

Σ beam asymmetry for η photoproduction off the proton at BGOOD

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on behalf of the BGOOD collaboration

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Outline:

- Brief introduction to Physics Motivations for meson photoproduction measurements
- Short description of BGOOD apparatus
- Analysis and η photoproduction events selection
- Beam Asymmetry Σ extraction method
- Σ measurement Results & comparison with Existing Data
- Summary & Conclusions

Why meson photoproduction ?

The understanding of the dynamics underlying the bound state of the nucleon and its excited spectrum still remain a crucial task since in this energy range QCD cannot be treated in perturbative mode.

Many models, based on different *degrees of freedom* descriptions, have been developed in order to describe the spectrum of excitation states and their features

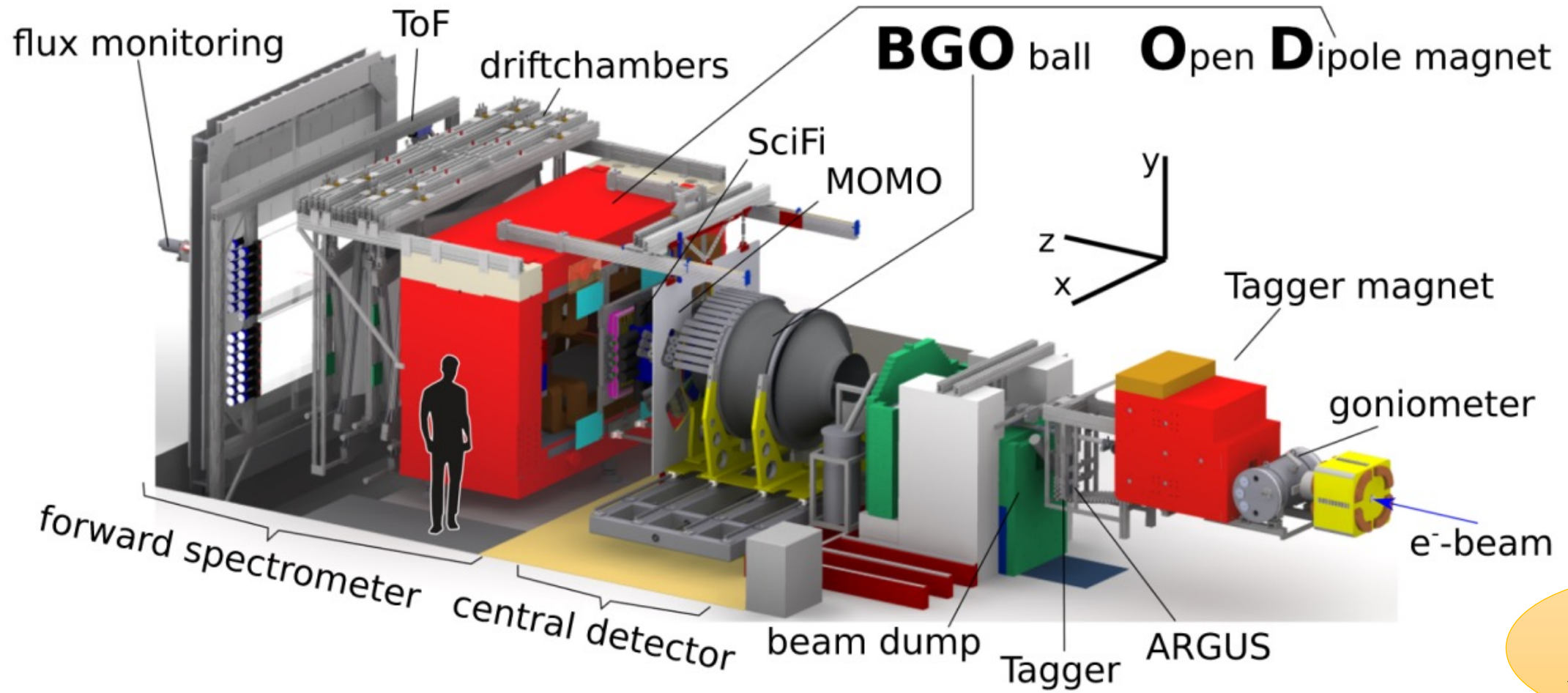
Meson photoproduction studies represent a strong tool for probing nucleon resonances:

- ✓ Access to resonance states coupled to photons which only weakly coupling to the πN processes (Missing Resonances problem)
- ✓ Access to Polarization Observables \Rightarrow Separation of overlapping resonances and characterization in terms of Spin and Parity, Constrains for unambiguous PWA
- Low e.m cross section \Rightarrow ✓ Overcome thanks to technological developments
- Non resonance contribution \Rightarrow ✓ Disentagled with polarization observables

η photoproduction isoscalar meson ($I=0$) \Rightarrow Isospin Filter \Rightarrow only $N^*(I=1/2)$ resonances as intermediate states

BGOOD Detector:

BGO calorimeter (central region) & Forward Spectrometer combination

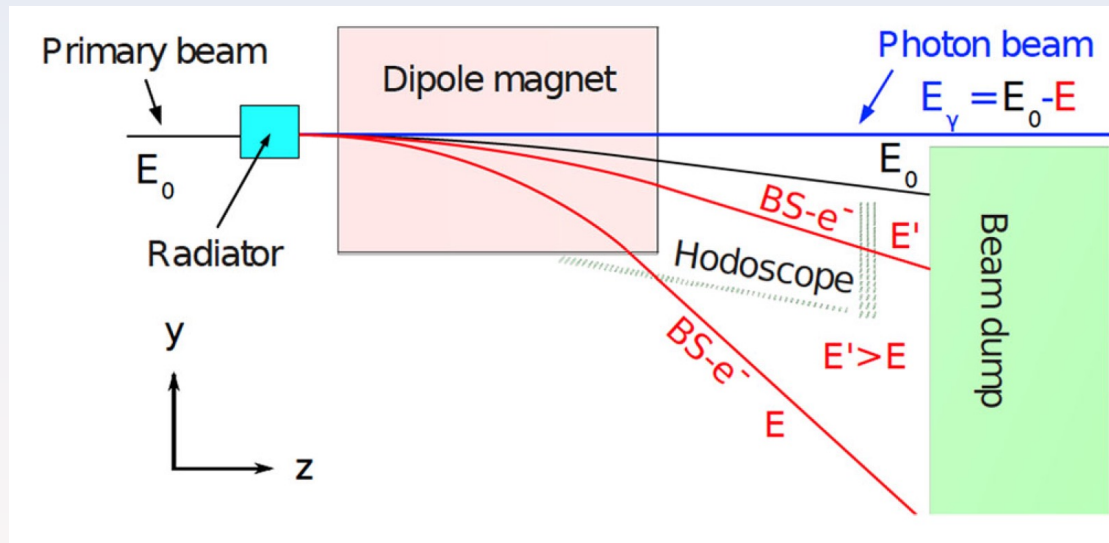


ELSA
Accelerator
 $E_{e_{MAX}} = 3.2 \text{ GeV}$

BGOOD Tagged and Polarized Bremsstrahlung Photon Beam

Tagging Detector

E_γ measured through the detection of the corresponding electrons in the tagging system.

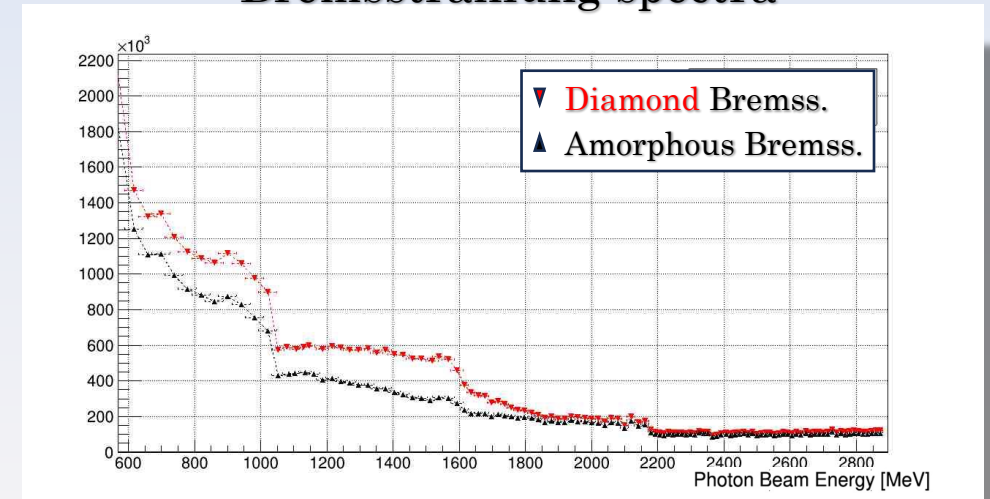


linearly polarised photon beams generated by coherent bremsstrahlung using a diamond crystal radiator.

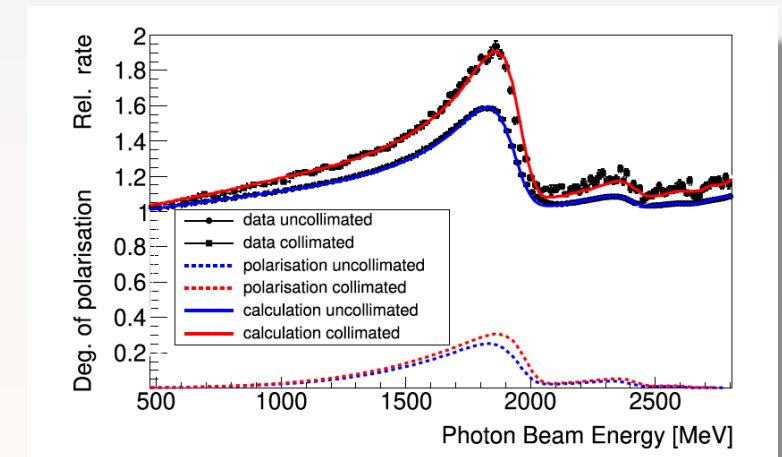
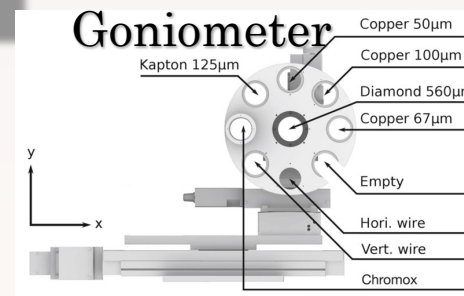
Cu Radiator → Incoherent Bremsstrahlung

Diamond Radiator → Coherent Bremss → Linearly Polarized γ beam

Bremsstrahlung spectra



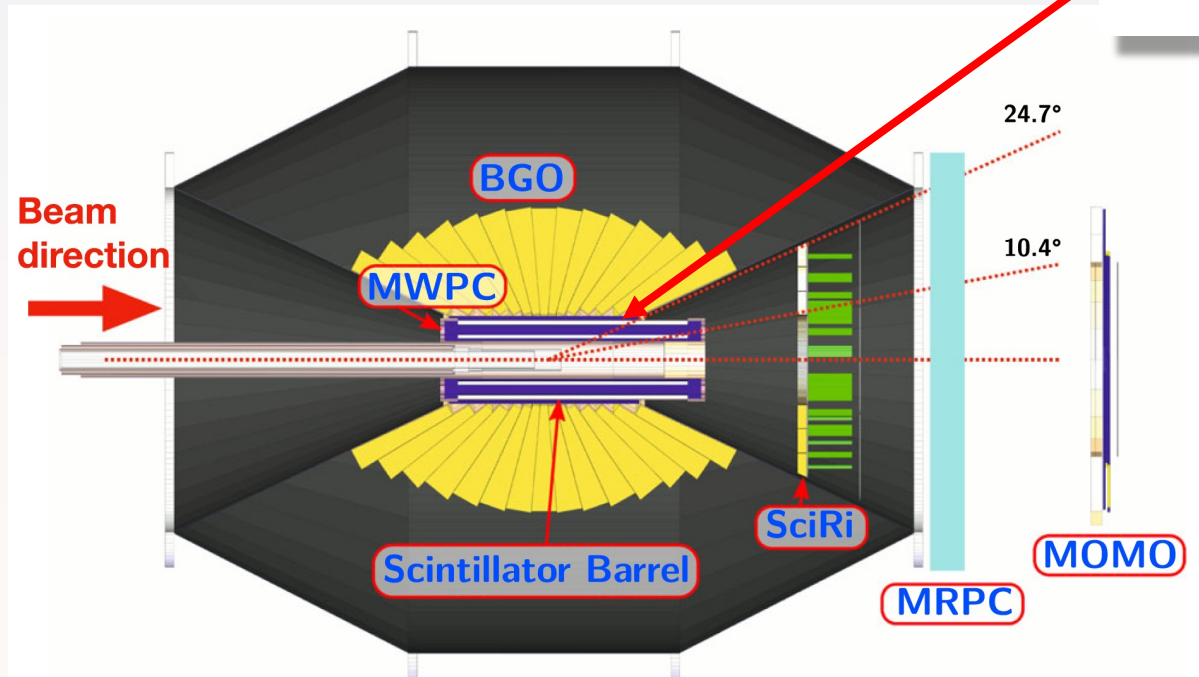
Normalized Diamond Spectra and Polarization



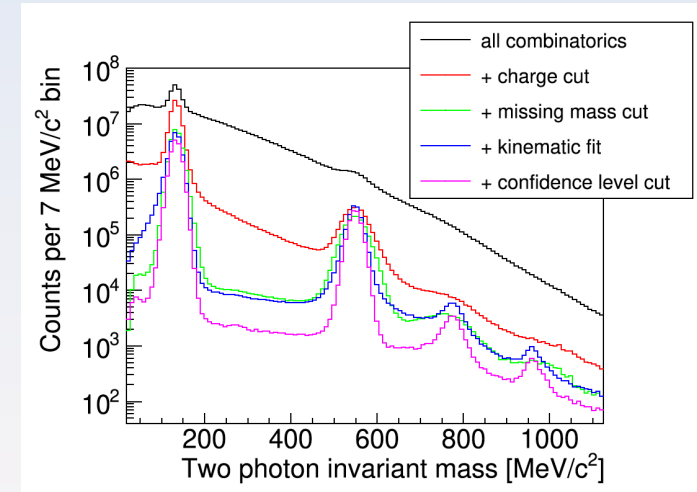
BGOOD Central Detectors:

large solid angle calorimeter:

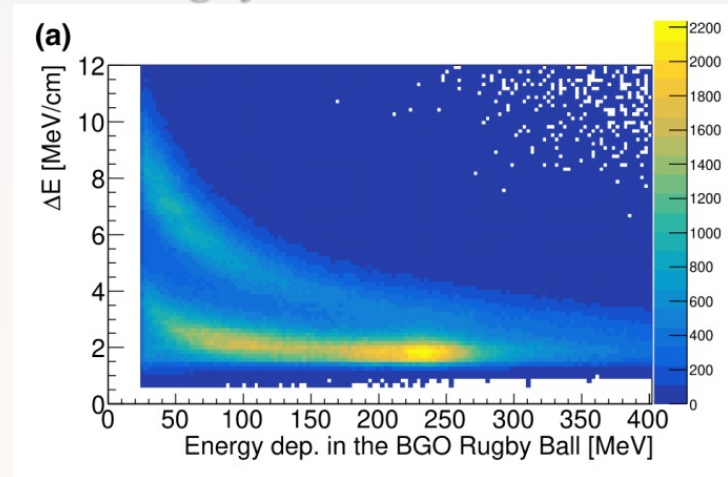
- excellent energy resolution for photons
- good detection efficiency for neutrons
- charged particle tracking and identification
- neutron/photon discrimination



Photons in BGO Rugby Ball

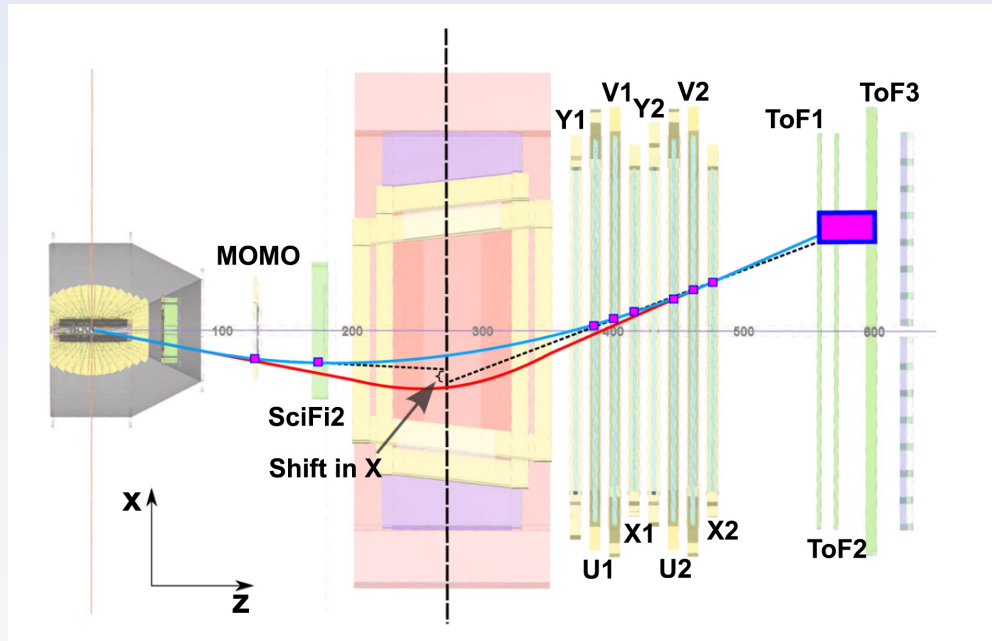


Pid BGORugbyBall-Plastic Scint. Barrel



BGOOD Forward Spectrometer

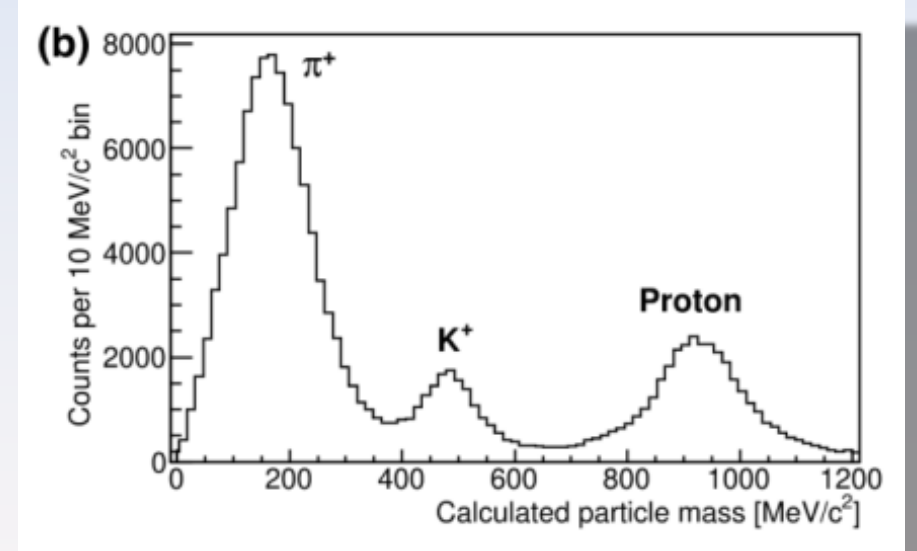
High momentum resolution forward tracking



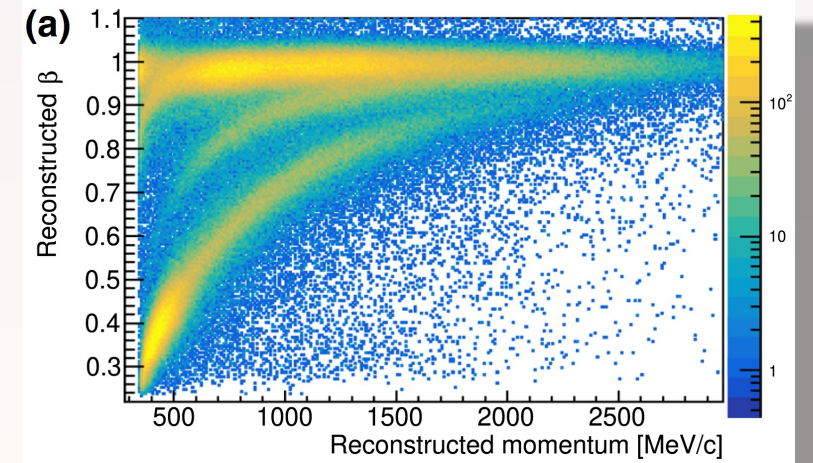
- Charged particles tracking in front of the magnet by means of two scintillating fibre detectors
- Behind the magnet, particle trajectories are determined through eight double layers of drift chambers

Particle identification through velocity measurements with the ToF Walls

Mass from ToF Walls



β vs Momentum in Forw Spectrometer

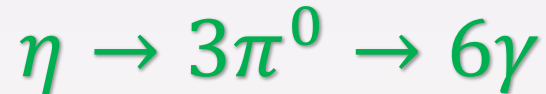
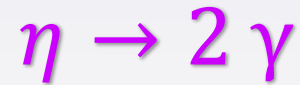


Σ Beam Asymmetry of η photoproduction on the proton



Energy Range: $E_\gamma = 1.2 \div 1.7 \text{ GeV}$

Analyzed simultaneously all main η decay channels:



for

4 data taking periods

- With different Polarization degrees
- With different detection and reconstruction efficiencies

$\vec{\gamma} + p \rightarrow \eta + p$ Events Selection:

1) $\eta \rightarrow 2\gamma$

2 γ detected in the BGO + 1 proton in whole apparatus

2) $\eta \rightarrow 3\pi^0$

6 γ detected in the BGO + 1 proton in whole apparatus

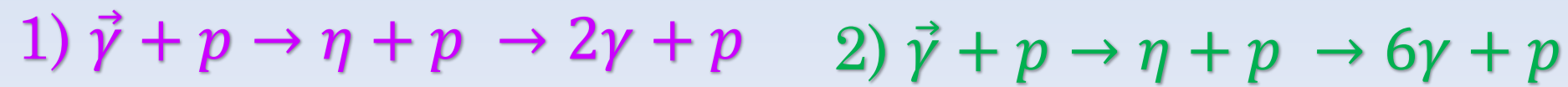
3) $\eta \rightarrow \pi^+ + \pi^- + \pi^0$

2 γ detected in BGO + 1 proton + $\pi^+ + \pi^-$ in whole apparatus

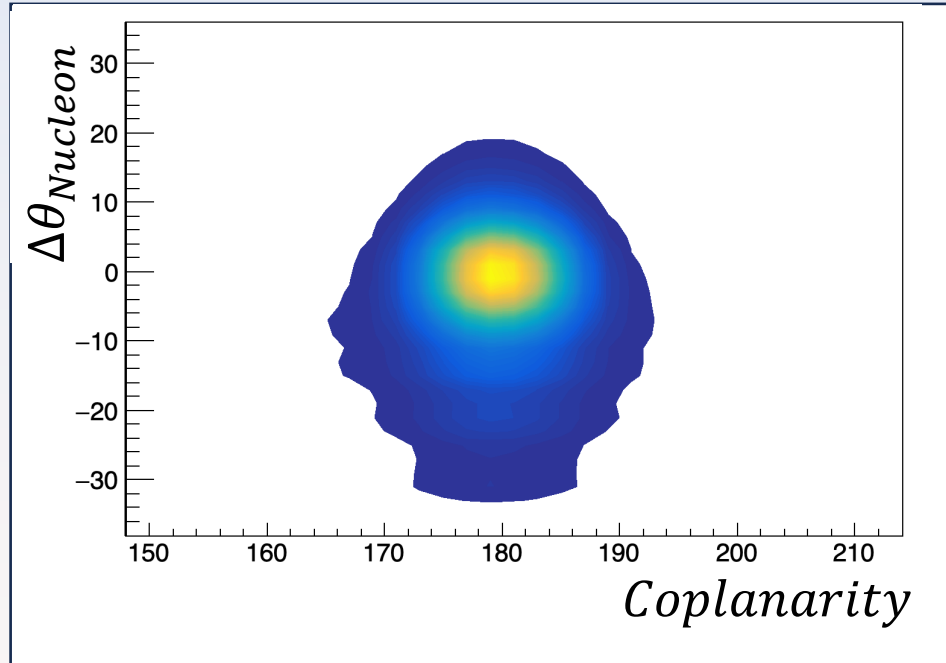
for this last case:

$|p_p|$ $|p_{\pi^+}|$ $|p_{\pi^-}|$ reconstructed from momentum conservation between Initial and Final State
Particles with NO hypothesis on the decaying meson

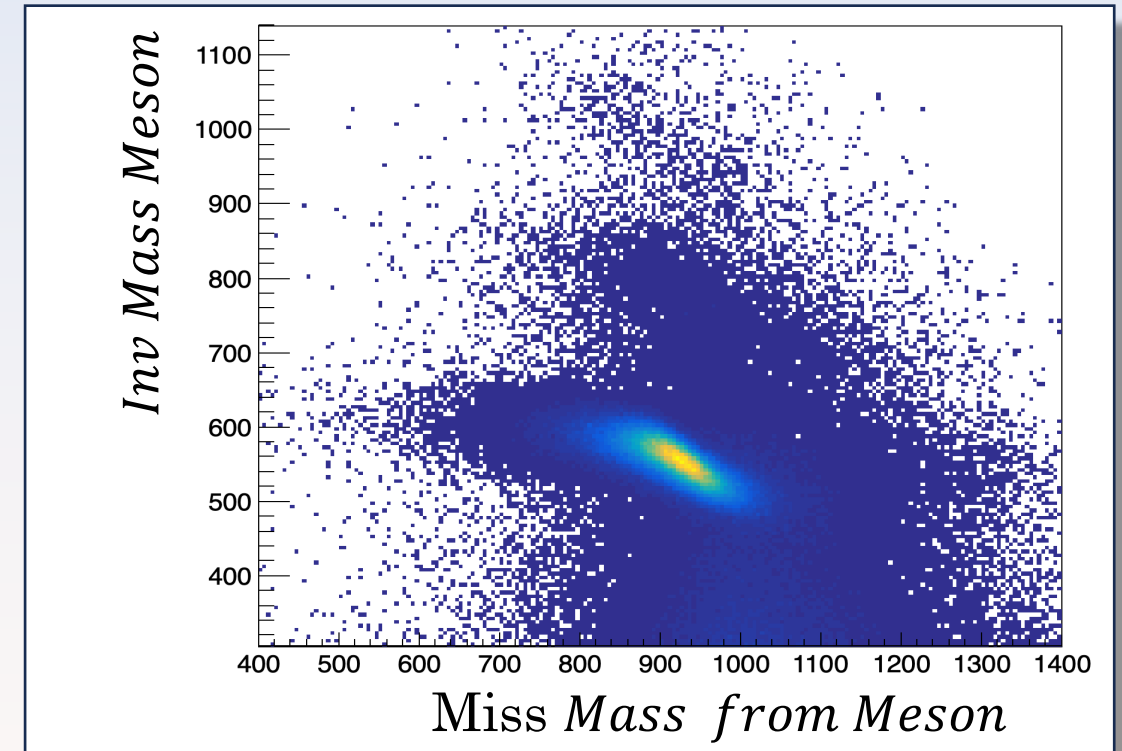
Measured quantities: => *Beam Energy* E_γ
 p and π^\pm angles (θ_p, φ_p) $(\theta_{\pi^+}, \varphi_{\pi^+})$ $(\theta_{\pi^-}, \varphi_{\pi^-})$
2 γ 's energies and angles $(E_{\gamma_1}, \theta_{\gamma_1}, \varphi_{\gamma_1})$ $(E_{\gamma_2}, \theta_{\gamma_2}, \varphi_{\gamma_2})$



$\theta_{missP}^{calc} - \theta_P^{meas}$ % Coplanarity



Meson *InvMass* % *MissMass* from Meson

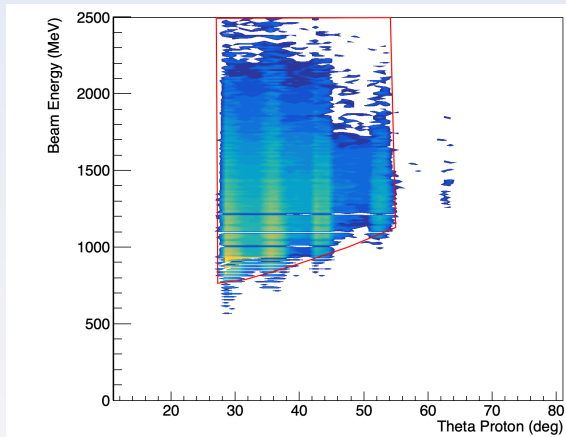


2-body reaction completely defined by the meson reconstructed from the 2 γ 's or 6 γ 's in BGO and the proton detected in whole apparatus.

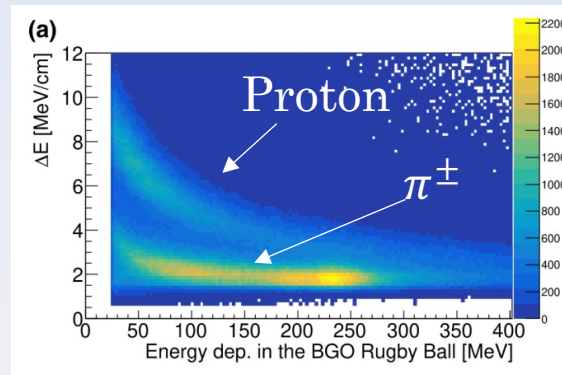
Redundancy of Measured variables => **Clean Event Selection by means of 2D Graphical Cuts**



Preselection

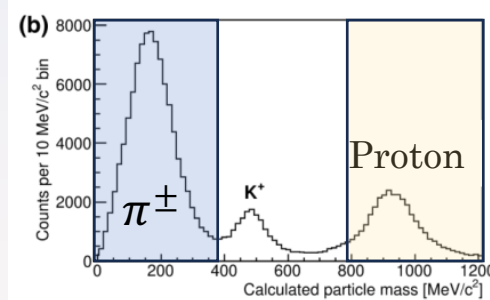


$E_\gamma \% \theta_{PLAB}$

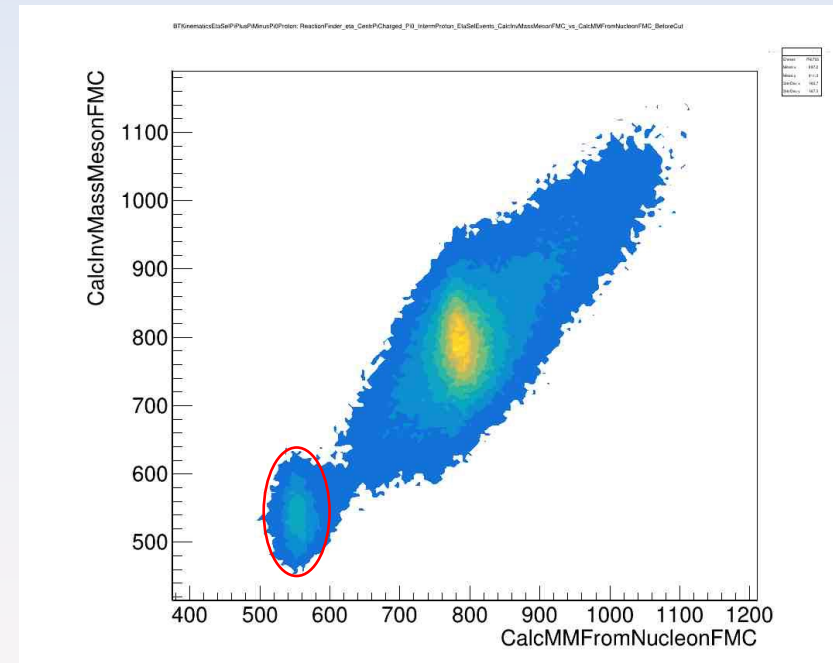


$\Delta E \% E$

PID



Mass from ToF



Preselection \Rightarrow Kinematical cut on $E_\gamma \% \theta_{PLAB}$

Cut on 2γ 's InvMass $\Rightarrow \pi^0$ selection

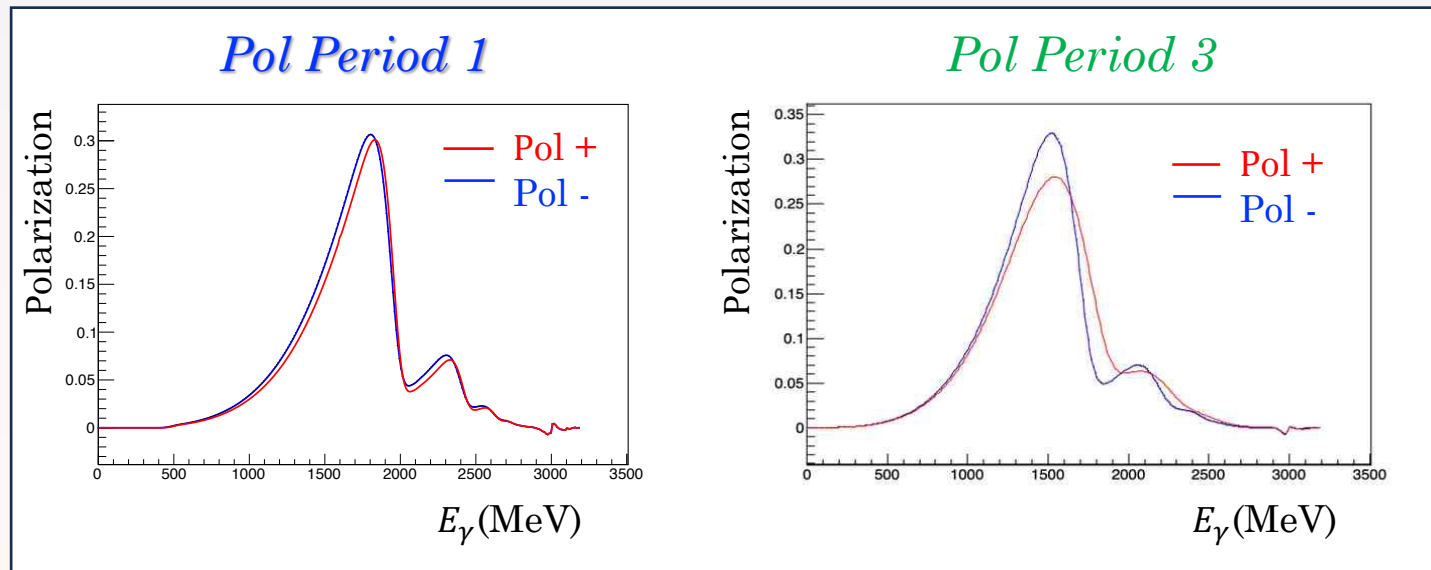
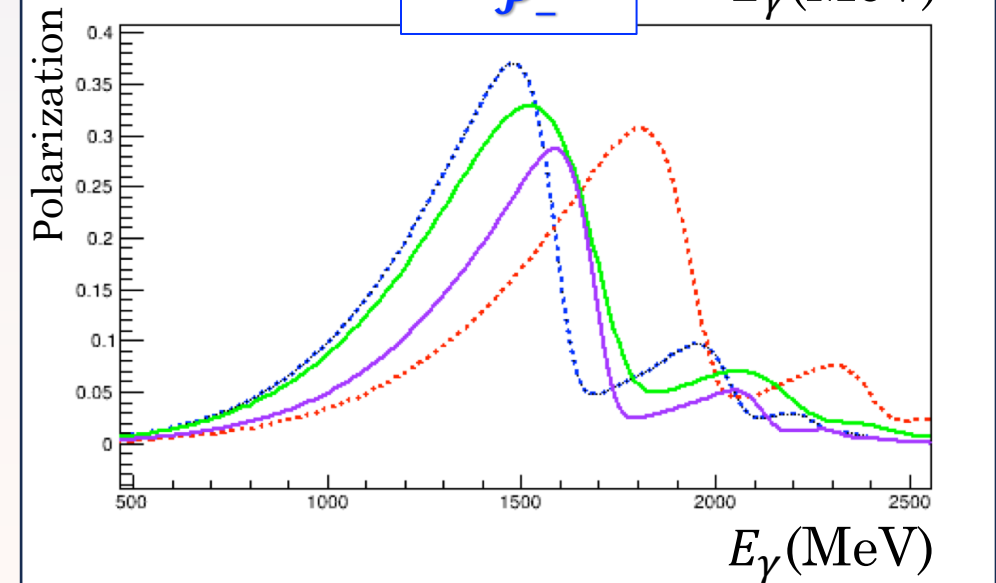
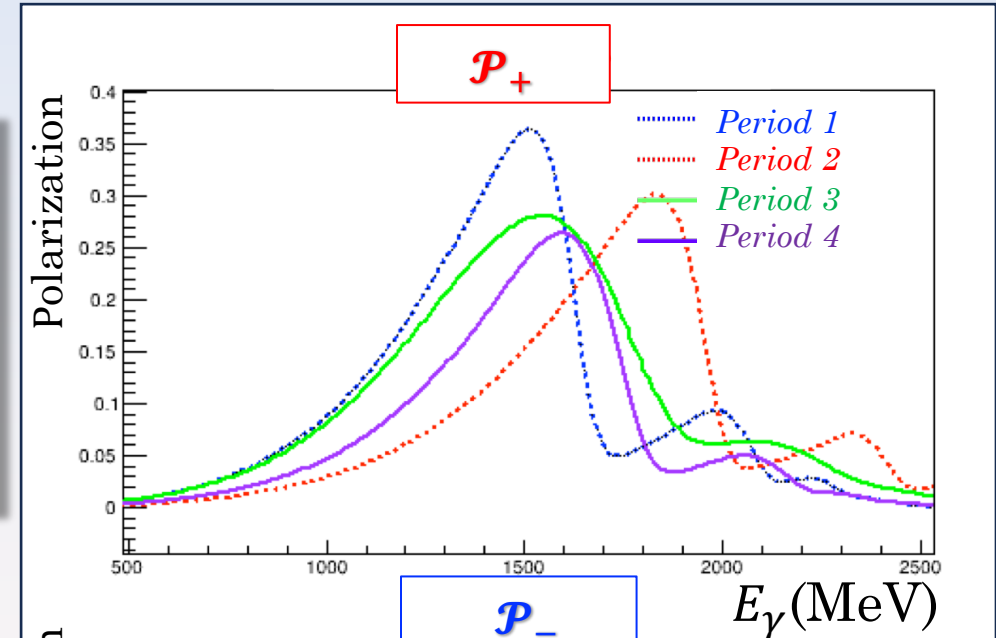
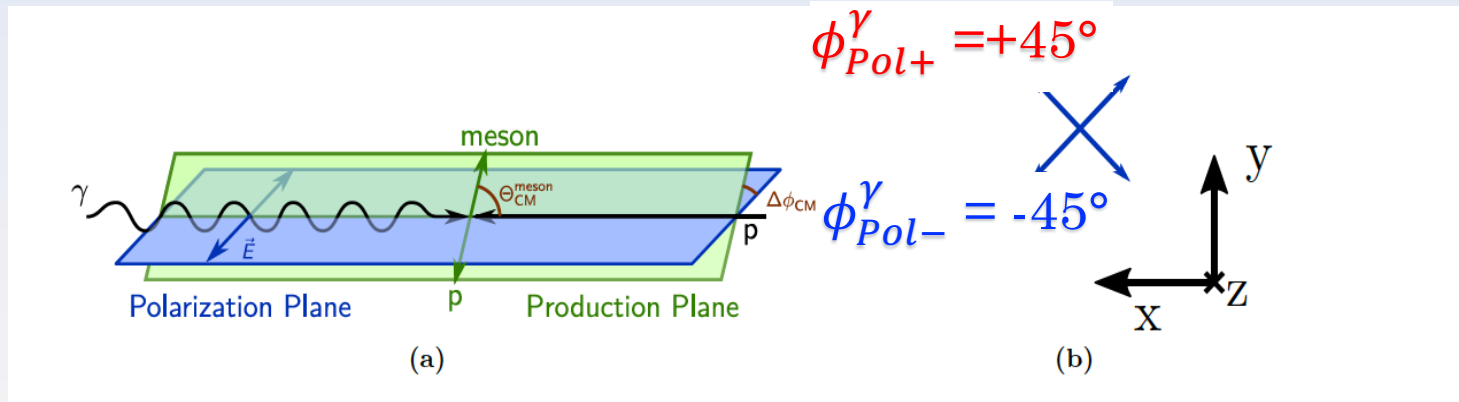
PID Proton - $\pi^\pm \Rightarrow \Delta E \% E$ in BGO-Barrel detectors

Mass from ToF in the Spectrometer

Events selection \Rightarrow 2D-Graph. Cut on Calc Meson Inv Mass % Calc Missing Mass from Proton

Four Data taking periods with different polarized E_γ spectra

During each data period two different polarization planes



«Standard» method for Σ beam asymmetry extraction:

In case of One Data Taking period P with polarization $\mathcal{P}_+^P = \mathcal{P}_-^P = \mathcal{P}^P$

- For fixed $(E_\gamma, \theta_\eta^{cm})$ bins the number of polarized events normalized to the flux is modulated in $\sin(2\varphi_\eta)$

$$\frac{N_\pm^P}{F_\pm^P} = \left(\frac{d\sigma}{d\Omega}\right)_{UNP} \varepsilon^P(\varphi) N_{SC} (1 \mp \mathcal{P}^P \Sigma \cdot \sin(2\varphi_\eta))$$

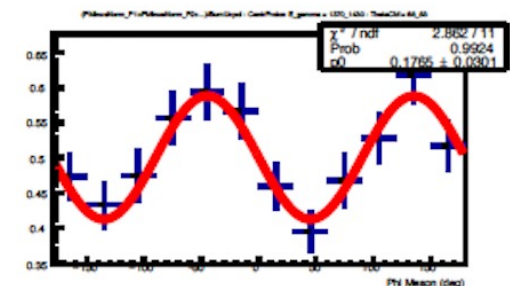
$$** \cos 2(\phi_{pol} - \varphi_\eta^{cm}) \xrightarrow{\phi_{pol} = \pm 45^\circ} \mp \sin(2\varphi_\eta^{cm}) = \mp \sin(2\varphi_\eta)$$

Since the detection and reconstruction efficiency $\varepsilon^P(\varphi)$ can be assumed to be the same for the two polarization states

$$N_{UNP,norm}^P = \frac{N_+^P}{F_+^P} + \frac{N_-^P}{F_-^P} \propto \varepsilon^P(\varphi) \cdot (1 - \cancel{\mathcal{P}^P \Sigma \cdot \sin(2\varphi_\eta)} + 1 + \cancel{\mathcal{P}^P \Sigma \cdot \sin(2\varphi_\eta)}) \Rightarrow N_{UNP,norm}^P \propto 2\varepsilon^P(\varphi_\eta)$$

- For each $(E_\gamma, \theta_\eta^{cm}) \Rightarrow$ The term $\mathcal{P}^P \Sigma$ can be extracted from an azimuthal fit of the ratio

$$\frac{\frac{N_\pm^P}{F_\pm^P}}{N_{UNP,norm}^P} = \frac{\frac{N_\pm^P}{F_\pm^P}}{\frac{N_+^P}{F_+^P} + \frac{N_-^P}{F_-^P}} = \frac{1}{2} (1 \mp \mathcal{P}^P \Sigma \sin(2\varphi_\eta))$$



«Standard» method for Σ beam asymmetry extraction (continued):

In case of a single period of data taking, but with $\mathcal{P}_+^P \neq \mathcal{P}_-^P$

In the ratio:

$$\frac{\frac{N_{\pm}^P}{F_{\pm}^P}}{N_{UNP,norm}^P} = \frac{1}{2} (1 \mp \mathcal{P}_{\pm}^P \Sigma \sin(2\varphi_{\eta}))$$

it is necessary to redefine the expression of the unpolarized term $N_{UNP,norm}$ because the form $N_{UNP,norm}^P = \frac{N_+^P}{F_+^P} + \frac{N_-^P}{F_-^P}$ has a dependence on $\sin(2\varphi)$.

$N_{UNP,norm}^P$ has to be defined as

$$N_{UNP,norm}^P = \frac{2}{\mathcal{P}_+^P + \mathcal{P}_-^P} \left(\mathcal{P}_-^P \frac{N_+^P}{F_+^P} + \mathcal{P}_+^P \frac{N_-^P}{F_-^P} \right) \Rightarrow N_{UNP,norm}^P \propto 2\varepsilon^P(\varphi_{\eta})$$

It is simple to show that this quantity depends on φ_{η} only via the efficiency $\varepsilon^P(\varphi_{\eta})$

Different Polarization spectra for the two polarization states and several periods

For each $(E_\gamma, \theta_\eta^{cm})$ bin we start from the ratio:

$$1) \frac{\sum_i N_{+,norm}^{Pi}}{\sum_i N_{UNP,norm}^{Pi}} = \frac{1}{2} (1 - \mathcal{P}_+^* \Sigma \sin(2\varphi))$$

Where:

- $P_i \Rightarrow$ i-th data taking period (i=1,2,3,4)

- $N_{\pm,norm}^{Pi} = \frac{N_{\pm}^{Pi}}{F_{\pm}^{Pi}} \Rightarrow$ yield of polarized events, normalized to flux for the period P_i

- $N_{UNP,norm}^{Pi}(E_\gamma, \theta_\eta^{cm}, \varphi) = \frac{2}{\mathcal{P}_+^{Pi} + \mathcal{P}_-^{Pi}} \left(\mathcal{P}_-^{Pi} \frac{N_+^{Pi}}{F_+^{Pi}} + \mathcal{P}_+^{Pi} \frac{N_-^{Pi}}{F_-^{Pi}} \right) \Rightarrow N_{UNP,norm}^{Pi}(E_\gamma, \theta_\eta^{cm}, \varphi) \propto \varepsilon^{Pi}(\varphi)$

- $\mathcal{P}_+^* = \frac{\sum_i (\varepsilon^{Pi} \mathcal{P}_+^{Pi})}{\sum_i \varepsilon^{Pi}} \Rightarrow$ sum of polarizations weighted on the efficiencies of the respective periods

- NB !!!!! $\mathcal{P}_+^* = \mathcal{P}_+^*(\varphi) \Rightarrow$ the fit on the ratio 1) is not possible

Σ beam asymmetry extraction method considering several data periods (cont.)

Since $N_{UNP,norm}^{P_i} \propto \varepsilon^{P_i}(\varphi)$:

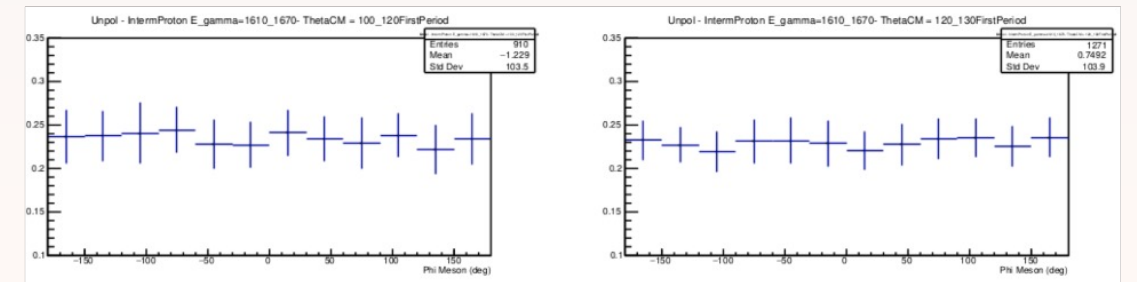
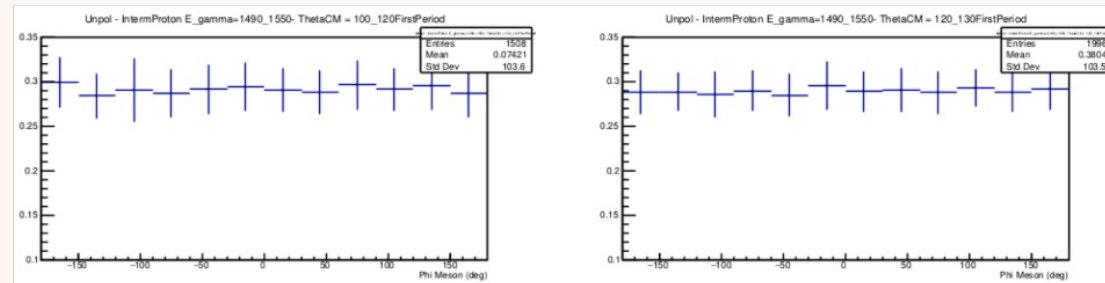
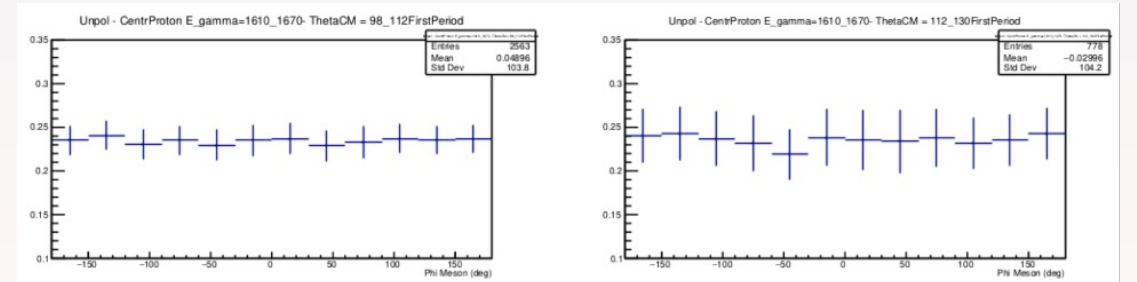
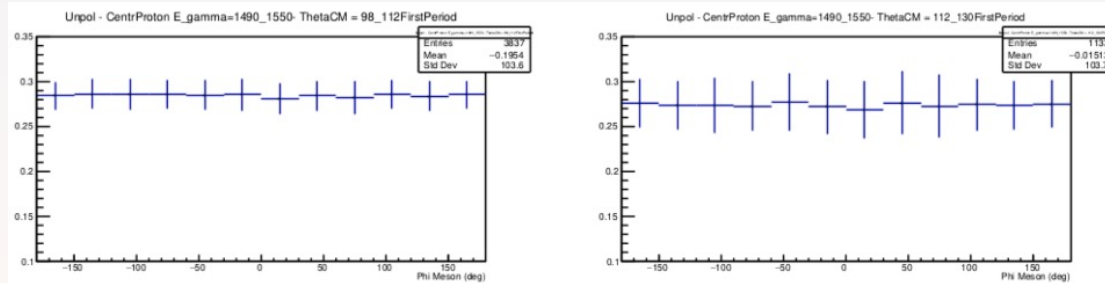
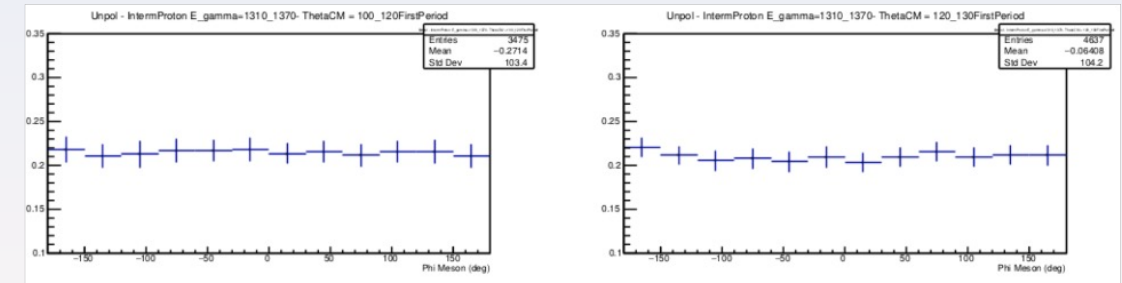
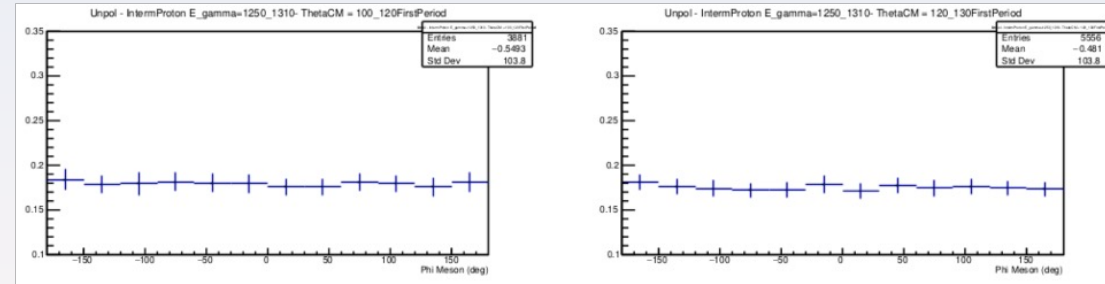
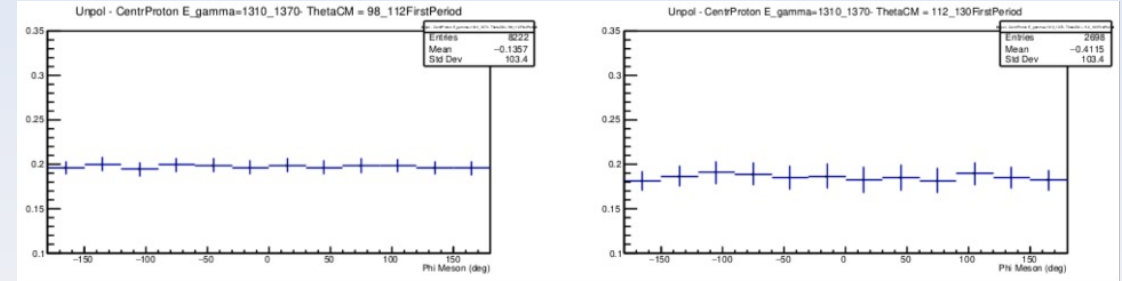
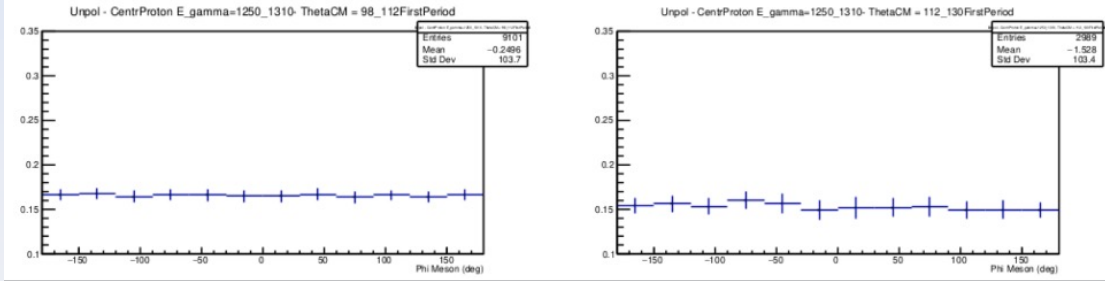
the quantity \mathcal{P}_+^* can be extracted directly from data replacing the efficiency ε^{P_i} by $N_{UNP,norm}^{P_i}$



$$\mathcal{P}_+^* = \frac{\sum_i (\varepsilon^{P_i} \mathcal{P}_+^{P_i})}{\sum_i \varepsilon^{P_i}} = \frac{\sum_i (N_{UNP,norm}^{P_i} \mathcal{P}_+^{P_i})}{\sum_i N_{UNP,norm}^{P_i}}$$

We still have the problem that \mathcal{P}_+^* depends on φ , but it is possible to show that the behavior of $\mathcal{P}_+^*(\varphi)$ is quite flat as function of the azimuthal angle φ for each $(E_\gamma, \theta_\eta^{cm})$

$$\mathcal{P}_+^* = \mathcal{P}_+(\varphi)$$



Σ beam asymmetry extraction method considering several data periods (cont.)

Since $N_{UNP,norm}^{P_i} \propto \varepsilon^{P_i}(\varphi)$:

the quantity \mathcal{P}_+^* can be extracted directly from data replacing the efficiency ε^{P_i} by $N_{UNP,norm}^{P_i}$



$$\mathcal{P}_+^* = \frac{\sum_i (\varepsilon^{P_i} \mathcal{P}_+^{P_i})}{\sum_i \varepsilon^{P_i}} = \frac{\sum_i (N_{UNP,norm}^{P_i} \mathcal{P}_+^{P_i})}{\sum_i N_{UNP,norm}^{P_i}}$$

We still have the problem that \mathcal{P}_+^* depends on φ , but it is possible to show that the behavior of $\mathcal{P}_+^*(\varphi)$ is quite flat as function of the azimuthal angle φ for each $(E_\gamma, \theta_\eta^{cm})$



It is possible to make the average of \mathcal{P}_+^* all over the φ bins.

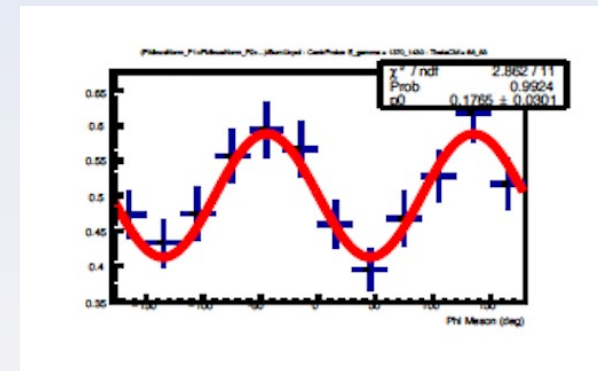
$$\overline{\mathcal{P}_+^*} = \frac{\sum_\varphi \mathcal{P}_+^*(\varphi)}{N_{\varphi \text{ bins}}}$$

Σ beam asymmetry extraction method considering several data periods (Cont.)

For each $(E_\gamma, \theta_\eta^{cm})$ bin we can then obtain *the Σ beam asymmetry* by:

1) extracting $\overline{\mathcal{P}}_{\pm}^* \Sigma$ from the azimuthal fit of the ratio:

$$\frac{\sum_i N_{\pm, norm}^{Pi}}{\sum_i N_{UNP, norm}^{Pi}} = \frac{1}{2} \left(1 \mp \overline{\mathcal{P}}_{\pm}^* \Sigma \sin(2\varphi) \right)$$



2) Dividing the value $\overline{\mathcal{P}}_{\pm}^* \Sigma$ by $\overline{\mathcal{P}}_{\pm}^*$

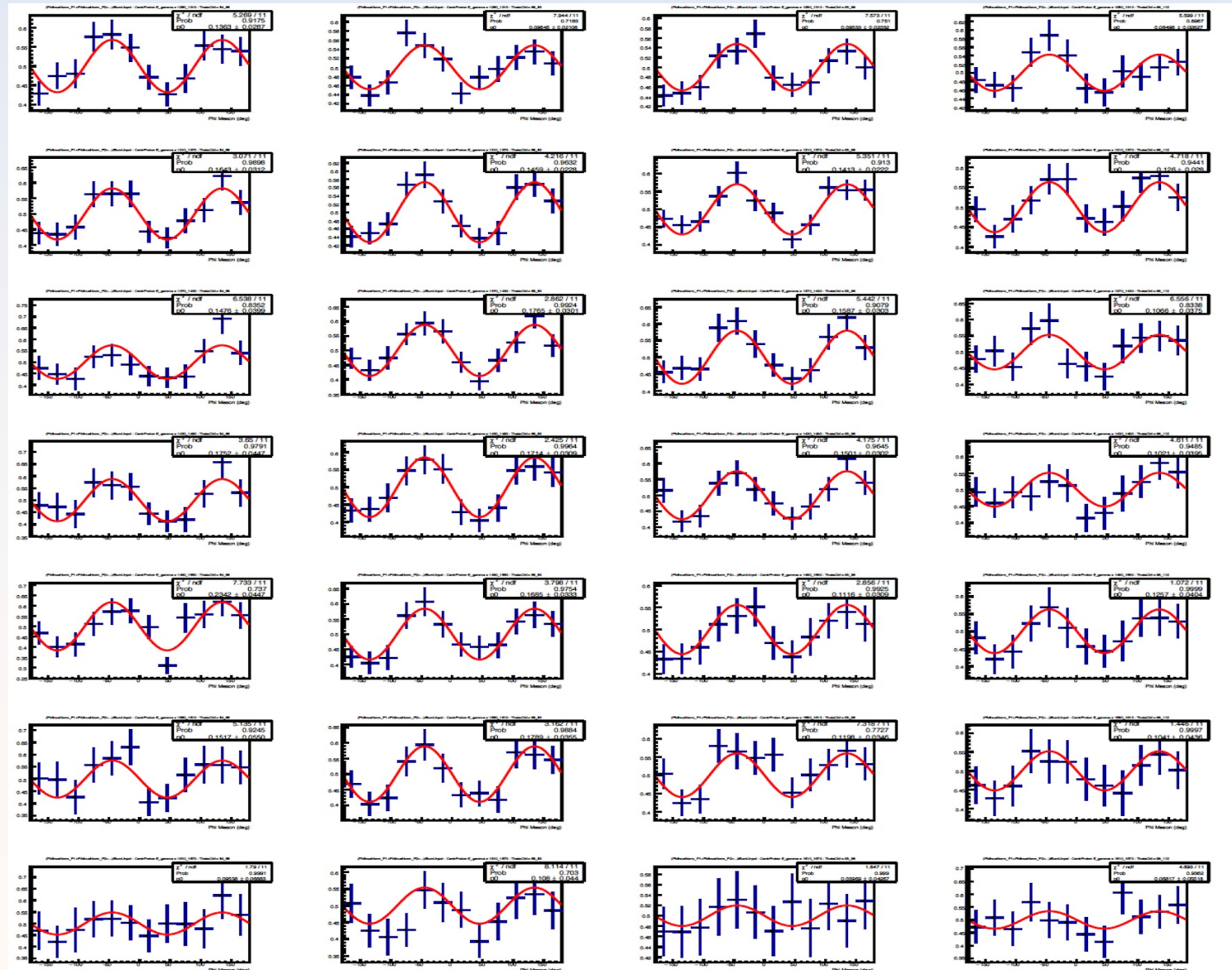
$$\Sigma = \frac{\overline{\mathcal{P}}_{\pm}^* \Sigma}{\overline{\mathcal{P}}_{\pm}^*}$$

NB: Since we analyze several η decay channels: $\eta \rightarrow 2\gamma$, $\eta \rightarrow 3\pi^0 \rightarrow 6\gamma$, $\eta \rightarrow \pi^+ + \pi^- + \pi^0$

$\varepsilon^{Pi}(\varphi)$ is the "global efficiency" \Rightarrow i.e the sum of the efficiencies of each channel

$$\varepsilon^{Pi}(\varphi) = \varepsilon^{Pi, C1}(\varphi) + \varepsilon^{Pi, C2}(\varphi) + \varepsilon^{Pi, C3}(\varphi)$$

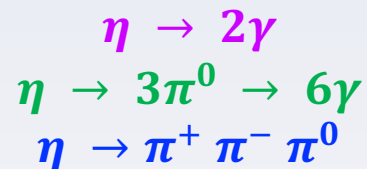
Azimuthal Fit of ratio $\sum_i N_{+,norm}^{Pi} / \sum_i N_{UNP,norm}^{Pi}$



Σ Beam Asymmetry Results (Preliminary)

8 E_γ bins (60 MeV)

11 θ_η^{cm} bins



4 ANALYZED PERIODS

selected events:Period1: 31747

selected events:Period2: 54598

selected events:Period3: 144879

selected events:Period4: 90579

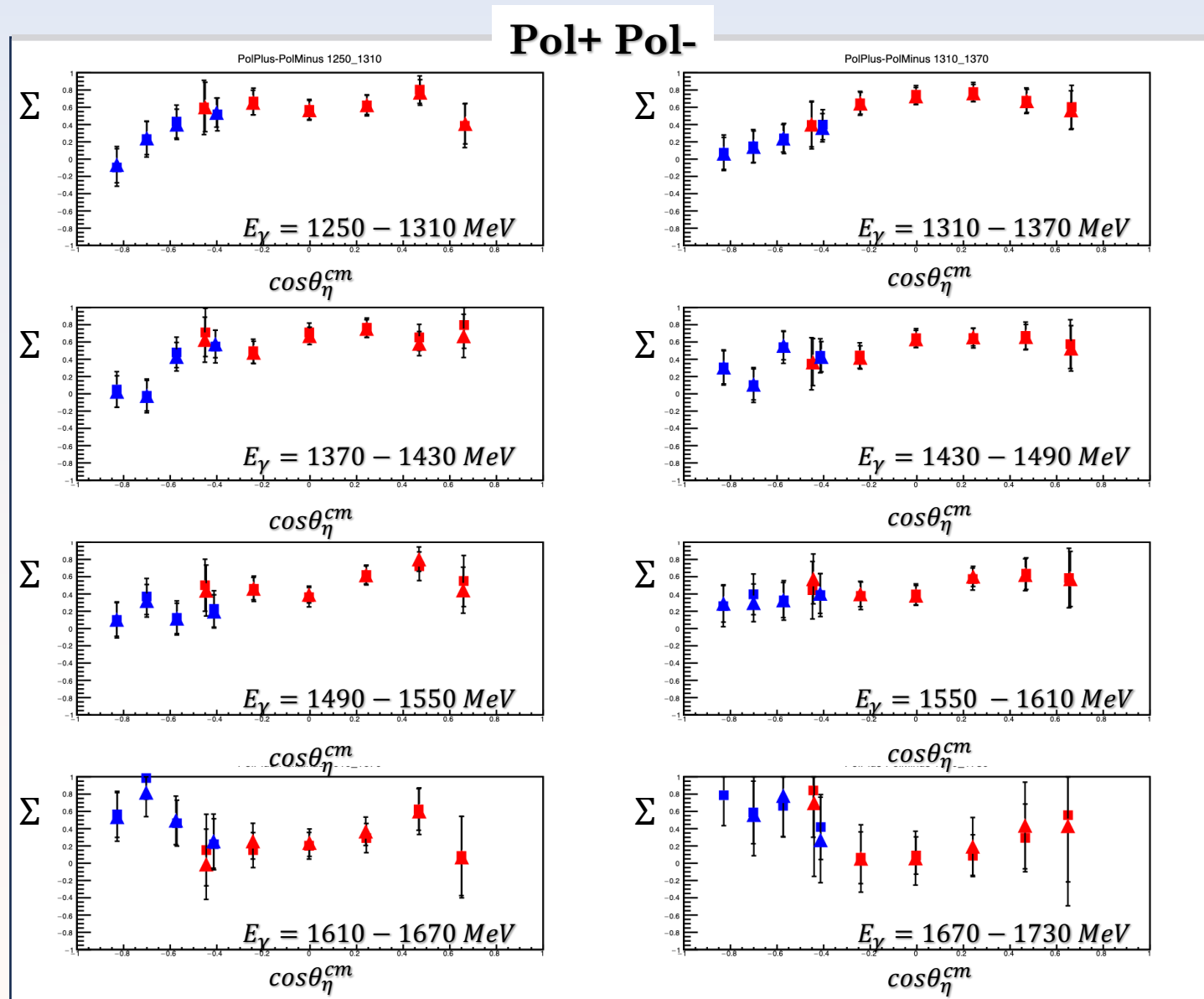
Total # selected events 321803

Pol- ▲ Proton in Intermediate Reg.

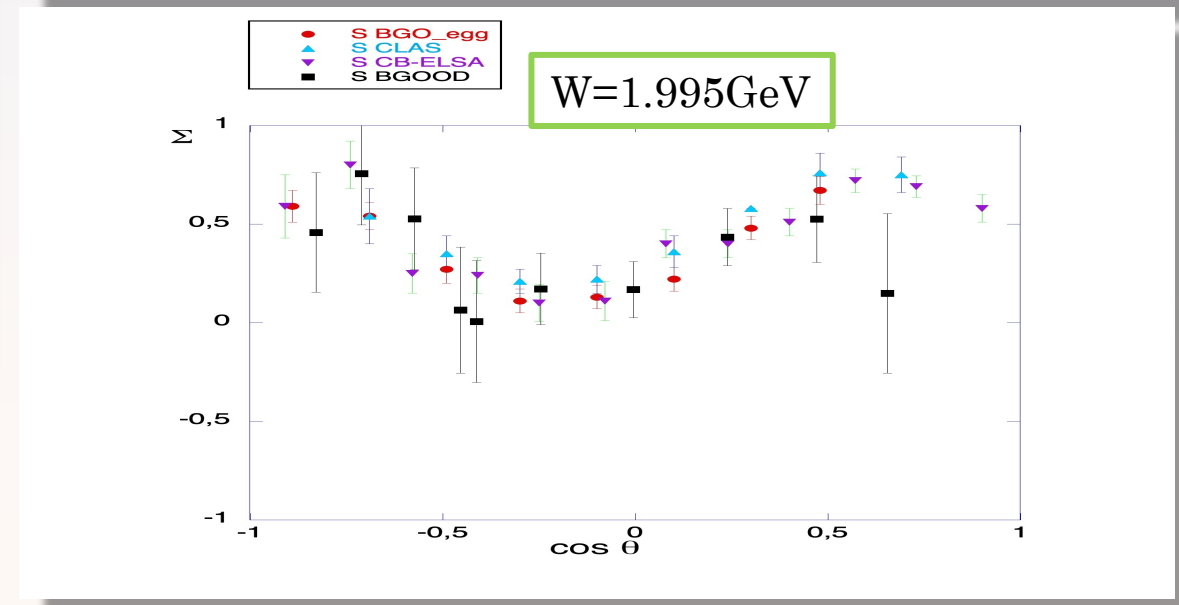
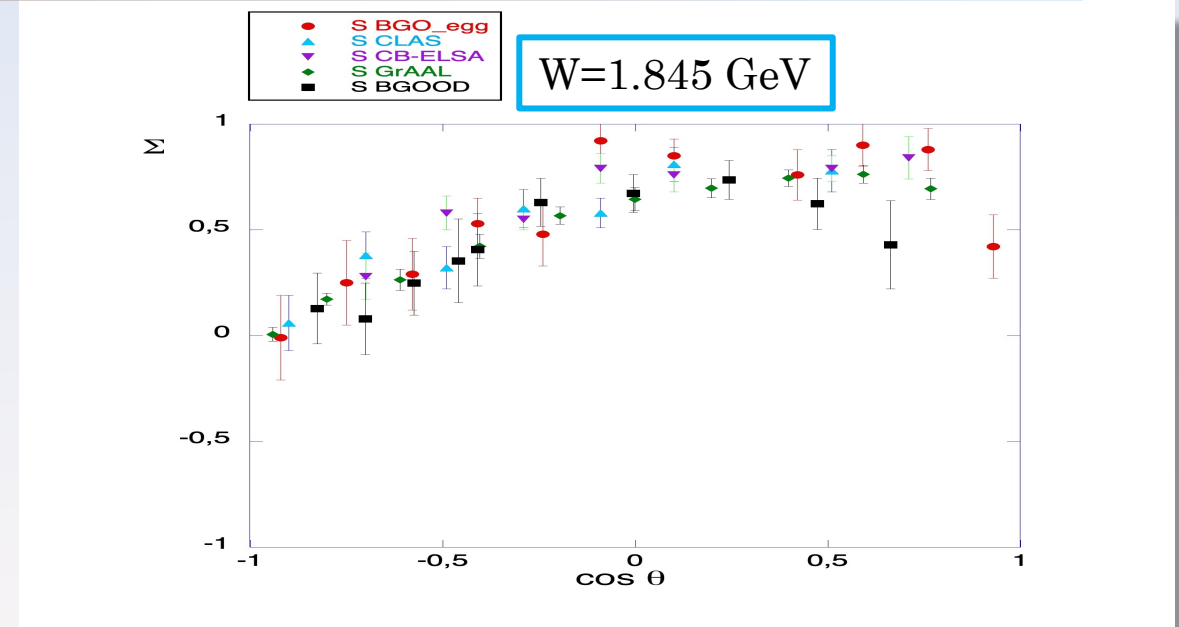
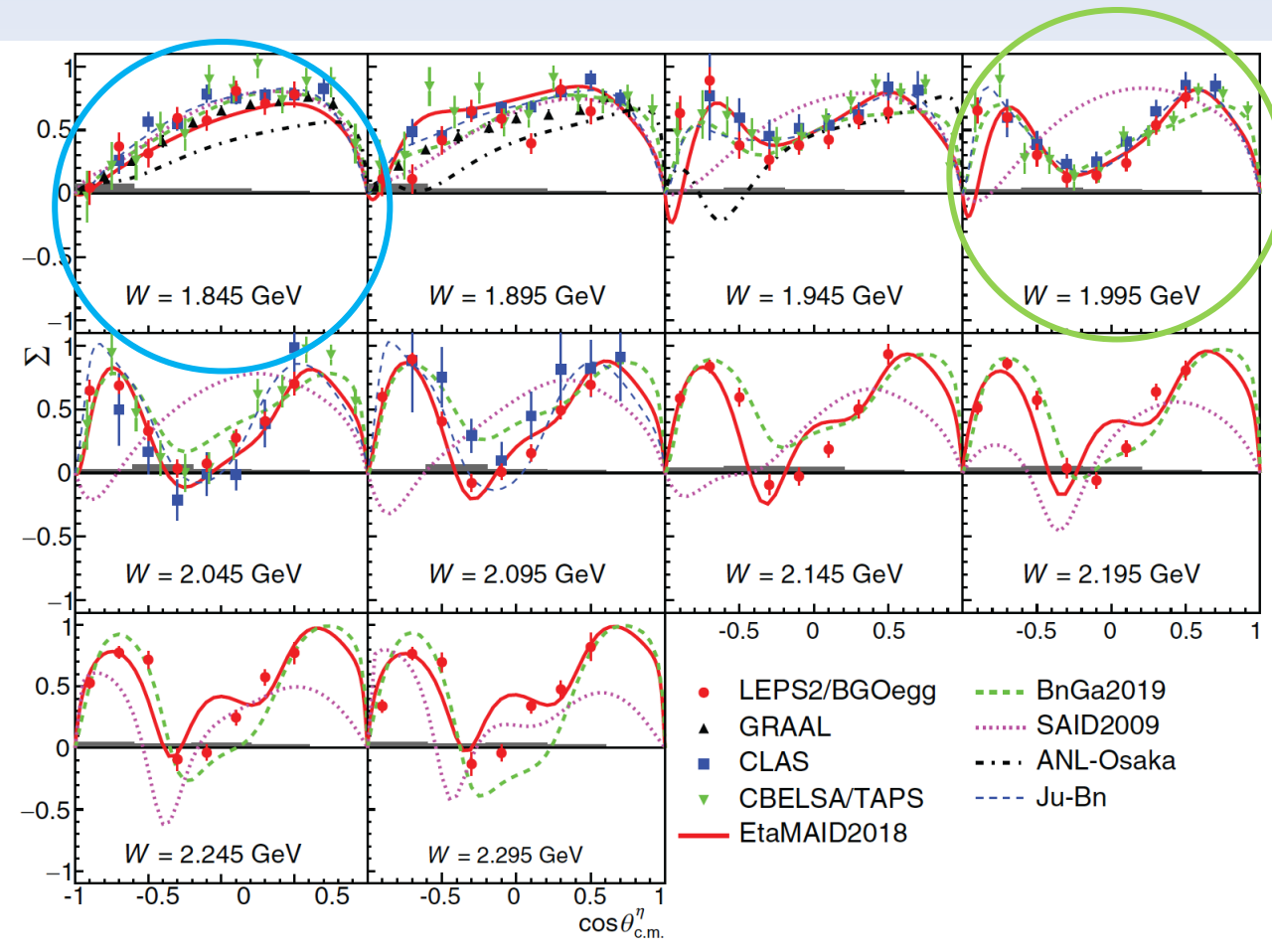
▲ Proton in BGO

■ Proton in Intermediate Reg.

Pol+ ■ Proton in BGO



Comparison with data from GRAAL-CLAS-CBELSA-LEP2



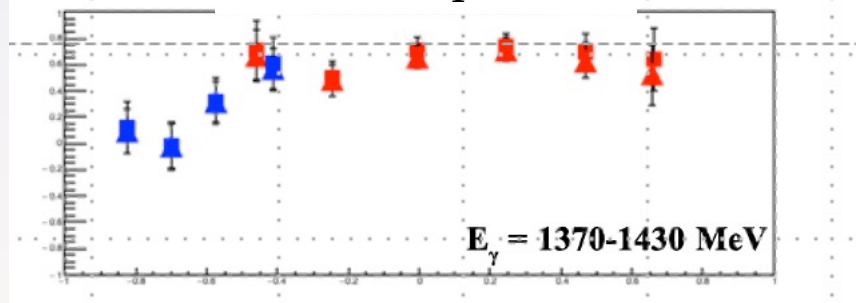
PHYSICAL REVIEW C 106, 035201 (2022)

T. Hashimoto et al. (LEPS2/BGOegg Collaboration)

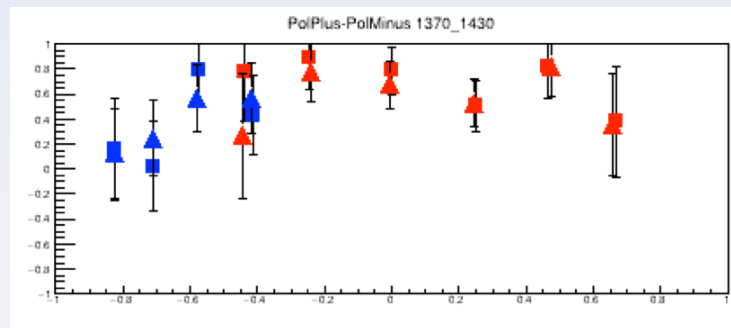
Comparison between Σ extracted analyzing 4 periods together or individually

$$E_\gamma = 1370 \div 1430 \text{ MeV}$$

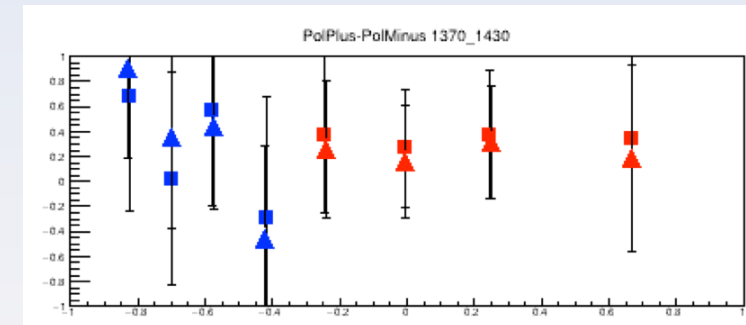
All the 4 periods



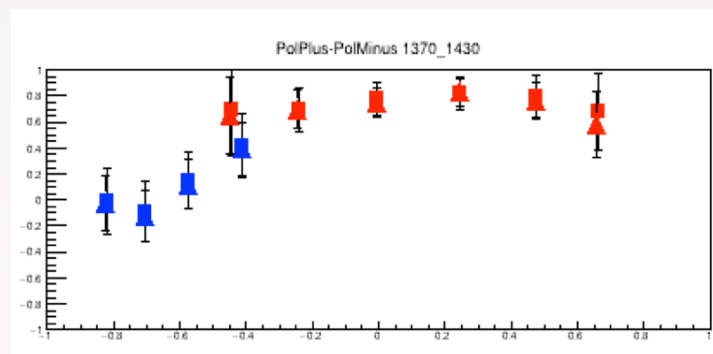
Period 1



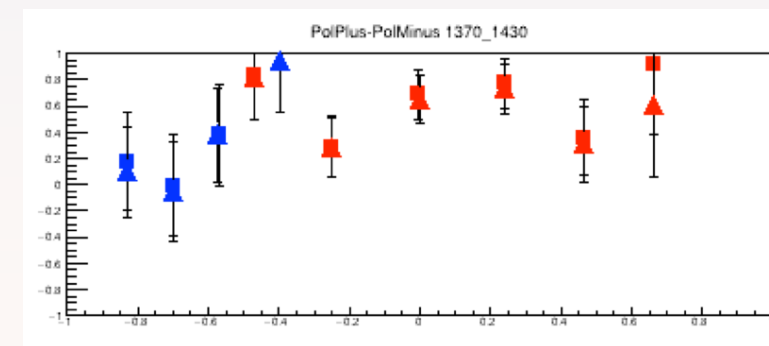
Period 2



Period 3



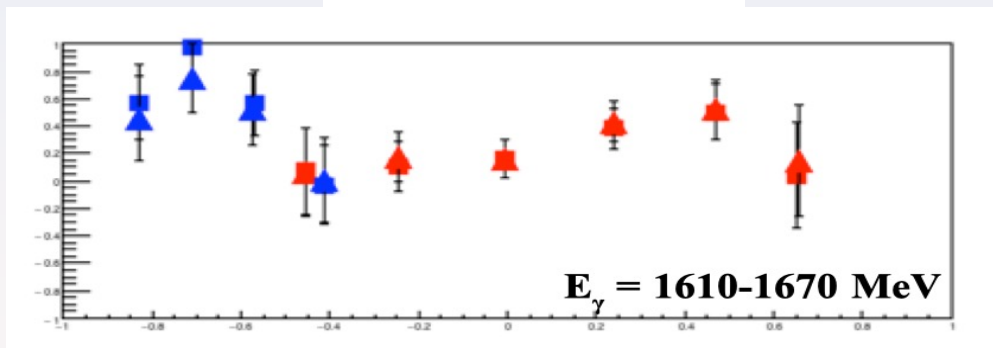
Period 4



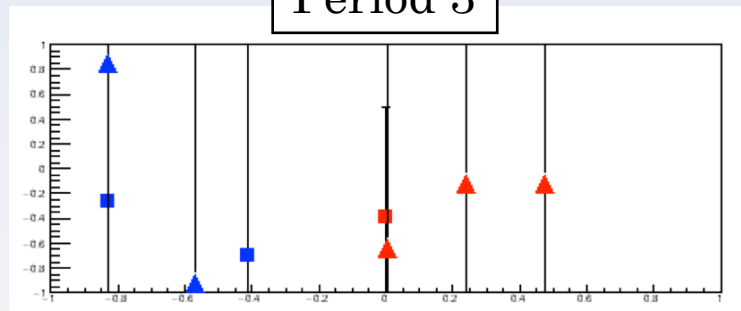
Comparison between Σ extracted analyzing 4 periods all together or individually

$$E_\gamma = 1610 \div 1670 \text{ MeV}$$

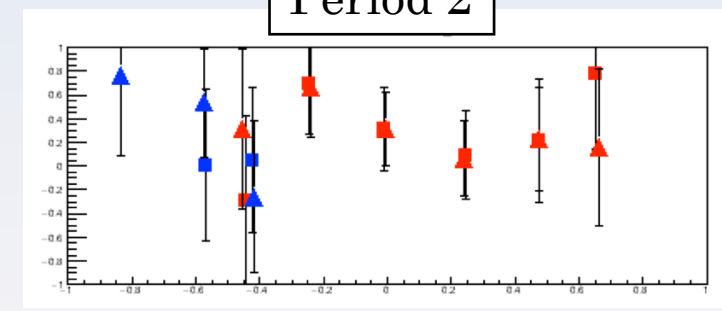
All the 4 periods



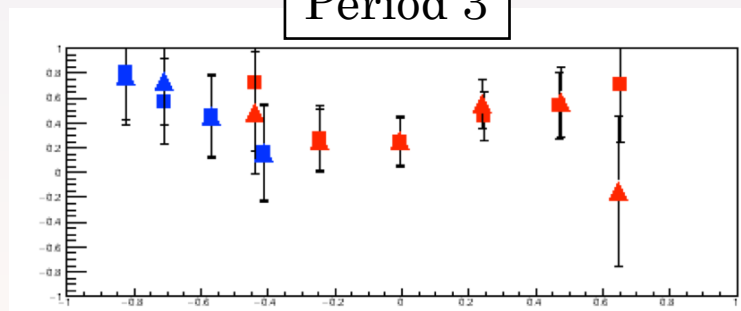
Period 3



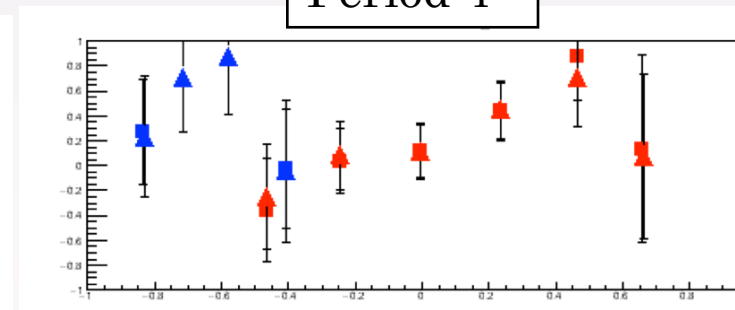
Period 2



Period 3



Period 4

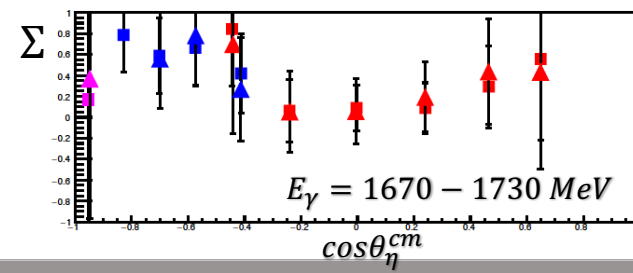
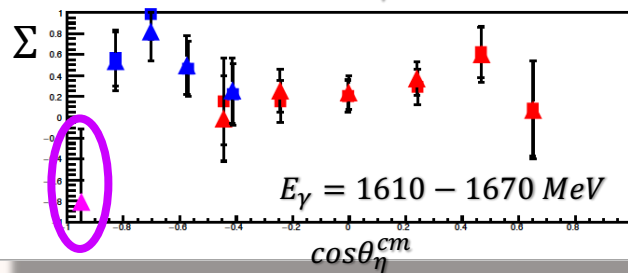
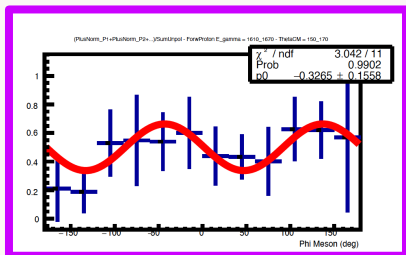
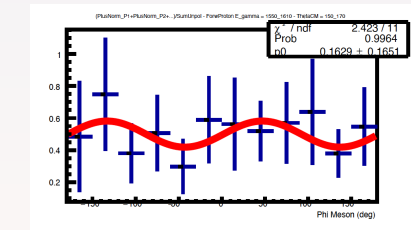
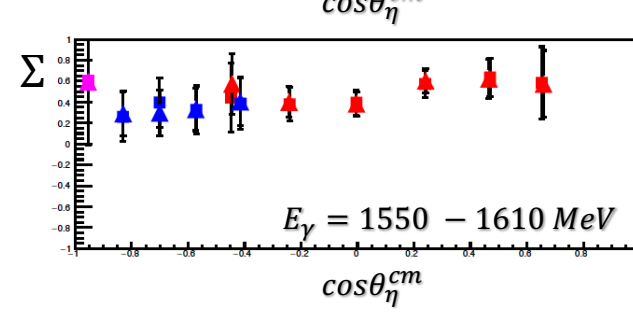
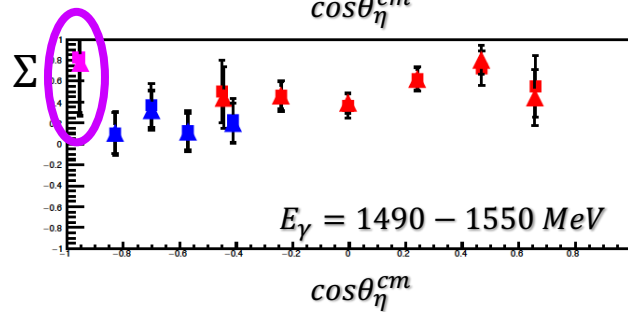
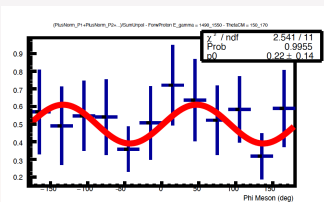
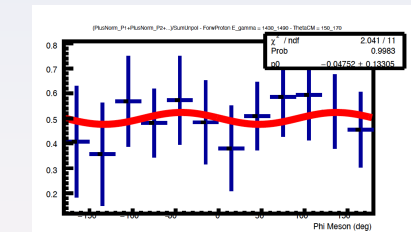
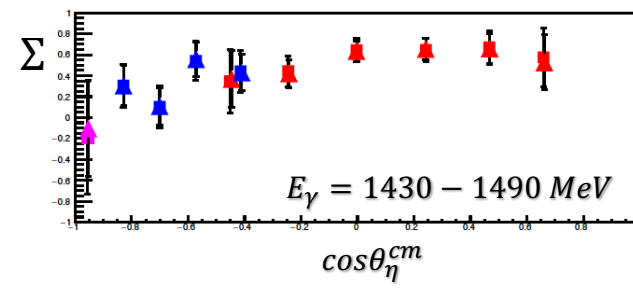
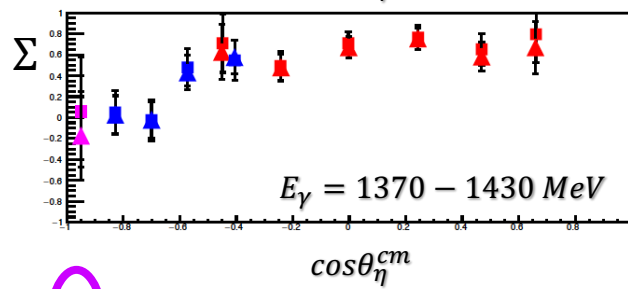
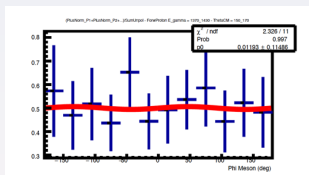
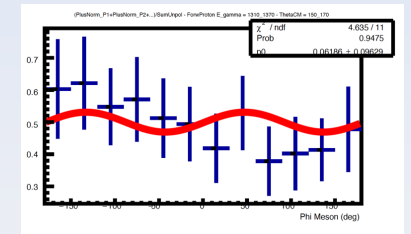
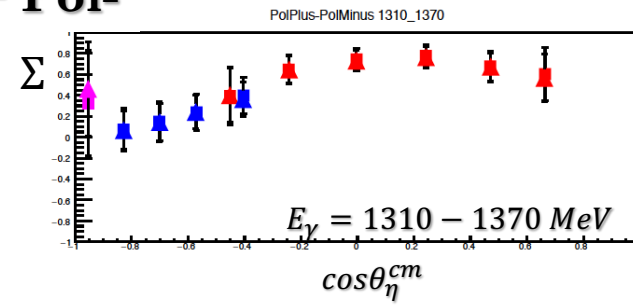
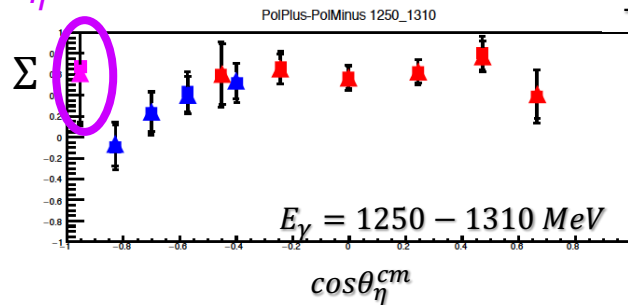
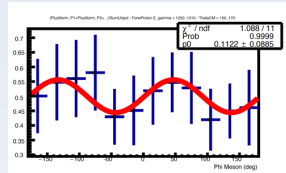


Σ Beam Asymmetry Results (VERY Preliminary)

- ▲ Proton in Forward region
- Proton in Forward region

$$\overline{\theta}_\eta^{cm} \approx 160^\circ$$

Pol+ Pol-

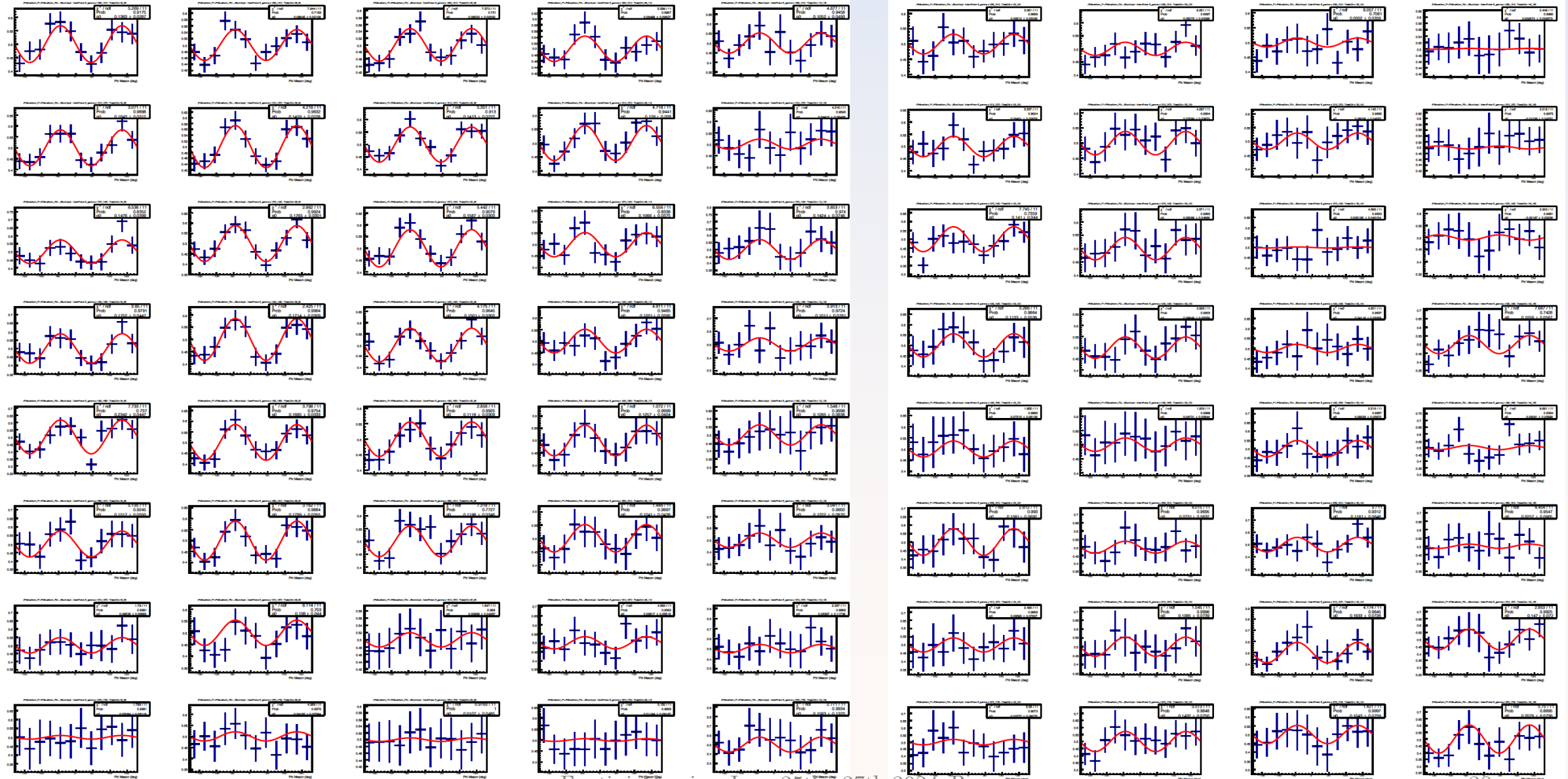


We extracted the Beam Asymmetry Σ for η photoproduction on the proton, with a technique which allows to analyze simultaneously data:

- With Different polarization Degrees
- With Different Detection and Reconstruction Efficiencies
- From Different Eta decay channels
- This technique permits to considerably improve the statistics and opens the possibility for us to analyze η photoproduction on neutron and η' photoproduction on proton.
- The work is in progress => systematic errors estimation & checks with simulation, other two periods have to be analyzed for η photoproduction on the proton and data taking is still in progress

- Backup

Σ beam asymmetry extraction method considering different data periods



$$\vec{\gamma} + p \rightarrow \eta + p \rightarrow \pi^+ + \pi^- + \pi^0 + p$$

2 γ in BGO (π^0) + 1 proton in Centr/Interm Det $\pi^+ \pi^-$ everywhere (most in BGOBarrel)

We miss the energies of pions and proton => apply momentum conservation along the three directions and we get a system of three equations:

$$0 = \begin{cases} 0 = p_P \sin\theta_P \cos\varphi_P + E_{\gamma_1} \sin\theta_{\gamma_1} \cos\varphi_{\gamma_1} + E_{\gamma_2} \sin\theta_{\gamma_2} \cos\varphi_{\gamma_2} + p_{\pi^+} \sin\theta_{\pi^+} \cos\varphi_{\pi^+} + p_{\pi^-} \sin\theta_{\pi^-} \cos\varphi_{\pi^-} \\ 0 = p_P \sin\theta_P \sin\varphi_P + E_{\gamma_1} \sin\theta_{\gamma_1} \sin\varphi_{\gamma_1} + E_{\gamma_2} \sin\theta_{\gamma_2} \sin\varphi_{\gamma_2} + p_{\pi^+} \sin\theta_{\pi^+} \sin\varphi_{\pi^+} + p_{\pi^-} \sin\theta_{\pi^-} \sin\varphi_{\pi^-} \\ E_\gamma = p_P \cos\theta_P + E_{\gamma_1} \cos\theta_{\gamma_1} + E_{\gamma_2} \cos\theta_{\gamma_2} + p_{\pi^+} \cos\theta_{\pi^+} + p_{\pi^-} \cos\theta_{\pi^-} \end{cases}$$

Where:

- Measured quantities:

proton and charged pions angles (θ_P, φ_P) ($\theta_{\pi^+}, \varphi_{\pi^+}$) ($\theta_{\pi^-}, \varphi_{\pi^-}$)

photons energies and angles ($E_{\gamma_1}, \varphi_{\gamma_1}$) ($E_{\gamma_2}, \varphi_{\gamma_2}$)

- Unknown quantities:

proton and charged pions momenta ($p_P, p_{\pi^+}, p_{\pi^-}$)

- NO HYPOTHESIS on the nature of the meson decaying into $\pi^+ \pi^- \pi^0$