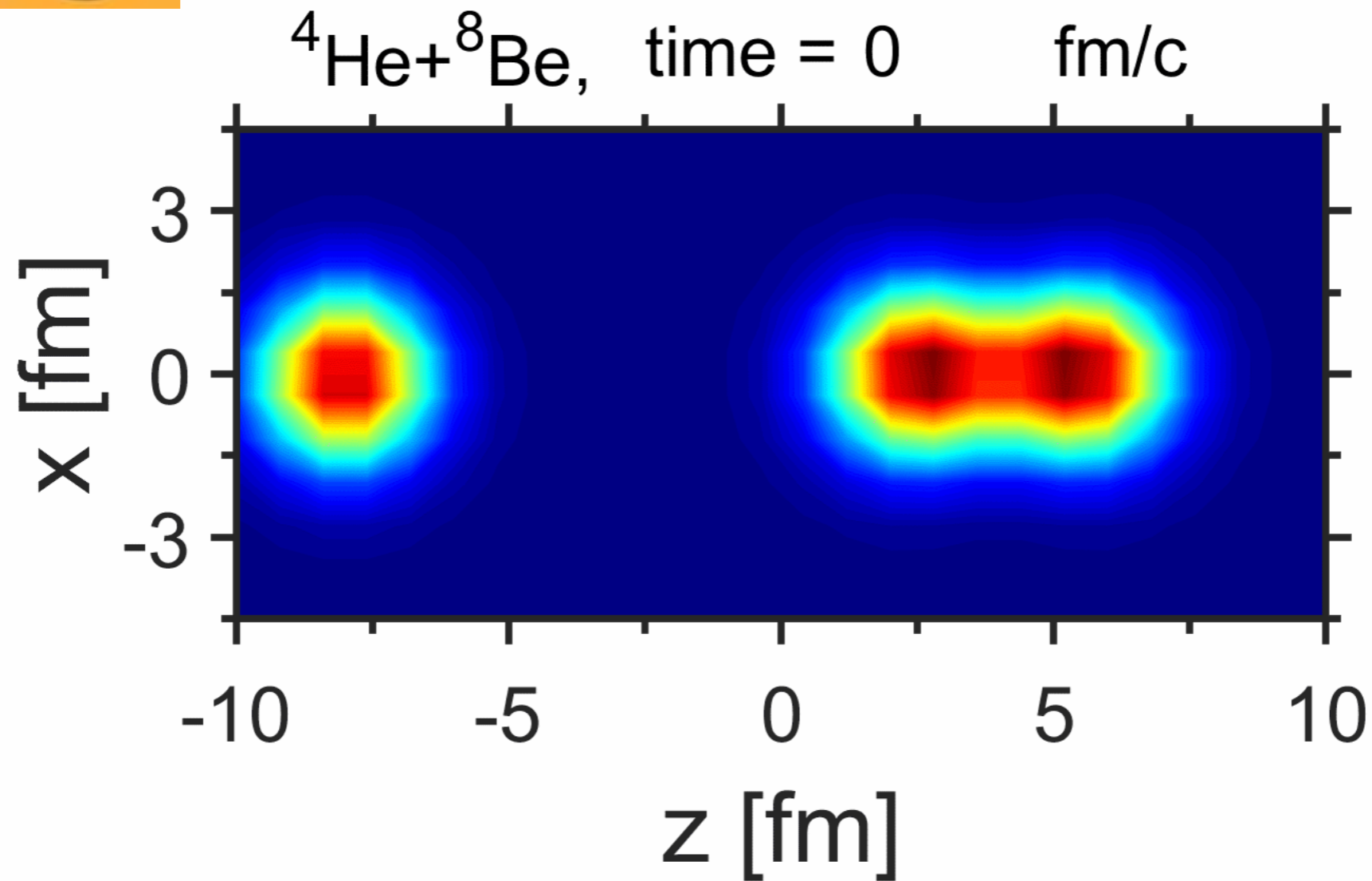


Rod shapes are generated as energy minima at a certain range of rotational frequencies.

Resonant scattering of ${}^4\text{He} + {}^8\text{Be}$

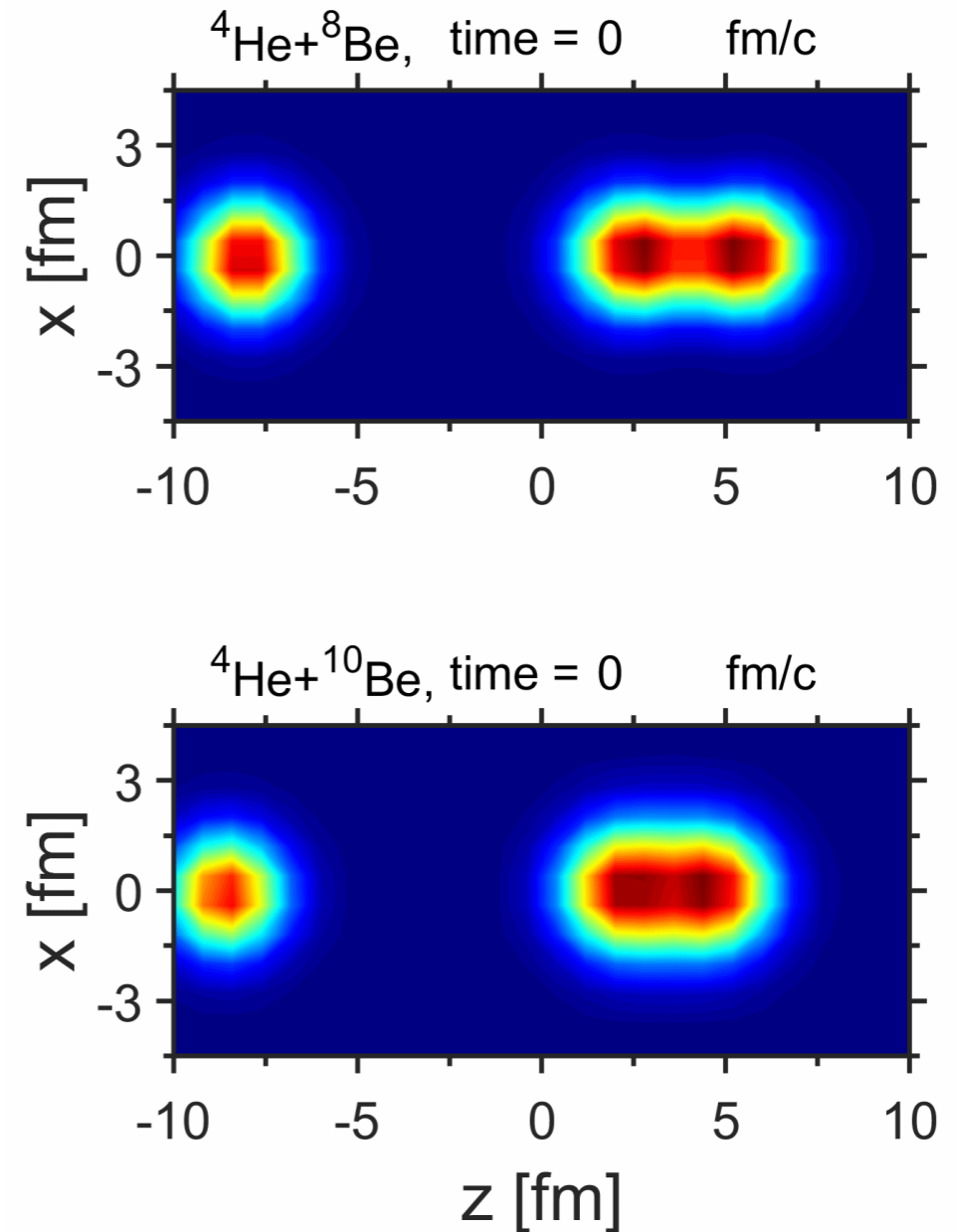
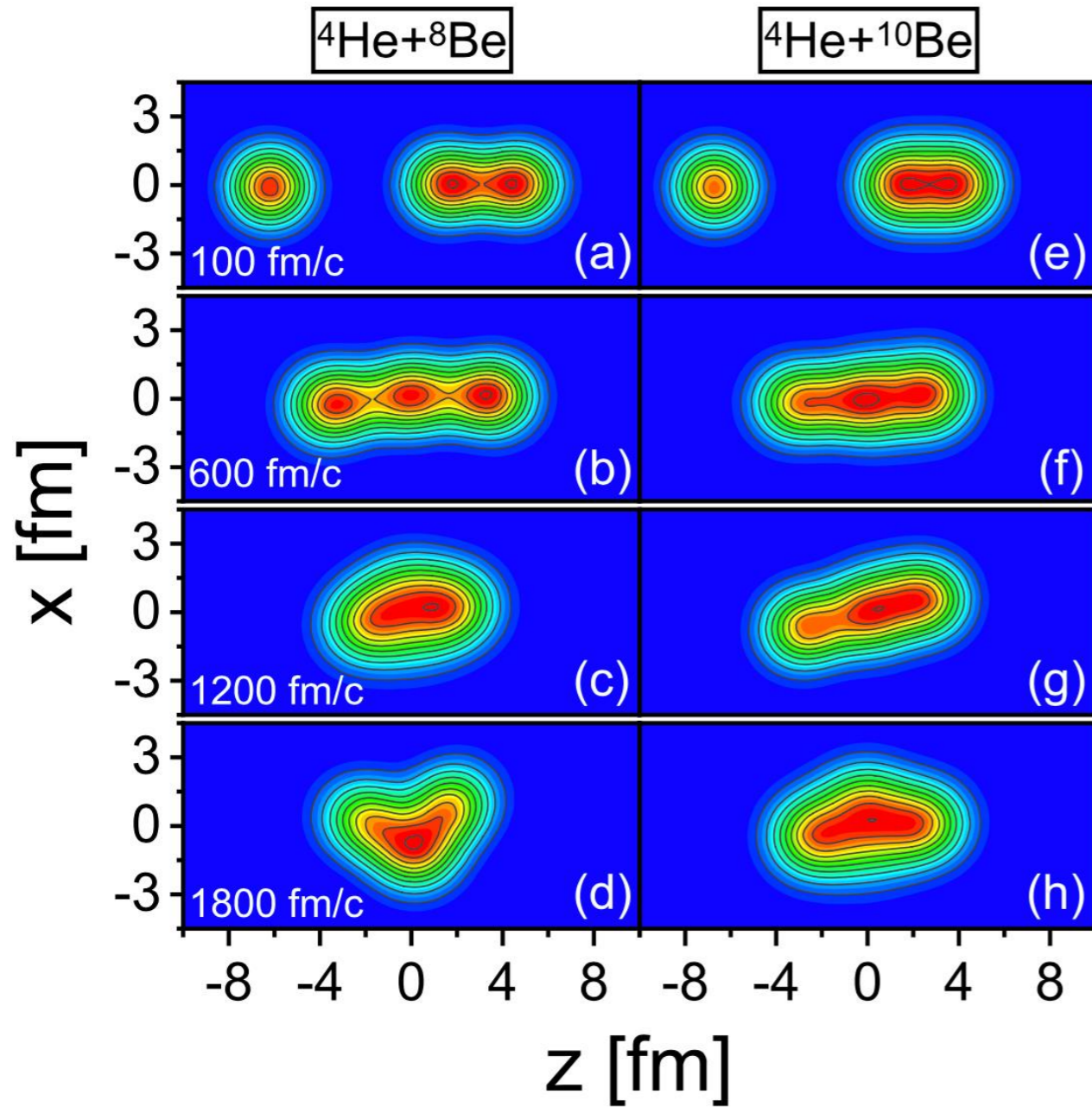


Ren, Zhao, Meng, PLB 801, 135194 (2020)



The linear-chain states are generated, and then evolve to a triangular configuration, and finally to a more compact shape.

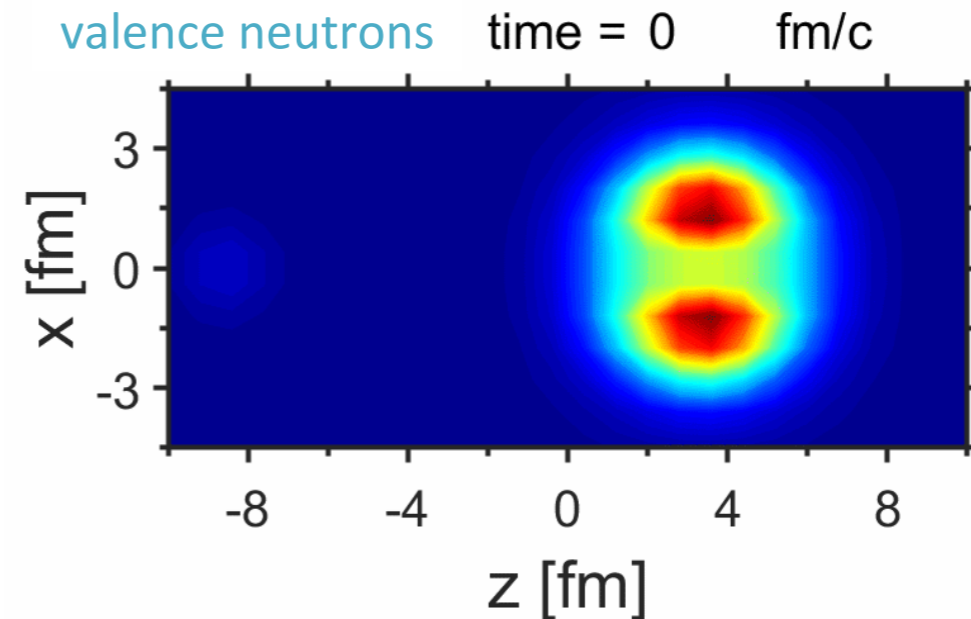
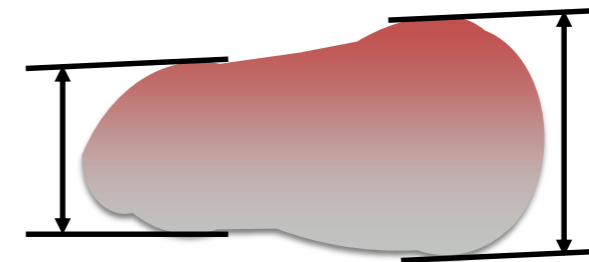
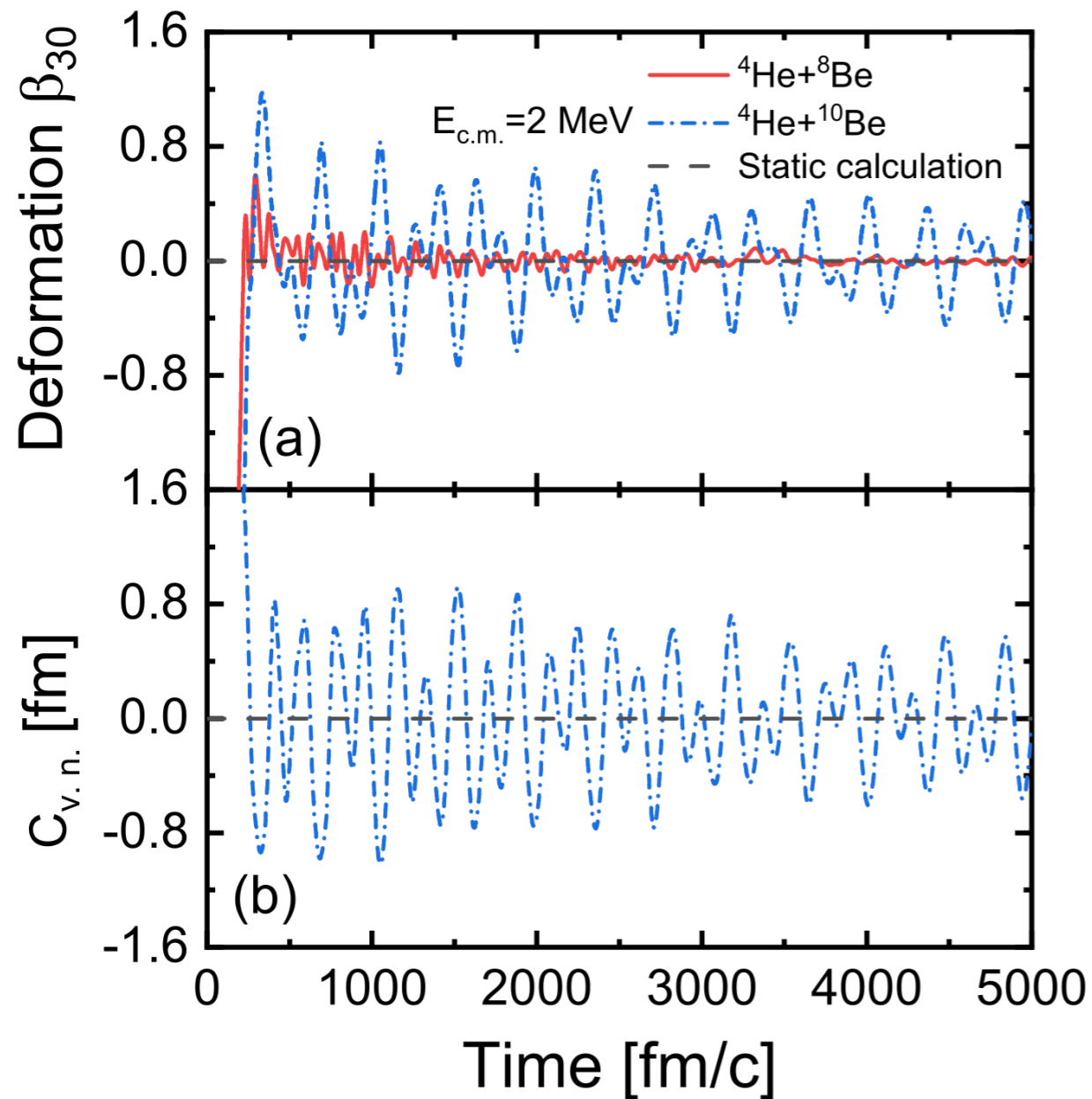
$^4\text{He} + ^{10}\text{Be}$



Ren, Zhao, Meng, PLB 801, 135194 (2020)

- ✓ The metastable linear chains can be formed in $^4\text{He} + ^8\text{Be}$ and $^4\text{He} + ^{10}\text{Be}$ collisions.
- ✓ During the time evolution of the linear-chain configuration, moving clusters can be found.

Octupole deformation

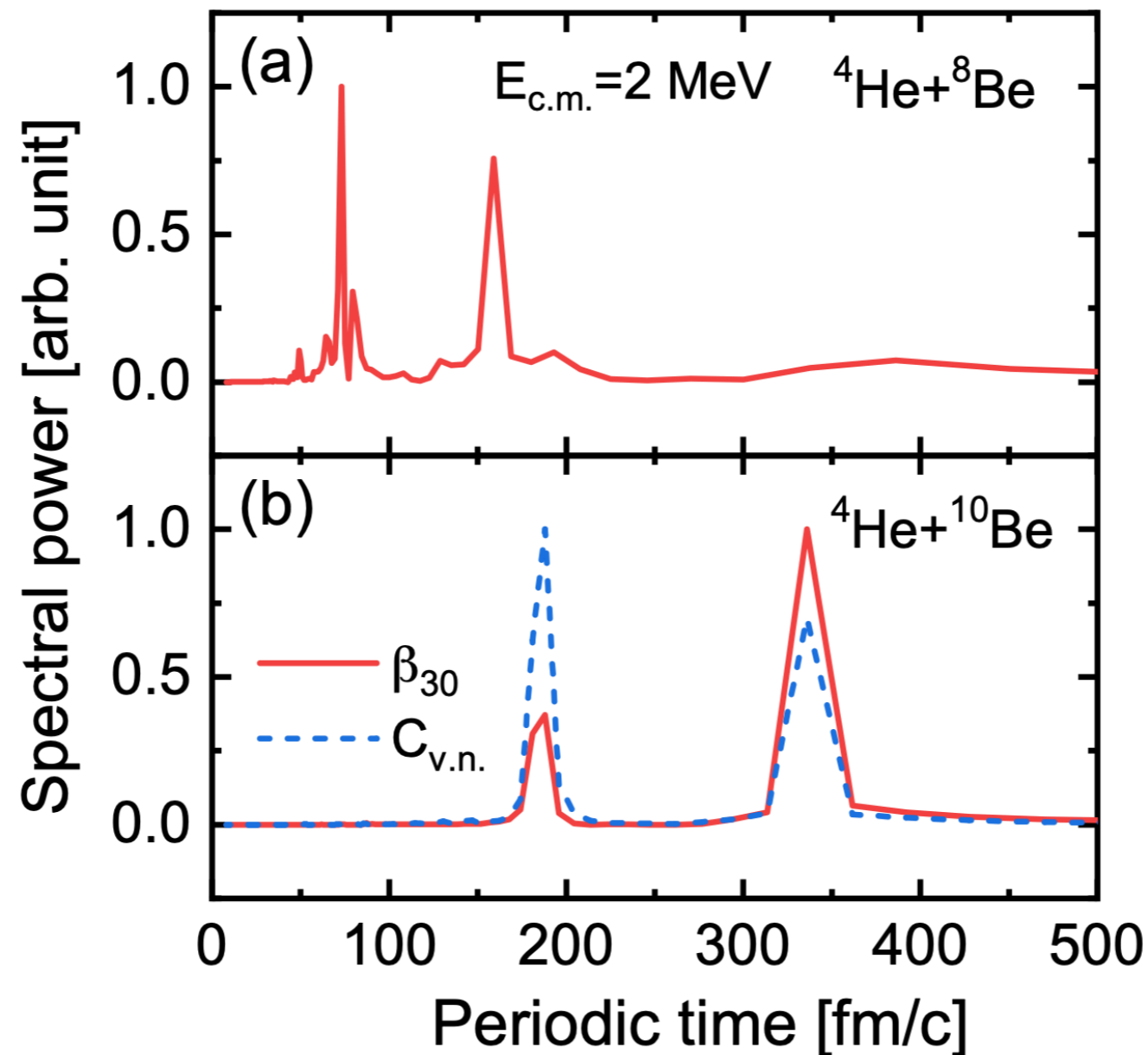


center of valence
neutrons

$$C_{v.n.} = \frac{\int d^3\mathbf{r} z (\rho_n - \rho_p)}{\int d^3\mathbf{r} (\rho_n - \rho_p)}$$

The oscillation of the two valence neutrons in the longitudinal direction induces the strong oscillation of the octupole deformation.

Dynamical isospin effects



Ren, Zhao, Meng, PLB 801, 135194 (2020)

Dynamical isospin effects: slowing down the longitudinal oscillations by the two valence neutrons.

Chiral dynamics

Nuclear spin-chirality

The **aplanar (3D-) rotation** of a **triaxial nucleus** could present **chiral geometry**.

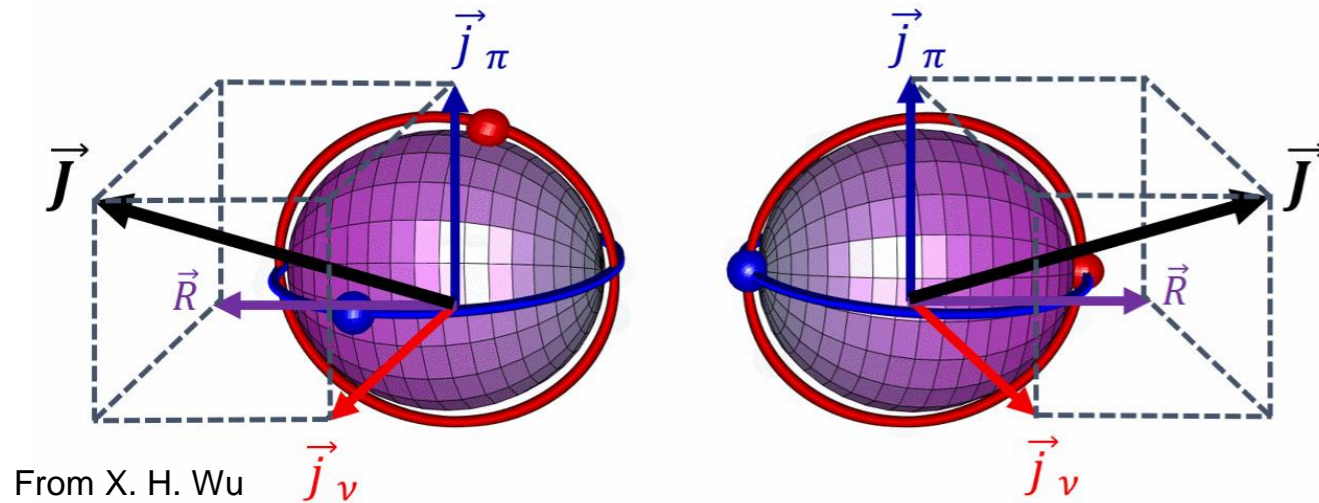
Frauendorf and Meng, Nucl. Phys. A **617**, 131 (1997)

Intrinsic frame :

Chiral Symmetry breaking

$$\hat{\chi} = \hat{T} \hat{R}_y(\pi)$$

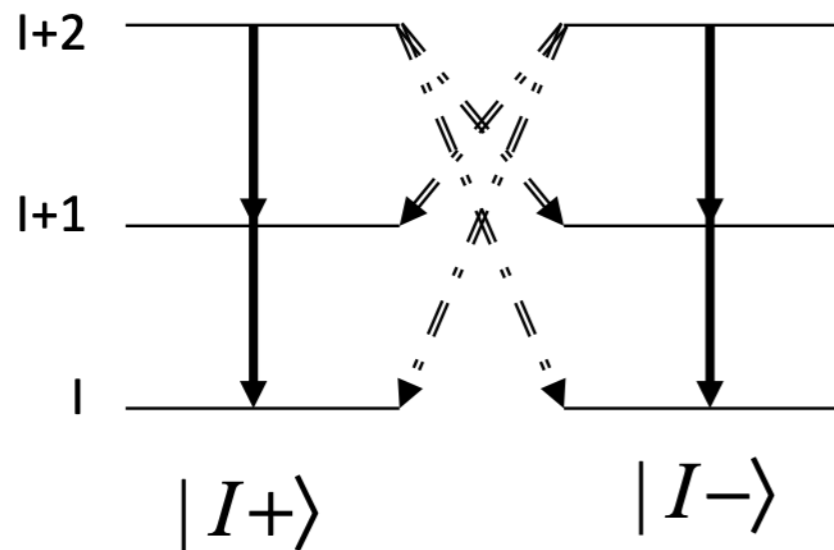
$$\hat{\chi} |\mathcal{L}\rangle = |\mathcal{R}\rangle \quad \hat{\chi} |\mathcal{R}\rangle = |\mathcal{L}\rangle$$



From X. H. Wu

Left-handed $|\mathcal{L}\rangle$

Right-handed $|\mathcal{R}\rangle$



Lab. frame :

Chiral Symmetry restoration

$$|I+\rangle = \frac{1}{\sqrt{2}}(|\mathcal{L}\rangle + |\mathcal{R}\rangle)$$

$$|I-\rangle = \frac{i}{\sqrt{2}}(|\mathcal{L}\rangle - |\mathcal{R}\rangle)$$

Exp. signal : **Two near degenerate $\Delta I = 1$ bands**, called **chiral doublet bands**



北京大学

Timeline for nuclear chirality

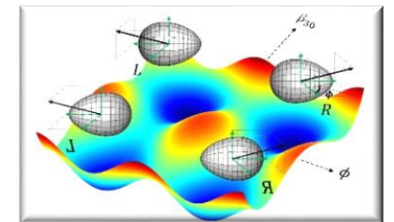
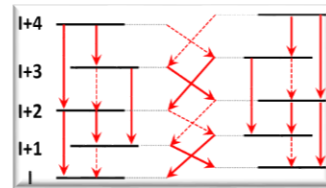
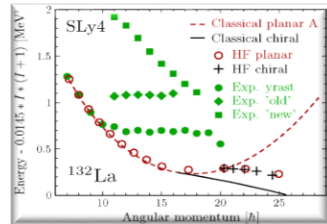
NPA 617, 131

PRL 93, 052501

PRL 93, 172502

PRC 73, 037303

Sci. Bull. 65, 2001



Prediction

Critical Frequency

EM transition Selection Rule

$M_{\chi D}$

Chirality-Parity Violation

1997

2001

2004

2006

2013

2014

2016

2018

2020

^{130}Cs ^{132}La
 ^{134}Pr ^{136}Pm
First Evidence

PRL 86, 971

^{128}Cs
Lifetime Analysis

PRL 97, 172501

^{133}Ce
First $M_{\chi D}$ Evidence

PRL 110, 172504

^{106}Ag
Chiral Conundrum Resolution

PRL 112, 202502

^{78}Br $M_{\chi D}$ with Octupole Correlations

PRL 116, 112501

^{128}Cs g factor

PRL 120, 022502

By Yiping Wang et al

Applications of TAC-CDFT for chiral nuclei

$^{102-107}\text{Rh}$, ^{106}Ag , ^{135}Nd , ^{136}Nd

Zhao PLB 773, 1-5 (2017)

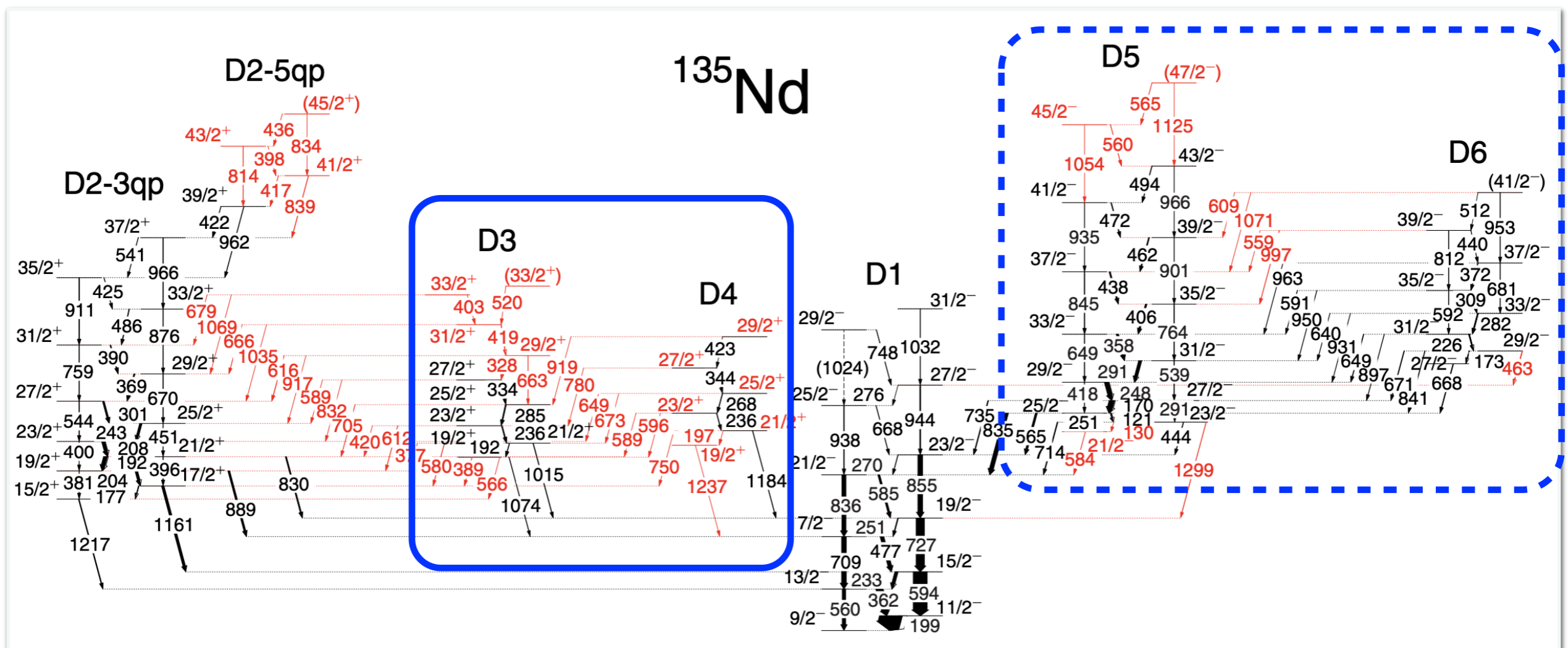
Zhao, Wang, and Chen, PRC 99, 054319 (2019)

Peng and Chen, PLB 810, 135795 (2020)

Peng and Chen, PRC 105, 044318 (2022)

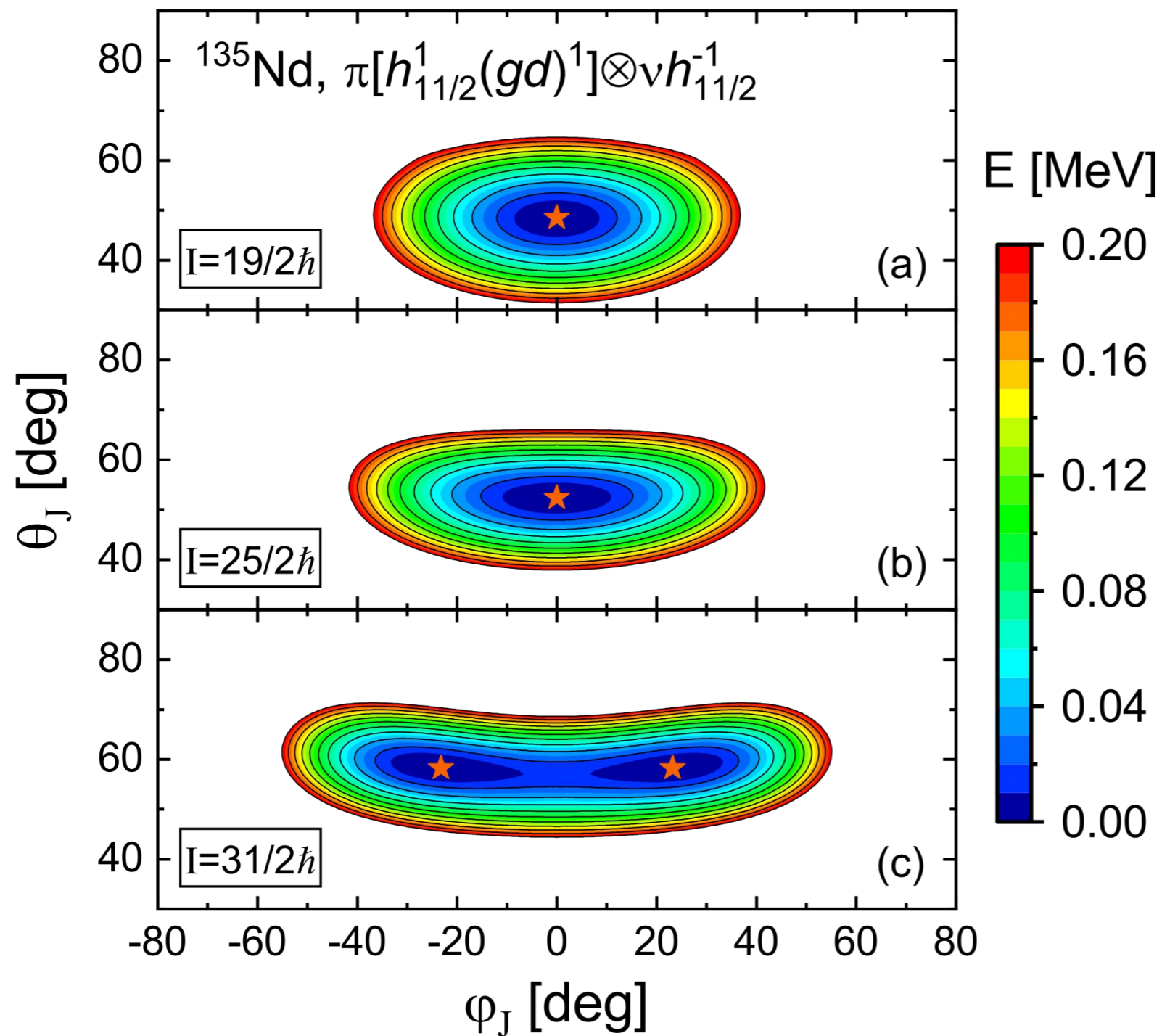
D3-D4 $\pi[h_{11/2}^1(gd)^1] \otimes \nu h_{11/2}^{-1}$

D5-D6 $\pi h_{11/2}^2 \otimes \nu h_{11/2}^{-1}$

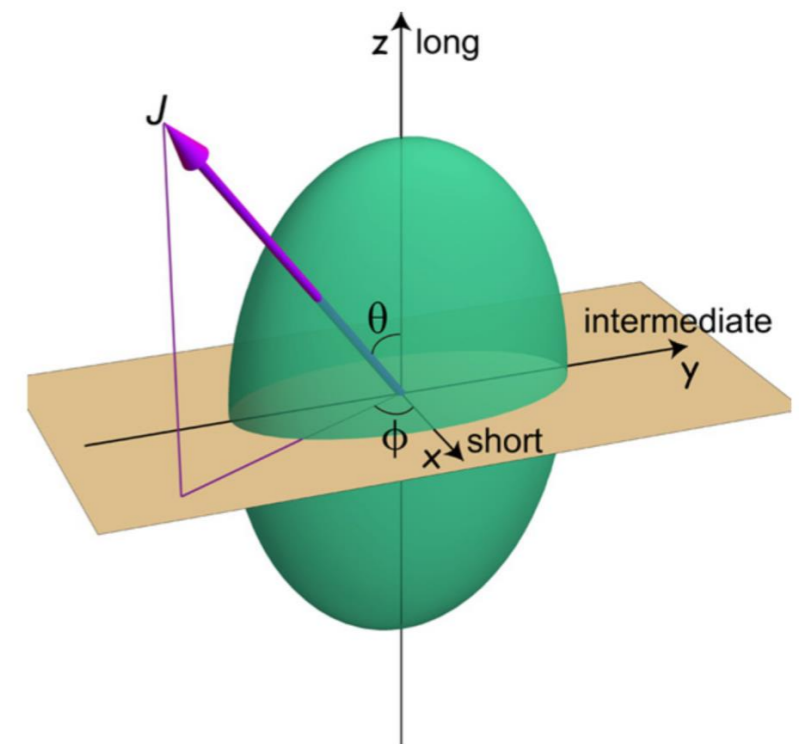


Experimental Data from Lv et al., Phys. Rev. C 100, 024314 (2019)

Energy surface against tilted angles



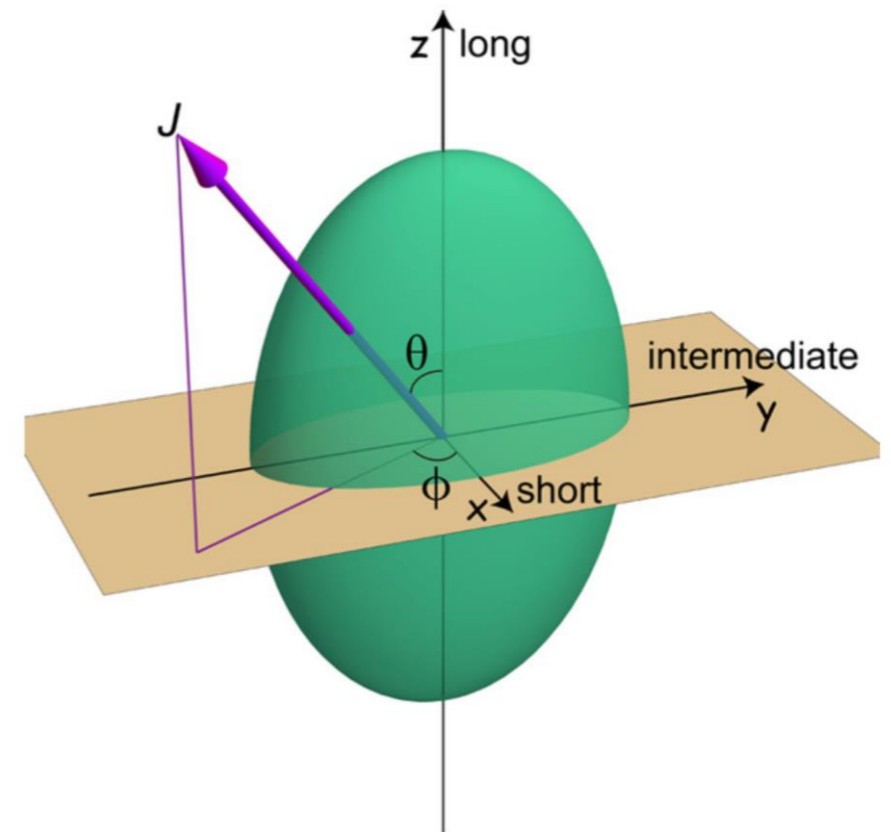
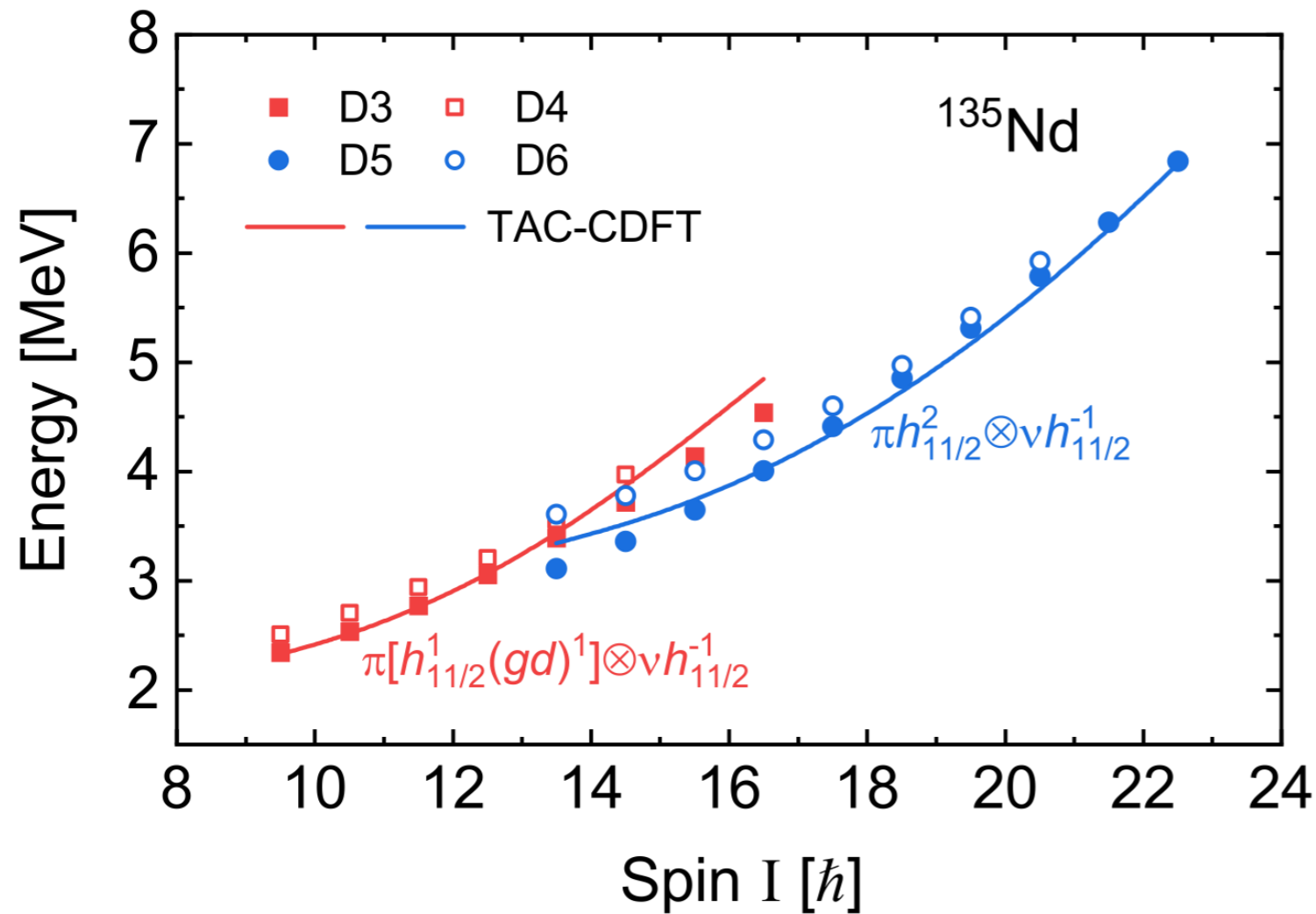
polar angle θ
azimuth angle φ



Chiral geometry is clearly seen.

Chiral rotation

The **lower bands** are well reproduced.



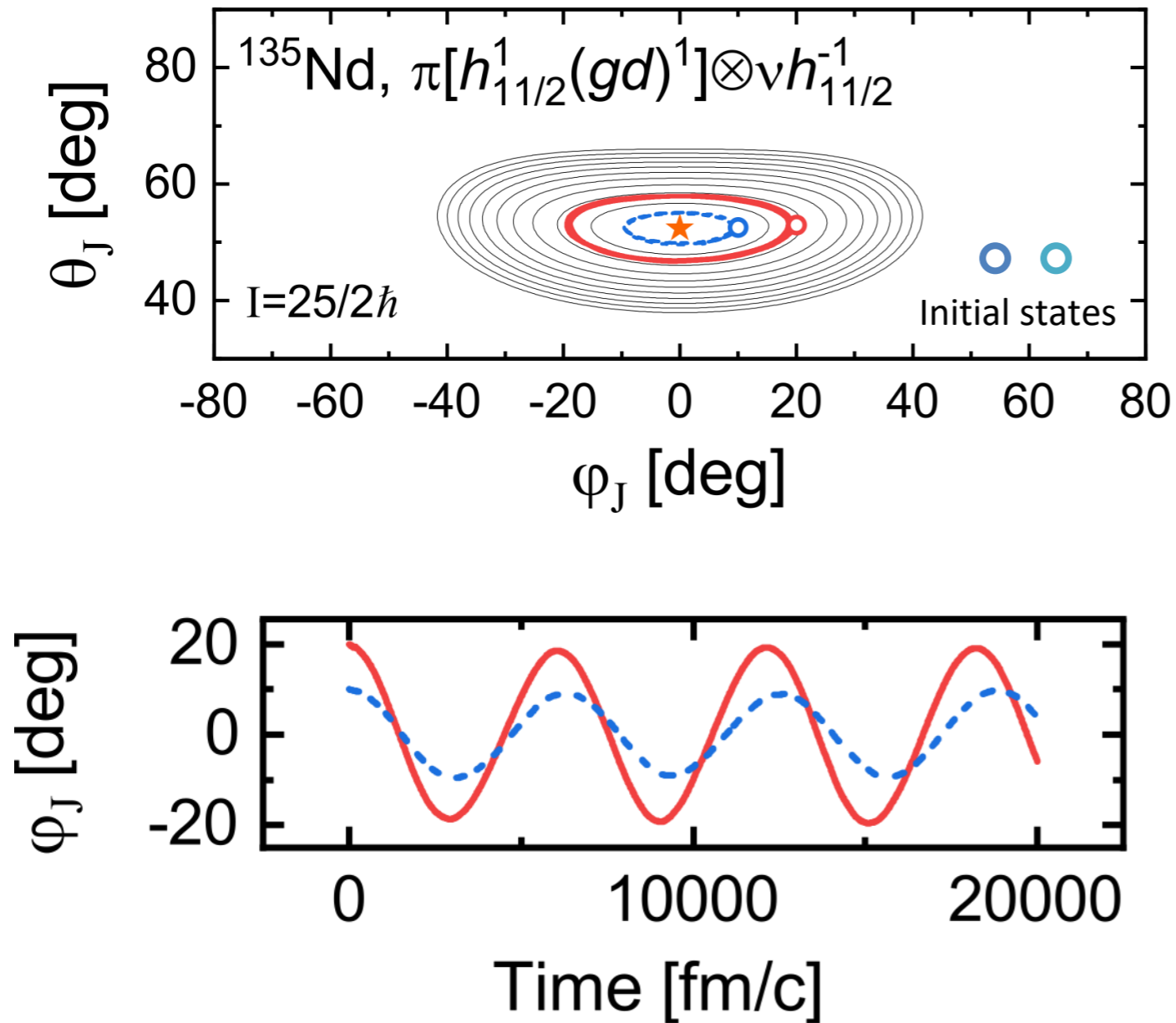
One does **NOT** get the **upper bands** !

Zhao, PLB **773**, 1 (2017)

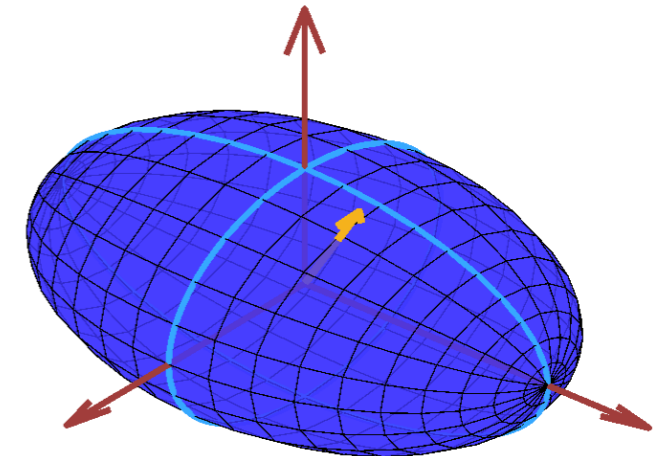
Ren, Zhao, Meng, PRC, **105**, L011301 (2022)

Chiral precession in triaxial nuclei

Chiral Precession across left and right sectors

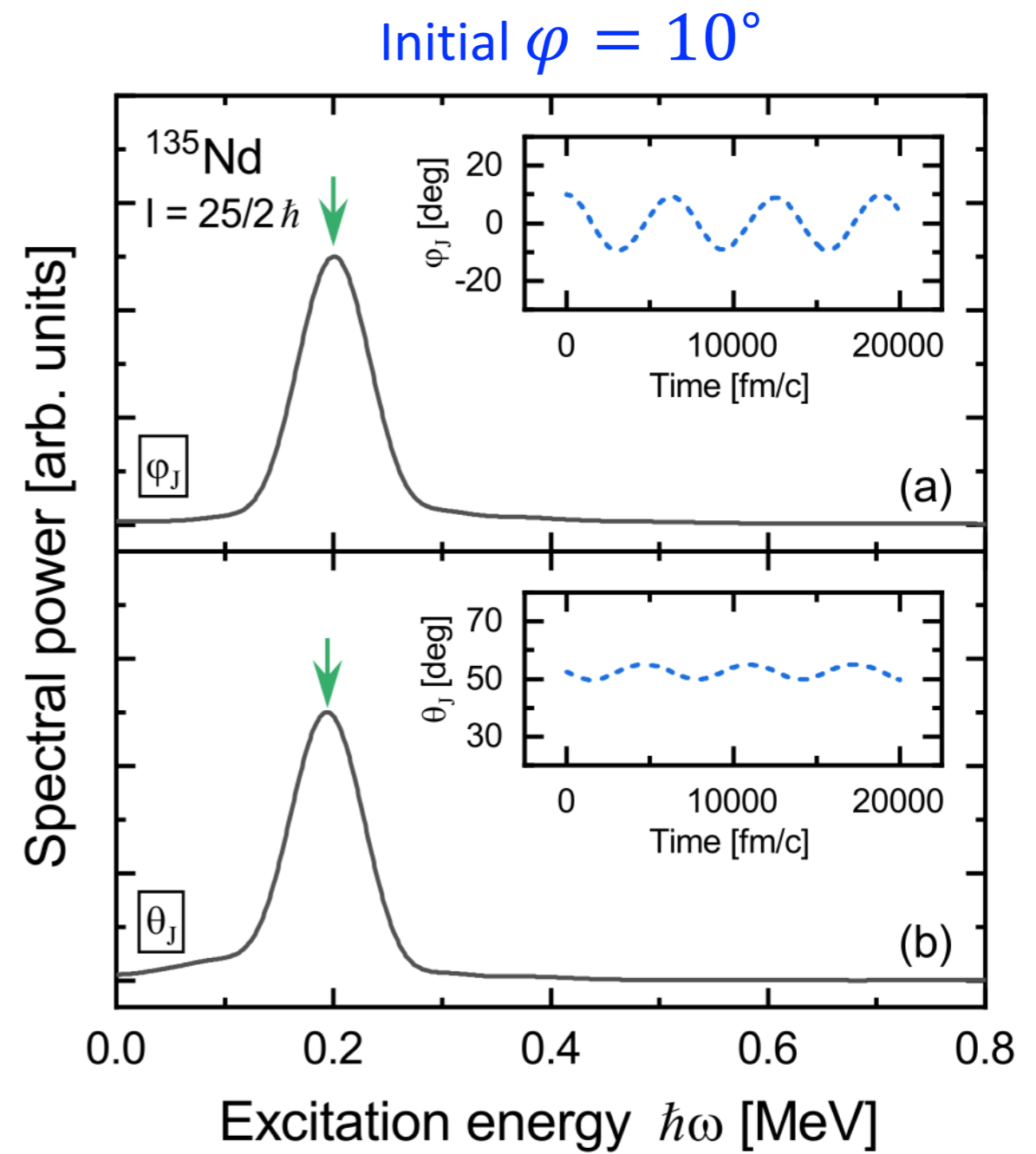
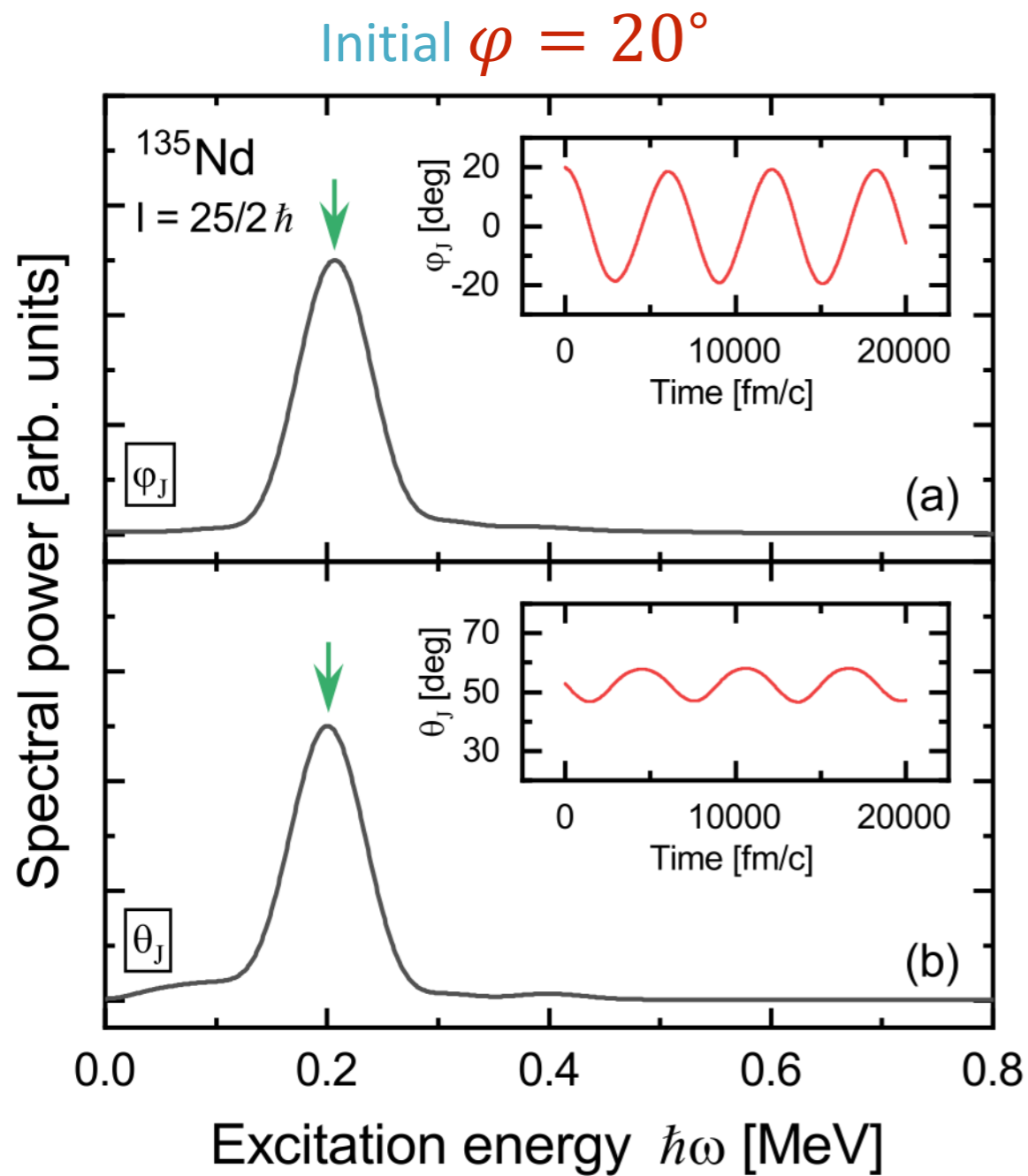


Body-fixed frame



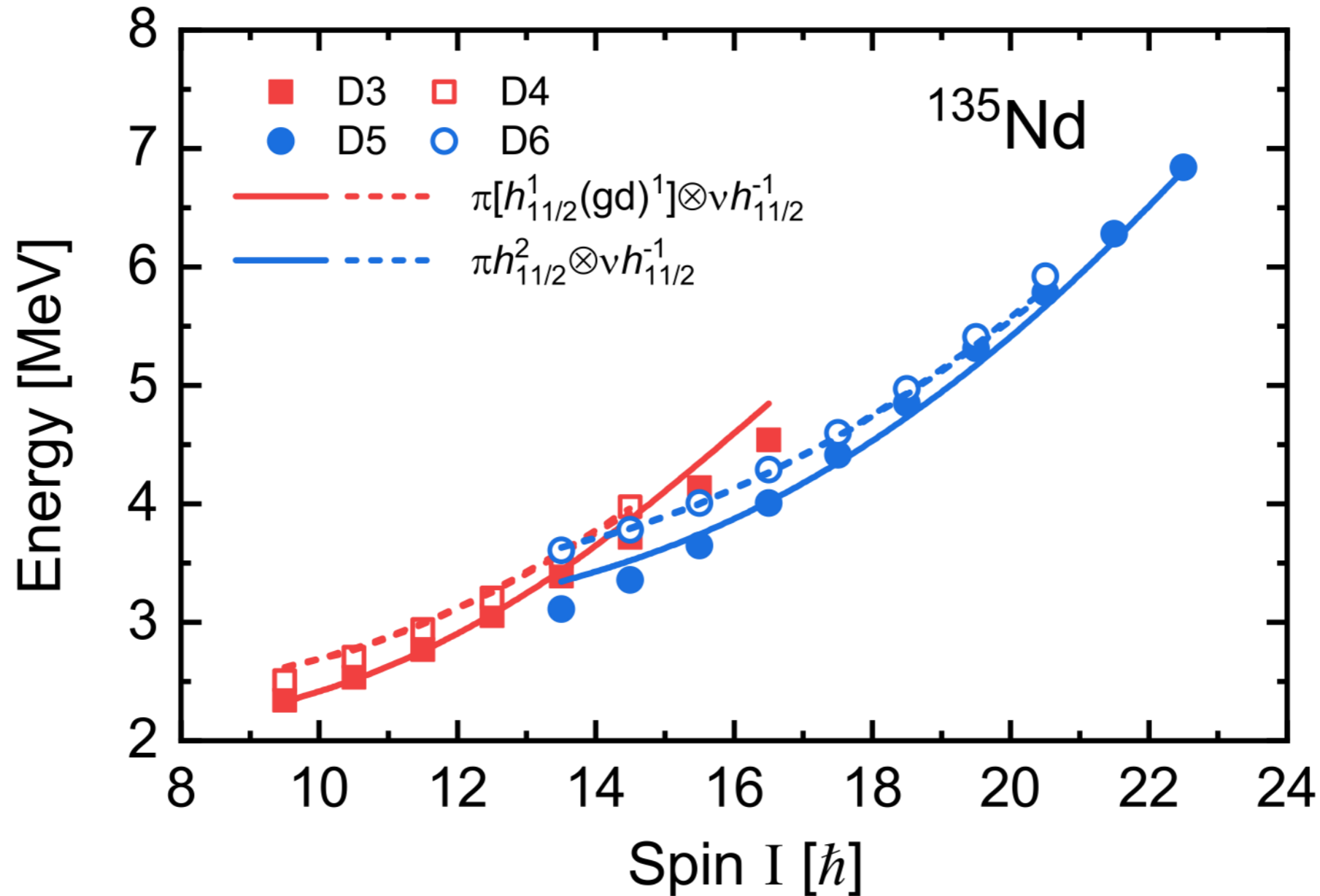
Precession periods are roughly **identical**.

Chiral excitation energies via Fourier analyses



Energy spectrum

Ren, Zhao, Meng, PRC, **105**, L011301 (2022)



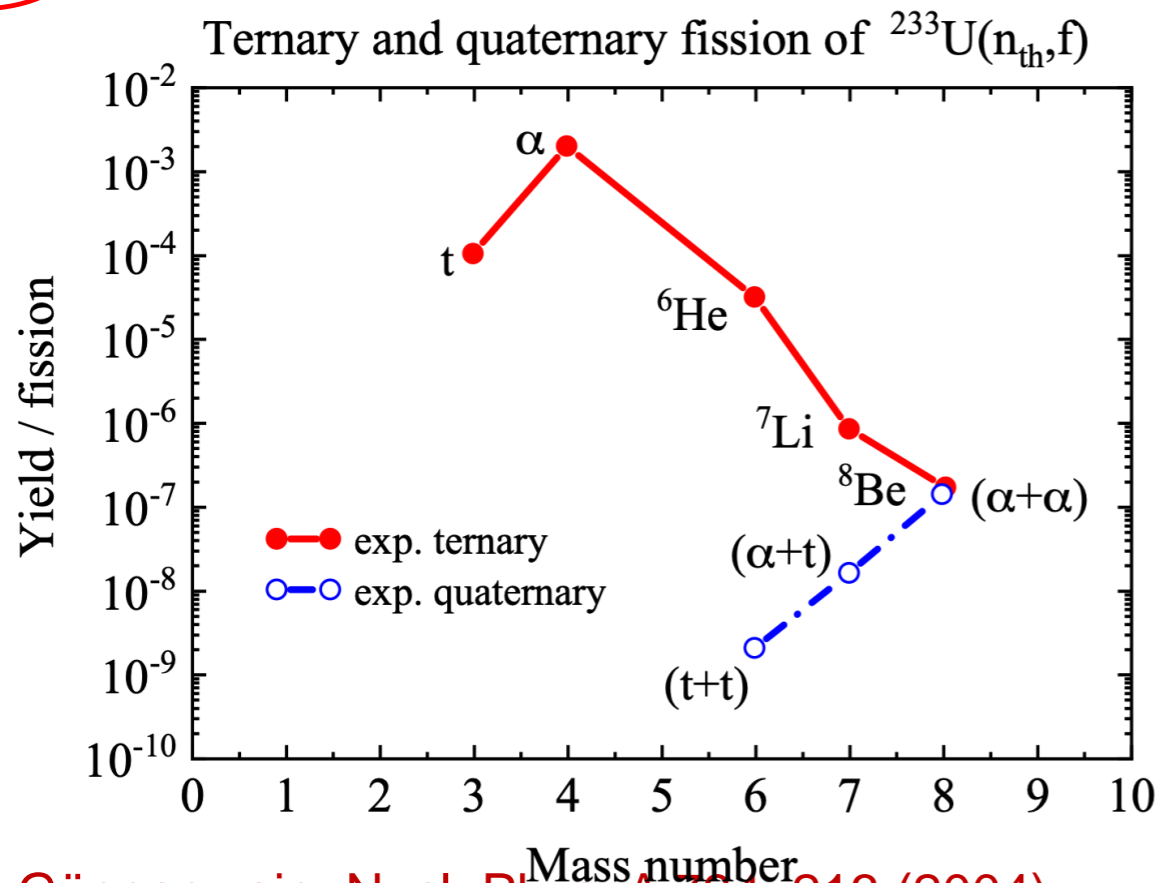
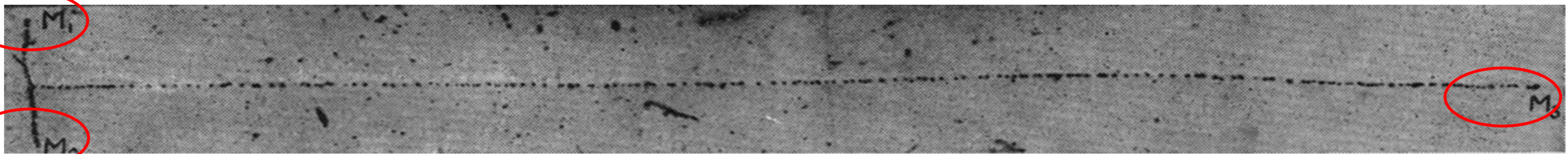
The first fully microscopic and self-consistent description for the *chiral twin bands* in the framework of DFTs.

Fission

Ternary and quaternary fission

- The ternary and quaternary fission were discovered by Tsien San-Tsiang, Ho Zah-Wei *et al.* in 1947.

Tsien San-Tsiang, Ho Zah-Wei, L. Vigneron, and R. Chastel, *Nature* 159, 773 (1947)



F. Gönnenwein, *Nucl. Phys. A* 734, 213 (2004)

Scission mechanism

- Statistical scission-point model
- Random neck rupture model
- Geometrical definitions
- Ratio between nuclear and Coul. Forces
- Quantum localization method
-

LETTERS TO THE EDITORS

The Editors do not hold themselves responsible for opinions expressed by their correspondents. No notice is taken of anonymous communications

Ternary and Quaternary Fission of Uranium Nuclei

AFTER our experimental proof of the existence of tripartition and quadripartition (ternary and quaternary fission) of U^{235} by means of photographic emulsion^{1,2}, a systematic study of the mass and kinetic energy of fission fragments has been made with the Ilford Nuclear Research C_2 plate. The experimental conditions were similar to those of previous work. C_2 plates soaked with 10 per cent solution of uranyl nitrate were bombarded with slow neutrons produced

Similar experiments have been made recently on the fission of uranium by fast neutrons. Short-range third light particles have been observed with the same frequency, but until now no long-range particle has been detected. Similar results have also been obtained in the case of fission of thorium by fast neutrons (observation made by the senior author in collaboration with Mrs. H. Faraggi).

TSIEN SAN-TSIANG
HO ZAH-WEI
L. VIGNERON
R. CHASTEL

Laboratoire de Chimie Nucléaire,
Collège de France, Paris.
March 15.

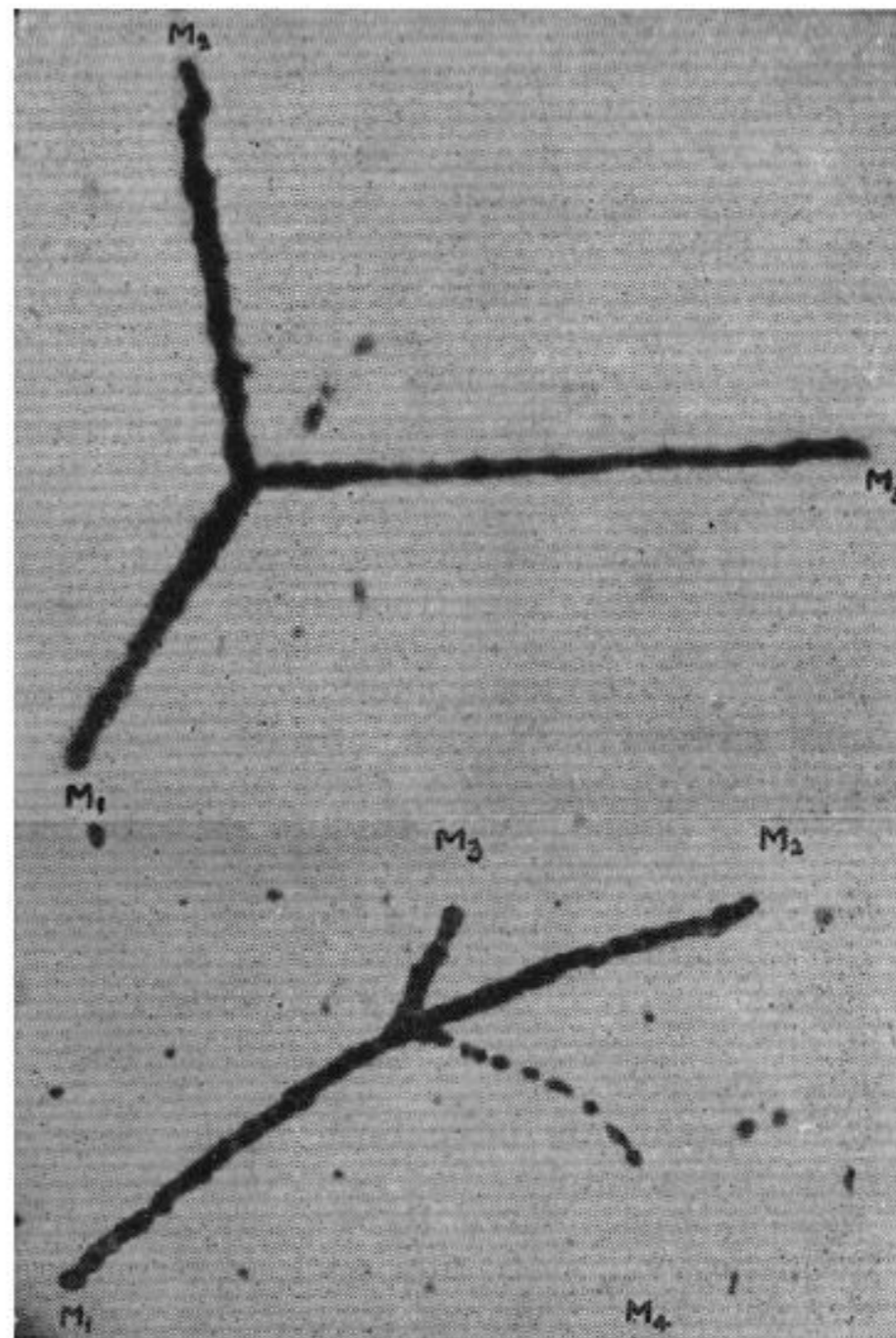


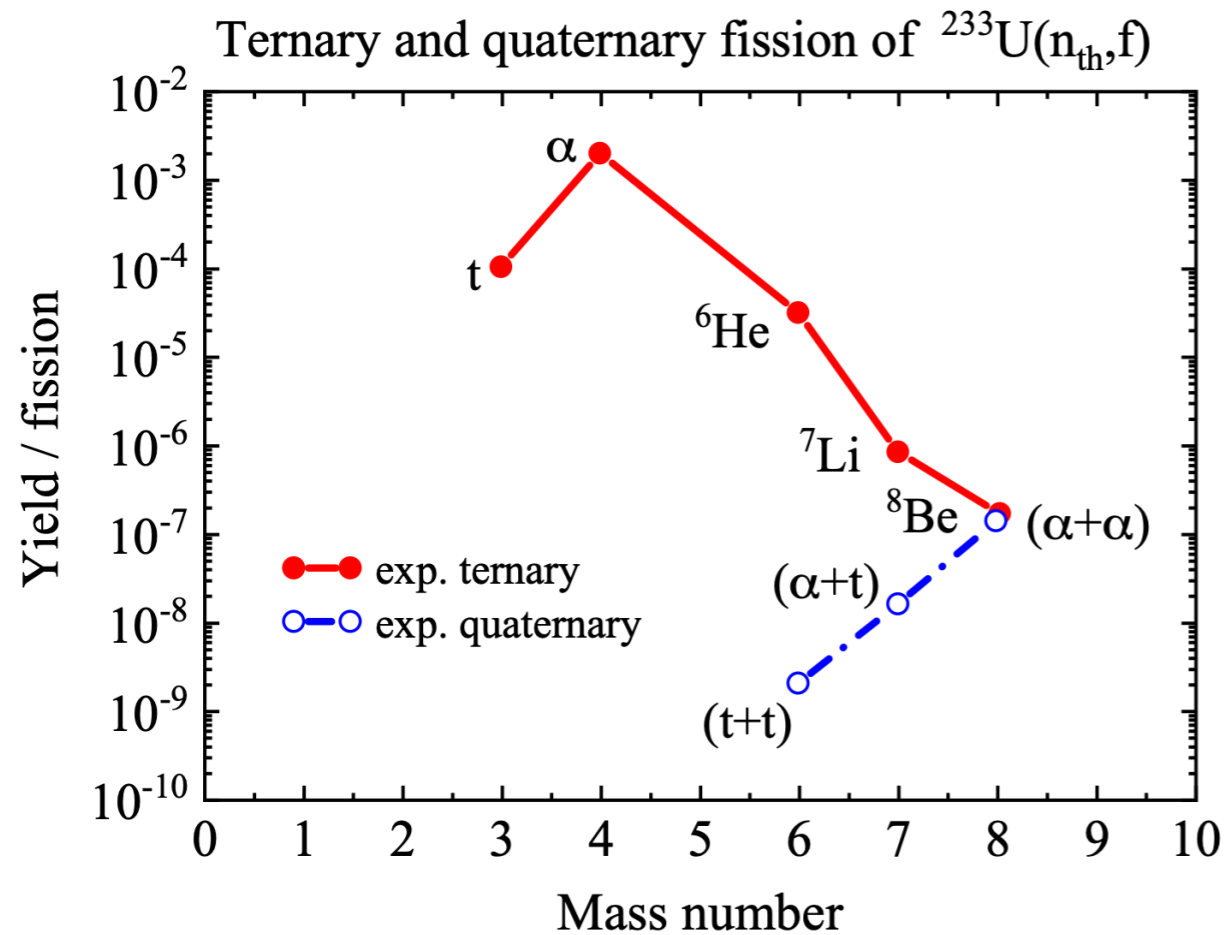
Fig. 2 (above); Fig. 3 (below)

Fig. 2. TRIPARTITION OF URANIUM NUCLEUS WITH THREE FRAGMENTS OF COMPARABLE MASSES: $m_1 = 127$; $m_2 = 77$; $m_3 = 32$

Fig. 3. QUADRIPARTITION OF URANIUM NUCLEUS

Ternary and quaternary fission

How are heavy nuclei ruptured?



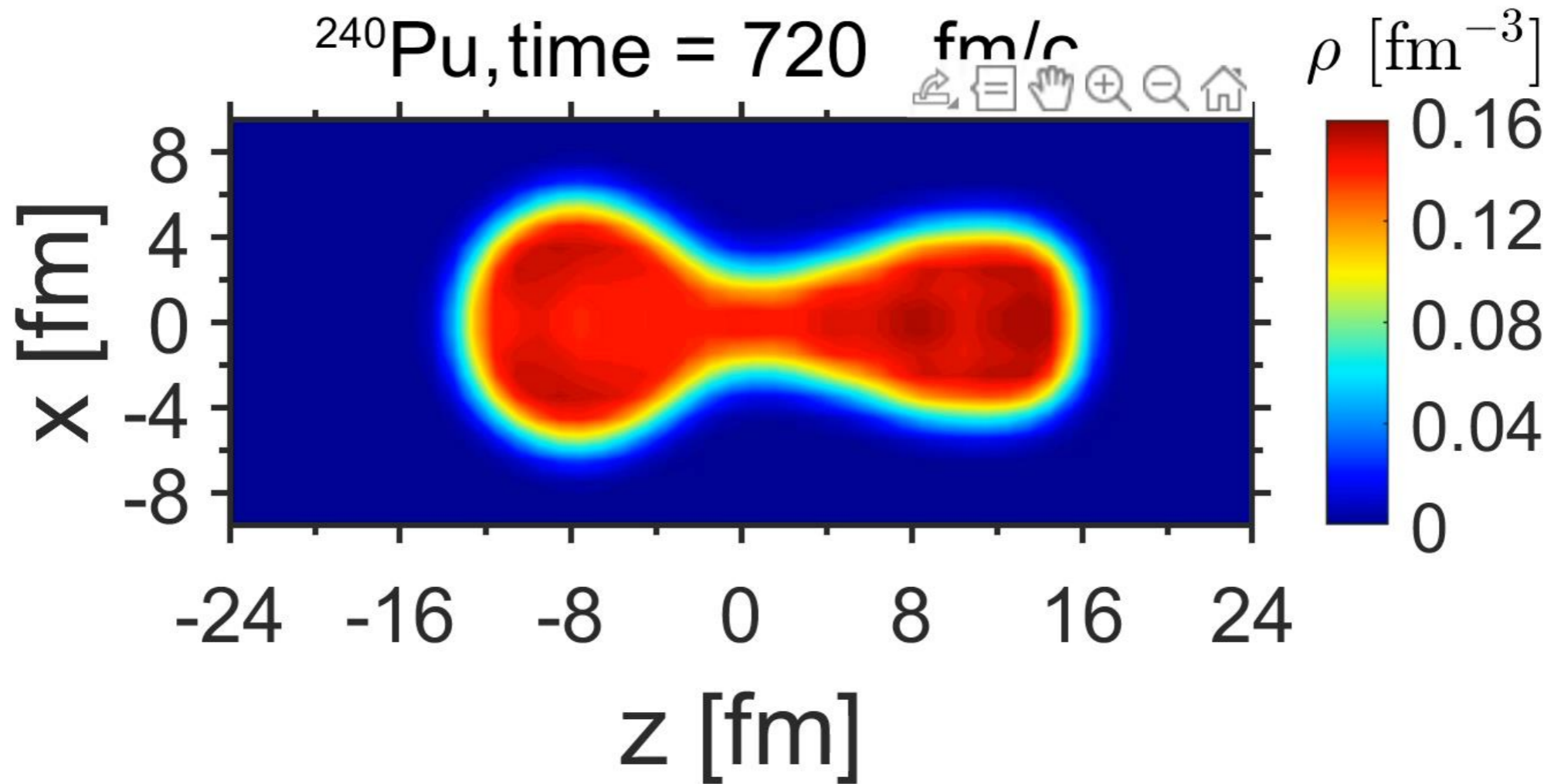
Scission mechanism

- Statistical scission-point model
- Random neck rupture model
- Geometrical definitions
- Ratio between nuclear and Coul. forces
- Quantum localization method
- ...

Tsien San-Tsiang, Ho Zah-Wei, L. Vigneron, and R. Chastel, *Nature (London)* 159, 773 (1947)

F. Gönnerwein, *NPA* 734 (2004) 213

^{240}Pu : Nuclear density



Ren, Vretenar, Nikšić, Zhao, Zhao, Meng, PRL 128, 172501 (2022)

Localization function

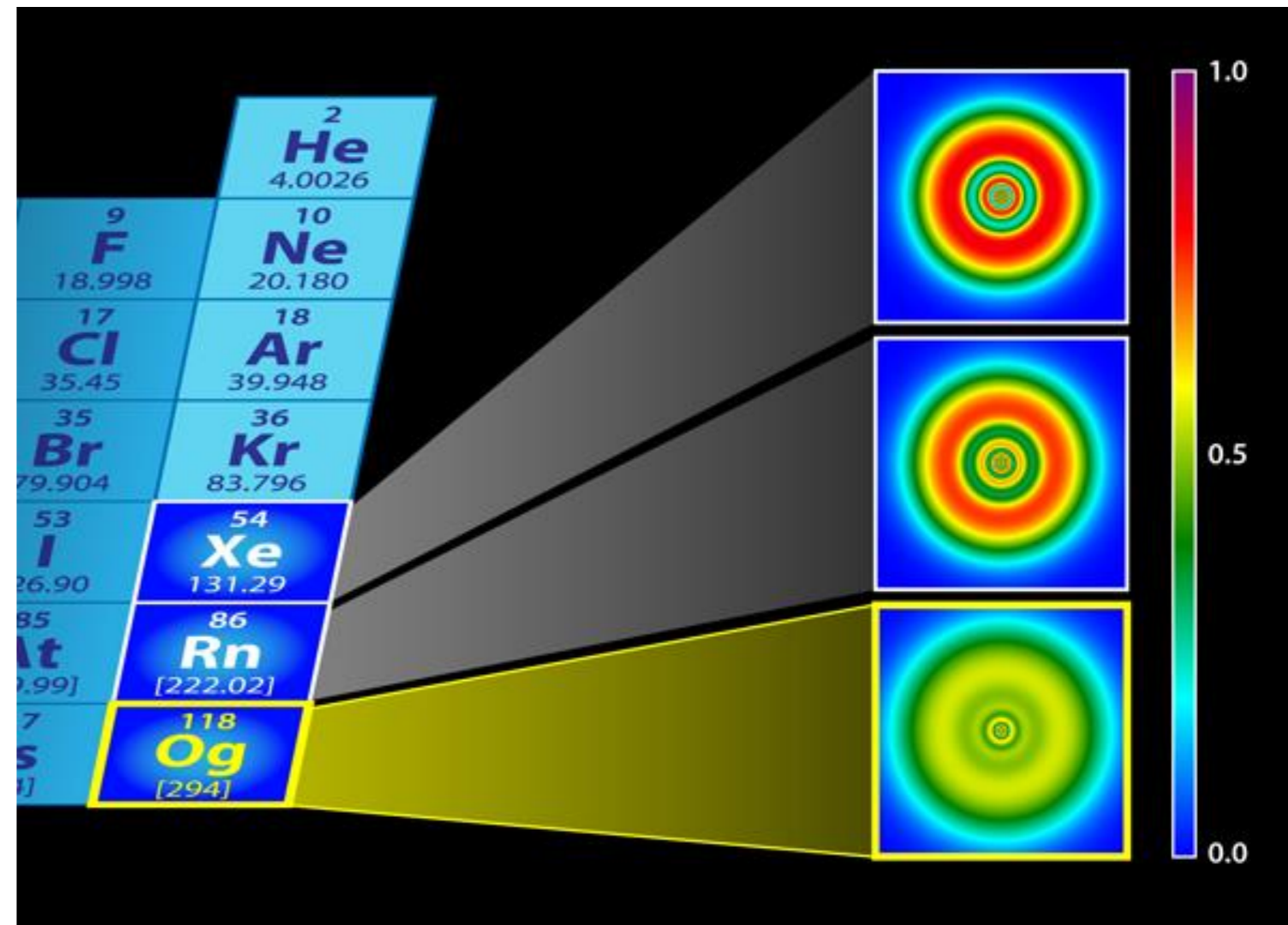
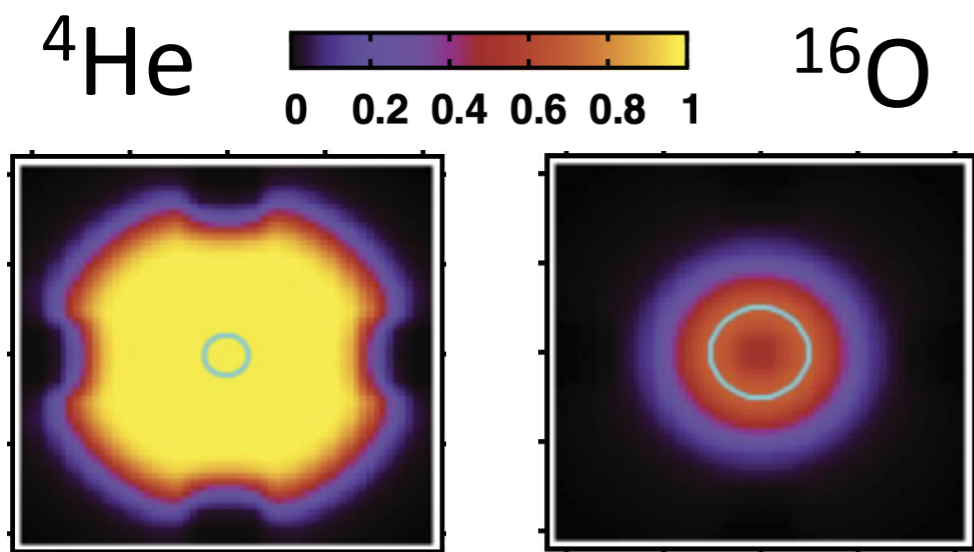
$$D_{q\sigma}(\mathbf{r}) = \left[\sum_{\alpha \in q} |\nabla \phi_{\alpha}(\mathbf{r}, \sigma)|^2 - \frac{|\sum_{\alpha \in q} \phi_{\alpha}^*(\mathbf{r}, \sigma) \nabla \phi_{\alpha}(\mathbf{r}, \sigma)|^2}{\rho_{q\sigma}(\mathbf{r})} \right]$$

the probability of finding two like-particles in the vicinity of each other \rightarrow Localization!

$$C_{q\sigma}(\mathbf{r}) = \left[1 + \left(\frac{D_{q\sigma}(\mathbf{r})}{\tau_{q\sigma}^{\text{TF}}(\mathbf{r})} \right)^2 \right]^{-1}$$

Becke and Edgecombe, J. Chem. Phys., **92**, 5397 (1990)

$C = 1/2$; Thomas-Fermi Gas
 $C = 1$; Highly Localized

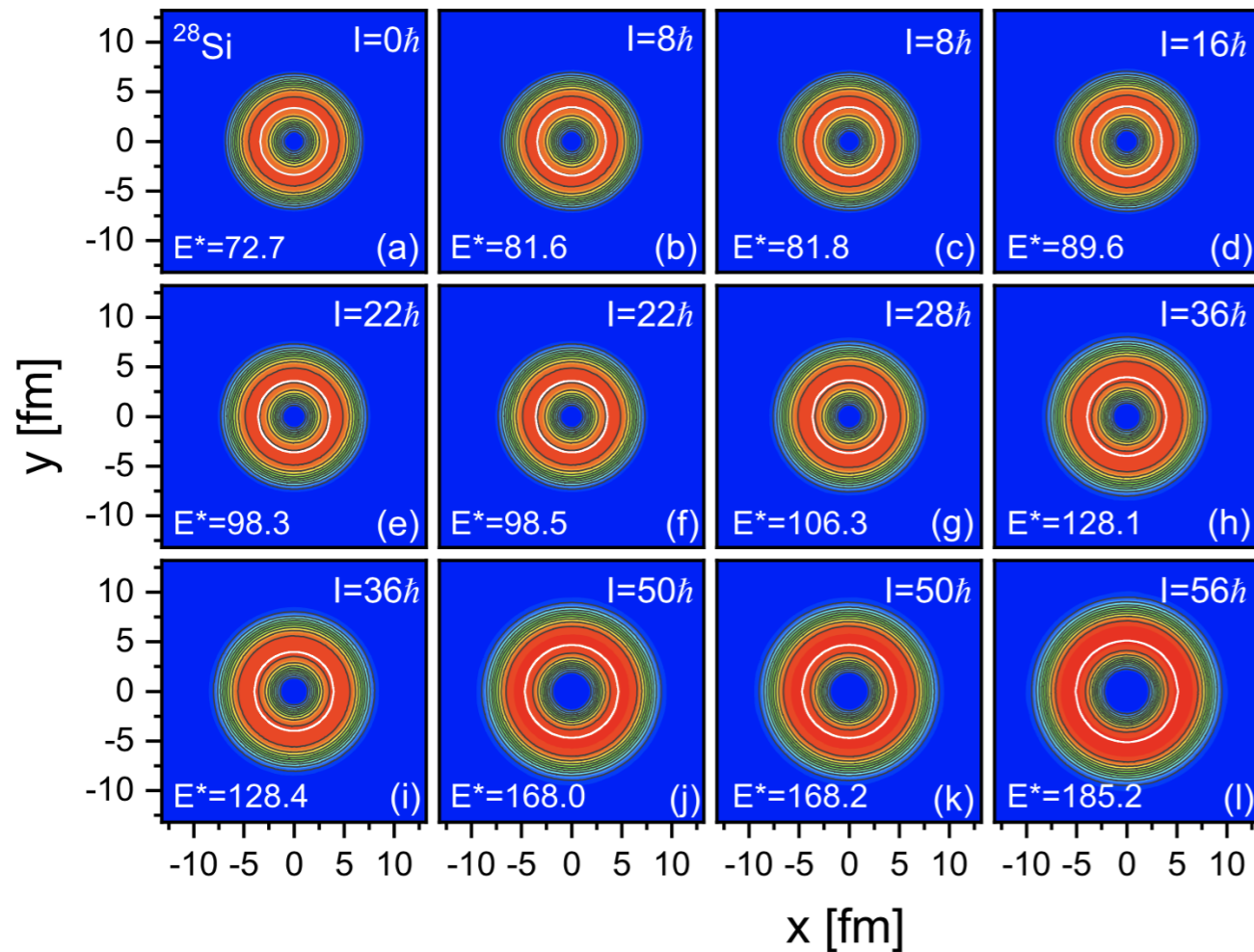


Reinhard, et al., Phys. Rev. C **83**, 034312 (2011)

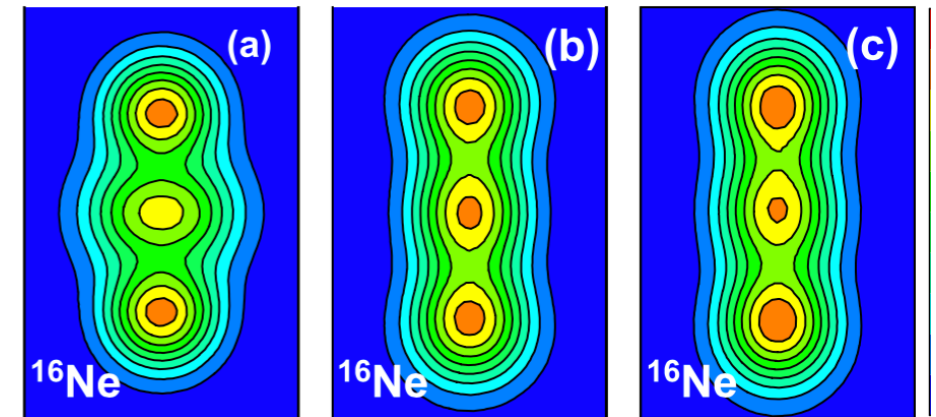
Jerabek et al., Phys. Rev. Lett. **120**, 053001 (2018)

Localization function

Toroidal states in ^{28}Si



molecular α -chain nuclei



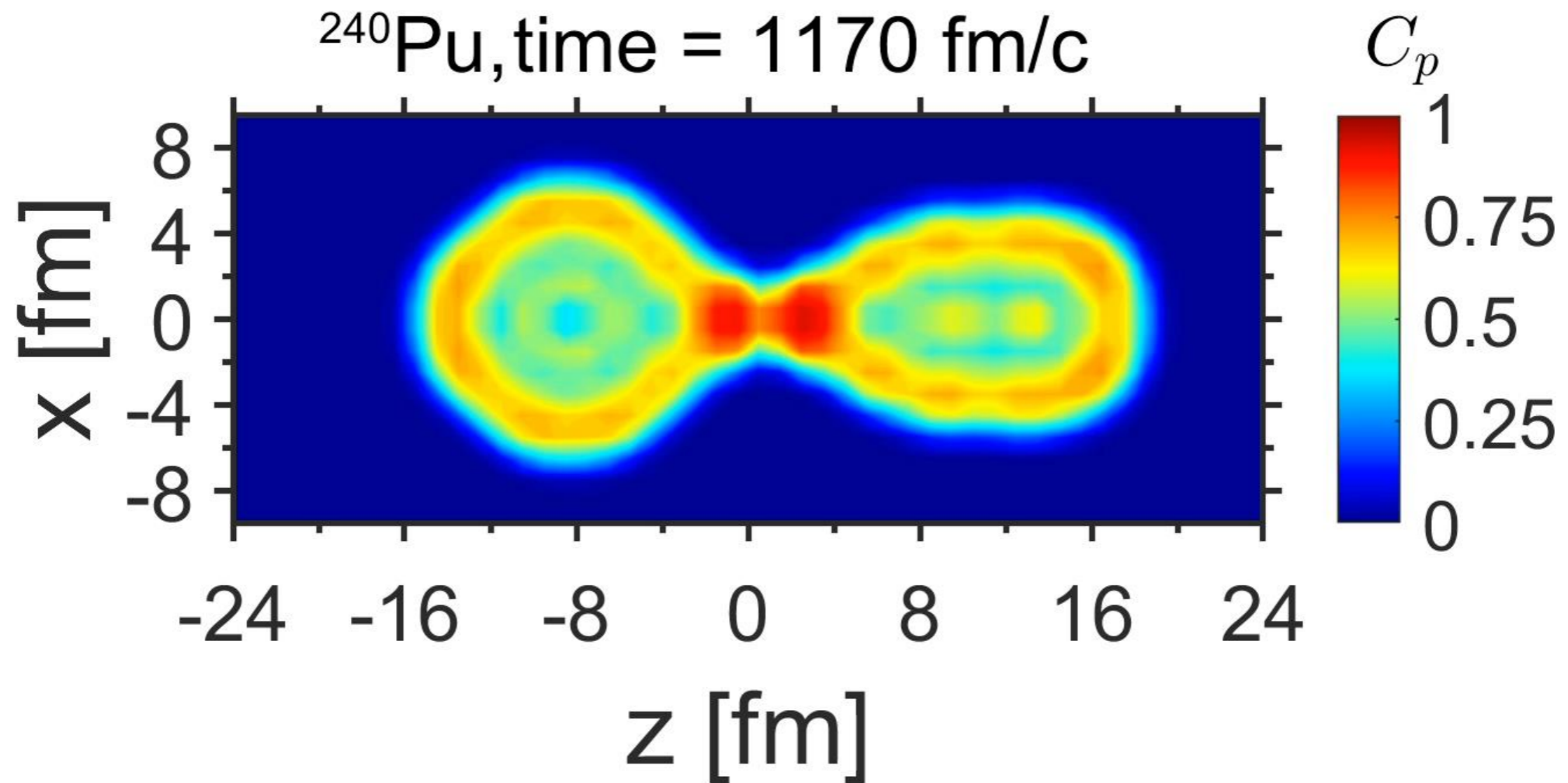
Zhang, Ren, Zhao, Vretenar, Nikšić, Meng, PRC, 105, 024322 (2022)

Cao et al., Phys. Rev. C, **99**, 014606 (2019)

Ren, Zhao, Zhang, Meng, Nucl. Phys. A, **996**, 121696, (2020)

^{240}Pu : Localization function

PRL 128, 172501 (2022)



Scission point between two alpha-like particles

Toward a comprehensive description of nuclear fission

TDGCM+GOA

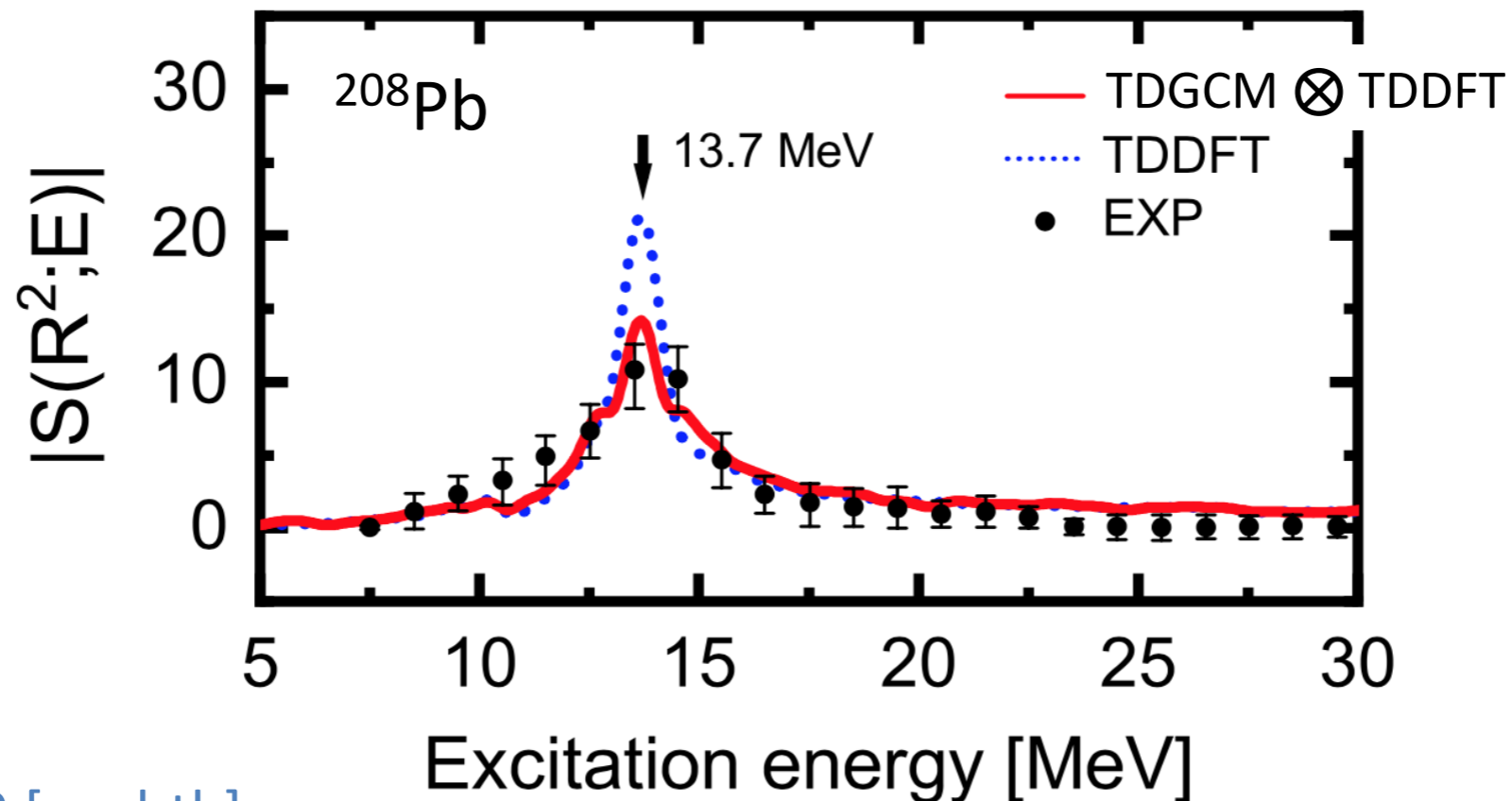
- collective degrees of freedom
- quantum correlations (beyond MF)
- no dissipation (adiabatic)

TD-CDFT

- independent nucleons
- no quantum correlations (MF)
- dissipative dynamics (nonadiabatic)

“TDGCM \otimes TDDFT”

$$|\Psi(t)\rangle = \int da f(a, t) |\Phi(a, t)\rangle$$

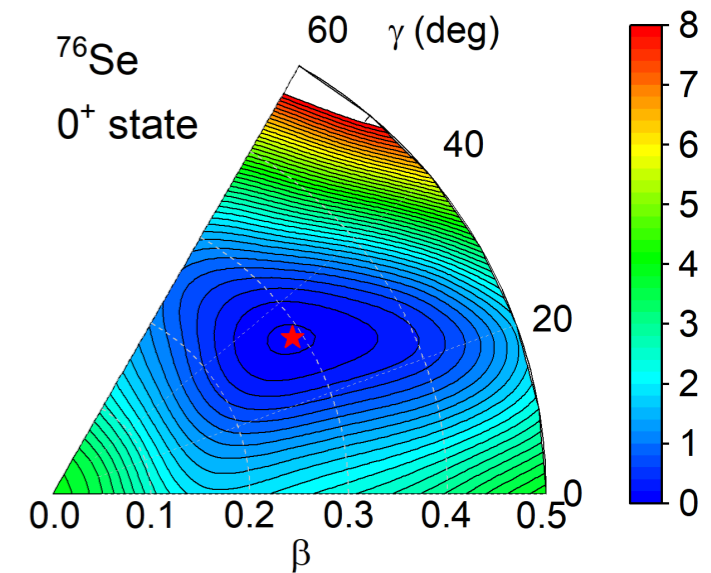
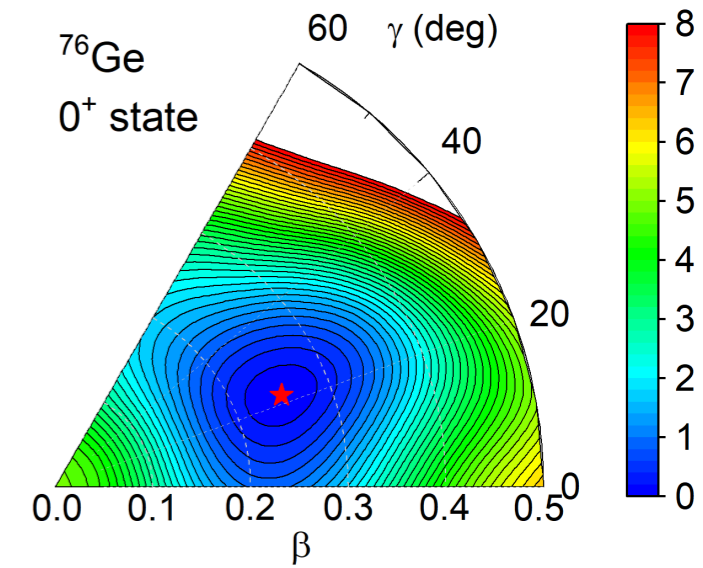
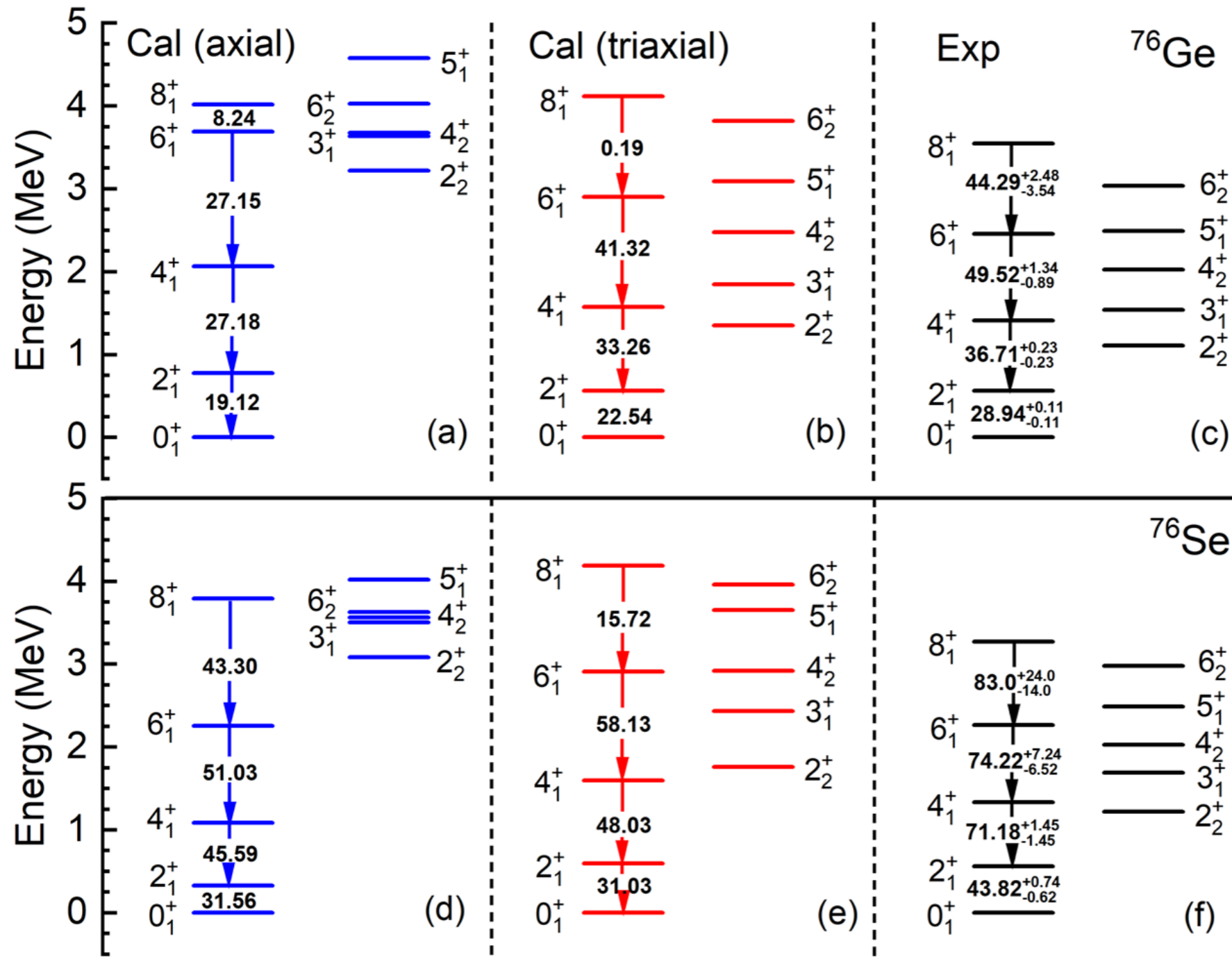


ReCD theory

Relativistic Configuration-interaction Density functional theory

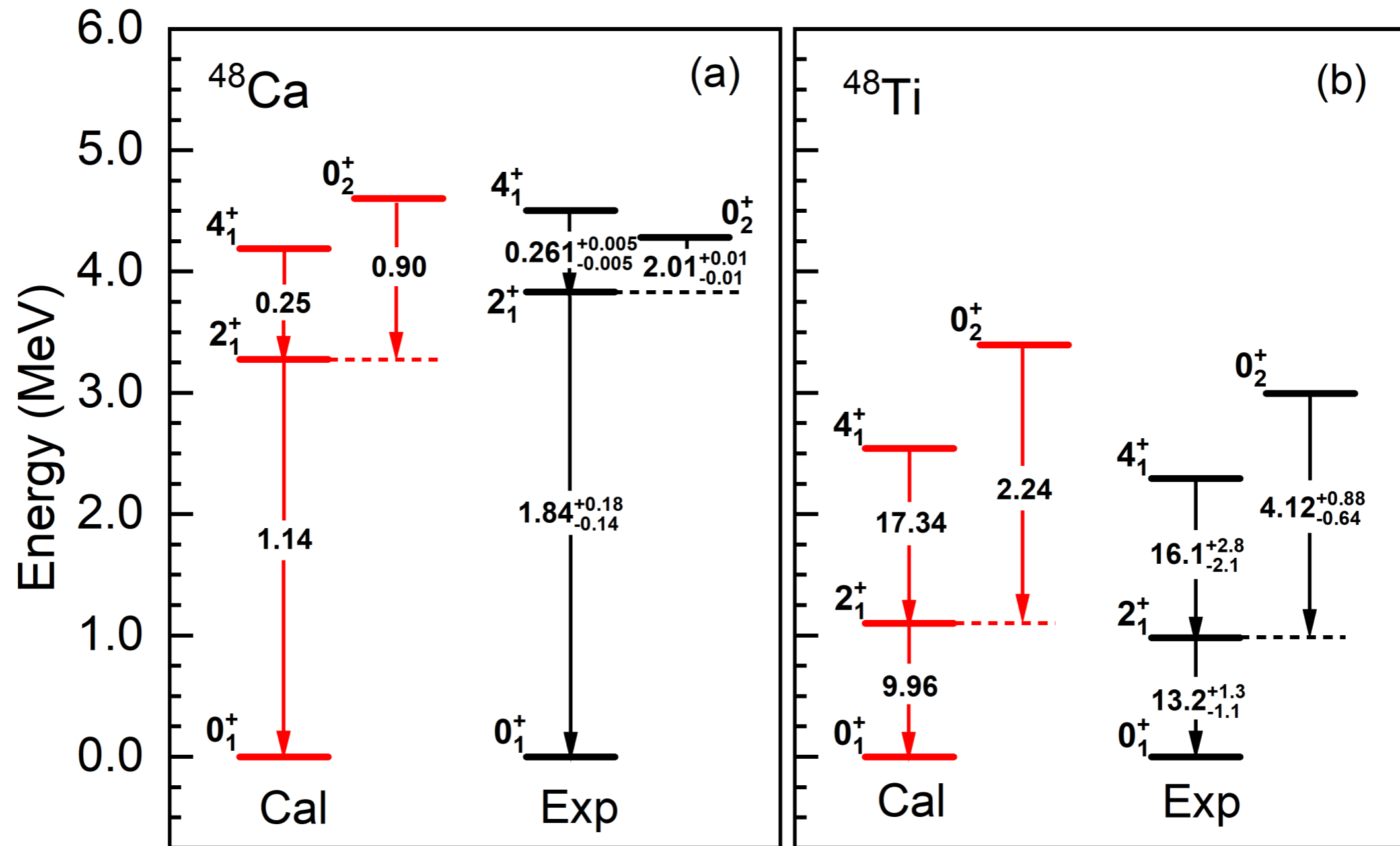
- 1. A self-consistent relativistic DFT calculation**
- 2. Construction of intrinsic configuration space**
- 3. Angular momentum projection**
- 4. Shell-model diagonalization**

ReCD: nuclear spectroscopy



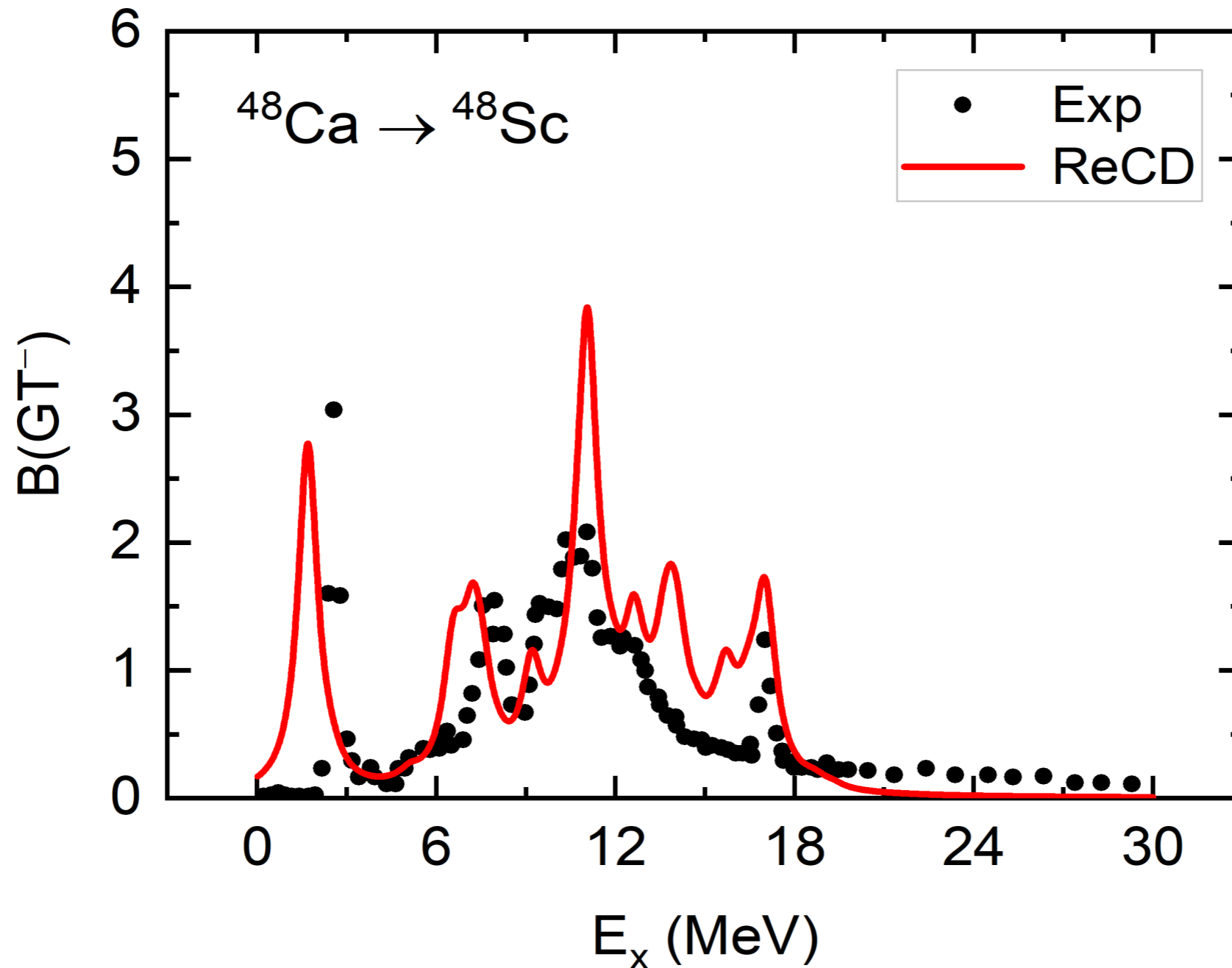
□ The energy spectra and the E2 transition probabilities are reproduced satisfactorily

ReCD: nuclear spectroscopy



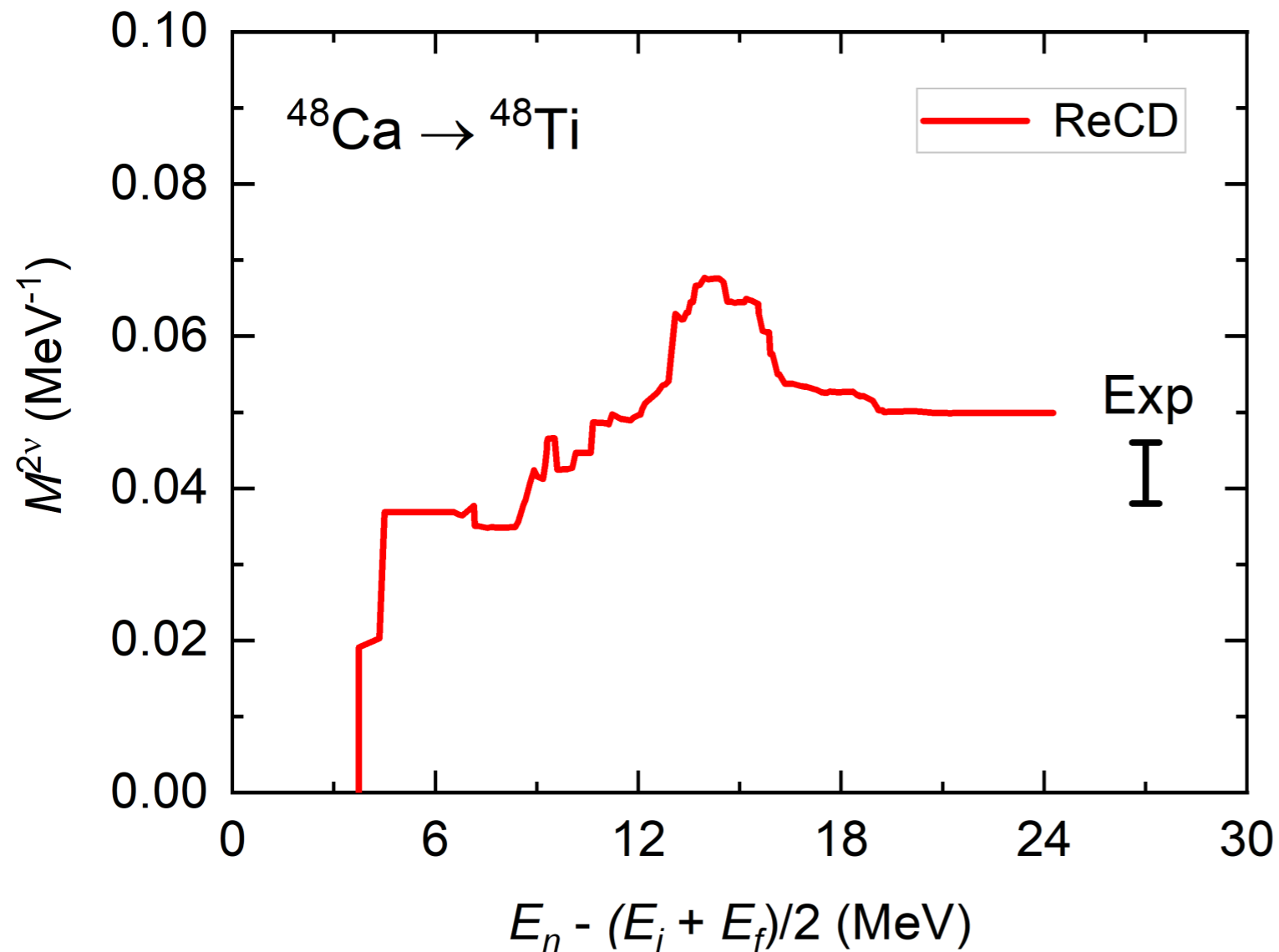
□ The energy spectra and the E2 transition probabilities are reproduced satisfactorily

ReCD: spin-isospin excitations



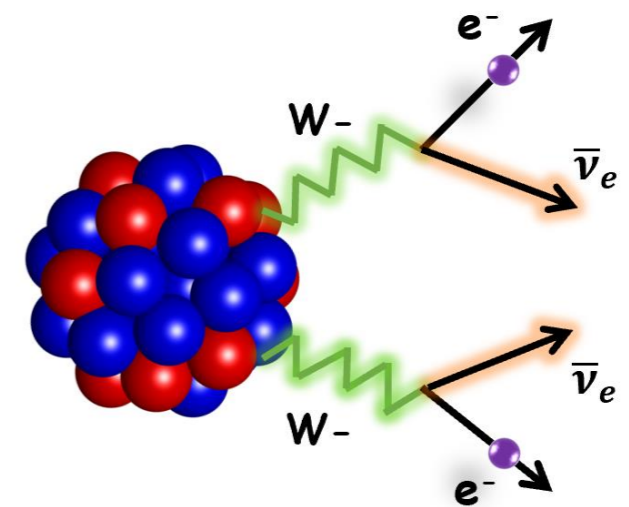
- The Gamow-Teller strength distribution around the **whole energy region** is reproduced satisfactorily

ReCD: nuclear weak decay



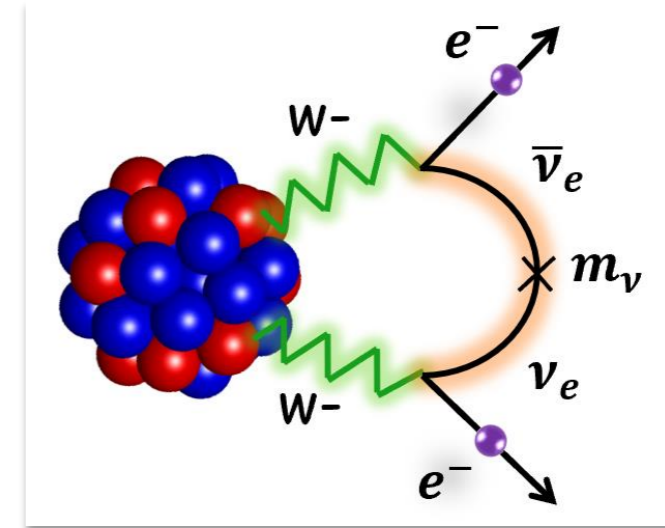
□ The $2\nu\beta\beta$ decay NME can be reproduced without introducing the quenching factor to the axial-vector coupling constant g_A

$$M^{2\nu} = q^2 \sum_n \frac{\langle f | \sum_a \hat{\sigma}_a \tau_a^+ | 1_n^+ \rangle \langle 1_n^+ | \sum_b \hat{\sigma}_b \tau_b^+ | i \rangle}{E_n - 1/2(E_i + E_f)}$$



Nuclear $0\nu\beta\beta$ decay

- ❑ Violation of lepton number
- ❑ Majorana nature of neutrinos
- ❑ Matter dominance in the Universe
- ❑ Neutrino mass scale and hierarchy



Avignone, Elliott, Engel, Rev. Mod. Phys. 80, 481 (2008) $(A, Z) \rightarrow (A, Z + 2) + e^- + e^-$

Isotopes	$T_{1/2}^{0\nu}$ (yr)	Collaborations
^{48}Ca	$> 5.8 \times 10^{22}$	ELEGANT VI
^{76}Ge	$> 1.8 \times 10^{26}$	GERDA, MAJORANA, CDEX
^{82}Se	$> 3.5 \times 10^{24}$	CUPID-0, N ν DEx
^{100}Mo	$> 1.5 \times 10^{24}$	CUPID-Mo
^{130}Te	$> 3.2 \times 10^{25}$	CUORE
^{136}Xe	$> 2.3 \times 10^{26}$	KamLAND-Zen, EXO-200, PandaX
^{150}Nd	$> 2.0 \times 10^{22}$	NEMO-3

- ✓ No $0\nu\beta\beta$ -decay signal has been observed so far.
- ✓ Current limit on the decay half-life ranges from 10^{22} yr to 10^{26} yr.

Agostini, Benato, Detwiler, Menéndez, Vissani, Rev. Mod. Phys. 95, 025002 (2023)

$0\nu\beta\beta$ -decay nuclear matrix elements

□ $0\nu\beta\beta$ decay rates (inversion of the decay half-life) :

$$[T_{1/2}^{0\nu}]^{-1} = G^{0\nu}(Q, Z) |M^{0\nu}|^2 f(\Lambda), \quad f(\Lambda) \Rightarrow \langle m_{\beta\beta} \rangle^2$$

Nuclear matrix element $M^{0\nu} = \langle \Psi_f | \hat{O}^{0\nu} | \Psi_i \rangle$

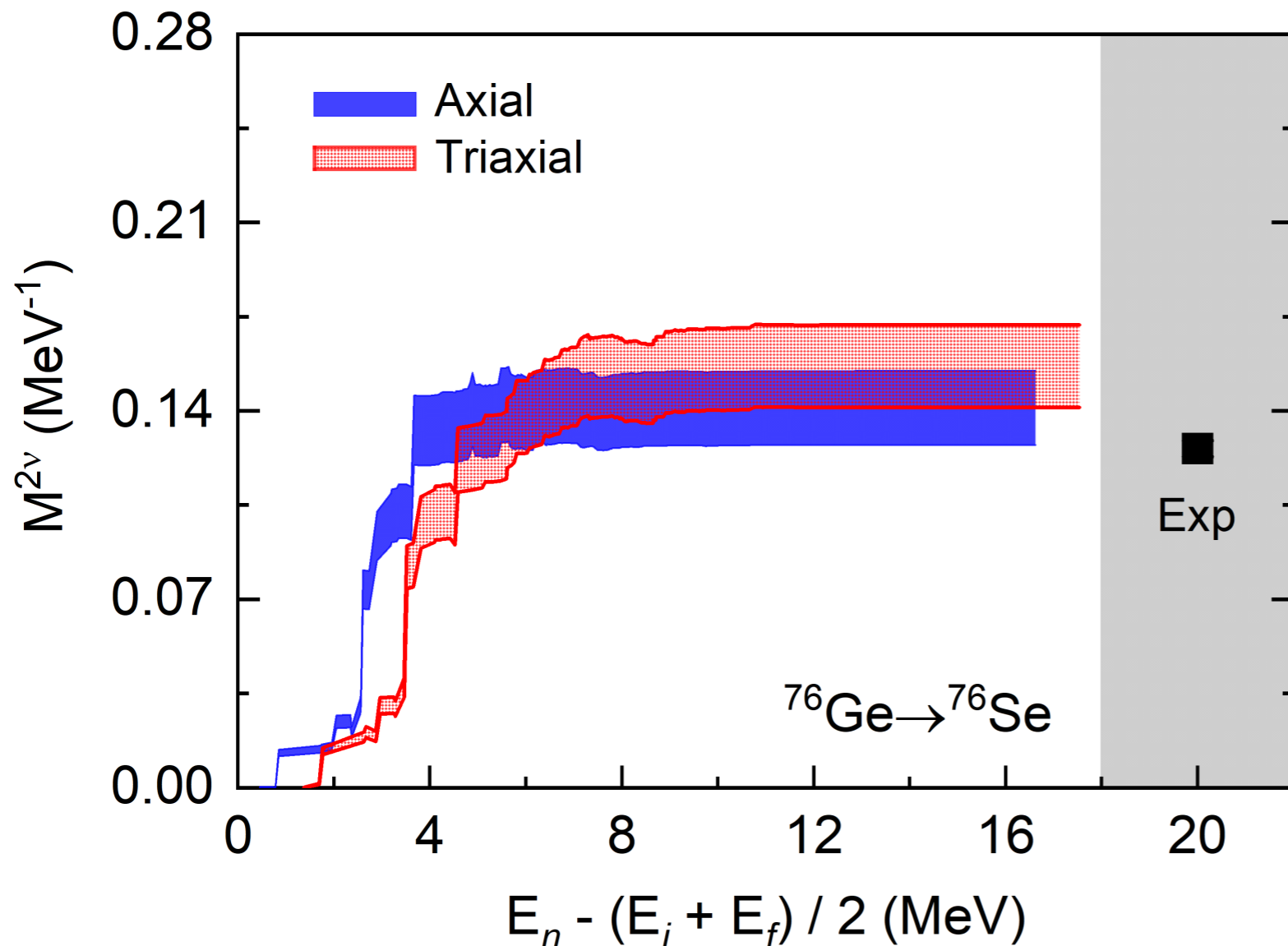
Theoretical studies on $\nu\beta\beta$ decay build a bridge between **data** and **underlying new physics**

PHYSICAL REVIEW LETTERS **123**, 102501 (2019)

Evidence for Rigid Triaxial Deformation in ^{76}Ge from a Model-Independent Analysis

A. D. Ayangeakaa^{1,*}, R. V. F. Janssens^{2,3,†}, S. Zhu^{4,‡}, D. Little^{2,3}, J. Henderson⁵, C. Y. Wu⁵, D. J. Hartley¹, M. Albers⁴, K. Auranen⁴, B. Bucher^{5,§}, M. P. Carpenter⁴, P. Chowdhury⁶, D. Cline⁷, H. L. Crawford⁸, P. Fallon⁸, A. M. Forney⁹, A. Gade^{10,11}, A. B. Hayes⁷, F. G. Kondev⁴, Krishichayan^{3,12}, T. Lauritsen⁴, J. Li⁴, A. O. Macchiavelli⁸, D. Rhodes^{10,11}, D. Seweryniak⁴, S. M. Stolze⁴, W. B. Walters⁹ and J. Wu⁴

NME of the $2\nu\beta\beta$ decay

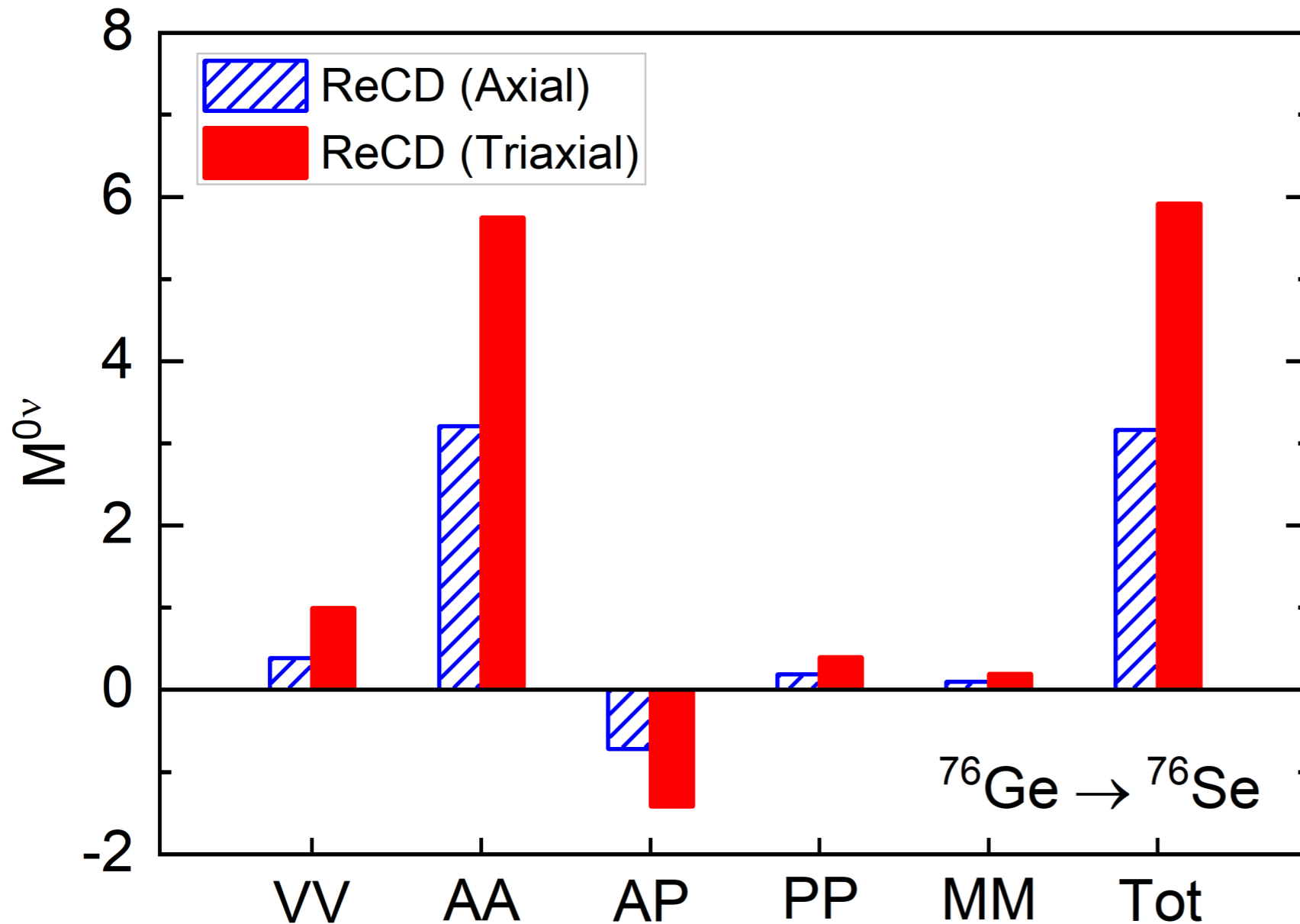


□ All possible odd-odd intermediate 1^+ states needs to be considered (around 200 1^+ states are included).

□ The two-body currents are not considered \Rightarrow quenching factor q ranging from **0.68** to **0.77** are adopted in our calculations.

$$M^{2\nu} = q^2 \sum_n \frac{\langle \Psi_f | \sum_a \sigma_a \tau_a^+ | \Psi_{1n^+} \rangle \langle \Psi_{1n^+} | \sum_b \sigma_b \tau_b^+ | \Psi^i \rangle}{E_n - (E_i + E_f) / 2}$$

NME: triaxial effects



$$[T_{1/2}^{0\nu}]^{-1} = G^{0\nu}(Q_{\beta\beta}, Z) |M^{0\nu}|^2 \left| \frac{\langle m_{\beta\beta} \rangle}{m_e} \right|^2$$

half-life decreases by
a factor of four

$$\propto \varepsilon \frac{i.a.}{A} Mt$$

Less detector materials
are needed

Y. K. Wang, P. W. Zhao, J. Meng, Science Bulletin, accepted

□ Triaxial deformation enhances the $0\nu\beta\beta$ -decay NME by a factor around two.

Double Gamow-Teller transition

- DGT cross section can be factorized as: **reaction factor**, and **DGT-transition NMEs** between initial state and the final state

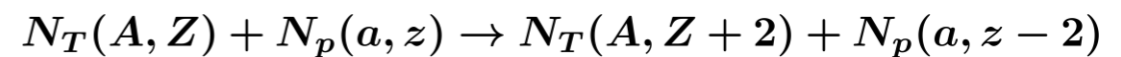
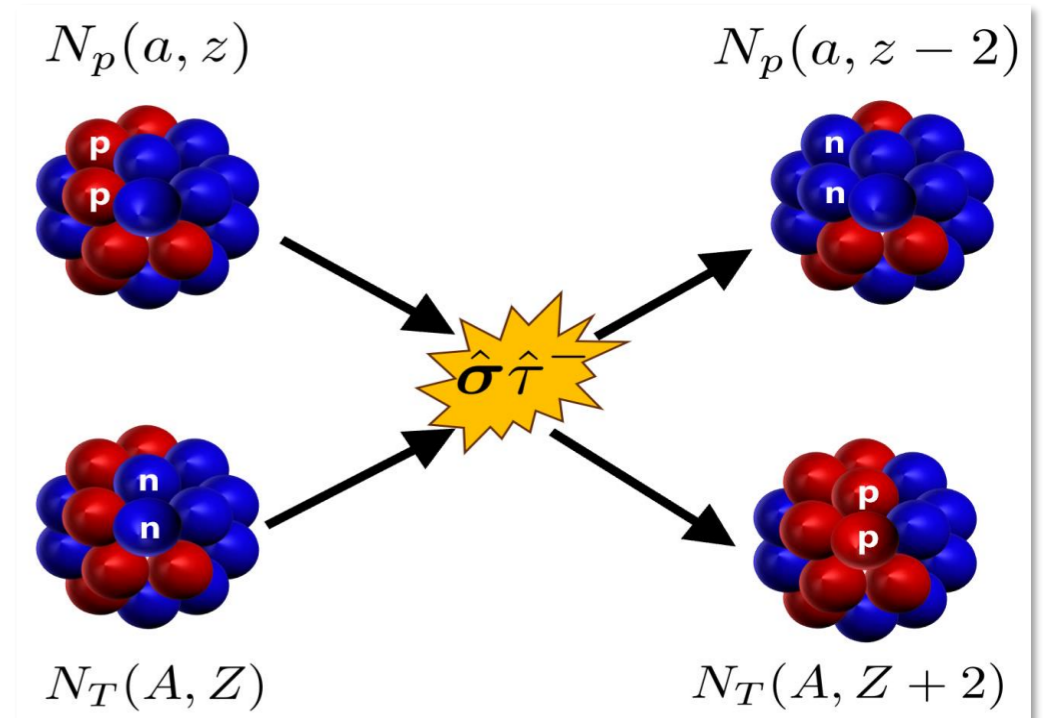
Santopinto et al., PRC 98, 061601(R) (2018)

Therefore, DGT-transition NMEs can be fixed by the **experimental cross section**

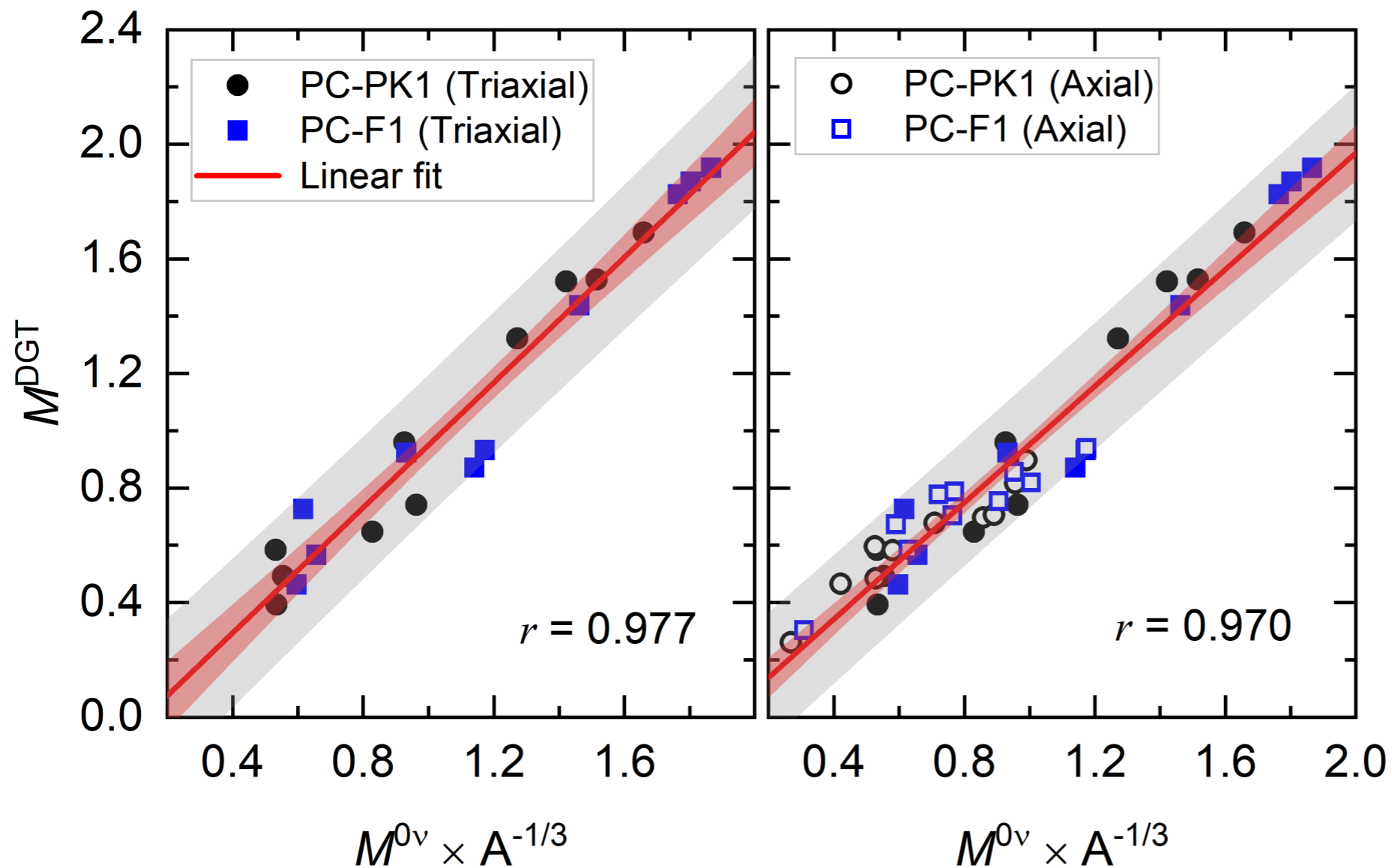
- Same initial and final states + similar decay operator \Rightarrow correlation between DGT transition and $0\nu\beta\beta$ decay?

Rodríguez et al., PLB 719, 174 (2013), Cappuzzello et al., EPJA 51, 145 (2015)

Constraining the $0\nu\beta\beta$ decay NME by DGT transitions!



Correlation between $0\nu\beta\beta$ decay and DGT transition



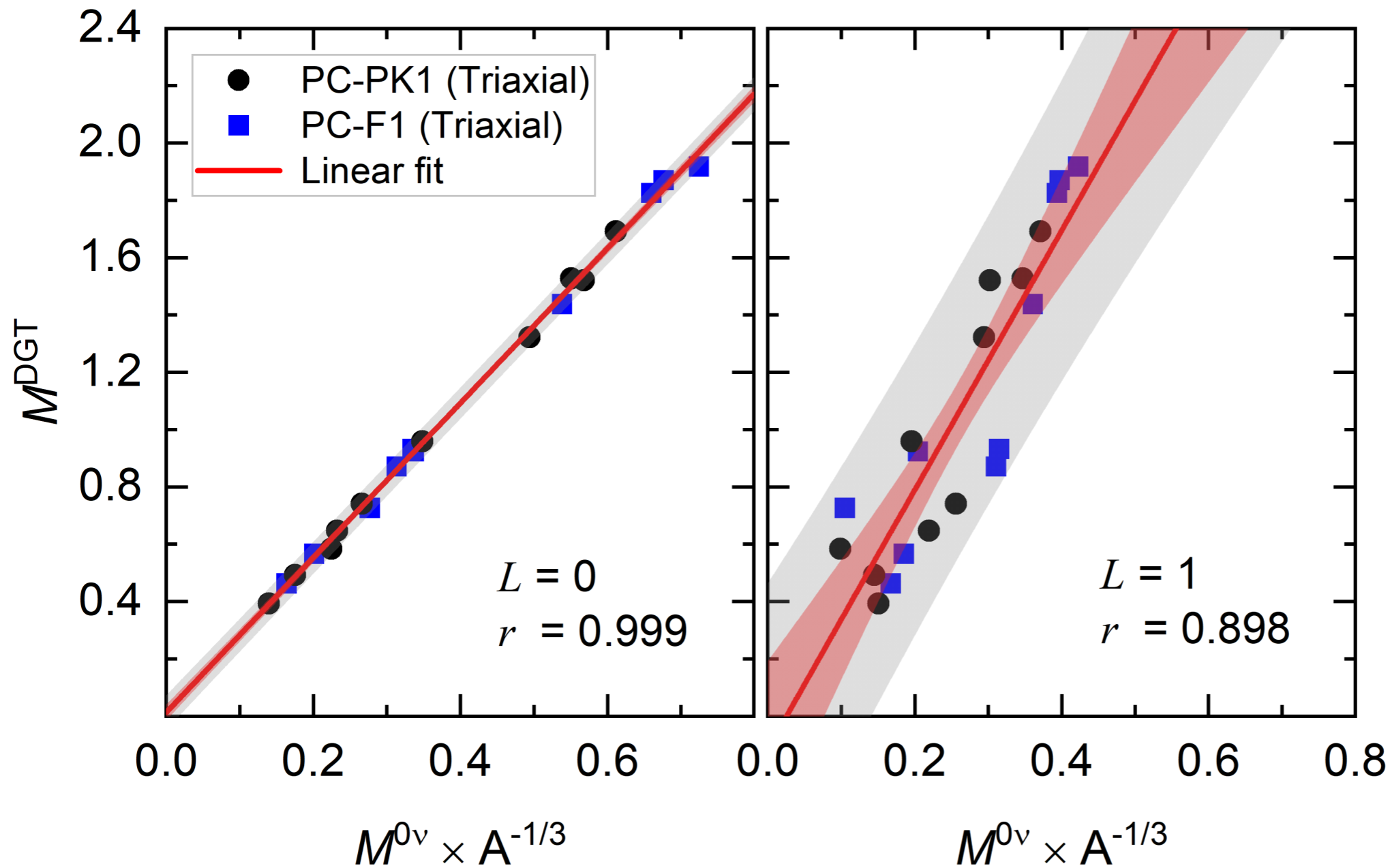
10 nuclei:

^{48}Ca , ^{76}Ge , ^{82}Se ,
 ^{96}Zr , ^{100}Mo , ^{116}Cd ,
 ^{124}Sn , ^{128}Te , ^{130}Te ,
 ^{136}Xe

Axial and triaxial
deformations
Full model space

- A strong linear correlation between $0\nu\beta\beta$ decay and DGT transition is demonstrated
- The linear correlation is robust against nuclear deformations

Decomposition of the NMEs

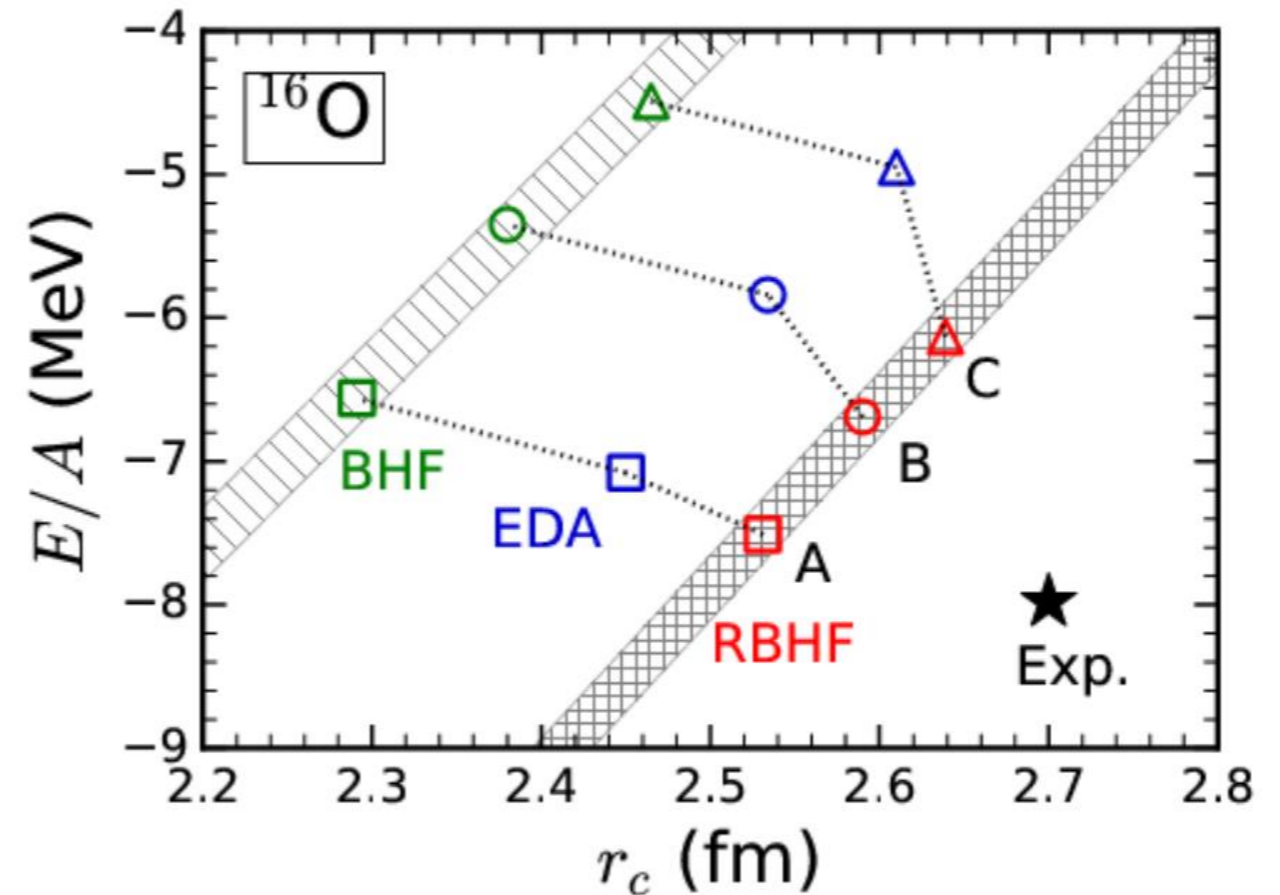
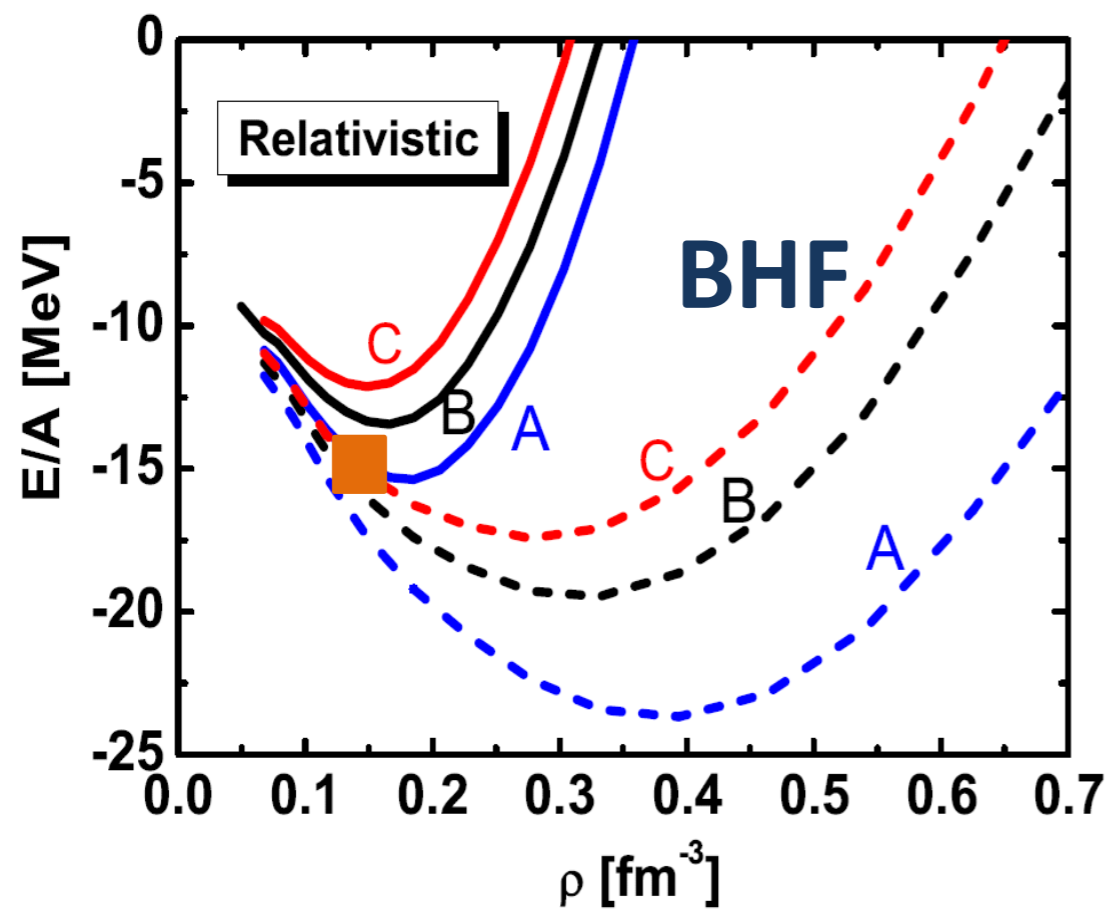


$$\hat{O}^{0\nu} = \sum_L \hat{O}_L^{0\nu}$$

$$M^{0\nu} = \sum_L M_L^{0\nu}$$

- The leading order $M_{L=0}^{0\nu}$ is more strongly correlated with M^{DGT} , while the correlation between $M_{L=1}^{0\nu}$ and M^{DGT} is weaker.

Toward Relativistic ab initio DFT





ab initio----- “from the beginning”

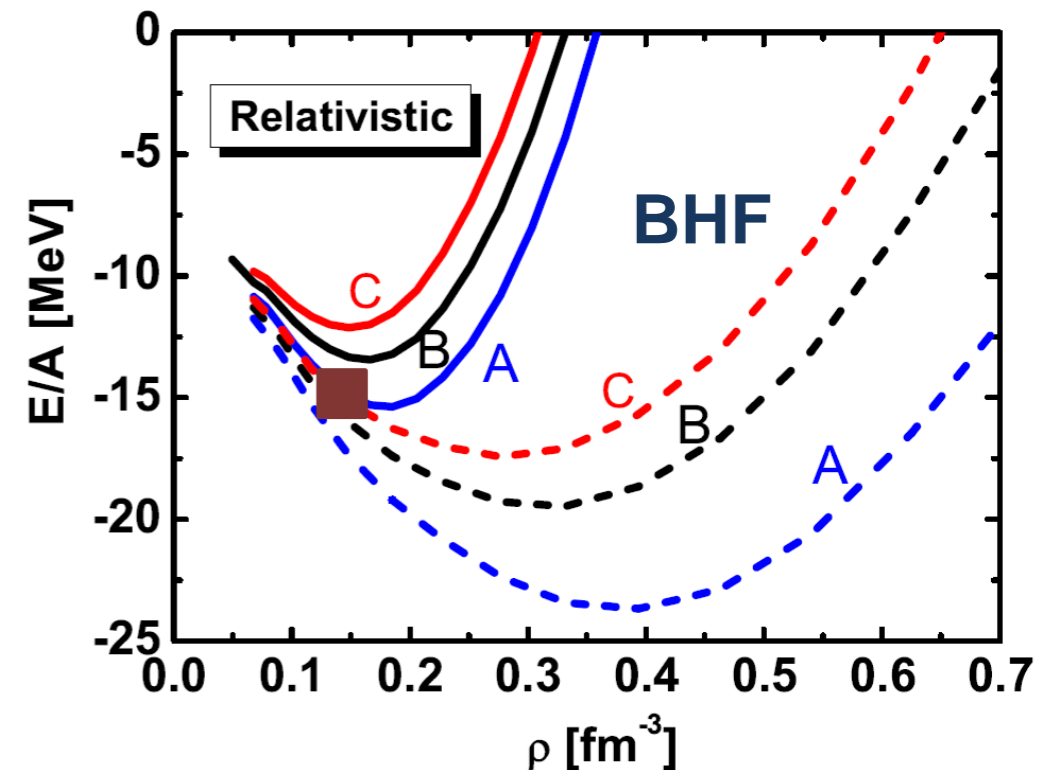
- without additional assumptions
- without additional parameters

ab initio in nuclear physics

- with **realistic** nucleon-nucleon interaction
- with **few-body** or **many-body** methods, such as Monte Carlo method, shell model and energy density functional theory

ab initio in nuclear matter

- Variational method Akmal PRC1998
- Green’s function method Dickhoff PPNP2004
- Chiral Perturbation theory Kaiser NPA2002
- Brueckner-Hartree-Fock (BHF) theory Baldo RPP2012
- Relativistic BHF theory Brockmann PRC1990
-





Progress in Particle and Nuclear Physics 109 (2019) 103713



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journal homepage: www.elsevier.com/locate/ppnp



Review

Towards an *ab initio* covariant density functional theory for nuclear structure



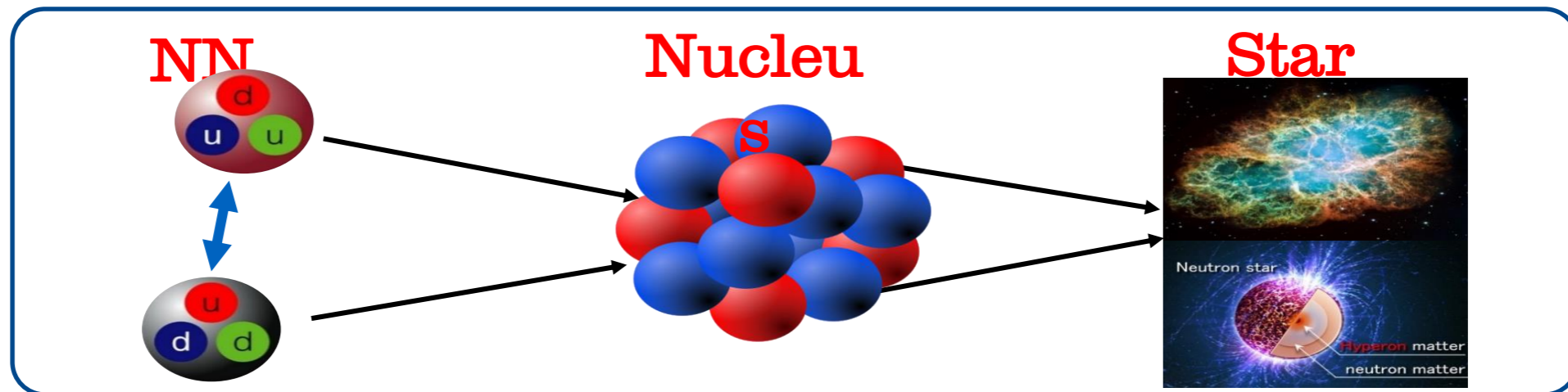
Shihang Shen^{a,b,c}, Haozhao Liang^{d,e}, Wen Hui Long^{f,g}, Jie Meng^{a,h,i,*}, Peter Ring^{a,j}

^a State Key Laboratory of Nuclear Physics and Technology, School of Physics, Peking University, Beijing 100871, China

^b Dipartimento di Fisica, Università degli Studi di Milano, Italy

^c INFN, Sezione di Milano, via Celoria 16, I-20133 Milano, Italy

^d RIKEN Nishina Center, Wako 351-0198, Japan



QCD/EFT

Brueckner

Relativistic *ab initio*

Effective

Relativistic DFT

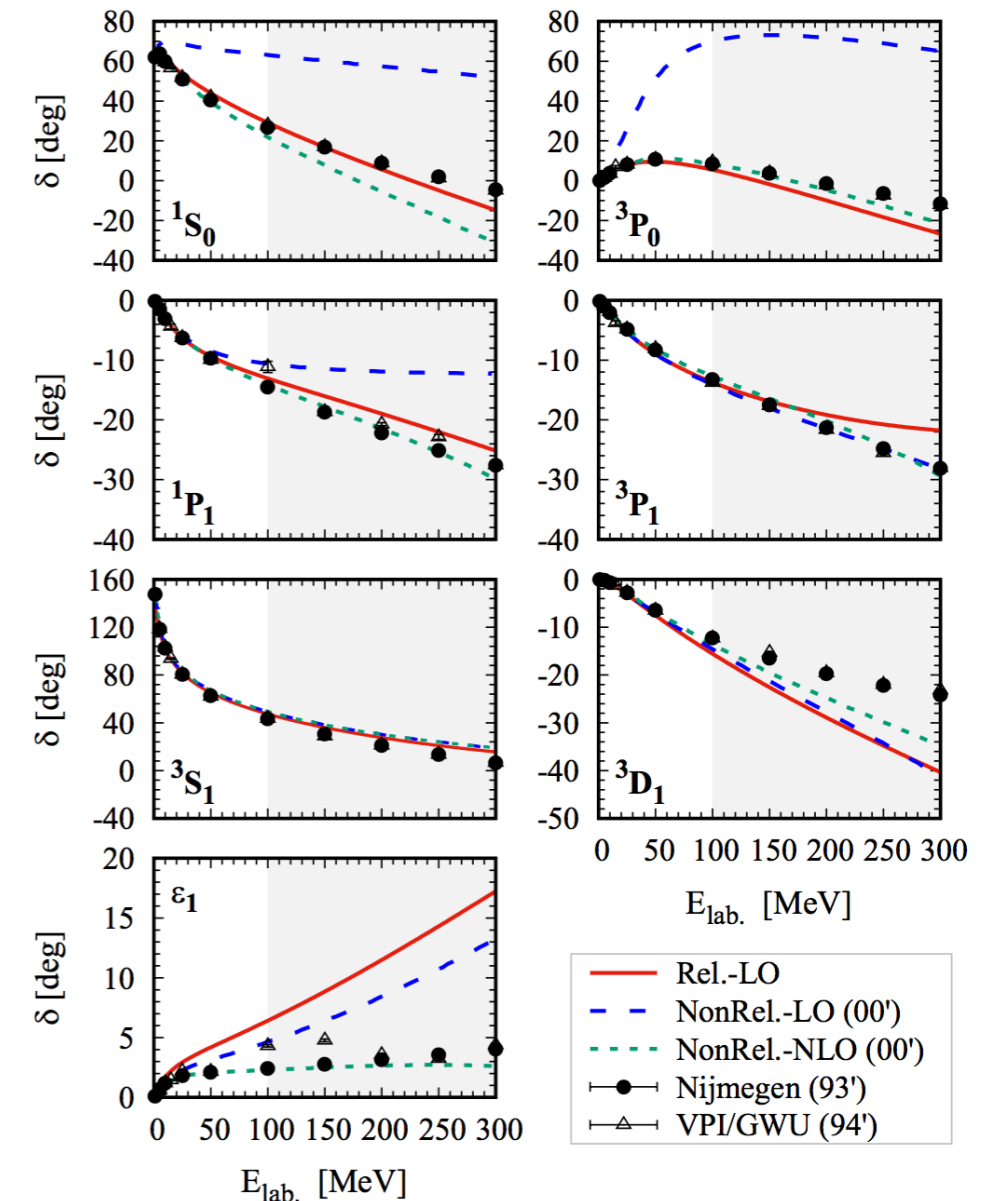


➤ Chiral Nucleon-Nucleon Interaction

X. L. Ren, K. W. Li, L. S. Geng, B. W. Long, P. Ring, and J. Meng,

Leading order relativistic chiral nucleon-nucleon interaction,

Chin. Phys. C 42, 014103(2018) J



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Accurate Relativistic Chiral Nucleon-Nucleon Interaction up to Next-to-Next-to-Leading Order

Jun-Xu Lu, Chun-Xuan Wang, Yang Xiao, Li-Sheng Geng, Jie Meng, and Peter Ring
Phys. Rev. Lett. **128**, 142002 – Published 6 April 2022

➤ Baryon Interaction by lattice QCD

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Accepted Paper

Dibaryon with highest charm number near unitarity from lattice QCD

Phys. Rev. Lett.
Yan Lyu, Hui Tong, Takuya Sugiura, Sinya Aoki, Takumi Doi, Tetsuo Hatsuda, Jie Meng, and Takaya Miyamoto

Accepted 2 July 2021

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Editors' Suggestion

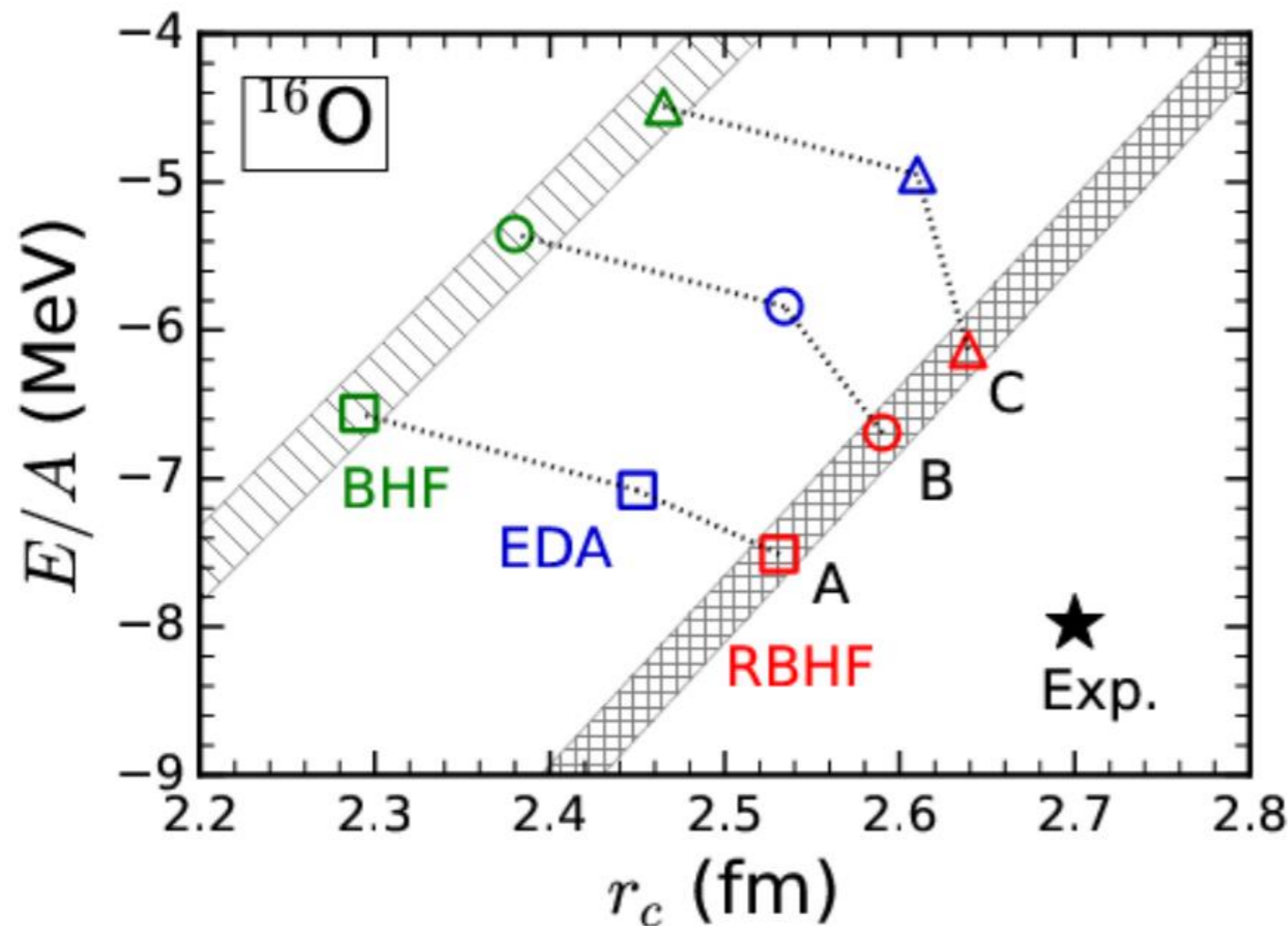
Open Access

Doubly Charmed Tetraquark T_{cc}^+ from Lattice QCD near Physical Point

Yan Lyu, Sinya Aoki, Takumi Doi, Tetsuo Hatsuda, Yoichi Ikeda, and Jie Meng
Phys. Rev. Lett. **131**, 161901 – Published 16 October 2023



Energies and charge radii of ^{16}O in RBHF in comparison with EDA and BHF



RBHF improves the description over EDA or BHF.

EDA and BHF taken from H. Mütter, R. Brockmann, and R. Machleidt, *PRC* **42**, 1981 (1990).

Shen, Hu, Liang, Meng, Ring, Zhang, *Chin. Phys. Lett.* **33** (2016) 102103

Shen, Liang, Meng, Ring, Zhang, *Phys. Rev. C* **96**, 014316 (2017)



PHYSICAL REVIEW C **98**, 054302 (2018)

Relativistic Brueckner-Hartree-Fock theory in nuclear matter without the average momentum approximation

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PHYSICAL REVIEW C **103**, 054319 (2021)

Nuclear matter in relativistic Brueckner-Hartree-Fock theory with Bonn potential in the full Dirac space

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Starting from the Bonn potential, the relativistic Brueckner-Hartree-Fock (RBHF) equations are solved for nuclear matter in the full Dirac space, which provides a unique way to determine the single-particle potentials and avoids the approximations applied in the RBHF calculations in the Dirac space with positive-energy states (PESs) only. The uncertainties of the RBHF calculations in the Dirac space with PESs only are investigated, and the importance of RBHF calculations in the full Dirac space is demonstrated. In the RBHF calculations in the full Dirac space, the empirical saturation properties of symmetric nuclear matter are reproduced, and the obtained equation of state agrees with the results based on the relativistic Green's function approach up to the saturation density.

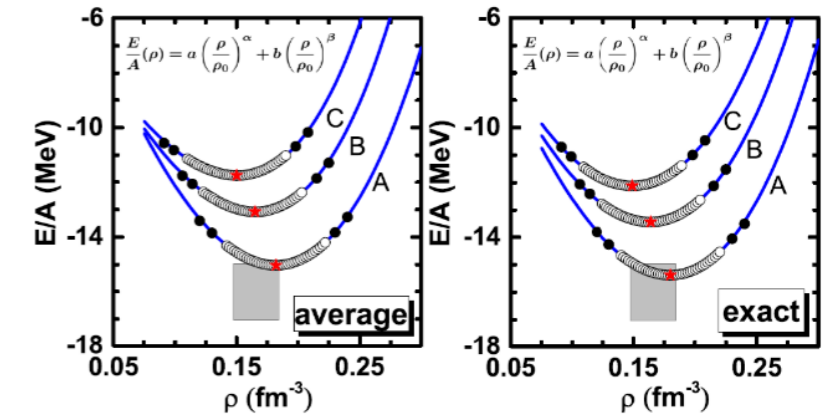
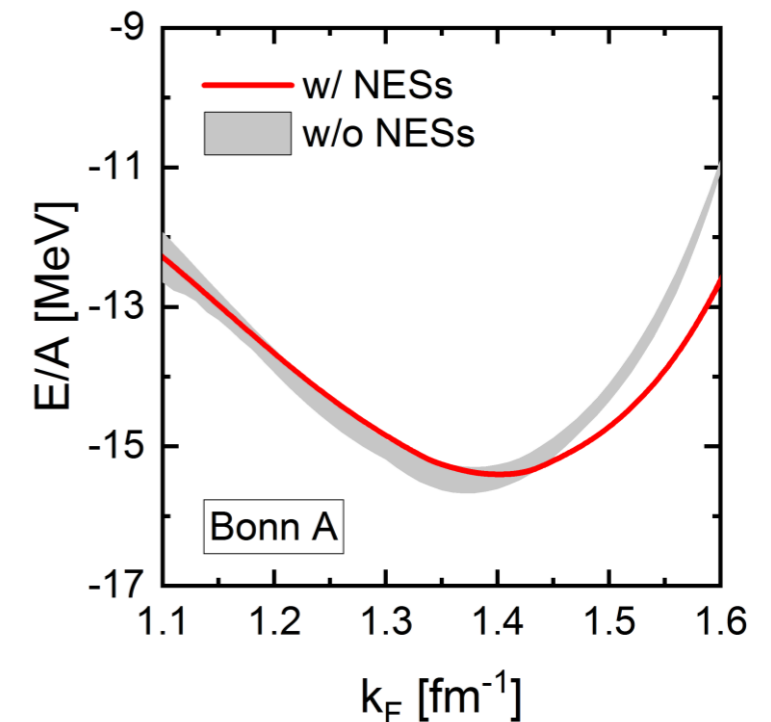


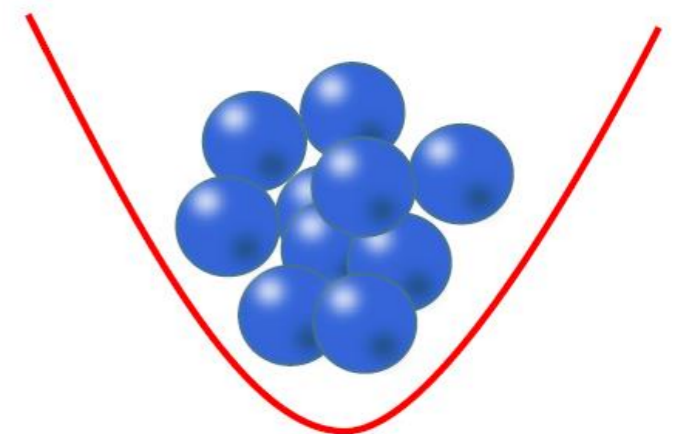
FIG. 1. Binding energy per nucleon for nuclear matter as a function of the total density ρ . Results for Bonn potentials A, B, C with (left panel) and without (right panel) c.m. momentum approximation are shown. The RBHF results are represented by open and solid circles, where open circles stand for the data used in the fit and solid circles indicate the validity of the results of the fit (solid curves). The red stars indicate the saturation points obtained from RBHF results.

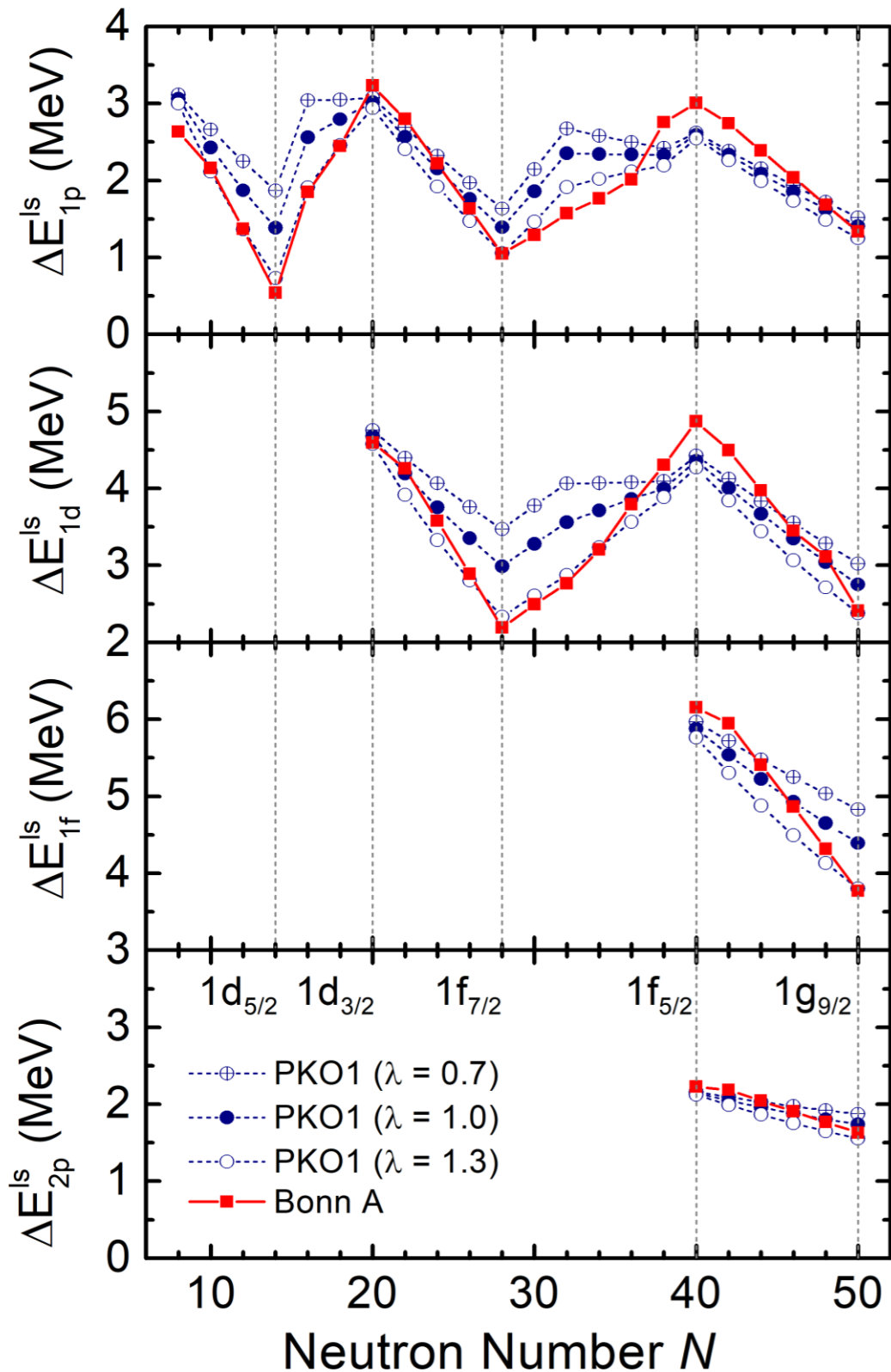




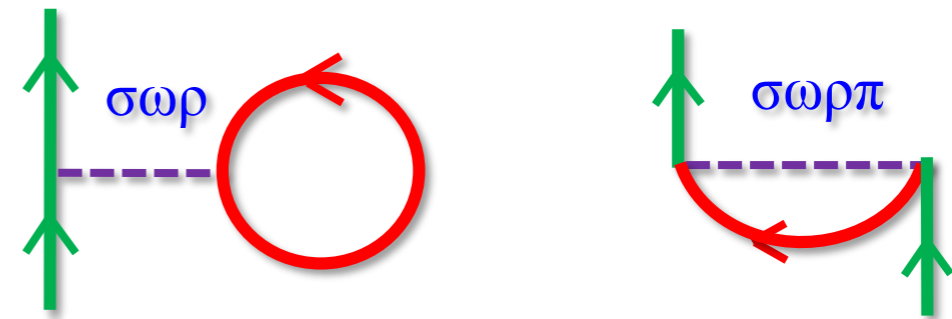
- Neutron drop is a neutron system confined in an external field. It is an ideal and simple system to investigate the neutron-rich environment by *ab initio* methods and phenomenological density functional theory. Pudliner et al., PRL 76, 2416 (1996). Gandolfi, Carlson, Pieper, PRL 106, 012501 (2011). Maris et al., PRC 87, 054318 (2013). Potter et al., PLB 739, 445 (2014). Zhao & Gandolfi, PRC 94, 041302(R) (2016).
-
- A neutron drop provides also an ideal and simple system to investigate the effects of tensor forces.
- From fully self-consistent relativistic Brueckner theory, a systematic and specific pattern due to the tensor forces is found in spin-orbit splitting in neutron drops, which forms a guide for the derivations of relativistic and nonrelativistic nuclear energy density functional.

Shen, Liang, Meng, Ring, Zhang,
Effects of tensor forces in nuclear spin-orbit
splittings from *ab initio* calculations.
Phys. Lett. B778 (2018) 344–348





➤ Comparison with the **phenomenological** relativistic Hartree-Fock (**RHF**) energy density functionals (EDF).



- **RHF** shows similar pattern, mainly contributed by **πNN tensor interaction**.
- Neither RBHF nor CDFT includes **beyond-mean-field effects** ➔ **a fair comparison!**

Shi-Hang Shen, Hao-Zhao Liang, Jie Meng, Peter Ring, Shuang-Quan Zhang,

Effects of tensor forces in nuclear spin-orbit splittings from ab initio calculations.

Phys. Lett. B 778 (2018) 344–348

Relativistic Brueckner-Hartree-Fock theory for neutron drops
Phys. Rev. C 97, 054312 (2018)



北京大學

Fully self-consistent relativistic Brueckner theory

Progress in Particle and Nuclear Physics 109 (2019) 103713



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Review

Towards an *ab initio* covariant density functional theory for nuclear structure

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- ✓ Origin of the heavy elements is a fundamental problems in modern science.
- ✓ Predictive power of PC-PK1 is demonstrated. Physics around the neutron drip line and $N=Z$ nuclei are discussed.
- ✓ Status of the DRHBc mass table collaboration is introduced. The effects of the continuum and deformation as well as the related interesting topics are discussed.
- ✓ Relativistic density functional theory solved in 3D lattice and its time-dependent version are introduced together with applications for Linear-chain, Chiral dynamics, Fission, etc.
- ✓ Strategy to build density functional based on QCD-spirited interaction and *ab initio* calculation are outlined.

Group members

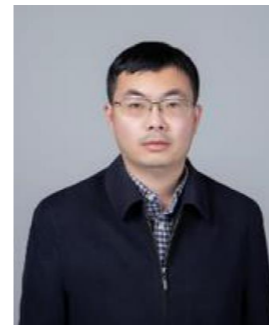
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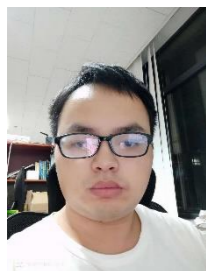
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Peng Guo



北京大學

**Thank you for your
attention!**

A decorative horizontal bar at the bottom of the slide, consisting of a blue segment on the left and a red segment on the right.