

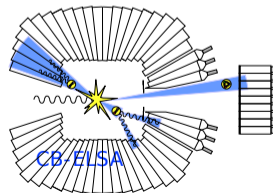
MULTI-MESON PHOTOPRODUCTION OFF THE PROTON

RECENT RESULTS FROM THE CBELSA/TAPS EXPERIMENT

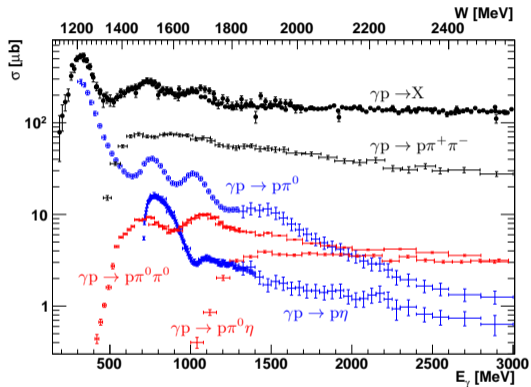
Nils Stausberg

25th June 2024

University of Bonn



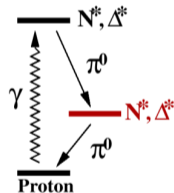
WHY MULTI-MESON PHOTOPRODUCTION?



Photoproduction as excellent tool for baryon spectroscopy!

Because...

- increasingly important at high E_γ
- sensitive to sequential decays



e.g. N^*/Δ^*

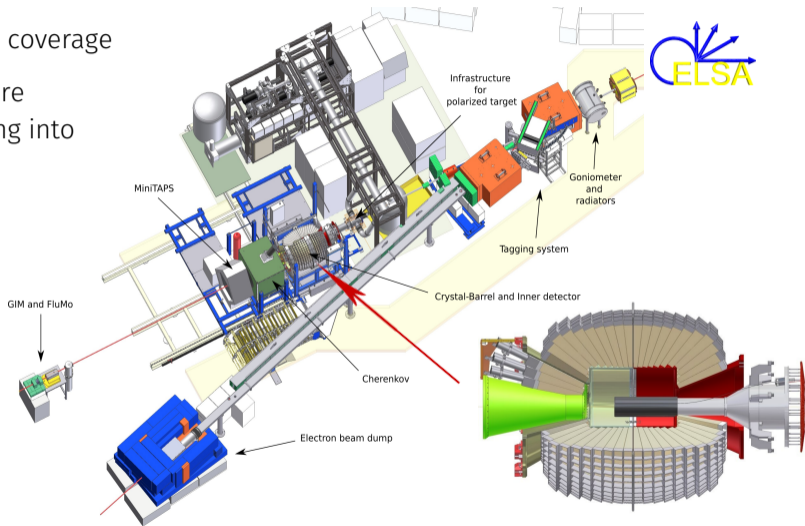
$$\rightarrow N(1520) \frac{3}{2} \pi^0$$

$$\rightarrow p \pi^0 \pi^0$$

Less background amplitudes for $p \pi^0 \pi^0$ and $p \pi^0 \eta$ compared to $p \pi^+ \pi^-$.

CBELSA/TAPS EXPERIMENT

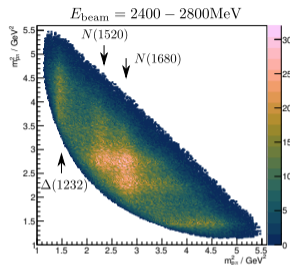
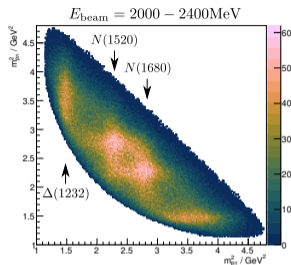
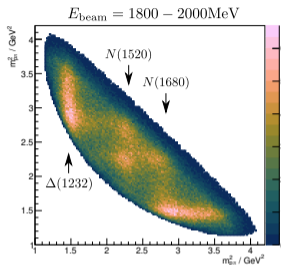
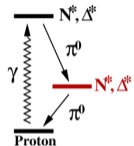
- Nearly 4π solid angular coverage
- Ideally suited to measure neutral mesons decaying into photons
- Polarized beam and target
→ Polarization observables



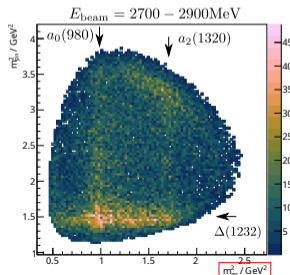
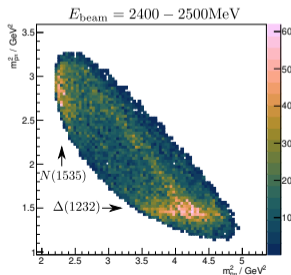
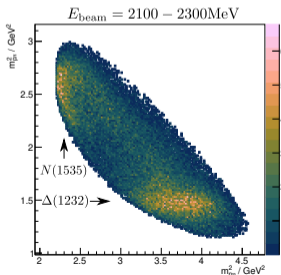
SEQUENTIAL DECAYS IN DALITZ PLOTS



[P. Mahlberg, Bonn]



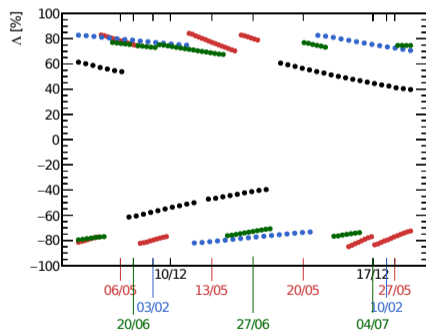
[G. Urff, Bonn]



MEASUREMENT WITH (DOUBLE) POLARIZATION

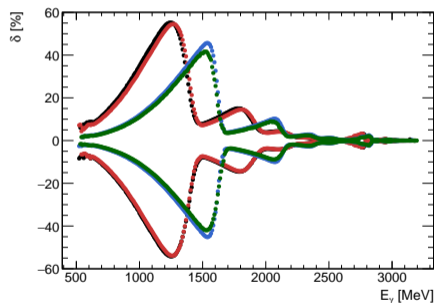
Transverse Pol. Target

Pol. using the Mainz-Dubna Frozen Spin target with Butanol as target material.



Linear Pol. Beam

Pol. via coherent bremsstrahlung with different settings of the diamond radiator.



Dec2017 May2018 Feb2019 Jun2021

POLARIZATION OBSERVABLES FOR $\gamma p \rightarrow p\pi^0\pi^0$

Polarization Observables (linear pol. beam + transverse pol. target)

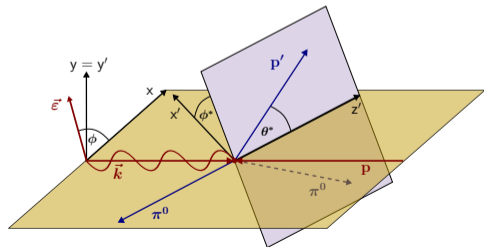
$$\frac{d\sigma}{d\Omega} = \frac{d\sigma_0}{d\Omega} \cdot \left(1 + \Lambda_x \cdot P_x + \Lambda_y \cdot P_y + \delta_l \sin(2\phi) \cdot I^S + \delta_l \cos(2\phi) \cdot I^C \right. \\ \left. + \Lambda_y \delta_l \sin(2\phi) \cdot P_y^S + \Lambda_x \delta_l \sin(2\phi) \cdot P_x^S \right. \\ \left. + \Lambda_x \delta_l \cos(2\phi) \cdot P_x^C + \Lambda_y \delta_l \cos(2\phi) \cdot P_y^C \right)$$

Target (\mathcal{T}), Beam (\mathcal{B}) and Beam + Target (\mathcal{BT}) polarization

The reaction can be described by 5 kinematic variables:

$$\underbrace{E_\gamma \quad \cos(\theta_{\pi^0})}_{2\text{-body}} \quad m_{p\pi^0} \quad \phi_{p\pi^0}^* \quad \theta_p^*$$

3-body



TARGET ASYMMETRY $P_y(\phi_{p\pi^0}^*)$ FOR BINS IN E_γ

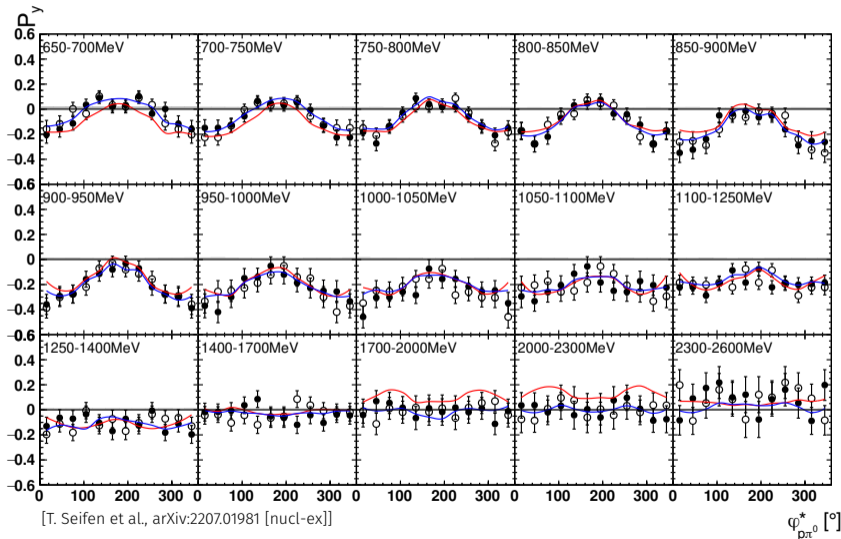
BnGa 2014-02

BnGa 2022-02 (fitted)

○ ●

Systematic check via
symmetry relation

$$P_y(2\pi - \phi^*) = P_y(\phi^*)$$



TARGET ASYMMETRY $P_y(\phi_{p\pi^0}^*)$ FOR BINS IN E_γ

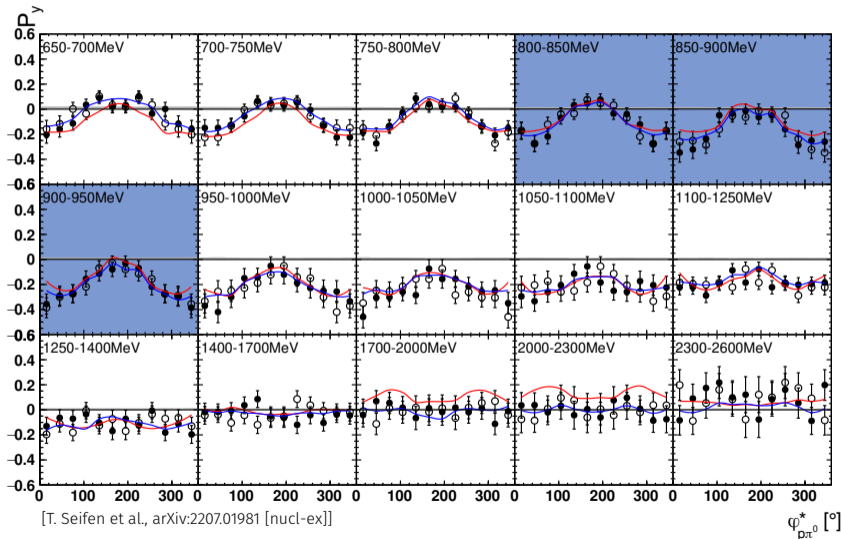
BnGa 2014-02

BnGa 2022-02 (fitted)

○ ●

Systematic check via
symmetry relation

$$P_y(2\pi - \phi^*) = P_y(\phi^*)$$



4D-TARGET ASYMMETRY $P_y(\phi_{\rho\pi^0}^*)$ FOR $E = 800 \text{ MeV} - 950 \text{ MeV}$

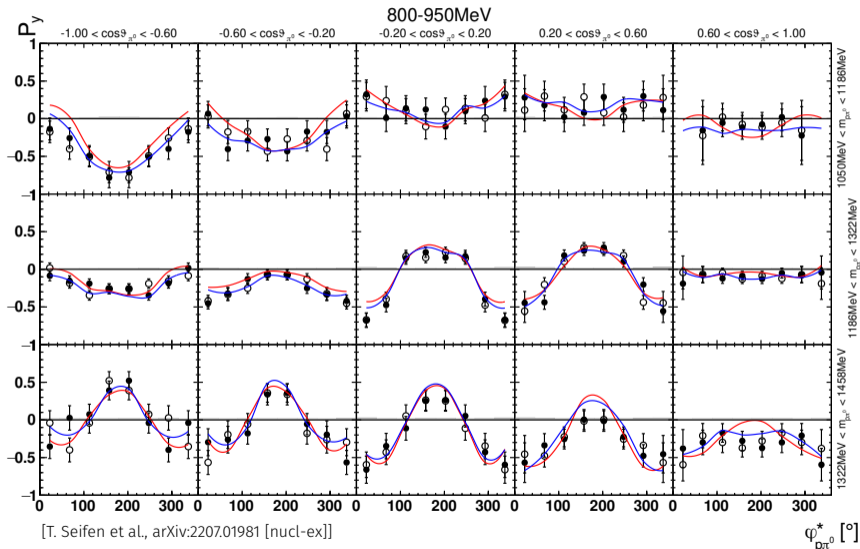
BnGa 2014-02

BnGa 2022-02 (fitted)



Systematic check via
symmetry relation

$$P_y(2\pi - \phi^*) = P_y(\phi^*)$$



4D-TARGET ASYMMETRY $P_y(\phi_{\rho\pi^0}^*)$ FOR $E = 800 \text{ MeV} - 950 \text{ MeV}$

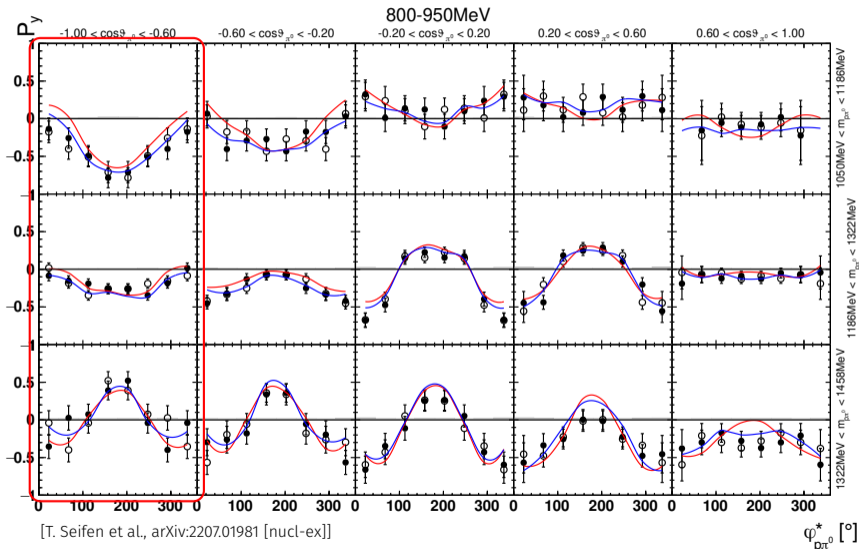
BnGa 2014-02

BnGa 2022-02 (fitted)



Systematic check via
symmetry relation

$$P_y(2\pi - \phi^*) = P_y(\phi^*)$$



4D-TARGET ASYMMETRY $P_y(\phi_{\rho\pi^0}^*)$ FOR $E = 800 \text{ MeV} - 950 \text{ MeV}$

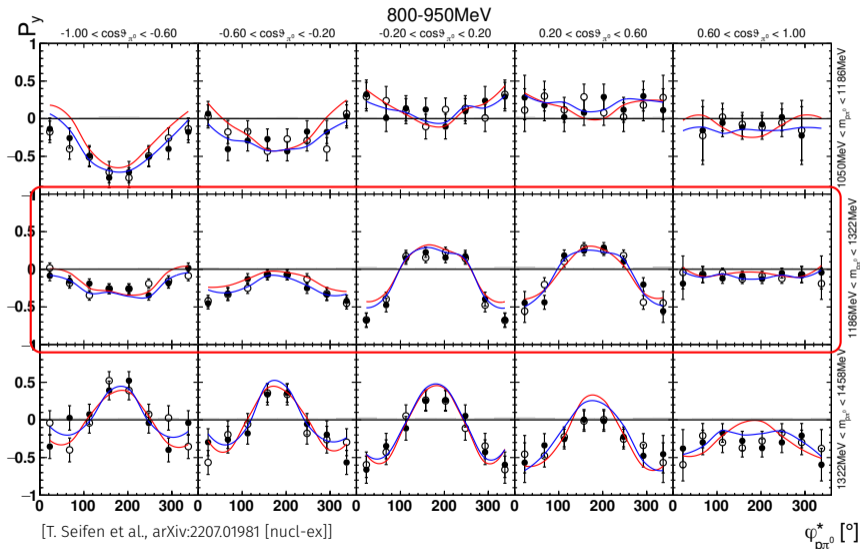
BnGa 2014-02

BnGa 2022-02 (fitted)

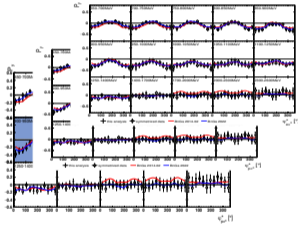


Systematic check via
symmetry relation

$$P_y(2\pi - \phi^*) = P_y(\phi^*)$$



FROM DATA TO RESONANCE PARAMETERS



Data

Partial Wave Analysis
(PWA)

		$N\pi$	L	$\Delta\pi$	$L < J$	$\Delta\pi$	$L > J$
1st shell, one oscillator	$N(1535)1/2^-$	46 ± 5	0	x		5 ± 3	2
		52 ± 5				2.5 ± 1.5	
	$N(1520)3/2^-$	61 ± 3	2	10 ± 4	0	10 ± 3	2
		61 ± 2		19 ± 4		9 ± 2	
	$N(1650)1/2^-$	48 ± 4	0	x		6 ± 3	2
		51 ± 4				12 ± 6	
$N(1700)3/2^-$	20 ± 8	2	66 ± 17	0	7 ± 4	2	
	15 ± 6		65 ± 15		9 ± 5		
$N(1675)5/2^-$	40 ± 1	2	19 ± 3	2	-	4	
	41 ± 2		30 ± 7		-		
$\Delta(1620)1/2^-$	30 ± 5	0	x		28 ± 15	2	
	28 ± 3				62 ± 10		
$\Delta(1700)3/2^-$	22 ± 6	2	16 ± 15	0	8 ± 6	2	
	22 ± 4		20 ± 15		10 ± 6		
2nd shell, one oscillator	$N(1440)1/2^+$	66 ± 3	1	x		10 ± 6	1
		63 ± 2				20 ± 7	
	$\Delta(1600)3/2^+$	17 ± 4	1	70 ± 6	1	< 2	3
		14 ± 4		77 ± 5		< 2	
	$N(1720)3/2^+$	13 ± 5	1	15 ± 7	1	6 ± 6	3
		11 ± 4		62 ± 15		6 ± 6	
	$N(1680)5/2^+$	68 ± 8	3	8 ± 4	1	8 ± 4	3
		62 ± 4		7 ± 3		10 ± 3	
	$\Delta(1910)1/2^+$	16 ± 6	1	x		17 ± 9	1
	12 ± 3				50 ± 16		
$\Delta(1920)3/2^+$	12 ± 6	1	5 ± 4	1	40 ± 20	3	
	8 ± 4		18 ± 10		98 ± 14		
$\Delta(1905)5/2^+$	13 ± 4	3	20 ± 12	1	-	3	
	13 ± 2		33 ± 10		-		
$\Delta(1950)7/2^+$	46 ± 4	3	5 ± 4	3	-	5	
	46 ± 2		5 ± 4		-		

Resonance Parameters

BRANCHING RATIOS & INTERPRETATION IN HARMONIC OSCILLATOR MODEL

$\Delta(1910)_{\frac{1}{2}}^{1+}, \Delta(1920)_{\frac{3}{2}}^{3+}, \Delta(1905)_{\frac{5}{2}}^{5+}, \Delta(1950)_{\frac{7}{2}}^{7+}$

BR into ground state: $(44 \pm 7) \%$
 $N\pi^0, \Delta(1232)\pi^0$

BR into excited states (with L=1): $(5 \pm 2) \%$
 $N(1520)\pi^0, N(1535)\pi^0, N\sigma$

$N(1880)_{\frac{1}{2}}^{1+}, N(1900)_{\frac{3}{2}}^{3+}, N(2000)_{\frac{5}{2}}^{5+}, N(1990)_{\frac{7}{2}}^{7+}$

BR into ground state: $(34 \pm 6) \%$
 $N\pi^0, \Delta(1232)\pi^0$

BR into excited states (with L=1): $(21 \pm 5) \%$
 $N(1520)\pi^0, N(1535)\pi^0, N\sigma$

BRANCHING RATIOS & INTERPRETATION IN HARMONIC OSCILLATOR MODEL

$$\Delta(1910)_{\frac{1}{2}^+}, \Delta(1920)_{\frac{3}{2}^+}, \Delta(1905)_{\frac{5}{2}^+}, \Delta(1950)_{\frac{7}{2}^+}$$

BR into ground state: $(44 \pm 7) \%$

$N\pi^0, \Delta(1232)\pi^0$

BR into excited states (with L=1): $(5 \pm 2) \%$

$N(1520)\pi^0, N(1535)\pi^0, N\sigma$

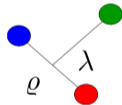
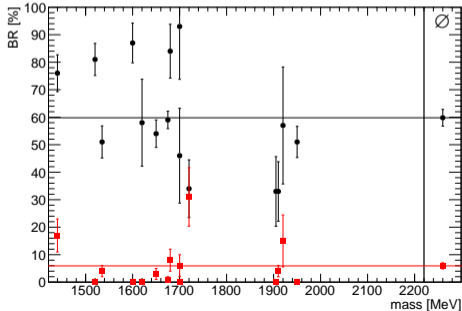
$$N(1880)_{\frac{1}{2}^+}, N(1900)_{\frac{3}{2}^+}, N(2000)_{\frac{5}{2}^+}, N(1990)_{\frac{7}{2}^+}$$

BR into ground state: $(34 \pm 6) \%$

$N\pi^0, \Delta(1232)\pi^0$

BR into excited states (with L=1): $(21 \pm 5) \%$

$N(1520)\pi^0, N(1535)\pi^0, N\sigma$



One-oscillator excitation \rightarrow ground state

One-oscillator excitation \rightarrow excited states

$$S = \frac{1}{2} ([0s \times 0d] + [0d \times 0s])^{(L=2)}$$

BRANCHING RATIOS & INTERPRETATION IN HARMONIC OSCILLATOR MODEL

$$\Delta(1910)_{\frac{1}{2}^+}, \Delta(1920)_{\frac{3}{2}^+}, \Delta(1905)_{\frac{5}{2}^+}, \Delta(1950)_{\frac{7}{2}^+}$$

BR into ground state: $(44 \pm 7) \%$

$N\pi^0, \Delta(1232)\pi^0$

BR into excited states (with L=1): $(5 \pm 2) \%$

$N(1520)\pi^0, N(1535)\pi^0, N\sigma$

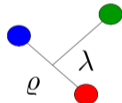
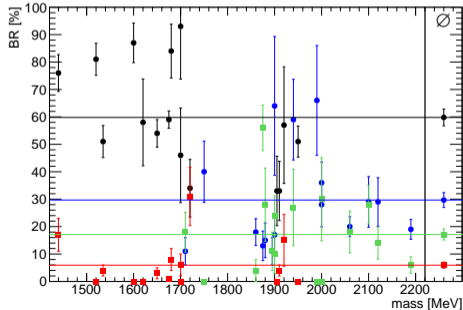
$$N(1880)_{\frac{1}{2}^+}, N(1900)_{\frac{3}{2}^+}, N(2000)_{\frac{5}{2}^+}, N(1990)_{\frac{7}{2}^+}$$

BR into ground state: $(34 \pm 6) \%$

$N\pi^0, \Delta(1232)\pi^0$

BR into excited states (with L=1): $(21 \pm 5) \%$

$N(1520)\pi^0, N(1535)\pi^0, N\sigma$



Mixed-oscillator excitations \rightarrow ground state

Mixed-oscillator excitations \rightarrow excited states

$$M_S = \frac{1}{2} ([0s \times 0d] - [0d \times 0s])^{(L=2)}$$

$$M_A = -\frac{1}{2} [0p \times 0p]^{(L=2)}$$

BRANCHING RATIOS & INTERPRETATION IN HARMONIC OSCILLATOR MODEL

$\Delta(1910)_{\frac{1}{2}}^{1+}$, $\Delta(1920)_{\frac{3}{2}}^{3+}$, $\Delta(1905)_{\frac{5}{2}}^{5+}$, $\Delta(1950)_{\frac{7}{2}}^{7+}$

BR into ground state: $(44 \pm 7) \%$

$N\pi^0$, $\Delta(1232)\pi^0$

BR into excited states (with L=1): $(5 \pm 2) \%$

$N(1520)\pi^0$, $N(1535)\pi^0$, $N\sigma$

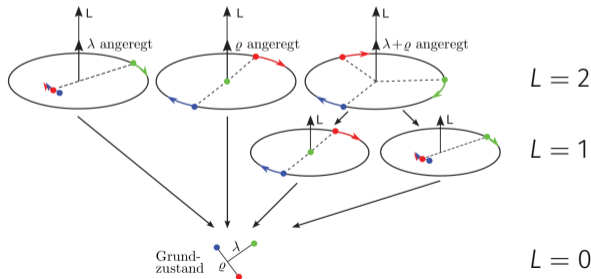
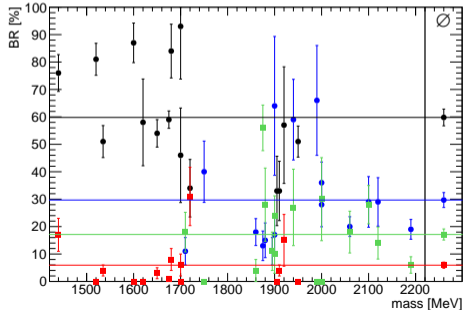
$N(1880)_{\frac{1}{2}}^{1+}$, $N(1900)_{\frac{3}{2}}^{3+}$, $N(2000)_{\frac{5}{2}}^{5+}$, $N(1990)_{\frac{7}{2}}^{7+}$

BR into ground state: $(34 \pm 6) \%$

$N\pi^0$, $\Delta(1232)\pi^0$

BR into excited states (with L=1): $(21 \pm 5) \%$

$N(1520)\pi^0$, $N(1535)\pi^0$, $N\sigma$



SINGLE POLARIZATION P_y - NEW DATA FOR $\gamma p \rightarrow p\pi^0\pi^0$

4D-Target Asymmetry

$P_y(\phi_{p\pi^0}^*)$ for

$E = 1400 \text{ MeV} - 1550 \text{ MeV}$

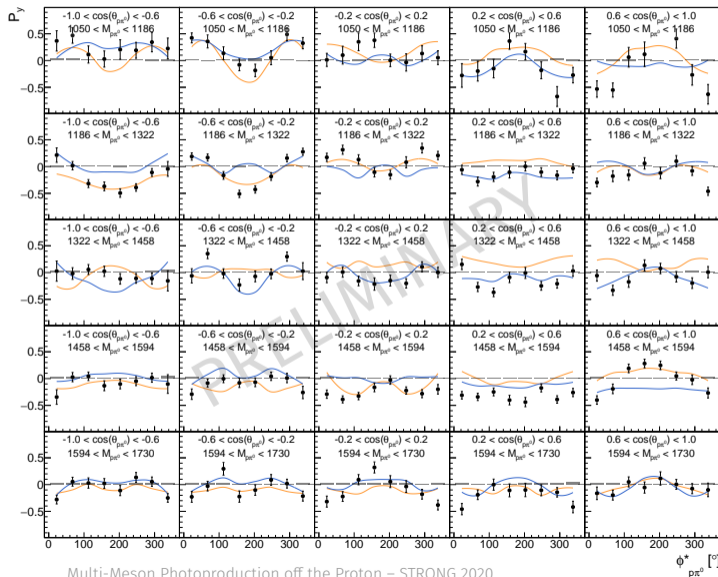
BnGa 2014-02

BnGa 2022-02

New data

Not fitted

[N. Stausberg, Bonn]



DOUBLE POLARIZATION P_y^c - NEW DATA FOR $\gamma p \rightarrow p\pi^0\pi^0$

linear pol. beam
+ transverse pol.
target

BnGa 2014-02

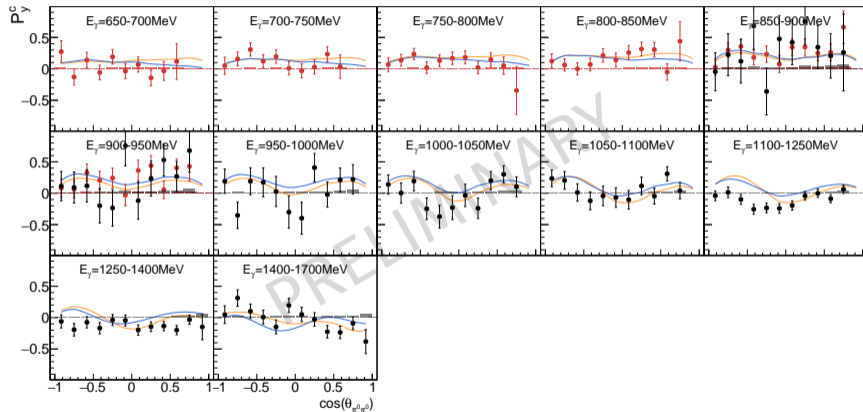
BnGa 2022-02

Seifen et al.

New data

Not fitted

[N. Stausberg, Bonn]



SINGLE POLARIZATION I^C - NEW DATA FOR $\gamma p \rightarrow p\pi^0\pi^0$

Attention!

Other beam time with liquid hydrogen target and different coherent edges.

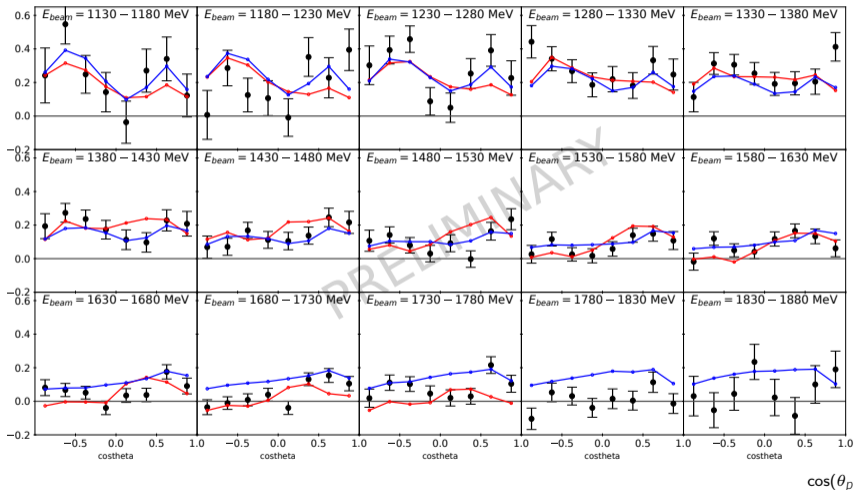
BnGa 2014-02

BnGa 2022-02

New data

Not fitted

[P. Mahlberg, Bonn]



Summary

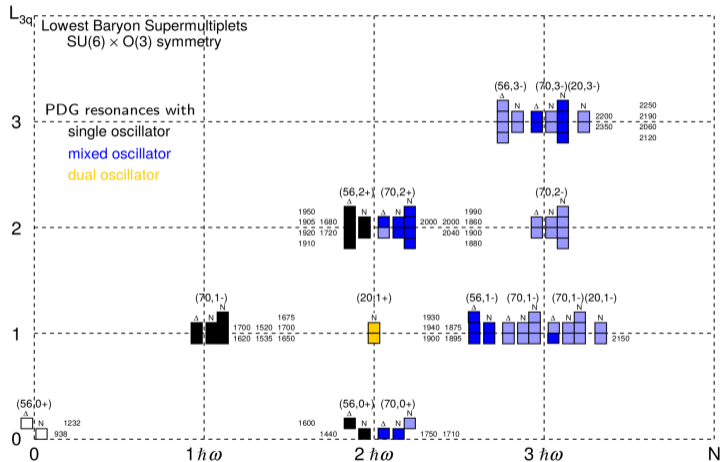
- Multidimensional determination of polarization observables for multi-meson final states essential
- Continuously achieving better precision of resonance parameters
 - BR shows systematics related to the wave function of baryons

Outlook

- Finalizing double polarization observables $P_x^C, P_y^C, P_x^S, P_y^S$ at high energies for $p\pi^0\pi^0$
- Additional data for linear pol. beam and unpol. target
 - Beam polarization observables I^S, I^C at high energies for $p\pi^0\pi^0$ and $p\pi^0\eta$

Thank You

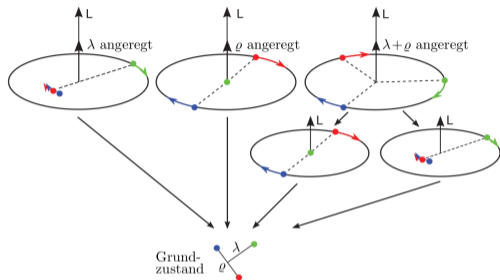
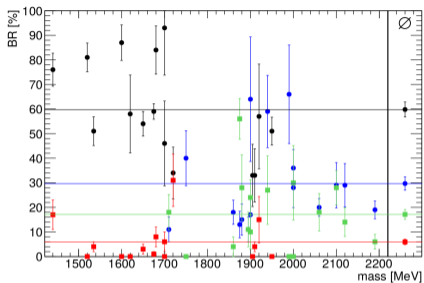
SU(6) × O(3) SUPERMULTIPLETS



$$N = 2n_\rho + 2n_\lambda + l_\rho + l_\lambda$$

$$L = l_\rho \otimes l_\lambda$$

BRANCHING RATIOS & INTERPRETATION IN HARMONIC OSCILLATOR MODEL



One-oscillator excitation → ground state

One-oscillator excitation → excited states

Mixed-oscillator excitations → ground state

Mixed-oscillator excitations → excited states

$$S = \frac{1}{2} ([0s \times 0d] + [0d \times 0s])^{(L=2)}$$

$$M_S = \frac{1}{2} ([0s \times 0d] - [0d \times 0s])^{(L=2)}$$

$$M_A = -\frac{1}{2} [0p \times 0p]^{(L=2)}$$

BRANCHING RATIOS

		$N\pi$	L	$\Delta\pi$	$L < J$	$\Delta\pi$	$L > J$	$N(1440)\pi$	L	$N(1520)\pi$	L	$N(1535)\pi$	L	$N(1680)\pi$	L	$N\sigma$	L
1st shell, one oscillator	$N(1535)1/2^-$	46±5	0	x		5±3	2	6±5	0	-	1	-	1	-	2	4±2	1
		52±5				2.5±1.5		12±8		-		-		-		6±4	
	$N(1520)3/2^-$	61±3	2	10±4	0	10±3	2	<1	2	-	1	-	1	-	2	<2	1
		61±2		19±4		9±2		<1		-		-		-		<2	
	$N(1650)1/2^-$	48±4	0	x		6±3	2	5±3	0	-	1	-	1	-	2	3±2	1
		51±4				12±6		16±10		-		-		-		10±8	
	$N(1700)3/2^-$	20±8	2	66±17	0	7±4	2	9±5	2	<2	1	<1	1	-	2	6±4	1
	15±6		65±15		9±5		7±4		<4		<1		-		8±6		
$N(1675)5/2^-$	40±1	2	19±3	2	-	4	-	2	-	1	-	3	-	0	1±1	3	
	41±2		30±7		-		-		-		-		-		5±2		
$\Delta(1620)1/2^-$	30±5	0	x		28±15	2	15±8	0	-	1	-	1	-	2	x		
	28±3				62±10		6±3		-		-		-				
$\Delta(1700)3/2^-$	22±6	2	16±15	0	8±6	2	3±2	2	<1	1	<1	1	-	2	x		
	22±4		20±15		10±6		<1		3±2		<1		-				
2nd shell, one oscillator	$N(1440)1/2^+$	66±3	1	x		10±6	1	-	1	-	0	-	2	-	3	17±6	1
		63±2				20±7		-		-		-		-		17±6	
	$\Delta(1600)3/2^+$	17±4	1	70±6	1	<2	3	<1	1	-	0	-	2	-	1	x	
		14±4		77±5		<2		22±5 [†]		-		-		-			
	$N(1720)3/2^+$	13±5	1	15±7	1	6±6	3	6±5	1	7±3	0	4±2	2	-	1	20±10	2
		11±4		62±15		6±6		<2		3±2		<2		-		8±6	
	$N(1680)5/2^+$	68±8	3	8±4	1	8±4	3	-	3	<1	2	-	2	-	1	8±4	2
		62±4		7±3		10±3		-		<1		-		-		14±5	
$\Delta(1910)1/2^+$	16±6	1	x		17±9	1	40±15	1	-	0	4±2	2	-	3	x		
	12±3				50±16		6±3		-		5±3		-				
$\Delta(1920)3/2^+$	12±6	1	5±4	1	40±20	3	9±6	1	10±8	0	5±5	2	-	1	x		
	8±4		18±10		58±14		<4		<5		<2		-				
$\Delta(1905)5/2^+$	13±4	3	20±12	1	-	3	-	3	-	2	<1	2	6±2	1	x		
	13±2		33±10		-		-		-		<1		10±5				
$\Delta(1950)7/2^+$	46±4	3	5±4	3	-	5	-	3	-	2	-	4	3±2	1	x		
	46±2		5±4		-		-		-		-		6±3				

[T. Seifen et al., arXiv:2207.01981 [nucl-ex]]

BRANCHING RATIOS

2nd shell, mixed oscillations	$N(1710)1/2^+$	5±3 5±3 5±3	1	x	6±4 7±4 25±10	1	22±12 30±10 <5	1	<2 <2 <2	2	4±4 - 15±6	0	- - -	3	14±6 55±15 10±5	0	
	$\Delta(1750)1/2^+$	18±5	1	x	22±10	1	49±26	1	-	0	<2	2	-	3	x		
	$\Delta(2000)3/2^+$	11±5	1	14±6	1	3±2	3	9±6	1	<1	0	<1	2	-	1	x	
	$N(2100)1/2^+$	17±7 16±5	1	x	12±6 10±4	1	- -	1	<2 <2	2	<1 30±4	0	- -	3	28±7 20±6	0	
	$N(1880)1/2^+$	11±6 6±3	1	x	4±2 30±12	1	- -	1	- -	2	6±3 8±4	0	- -	3	22±13 25±15	0	
	$N(1900)3/2^+$	4±3 3±2	1	9±6 17±8	1	4±3 33±12	3	9±6 <2	1	<2 15±8	0	15±6 7±3	2	- -	1	9±4 4±3	2
	$N(1860)5/2^+$	12±4 20±6	3	2±2 10±6 [†]	1	4±2 27±11 [†]	3	- -	3	- -	2	- -	2	- -	1	4±4 51±10 [†]	2
	$N(2000)5/2^+$	11±5 8±4	3	9±4 22±10	1	16±4 34±15	3	- -	1	2±2 21±10	2	- -	2	28±9 16±9	1	28±15 10±5	2
	$N(1990)7/2^+$	1±1 1.5±0.5	3	65±20 48±10	3	-	5	<2 <2	1	<2 <2	1	<2 <2	4	- -	1	- -	4
3rd shell, mixed oscillations	$N(1895)1/2^-$	6±4 2.5±1.5	0	x	5±3 7±4	2	2±2 8±8	0	- -	1	- -	1	- -	2	11±7 18±15	1	
	$N(1875)3/2^-$	3±2 4±2	2	6±4 14±7	0	4±3 7±5	2	11±5 5±3	2	4±4 <2	1	2±2 <1	1	- -	2	50±20 45±15	
	$\Delta(1900)1/2^-$	4±4 7±2	0	x	60±25 50±20	2	15±12 20±12	0	10±7 6±4	1	- -	1	- -	2	x		
	$\Delta(1940)3/2^-$	13±6 2±1	0	16±6 46±20	0	30±12 12±7	2	5±5 7±7	2	7±5 4±3	1	20±13 8±6	1	- -	2	x	
	$N(2120)3/2^-$	5±2 5±3	2	10±5 50±20	0	14±7 20±12	2	6±6 10±10	2	10±5 15±10	1	<2 15±8	1	- -	2	4±3 11±4	1
	$N(2060)5/2^-$	11±2 11±2	2	9±3 7±3	2	- -	4	8±5 9±5	2	13±7 15±6	1	- -	3	14±5 15±7	0	5±3 6±3	3
	$N(2190)7/2^-$	15±3 16±2	4	4±2 25±6	2	- -	4	- -	4	- -	3	- -	3	- -	2	6±3 5±3	3

[T. Seifen et al., arXiv:2207.01981 [nucl-ex]]

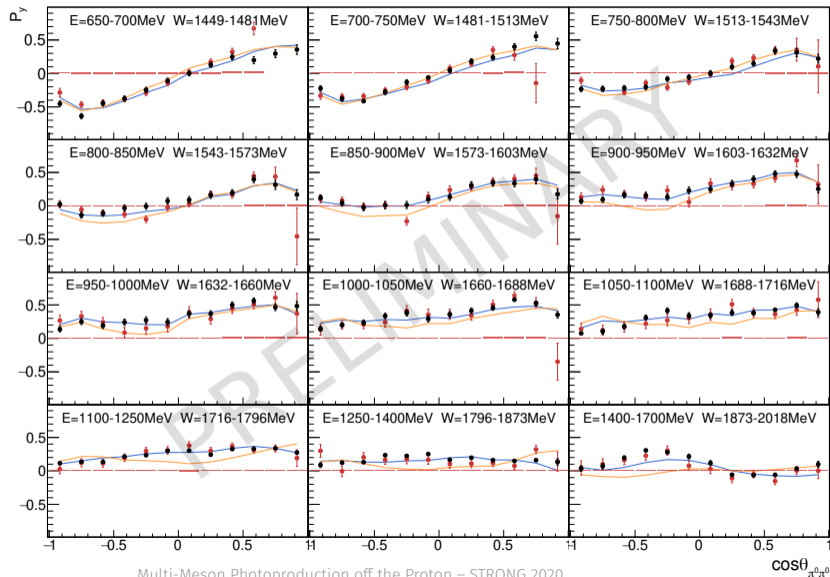
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New data



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