

Will be related to Lorenz's talk! arxiv:2407.07873

NETS

A Non-Equilibrium Transport Sampler

aka an Answer to Tej's question of the connection between flows and diffusions for sampling

aka a continuous time algorithm for what Alessandro is doing

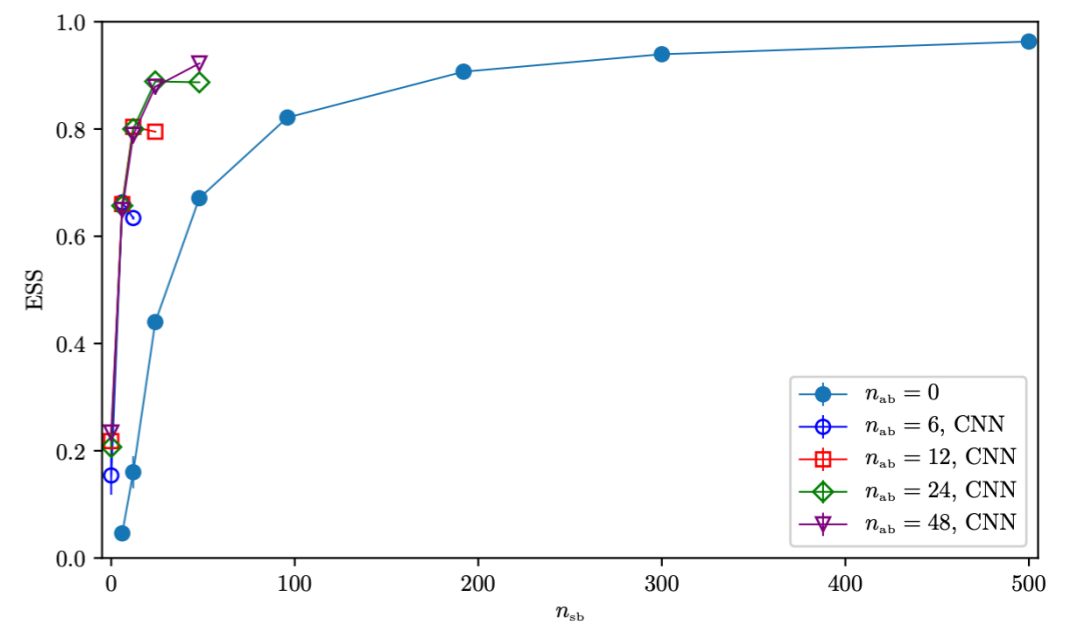
Michael Albergo Bonn, Germany October 24 2024

Remember this from yesterday?

Stochastic normalizing flows as non-equilibrium transformations

Michele Caselle^{1,2,*} Elia Cellini^{1,2,†} Alessandro Nada^{1‡} and Marco Panero^{1,2§}

Improved ESS by growing the # of discrete affine flows + stochastic steps



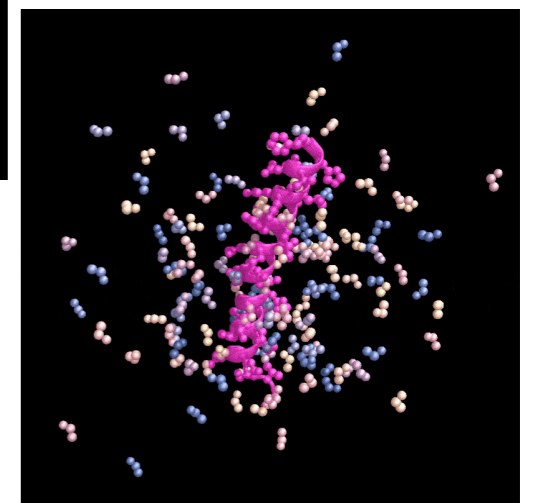
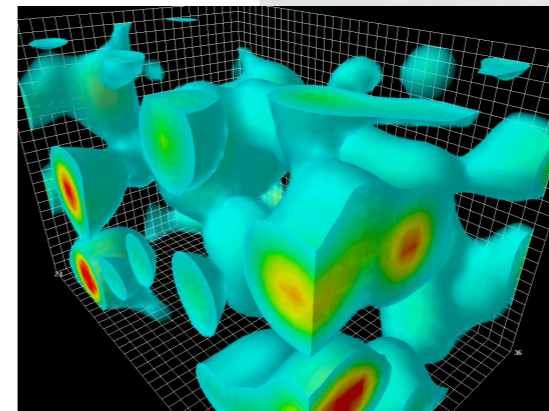
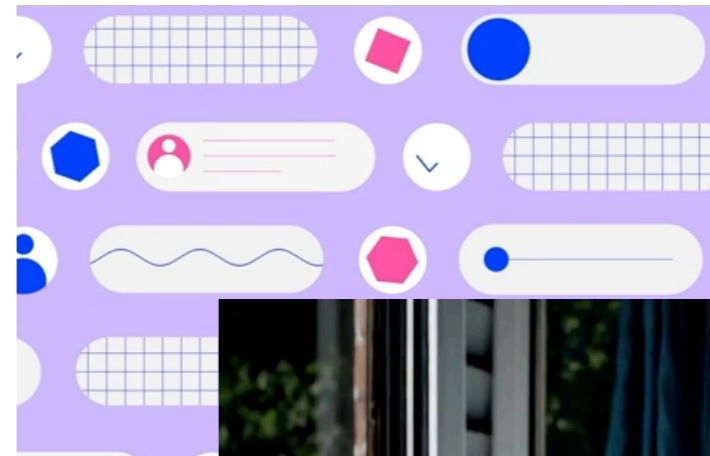
NETS is a continuous time limit of SNFs

- Can choose how many steps + diffusion *after training*
- Knob to explicitly get more performance from more compute

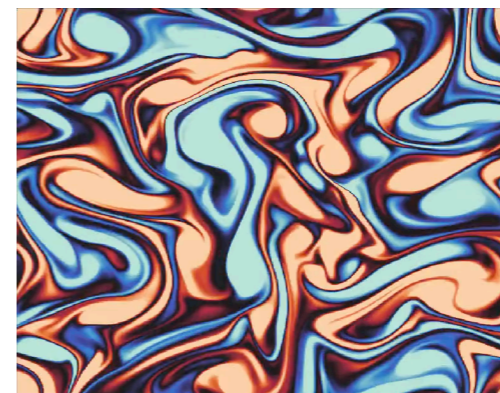
Advertisement: New Research group

In 2026 I will be starting a group at Harvard in Applied mathematics + Kempner Institute

- *Theme: Nature and Computation*
- *Interdisciplinary! Computationally inclined, mathematically inclined welcome*
- *Current undergraduates, master's students, graduating PhDs, and postdocs, please reach out if interested*
- *Advisors, please forward your students :)*



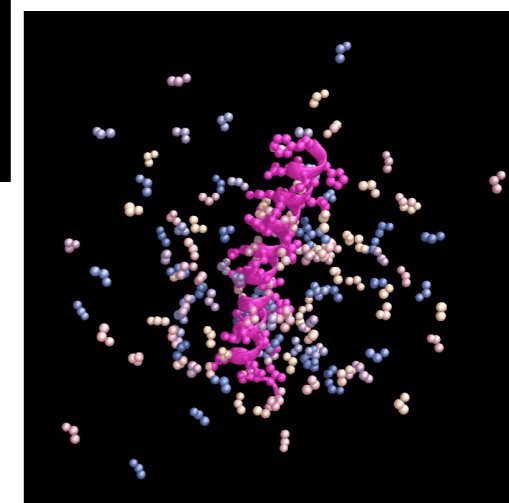
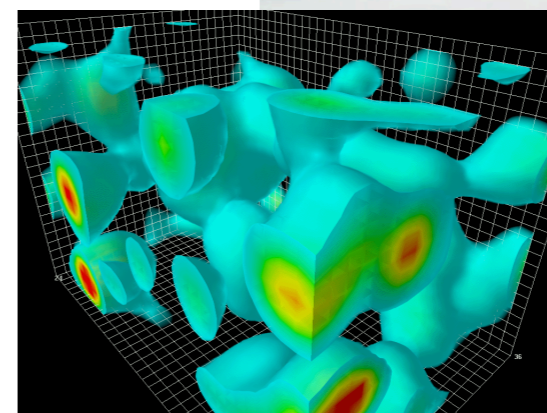
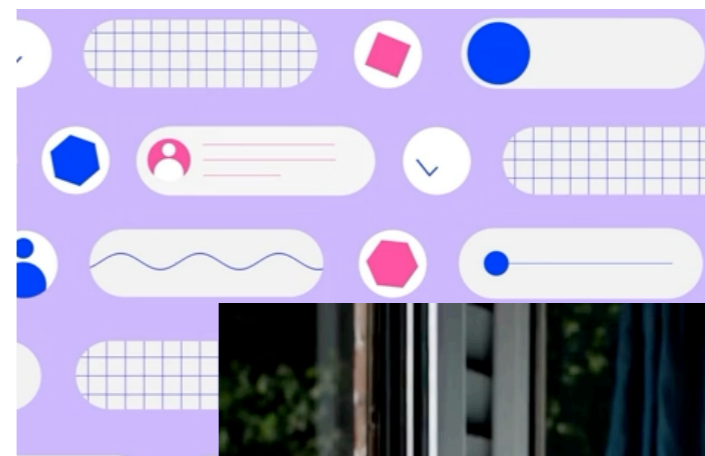
HARVARD
UNIVERSITY



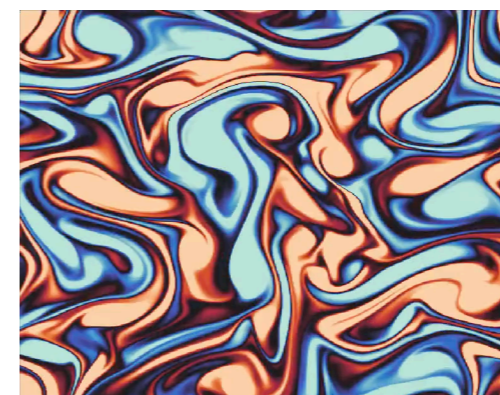
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Annealed Importance Sampling and Jarzynski's equality

Problem statement

Much related work!

Dynamical Measure Transport

Recent methods for learning maps between distributions

Combining the two!

New learning algorithms Applications, e.g. field theory

Annealed Importance Sampling and Jarzynski's equality

Problem statement

Main motivation for this work:

Can we explicitly get a machine learning-augmented sampling setup for which “when I pay more from using my model, I get more from my model”?

Dynamical Measure

Recent methods

Combining the two!

New learning algorithms

Applications, e.g. field theory

Agenda

Annealed Importance Sampling and Jarzynski's equality

Problem statement

Dynamical Measurements

Recent methods

Joint work with Eric
Vanden-Eijnden



Combining the two!

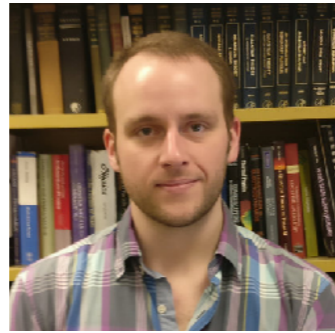
New learning algorithms

Applications, e.g. field theory

Thanks to all collaborators!



P. Shanahan



D. Hackett



F. Romero-



R. Abbott



D. Boyda



J. Urban



D. Rezende



S. Racanière



A. Razavi



A. Botev



M. Lindsey



K. Cranmer



N. Boffi



G. Kanwar



P. Lunts



A. Patel



M. Goldstein



R. Ranganath



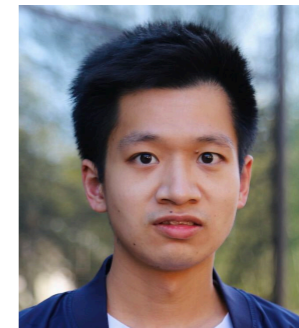
E. Vanden-Eijnden



Y. LeCun



S. Xie



Y. Chen



Problem Setup

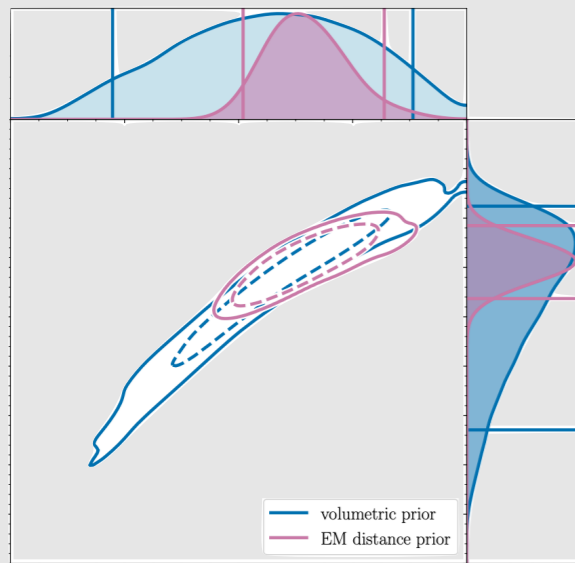
Goal: estimate the unknown *probability density function* $\rho_1 \in \mathcal{D}(\Omega)$ either through:

1. sample data $\{x_i\}_{i=1}^n$
2. **query access to the unnormalized log likelihood (energy function)**

Sampling problem ubiquitous!

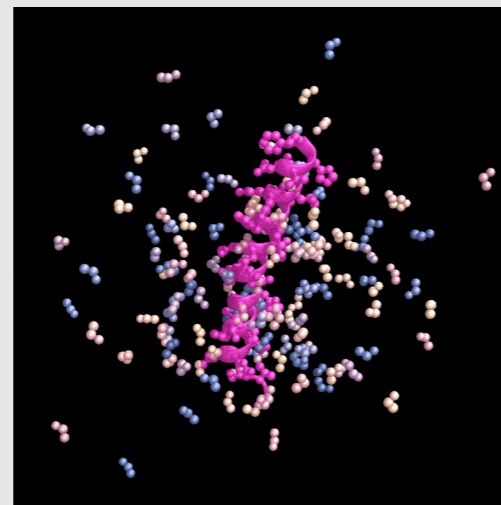
(obviously, to this audience)

energy function $U_1(x)$

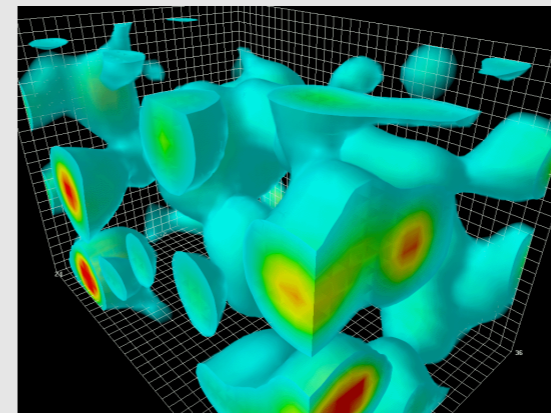


bayesian inference in GW astronomy

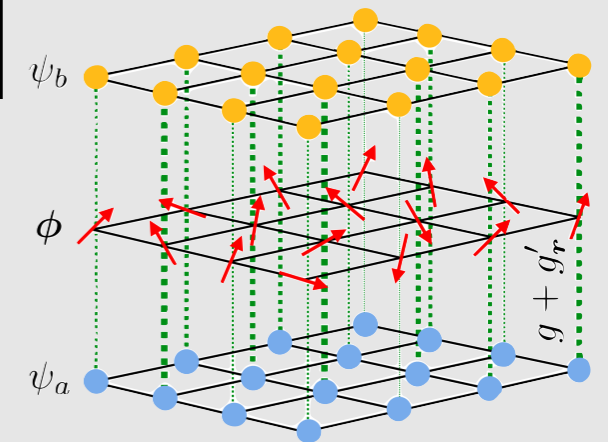
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MD simulations



quantum field theory



condensed matter

Problem Setup

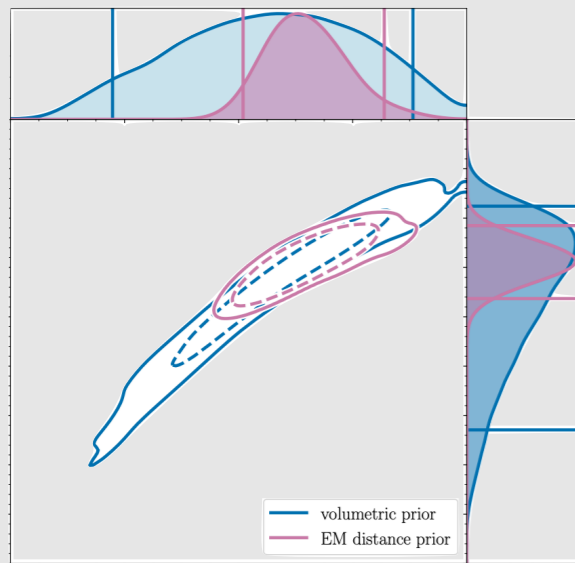
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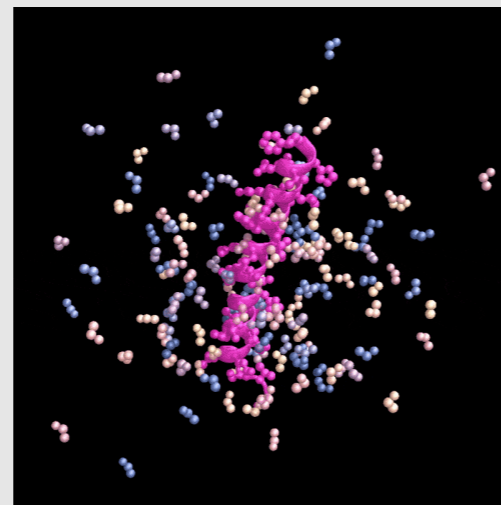
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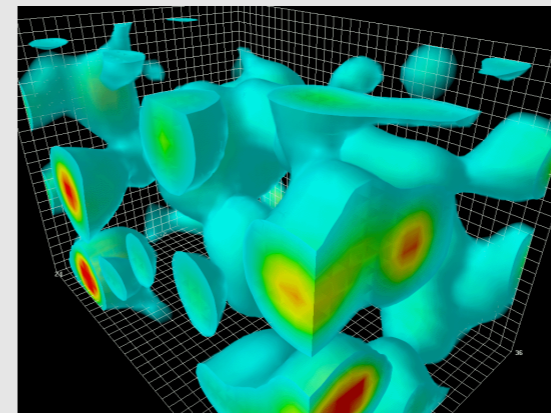


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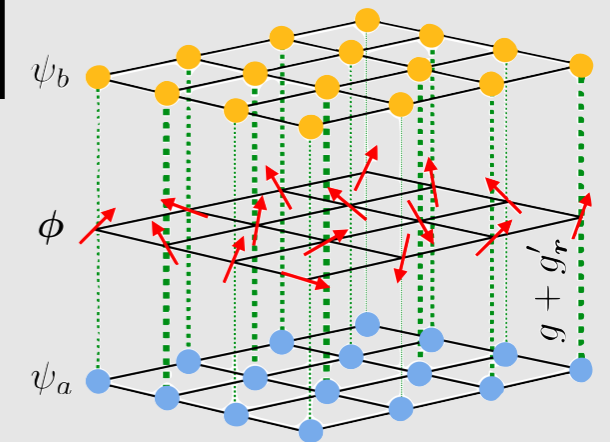
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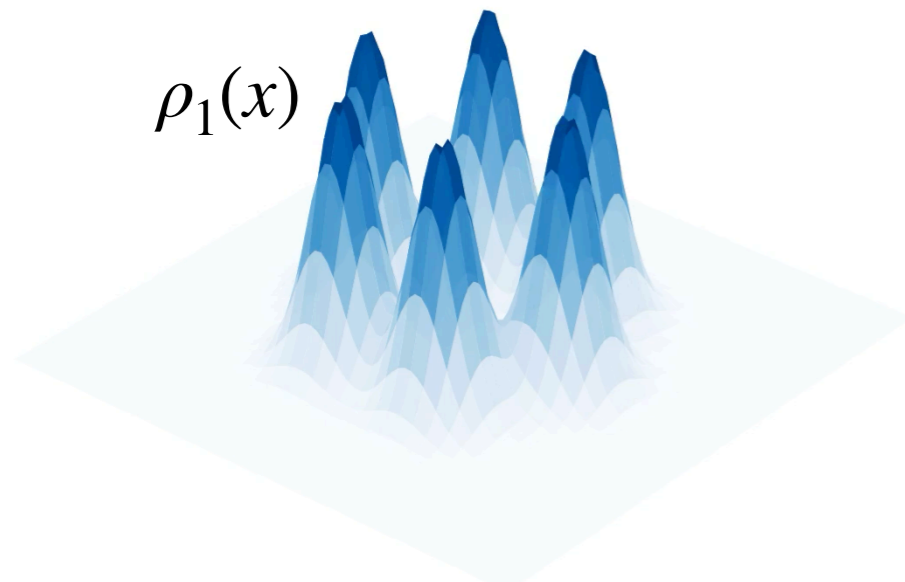
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Approaches to sampling

Markov Chain Monte Carlo build randomized sequence of samples $\{x_i\}_{i=1}^N$ so that



Langevin dynamics on 2-dimensional distribution

$$\lim_{N \rightarrow \infty} \mathbb{E}[h(x)]_N \rightarrow \mathbb{E}[h(x)]$$

Common tool: Langevin Dynamics

$$dX_t = -\epsilon \nabla U_1(X_t) dt + \sqrt{2\epsilon} dW_t$$

gradient drift

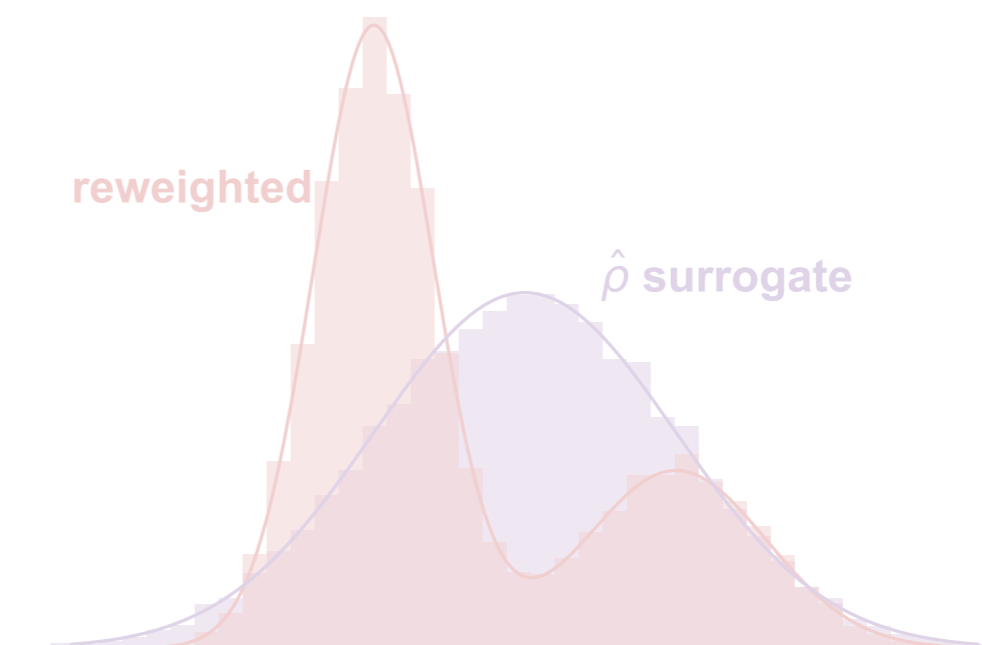
incremental brownian motion

Importance Sampling

Re-weight samples from cheap surrogate model

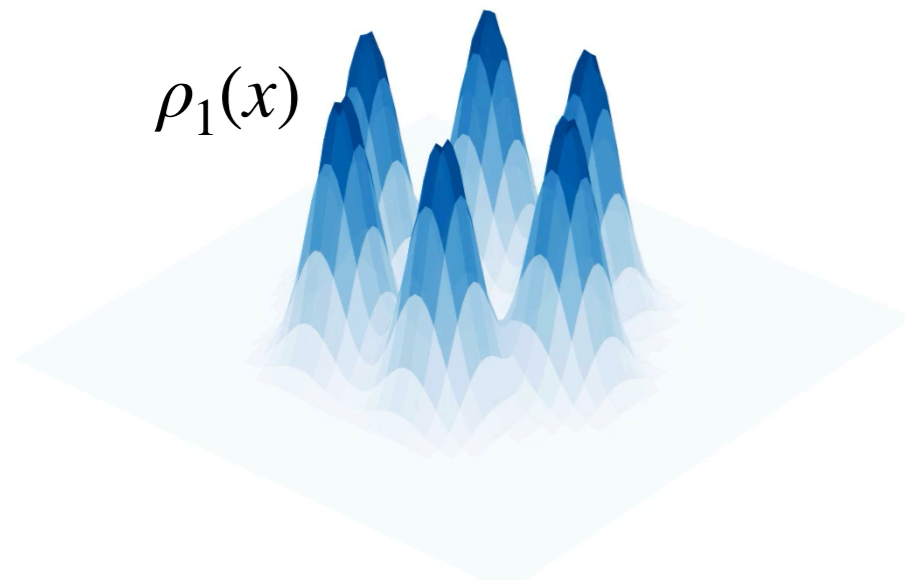
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Effective when $\rho_1, \hat{\rho}_1$ overlap



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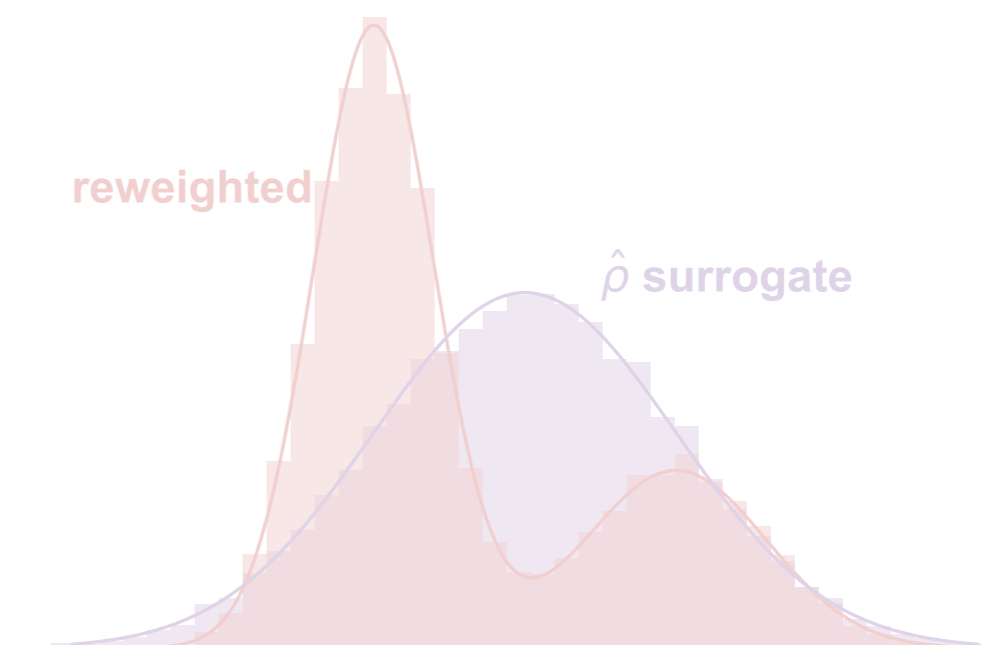
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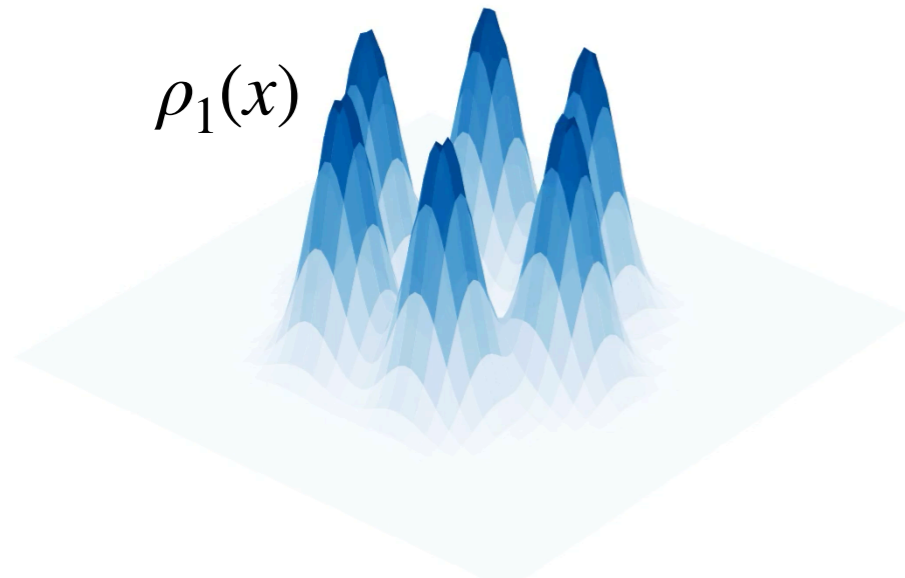
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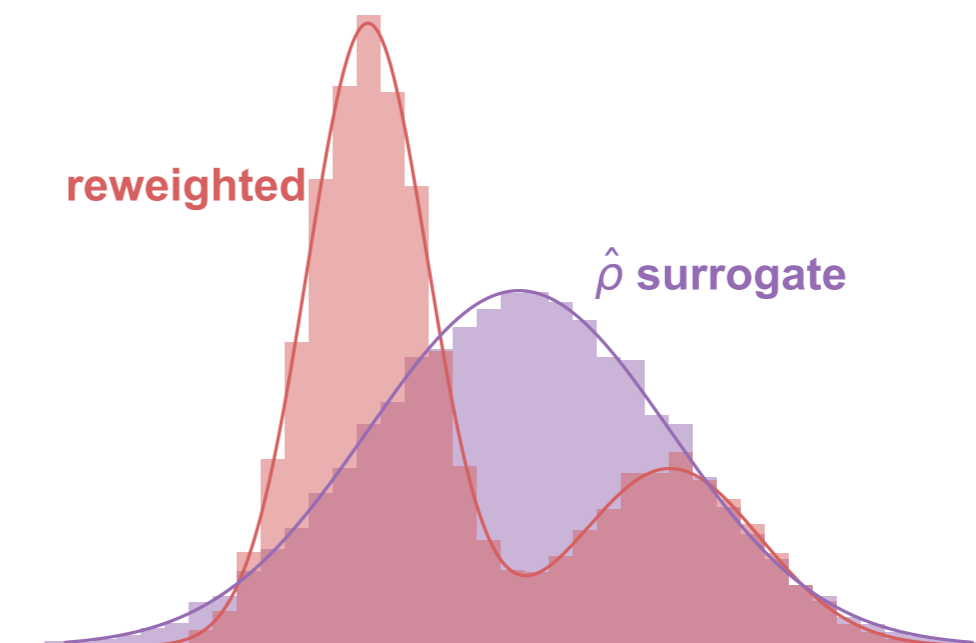
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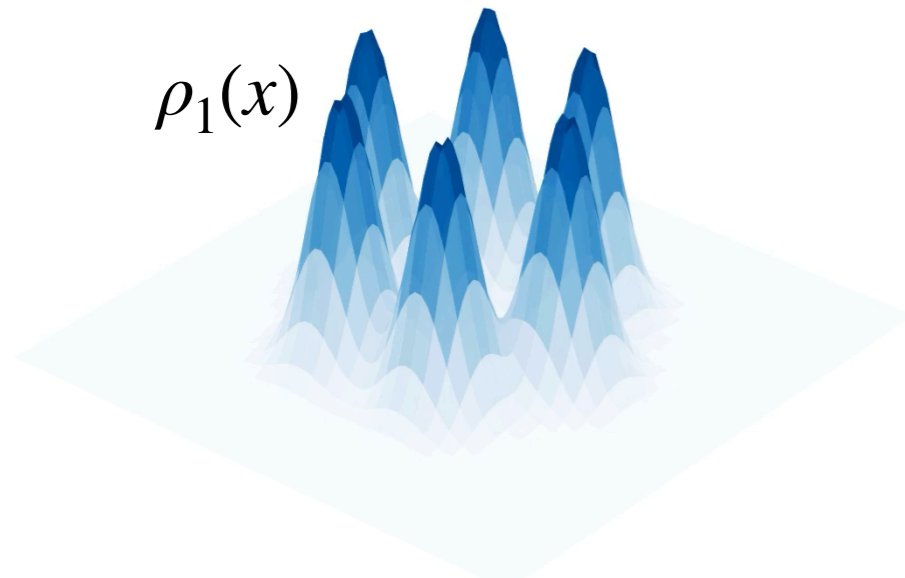
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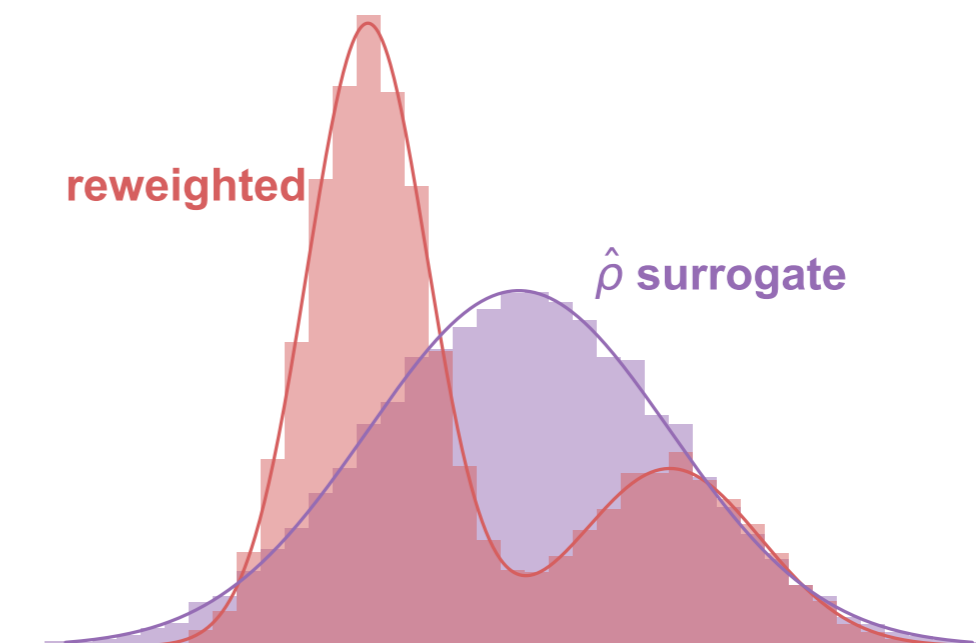
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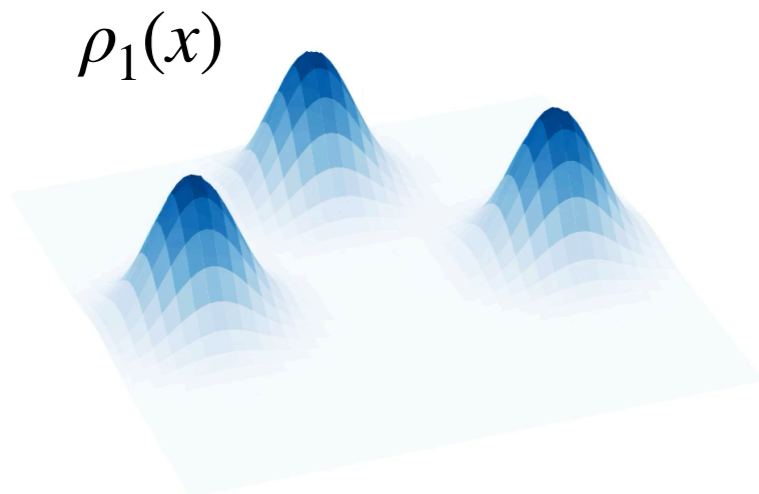
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Limitations of MCMC and IS

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Non-log concave target,
exponentially slow mixing

Convergence can be exponentially slow

Common tool: Langevin Dynamics

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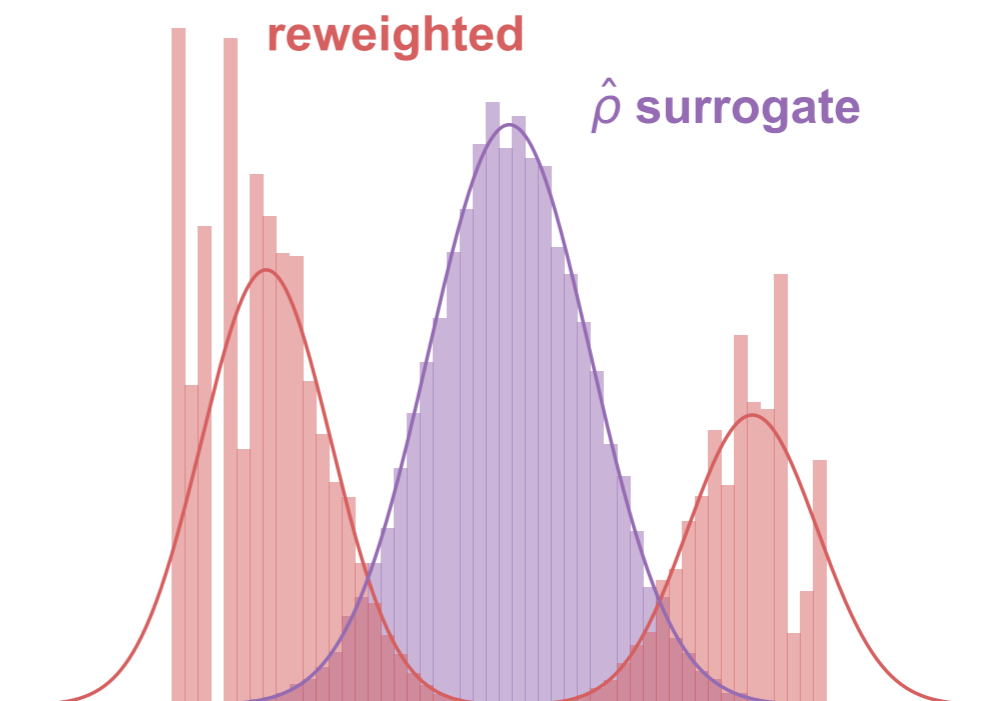
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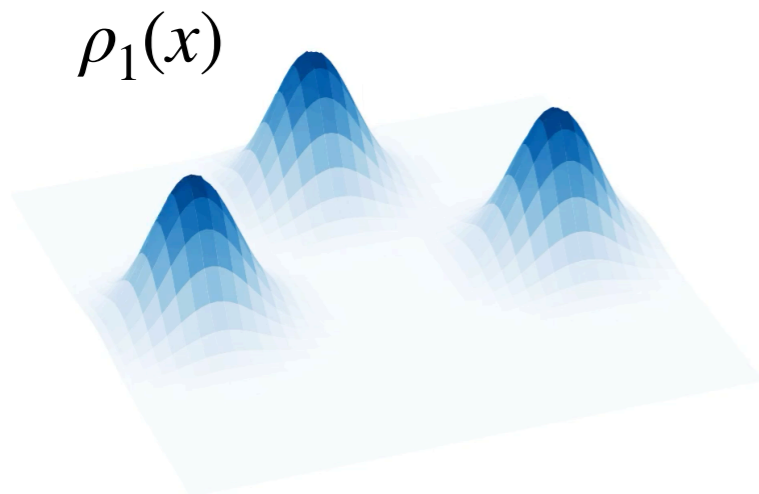
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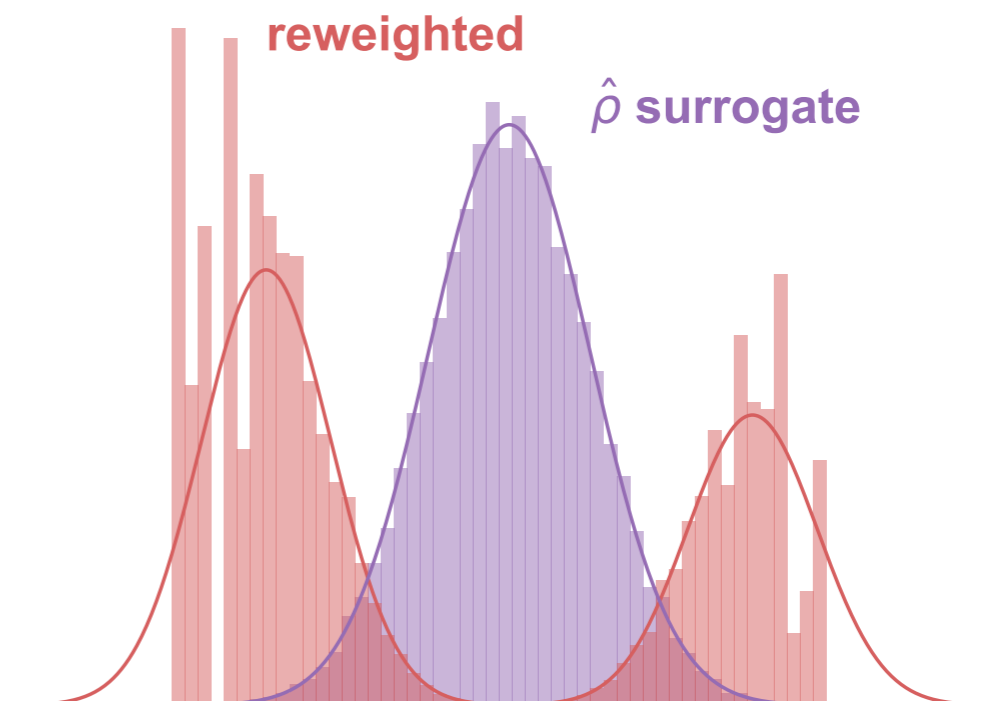
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Common Augmentation: Annealed Langevin Dynamics

Introduce dynamics which anneal to $U_1(x)$ from some $U_0(x)$

$$U_t(x) = (1 - t)U_0 + tU_1 \quad \text{PDF: } \rho_t(x) = e^{-U_t(x)+F_t}, \quad F_t = -\log Z_t$$

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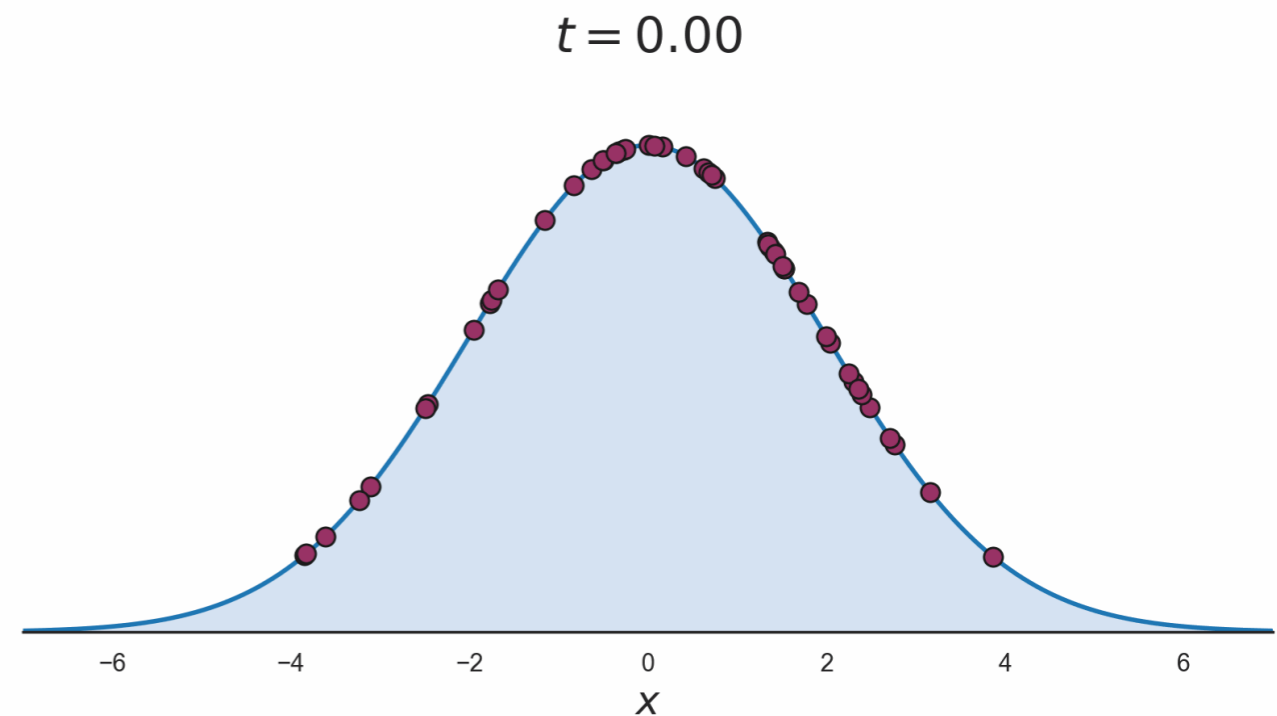
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SDE:

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- Time evolving potential
- ϵ_t sets speed of walkers per time step
- high temperature \rightarrow low temperature helps with multimodality



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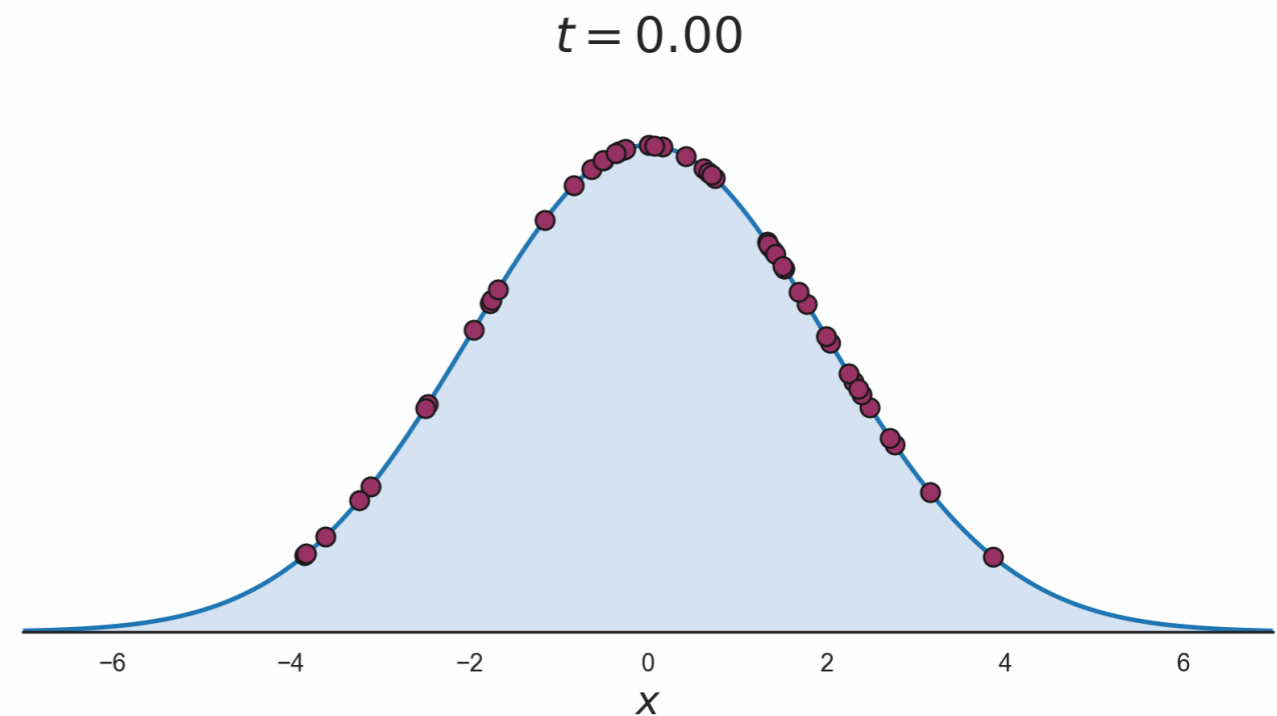
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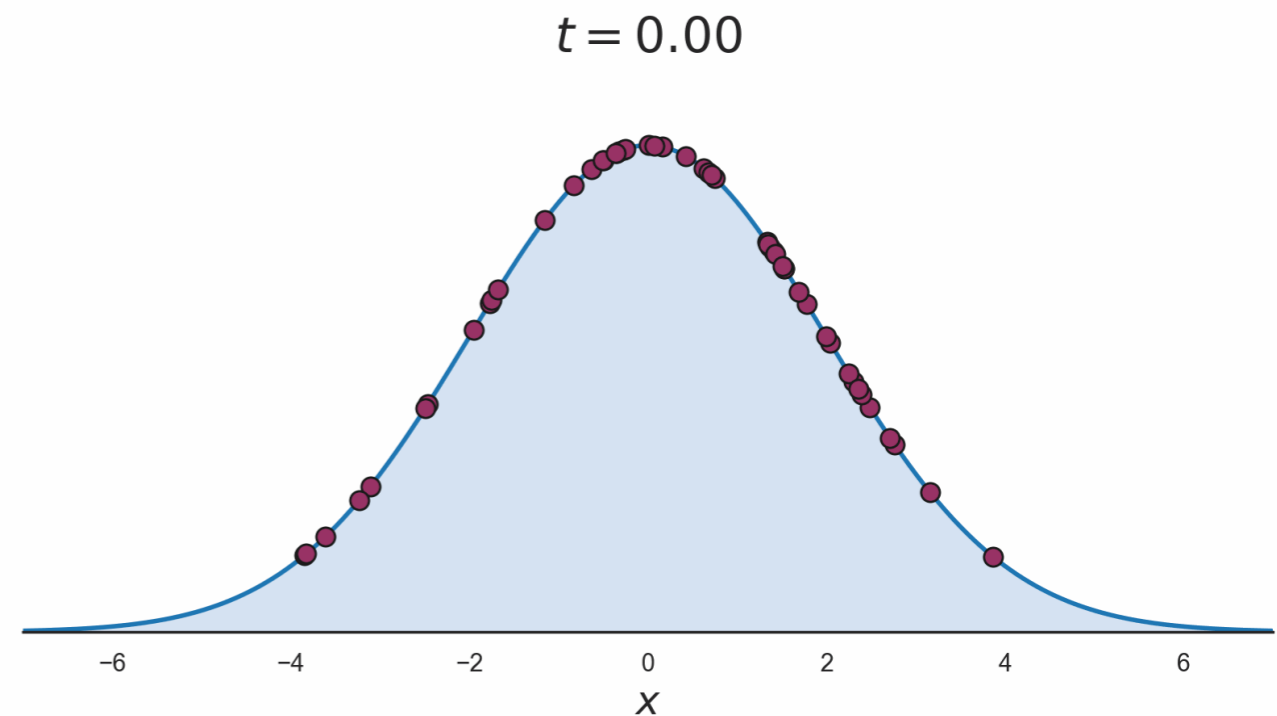
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Question: does the solution \tilde{X}_t to this SDE have ρ_t as its density?

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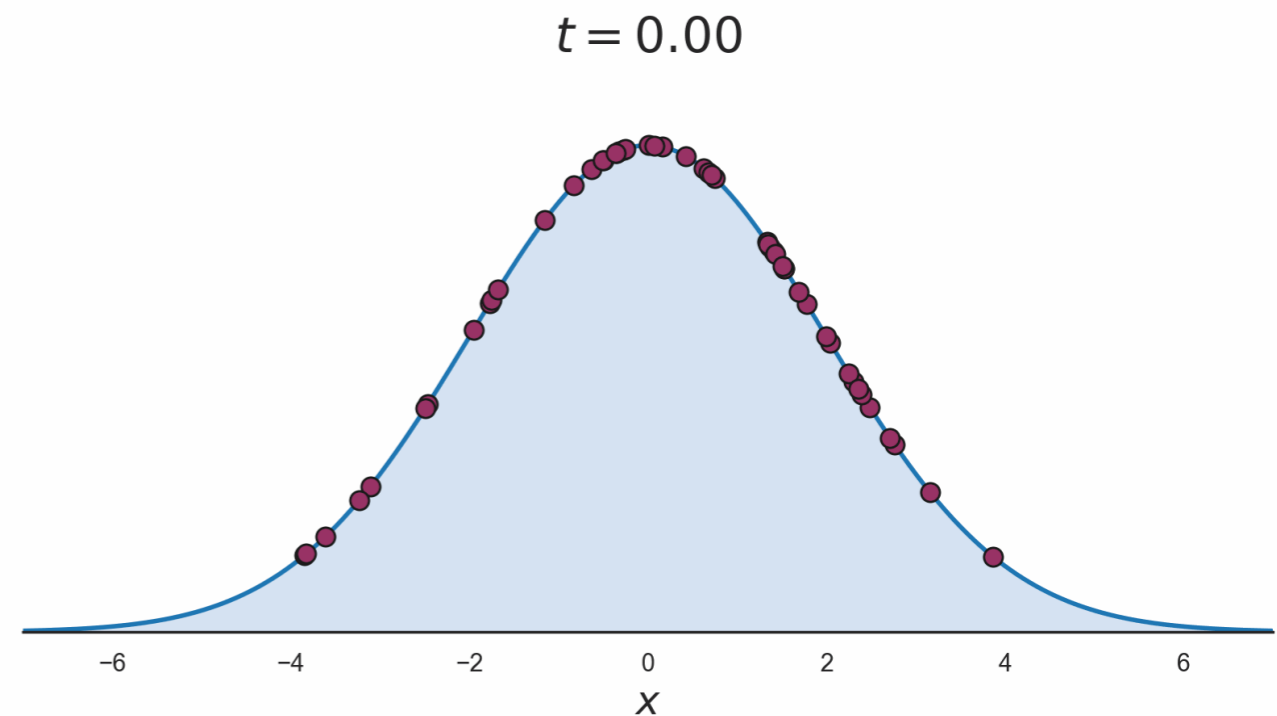
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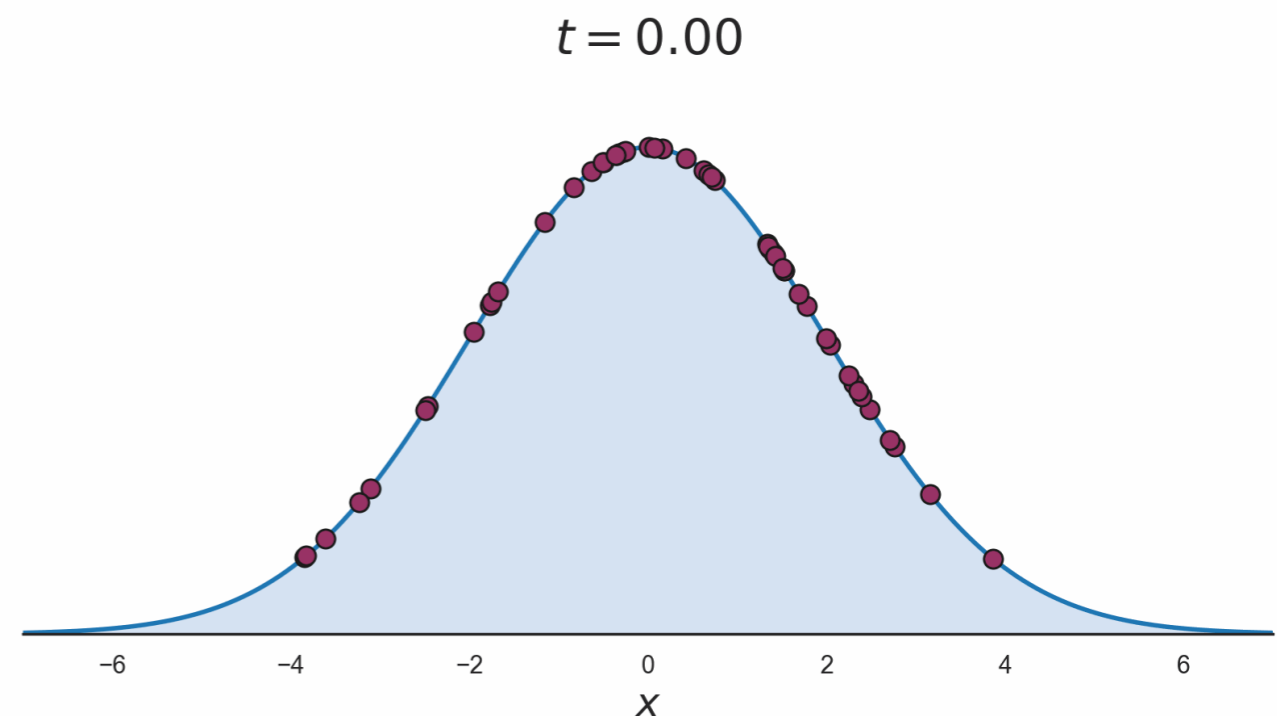
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NO! only if $\epsilon_t \rightarrow \infty$ and $dt \rightarrow 0$Why?

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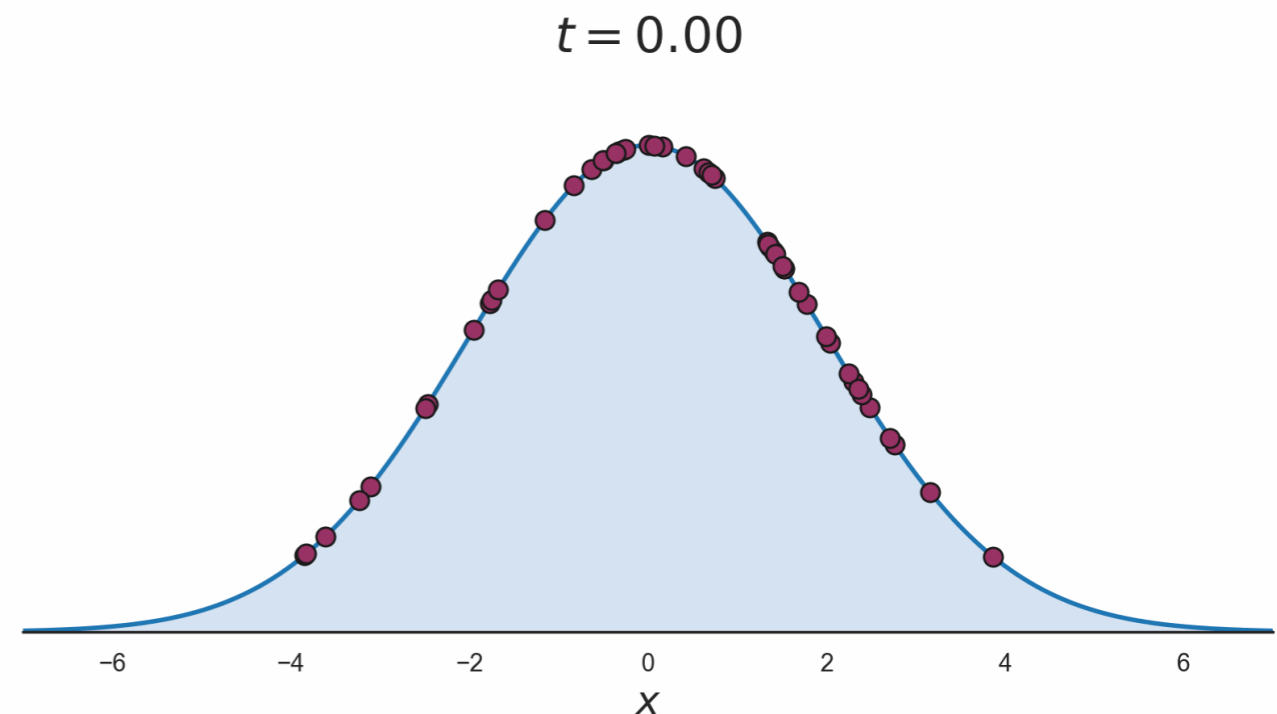
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$$\tilde{\rho}_t \neq \rho_t!$$

Compare the Fokker-Planck to $\partial_t \rho_t$

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FPE:

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Direct calculation:

$$\partial_t \rho_t = \frac{\partial}{\partial t} [e^{-U_t(x)+F_t}] - (\partial_t U_t - \partial_t F_t) \rho_t$$

$$= \epsilon_t \nabla \cdot (\nabla U_t \rho_t + \nabla \rho_t) + (\partial_t U_t - \partial_t F_t) \rho_t$$

since $\nabla \rho_t = -\nabla U_t \rho_t$

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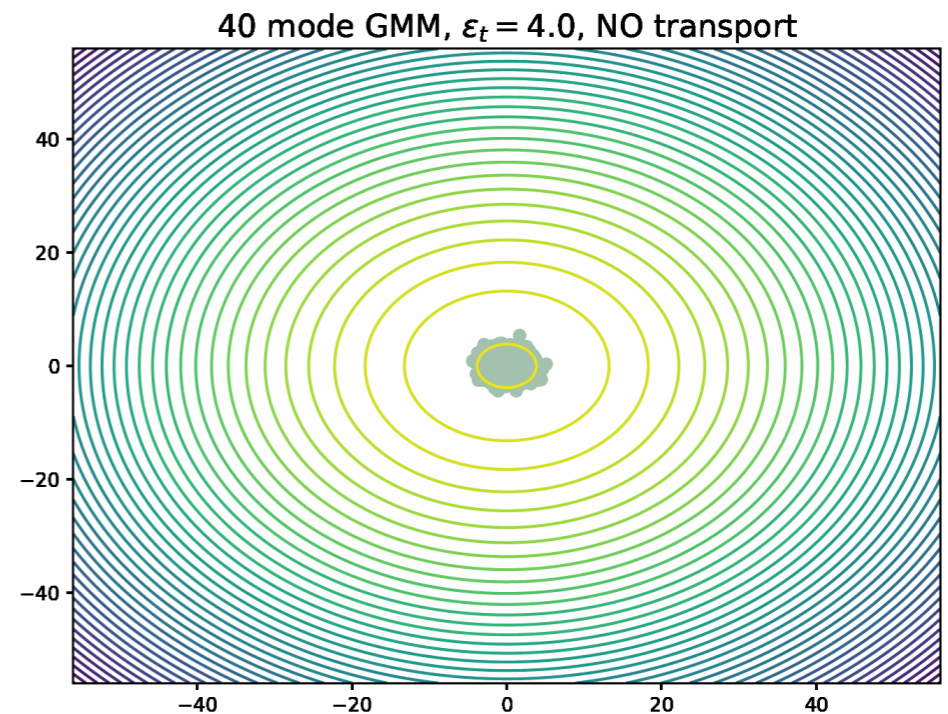
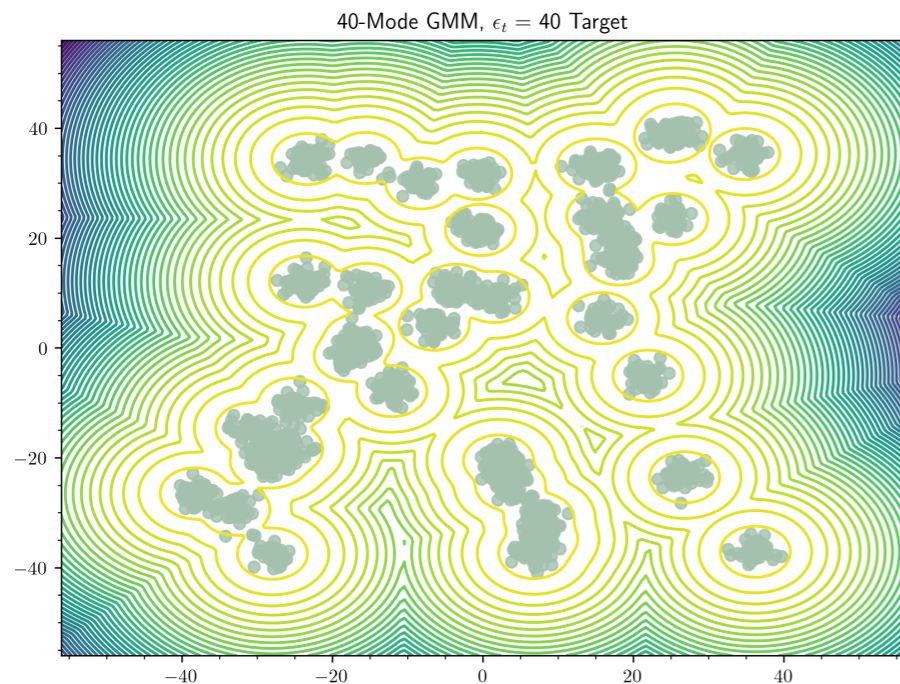
$\partial_t \rho_t$ and $\partial_t \tilde{\rho}_t$ differ by factor arising from time dynamics of U_t

In practice, the walkers \tilde{X}_t “lag behind” the intended evolution of ρ_t

The lag of \tilde{X}_t in practice

Compare the Fokker-Planck to $\partial_t \rho_t$

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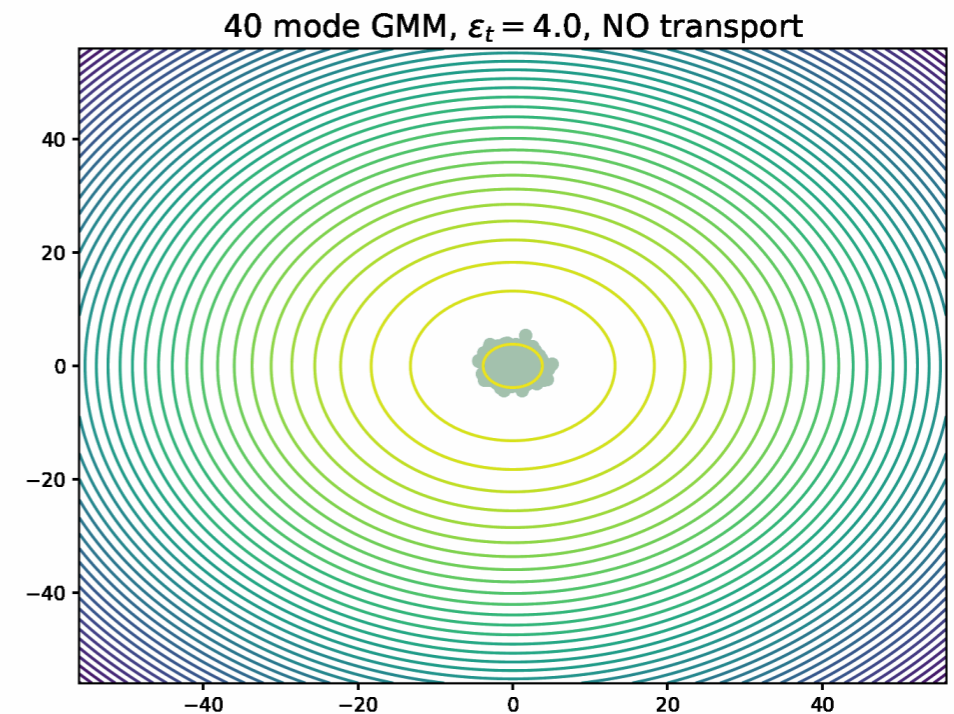
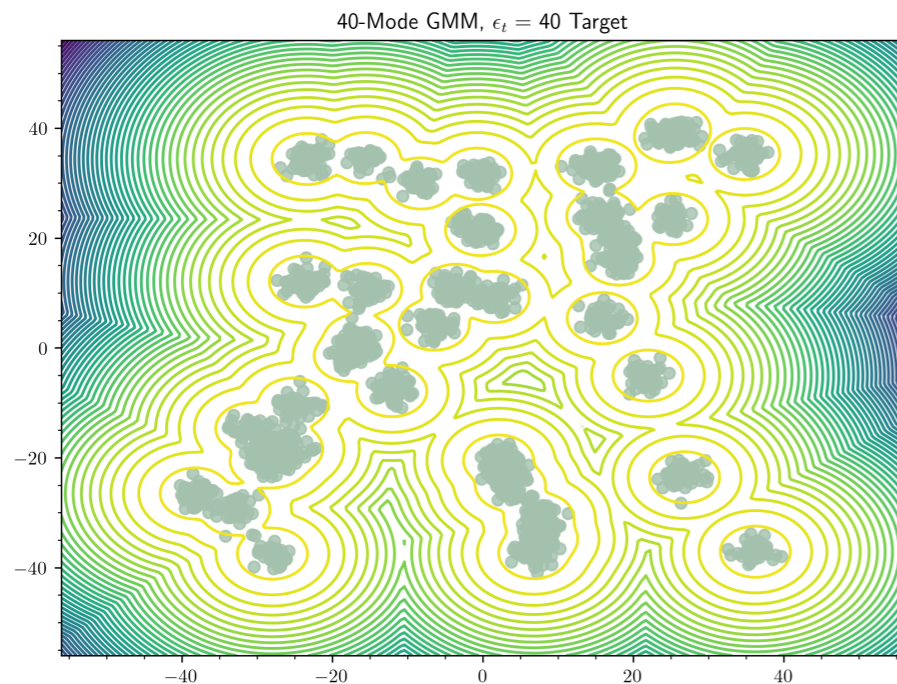
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This can in theory be fixed with re-weighting

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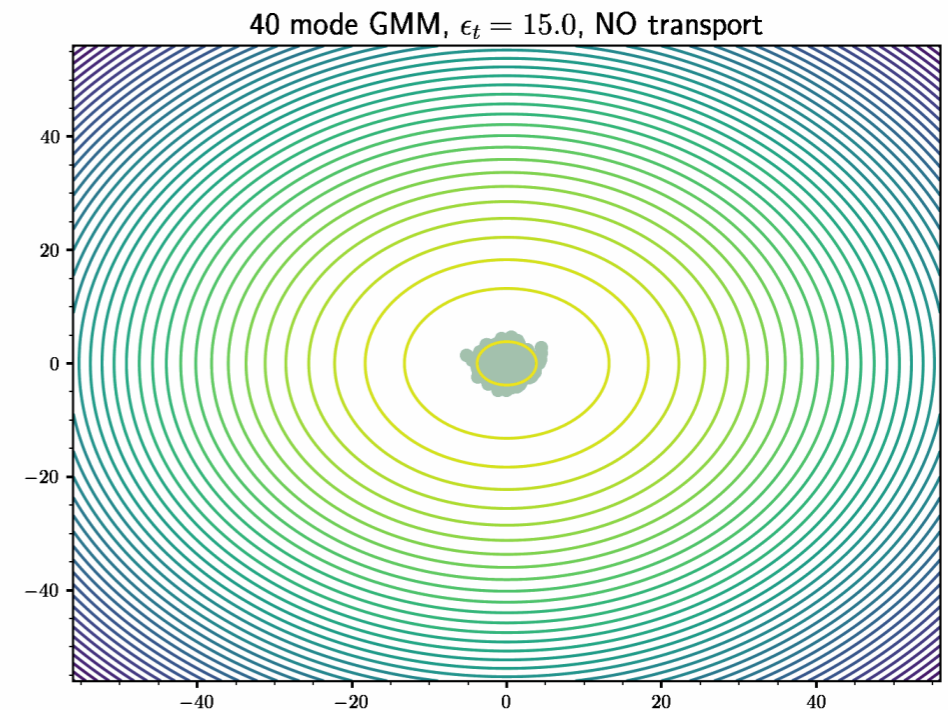
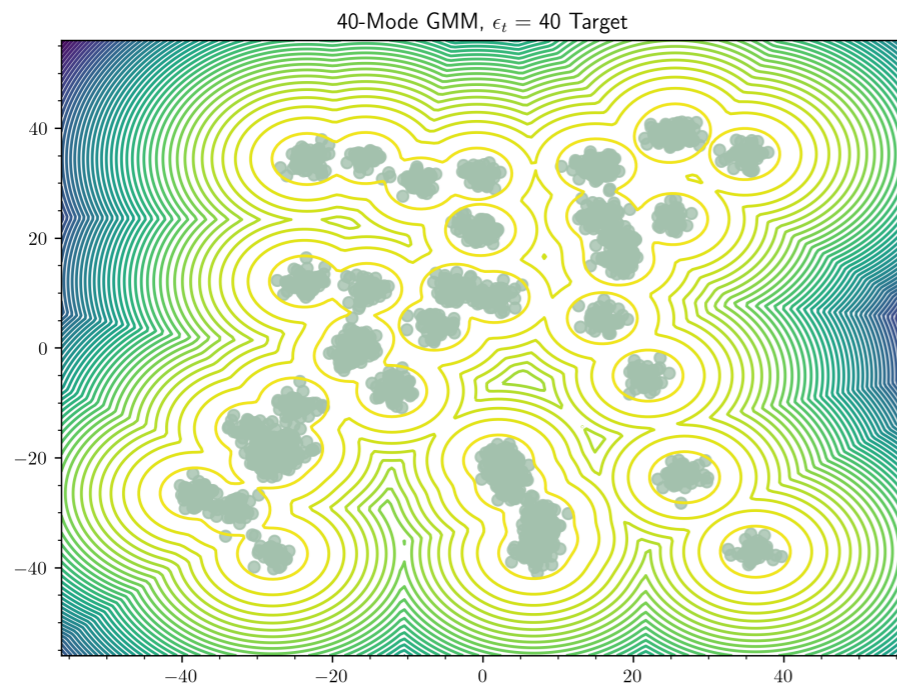
In practice, the walkers \tilde{X}_t “lag behind” the intended evolution of ρ_t

This can in theory be fixed with re-weighting

The lag of \tilde{X}_t in practice

Compare the Fokker-Planck to $\partial_t \rho_t$

$$\rho_t(x) = e^{-U_t(x)+F_t}$$



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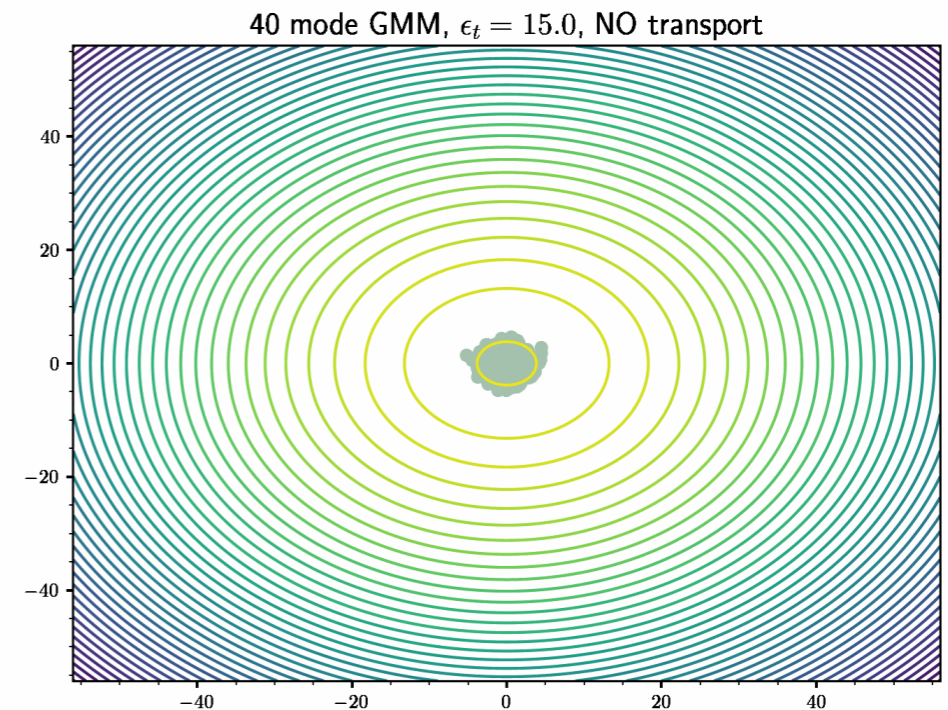
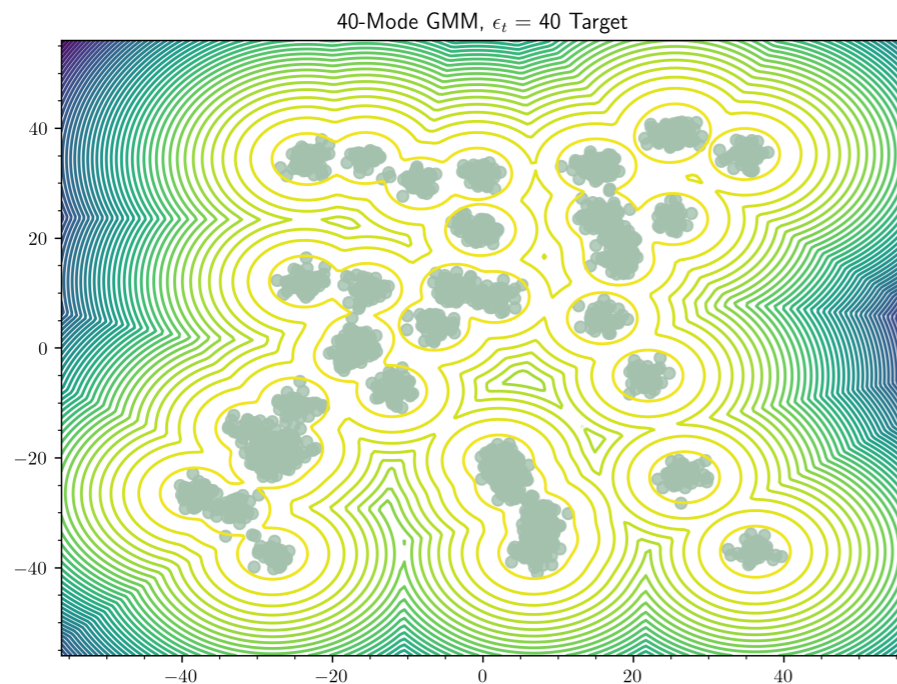
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Jarzynski Equality:

MSA & Vanden-Eijnden arXiv:2410.02711 (2024);
Jarzynski, PRL **78**, 2690 (1997)

Introduce weights A_t to account for the lag of the walkers

Proposition

Let (X_t, A_t) be the solution to the coupled SDE/ODE

$$\begin{aligned}dX_t &= -\epsilon_t \nabla U_t(X_t)dt + \sqrt{2\epsilon_t}dW_t, & X_0 &\sim \rho_0 \\dA_t &= -\partial_t U_t(X_t)dt & A_0 &= 0\end{aligned}$$

then for all test functions $h(x)$, we have

$$\int_{\mathbb{R}^d} h(x)\rho_t(x)dx = \frac{\mathbb{E}[e^{A_t}h(x)]}{\mathbb{E}[e^{A_t}]} \quad Z_t/Z_0 = e^{-F_t+F_0} = \mathbb{E}[e^{A_t}]$$

Jarzynski!

change in
free energy

average
work!

- Can be proven by looking at the FPE for the joint pdf $f_t(x, a) : \mathbb{R}^{d+1} \rightarrow \mathbb{R}$

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- **Problem:** variance of e^{A_t} may be so large that re-weighting not useful

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$$dX_t = -\epsilon_t \nabla U_t(X_t) dt + \sqrt{2\epsilon_t} dW_t, \quad X_0 \sim \rho_0$$

Can we fix this with measure transport?

- Can be proven by looking at the FPE for the joint pdf $f_t(x, a) : \mathbb{R}^{d+1} \rightarrow \mathbb{R}$
- **Problem:** variance of e^{A_t} may be so large that re-weighting not useful

\mathbb{R}^d

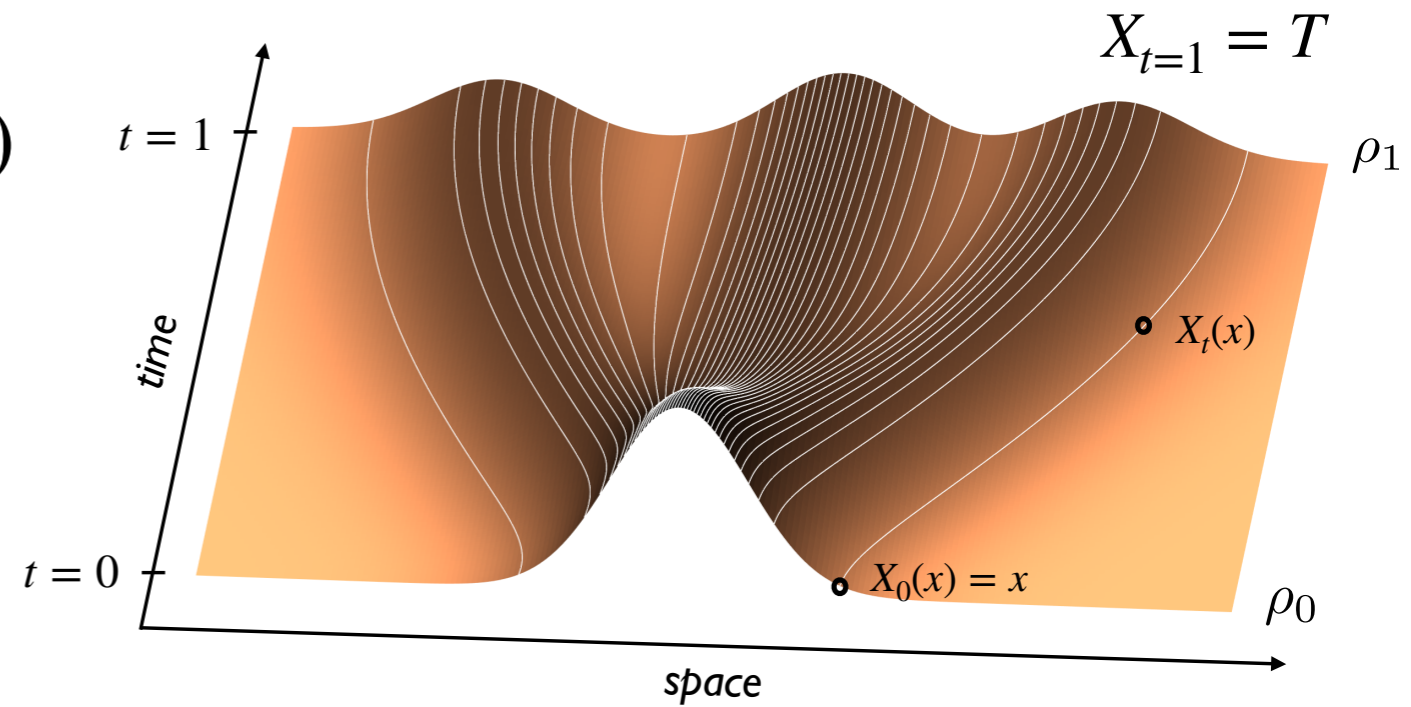
\mathbb{R}^d

Measure Transport

X_t flow map given by velocity field $b(t, x)$

$$X_{t=0}(x) = x \in \mathbb{R}^d$$

$$\dot{X}_t(x) = b_t(X_t(x))$$

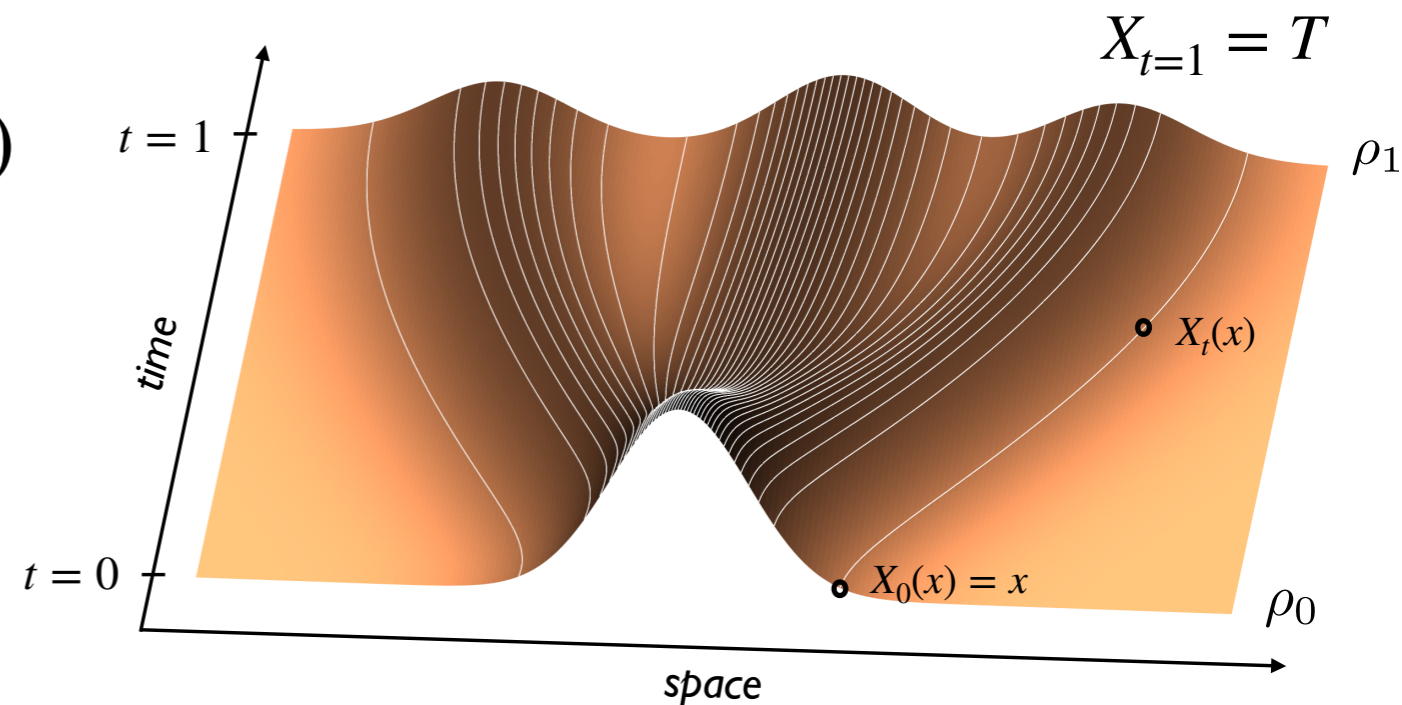


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At the level of the of the distribution, how does $\rho(t, x)$ evolve?

Transport equation

$$\partial_t \rho_t + \nabla \cdot (b_t \rho_t) = 0, \quad \rho_{t=0} = \rho_0$$

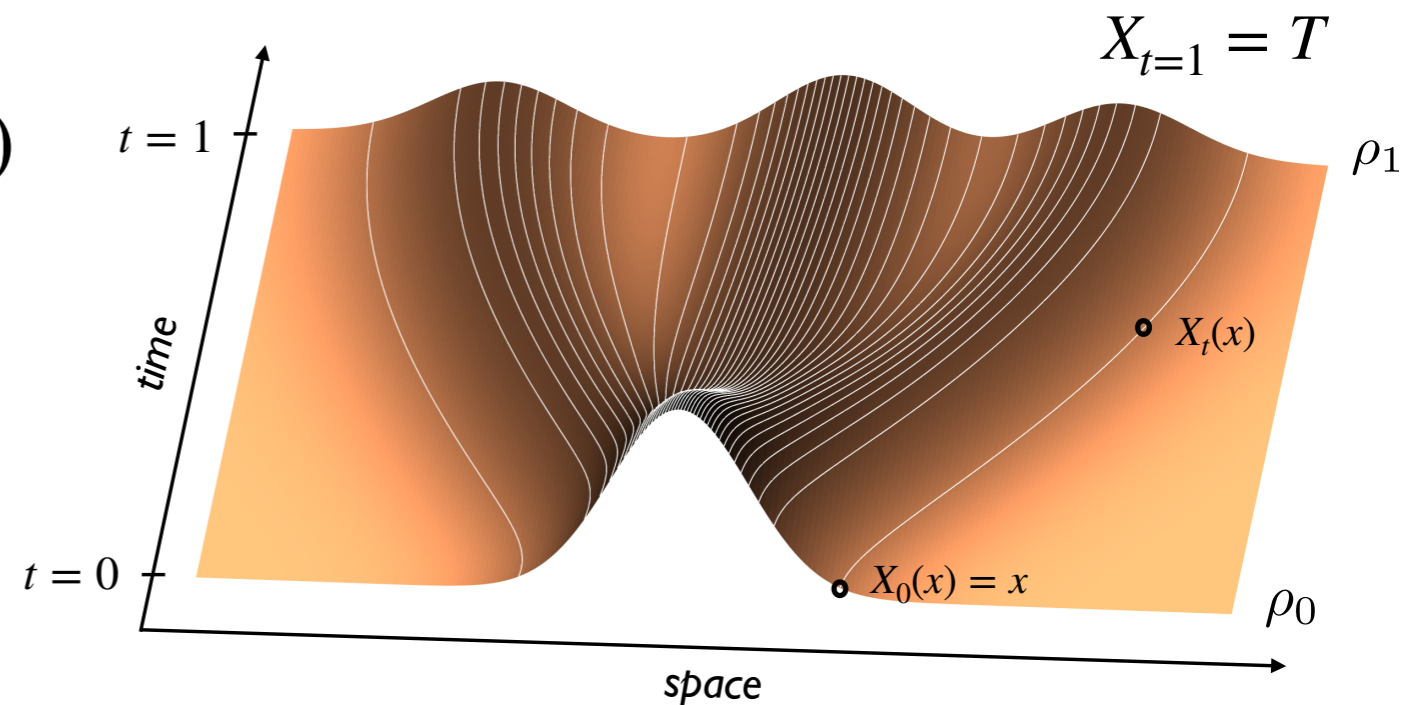
If $\rho(t)$ solves TE, **then** $\rho_{t=1} = \rho_1$

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Fokker-Planck Equation

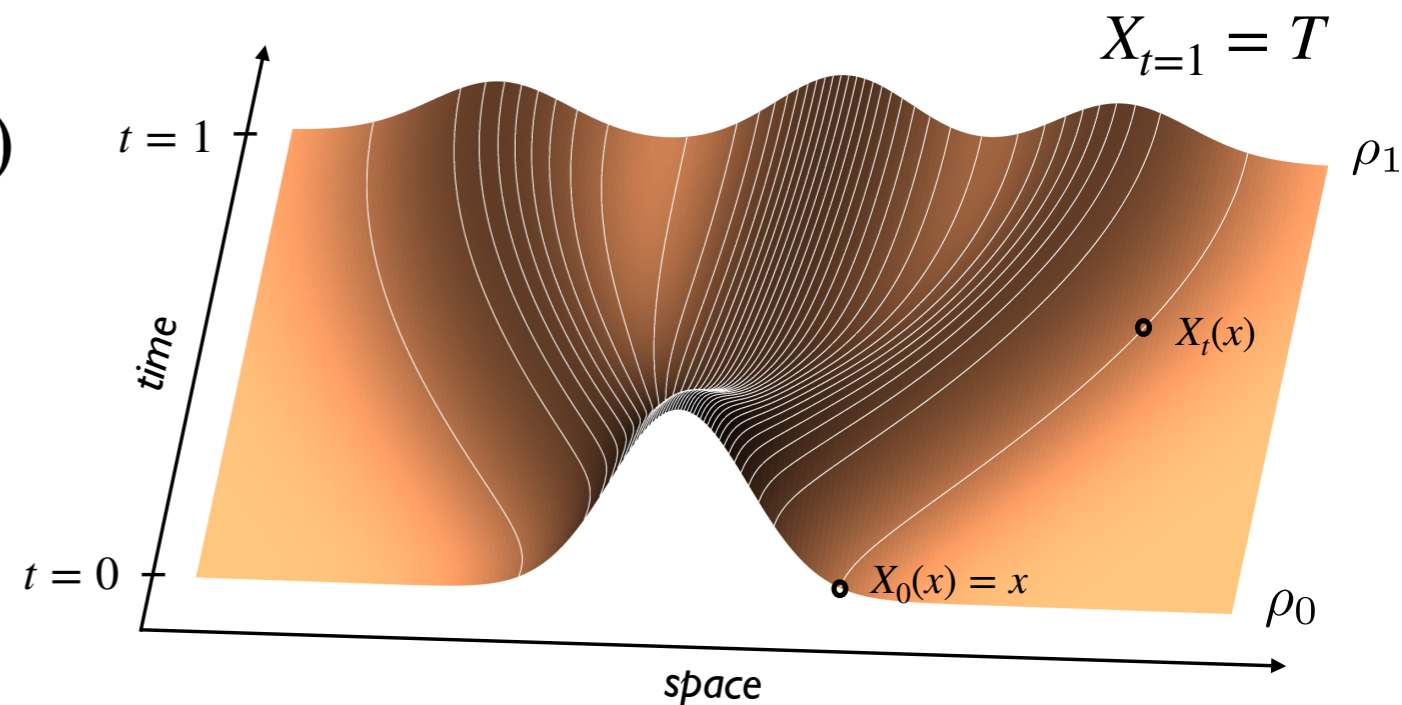
$$\partial_t \rho_t + \nabla \cdot (b_t \rho_t) = \epsilon \nabla \cdot (\nabla U_t \rho_t + \nabla \rho_t)$$

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Exact sampling

If X_t solves $dX_t = -\epsilon_t \nabla U_t(X_t)dt + b_t(X_t)dt + \sqrt{2\epsilon_t}dW_t$ **Then** $X_t \sim \rho_t$

Non-equilibrium transport sampler

What if you don't have the perfect b_t ?

MSA & Vanden-Eijnden *arXiv:2410.02711* (2024);
Vargas et al *ICLR* (2024);
Vaikuntanathan and Jarzynski, *PRL* **78**, 2690 (2008)

$$\text{Using } \nabla \cdot (\hat{b}_t \rho_t) = \nabla \cdot \hat{b}_t \rho_t - \nabla U_t \cdot \hat{b}_t \rho_t$$

FPE:

New non-eq term!

$$\partial_t \rho_t + \nabla \cdot (\hat{b}_t \rho_t) = \epsilon_t \nabla \cdot (\nabla U_t \rho_t + \nabla \rho_t) + (\nabla \cdot \hat{b}_t - \nabla U_t \cdot \hat{b}_t - \partial_t U_t + \partial_t F_t) \rho_t$$

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Let (X_t, A_t) be the solution to the coupled SDE/ODE

$$dX_t = -\epsilon_t \nabla U_t(X_t) dt + \sqrt{2\epsilon_t} dW_t, \quad X_0 \sim \rho_0$$

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Correctable dynamical transport for sampling

Valid for any diffusion ϵ_t which we will exploit

Strict augmentation of annealed Langevin dynamics



Learning b:

FPE:

$$\underbrace{\partial_t \rho_t + \nabla \cdot (\hat{b}_t \rho_t)}_{=0} = \epsilon_t \nabla \cdot (\nabla U_t \rho_t + \nabla \rho_t) + \underbrace{(\nabla \cdot \hat{b}_t - \nabla U_t \cdot \hat{b}_t - \partial_t U_t + \partial_t F_t)}_{=0} \rho_t$$

Need either

solves the transport

removes the non-equilibrium lag

Learning b: Physics Informed Neural Network Loss

MSA & Vanden-Eijnden arXiv:2410.02711 (2024);

Tian et. al ICML (2024);

FPE:

$$\underbrace{\partial_t \rho_t + \nabla \cdot (\hat{b}_t \rho_t)}_{=0} = \epsilon_t \nabla \cdot (\nabla U_t \rho_t + \nabla \rho_t) + \underbrace{(\nabla \cdot \hat{b}_t - \nabla U_t \cdot \hat{b}_t - \partial_t U_t + \partial_t F_t)}_{=0} \rho_t$$

Need either

solves the transport

removes the non-equilibrium lag

PINN Loss

All minimizers (b_t, F_t) of the objective

*Valid for any $\hat{\rho}_t$!
Controls the KL!*

$$L_{PINN}[\hat{b}, \hat{F}] = \int_0^1 \int_{\mathbb{R}^d} \left| \nabla \cdot \hat{b}_t(x) - \nabla U_t(x) \cdot \hat{b}_t(x) - \partial_t U_t(x) + \partial_t \hat{F}_t \right|^2 \hat{\rho}_t(x) dx dt$$

are such that $L_{PINN}[b, F] = 0$, F_t is the free energy, and b_t solves the transport

Learning b: Action Matching Loss

FPE:

$$\underbrace{\partial_t \rho_t + \nabla \cdot (\hat{b}_t \rho_t)}_{=0} = \epsilon_t \nabla \cdot (\nabla U_t \rho_t + \nabla \rho_t) + \underbrace{(\nabla \cdot \hat{b}_t - \nabla U_t \cdot \hat{b}_t - \partial_t U_t + \partial_t F_t)}_{=0} \rho_t$$

Need either

solves the transport

removes the non-equilibrium lag

Action matching loss

The minimizer $b_t = \nabla \phi_t$ of the objective

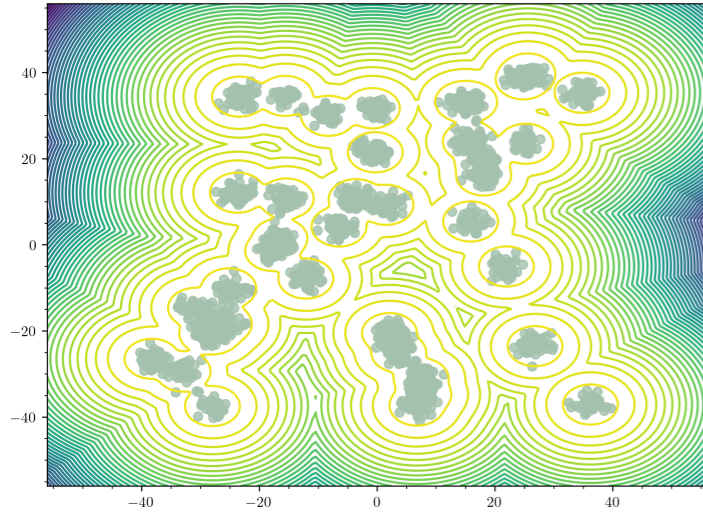
$$L_{AM}^T[\hat{\phi}] = \int_0^T \int_{\mathbb{R}^d} \left[\frac{1}{2} \left| \nabla \hat{\phi}_t(x) \right|^2 + \partial_t \hat{\phi}_t(x) \right] \rho_t(x) dx dt + \int_{\mathbb{R}^d} \left[\hat{\phi}_0(x) \rho_0(x) - \hat{\phi}_T(x) \rho_T(x) \right] dx$$

Needs reweighted samples from ρ_t

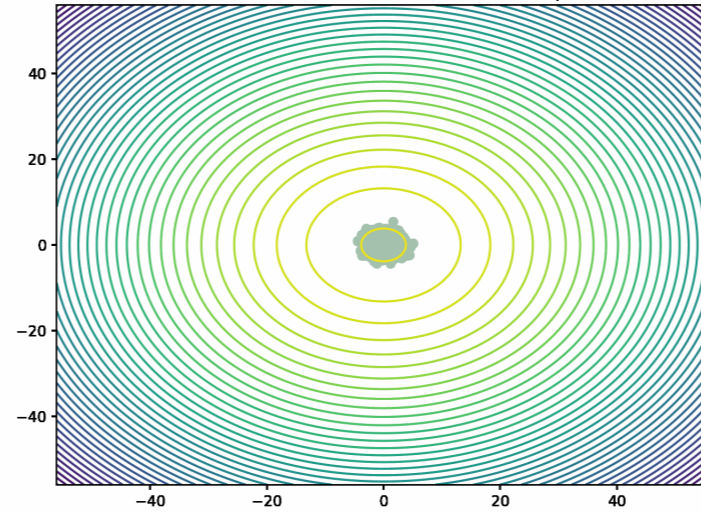
is unique up to a constant, and solves the transport.

Numerical Example: Painfully multimodal GMM

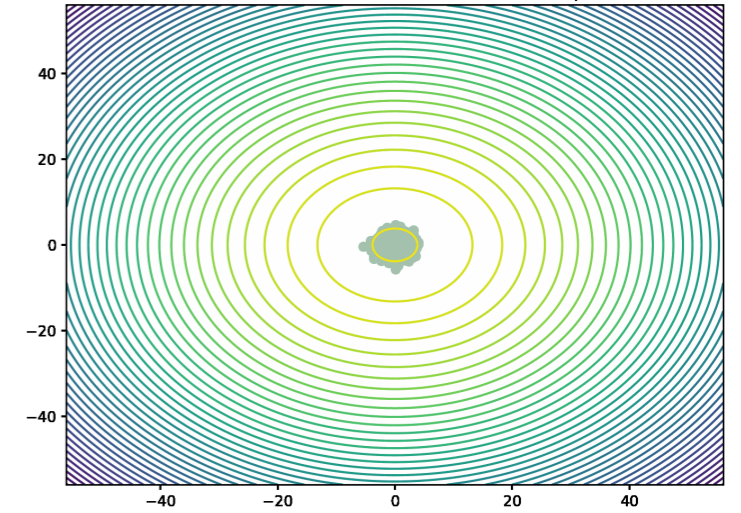
40-Mode GMM, $\epsilon_t = 40$ Target



40 mode GMM, $\epsilon_t = 4.0$, NO transport

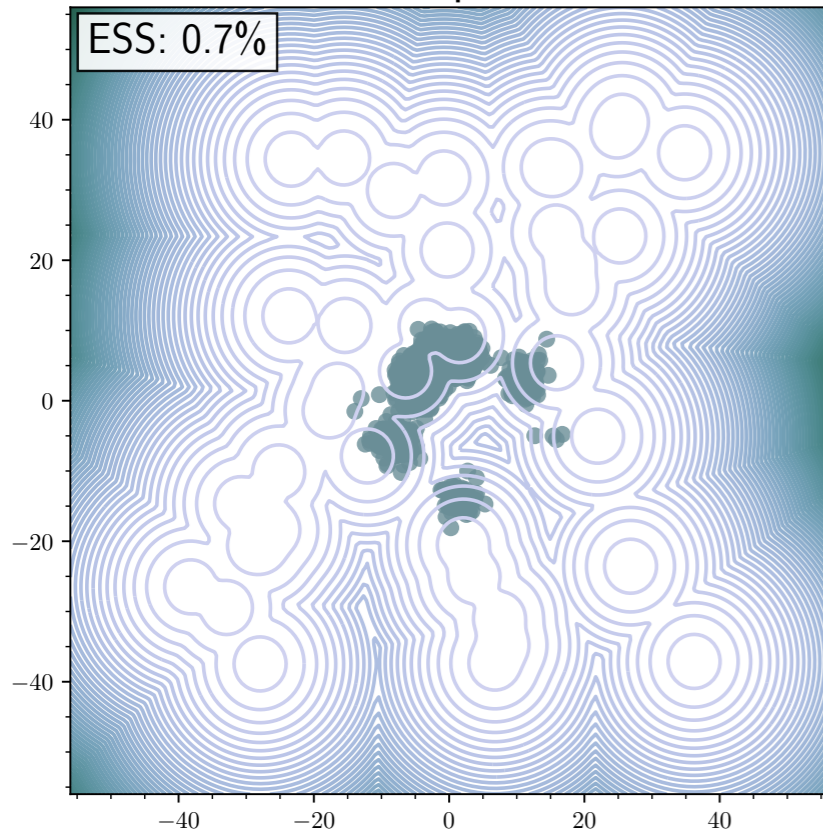


40 mode GMM, $\epsilon_t = 4.0$, with transport

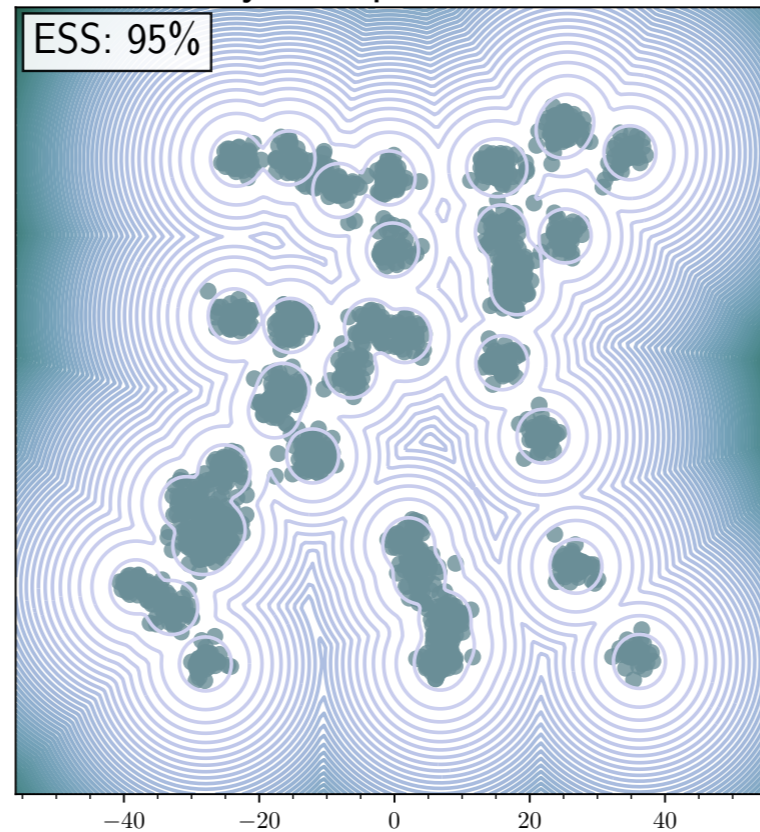


Turning on the diffusion improves ESS

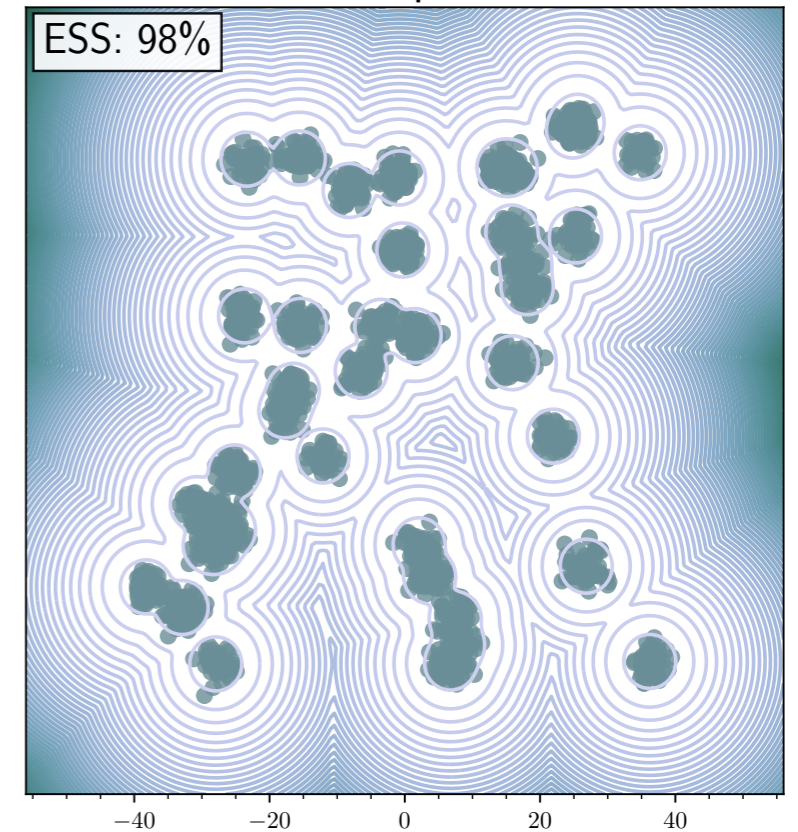
AIS, no transport, $\epsilon = 4.0$



Only transport, $\epsilon = 0.0$

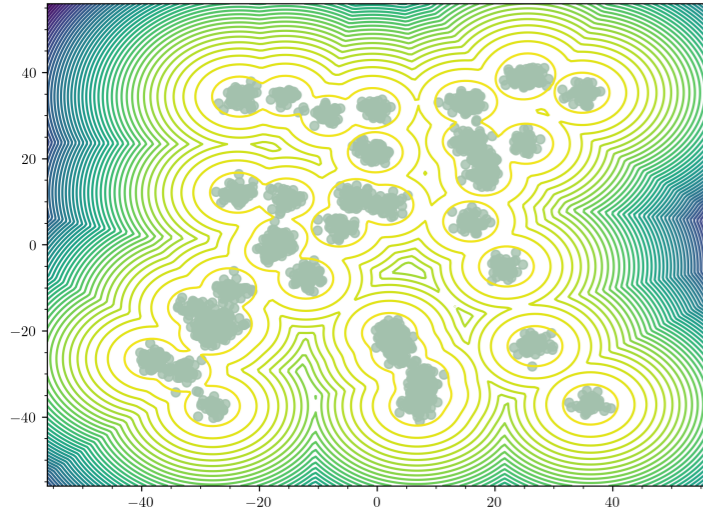


AIS and transport, $\epsilon = 4.0$

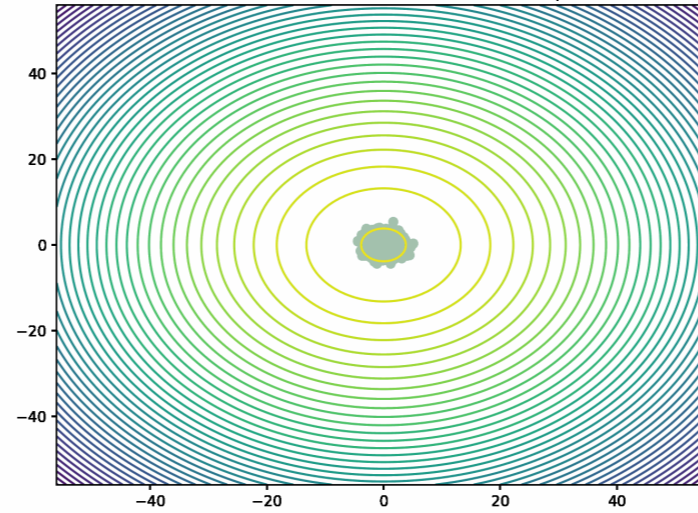


Numerical Example: Painfully multimodal GMM

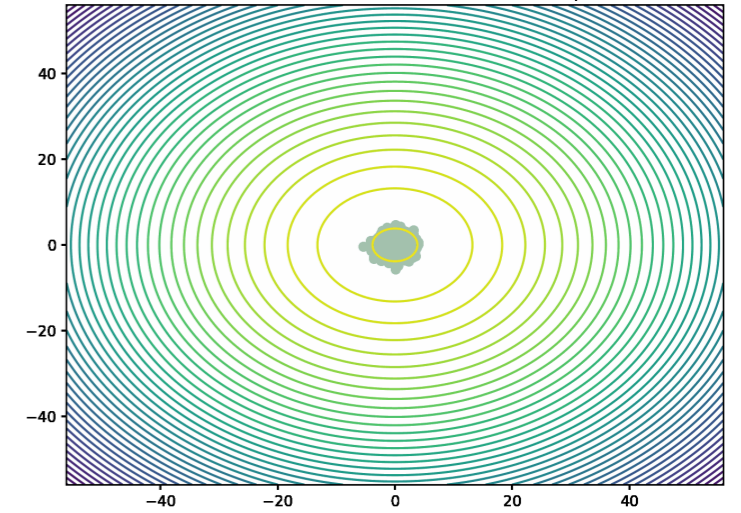
40-Mode GMM, $\epsilon_t = 40$ Target



40 mode GMM, $\epsilon_t = 4.0$, NO transport

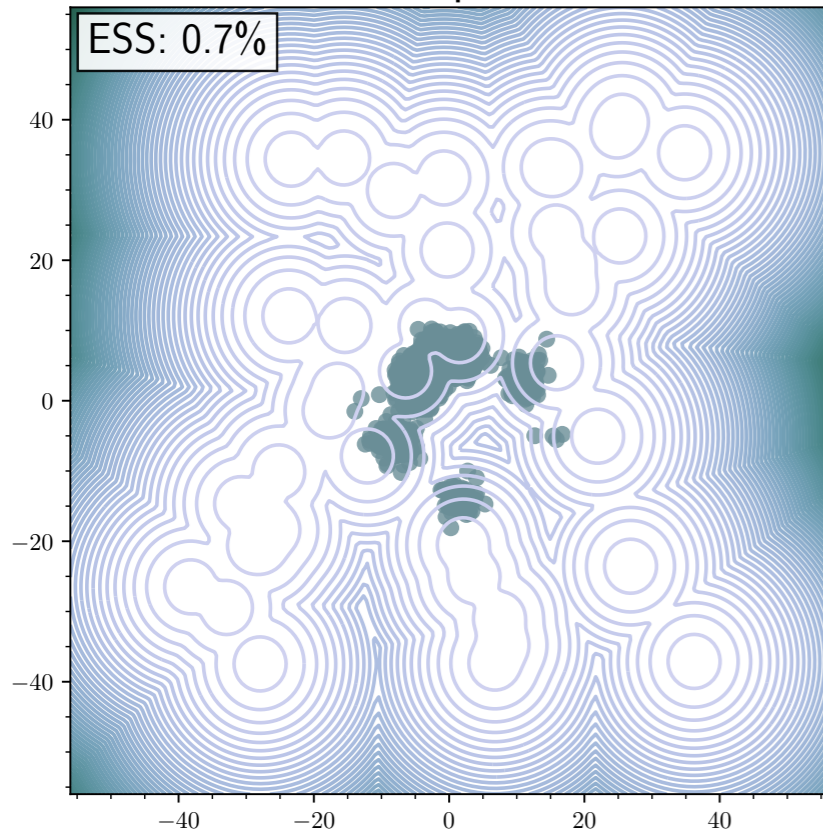


40 mode GMM, $\epsilon_t = 4.0$, with transport

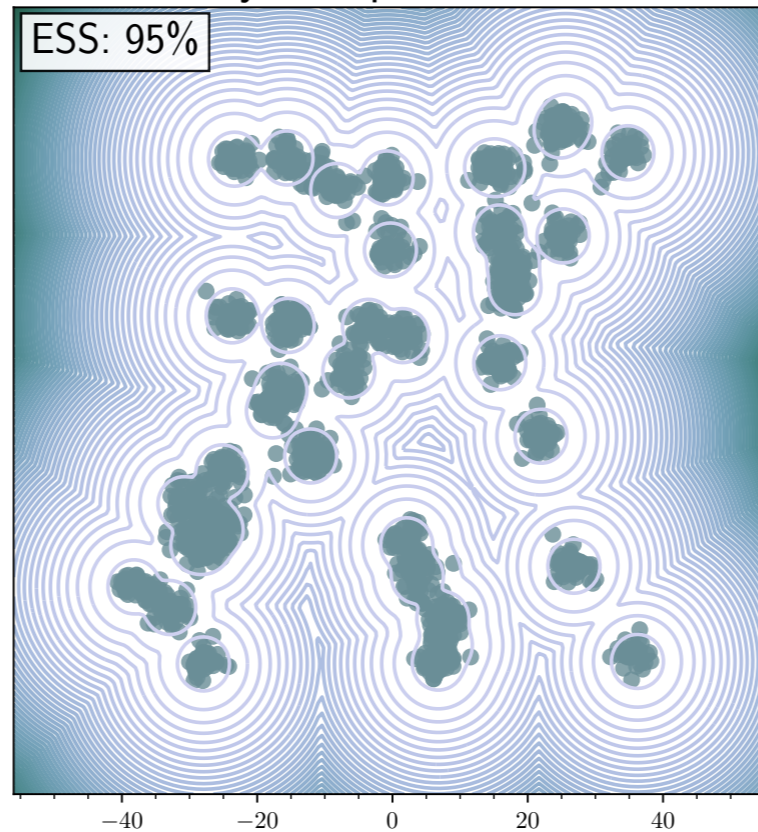


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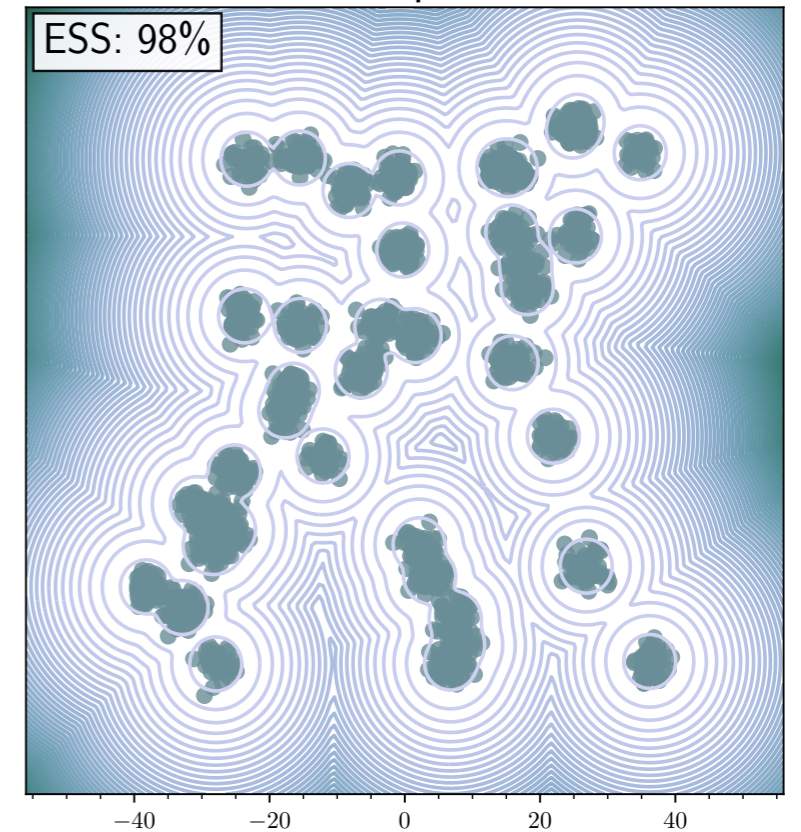
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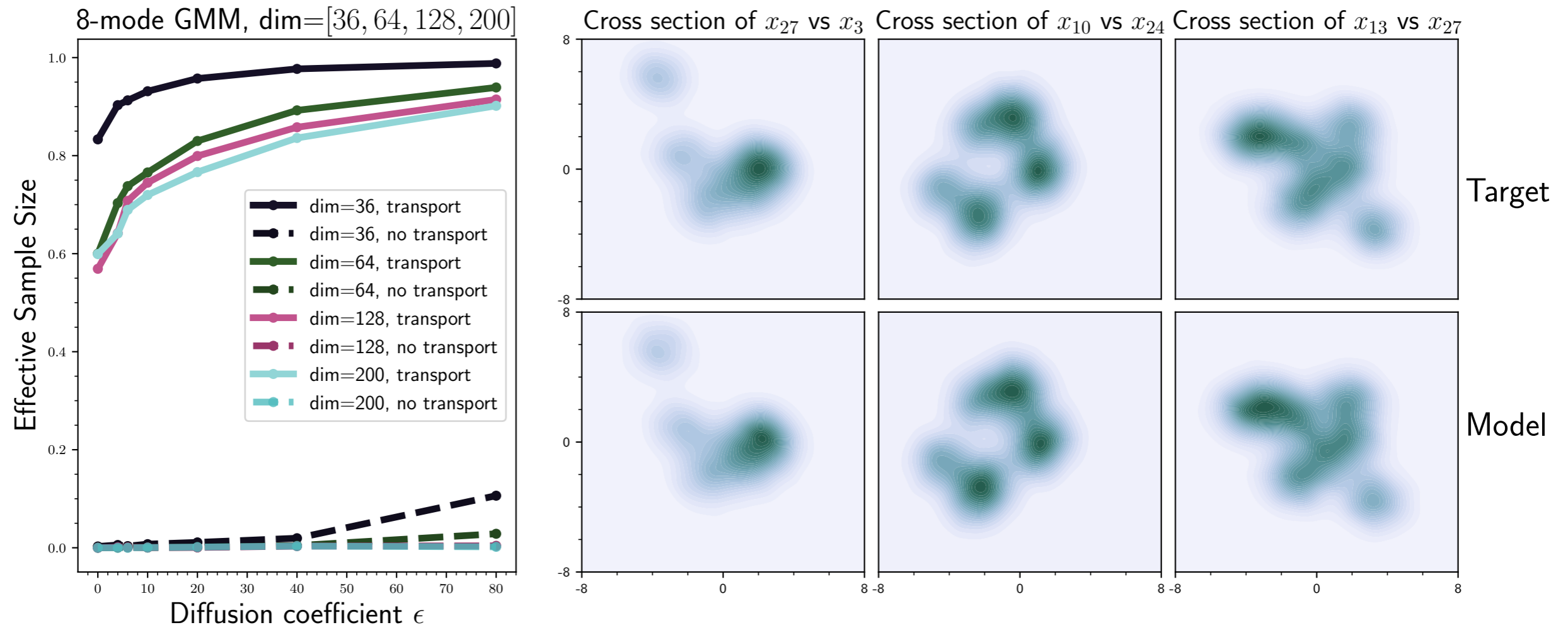


AIS and transport, $\epsilon = 4.0$



More diffusion helps more with transport than without

Scaling study: Use same neural network for multimodal GMM of growing dimension



Drop in ESS for deterministic flow $\epsilon_t = 0$ can be alleviated by growing ϵ_t

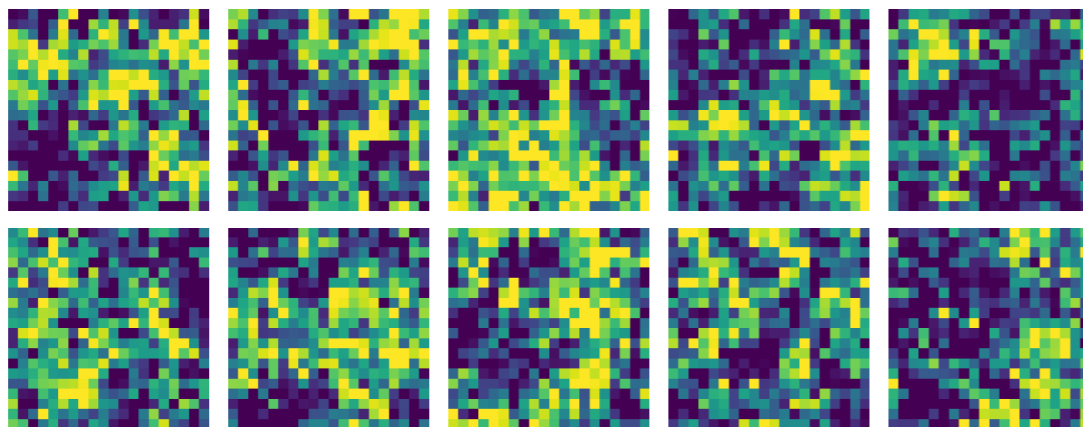
Less apparent in practice if you just use annealed Langevin along!

Standard test: ϕ^4 theory

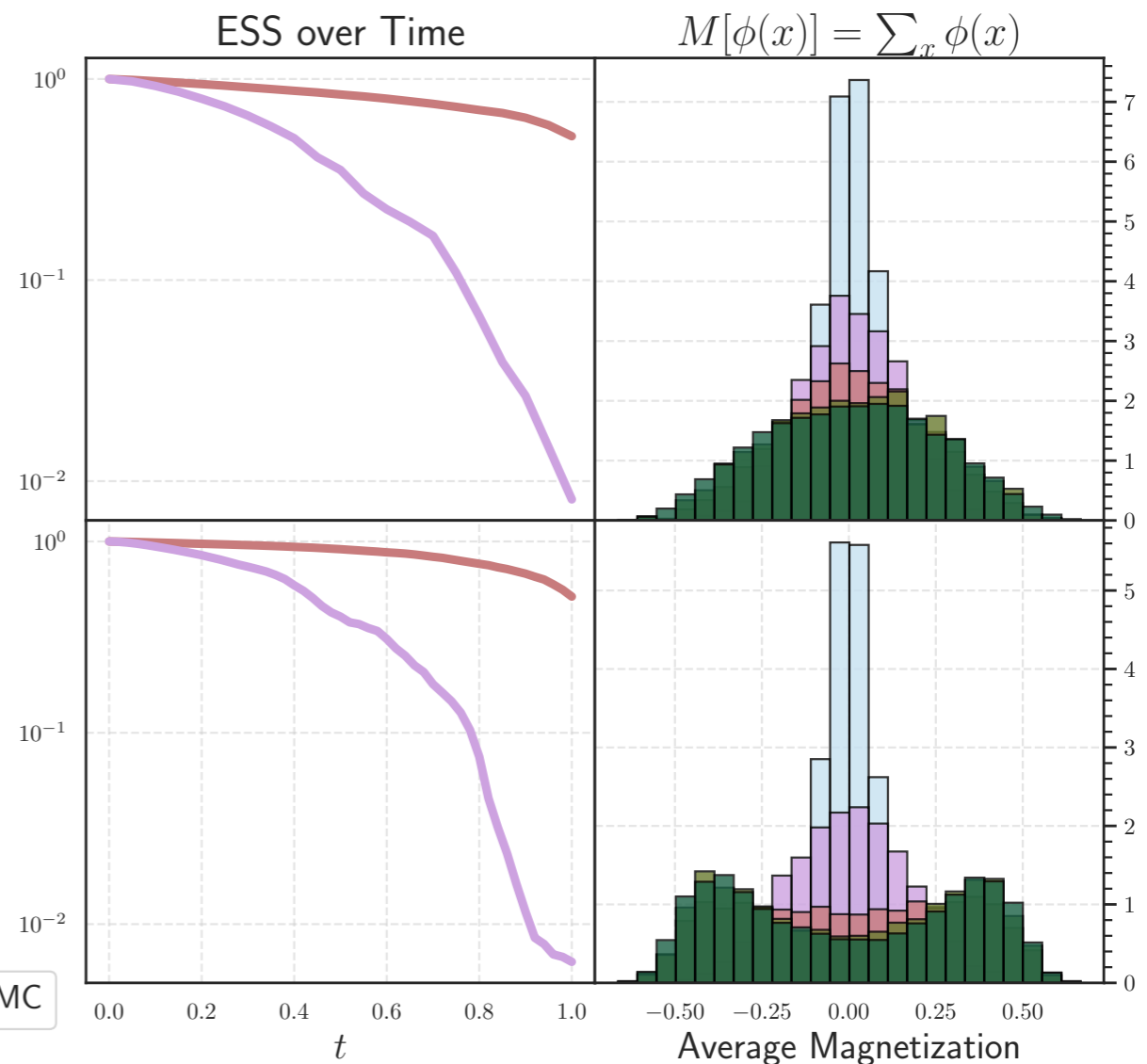
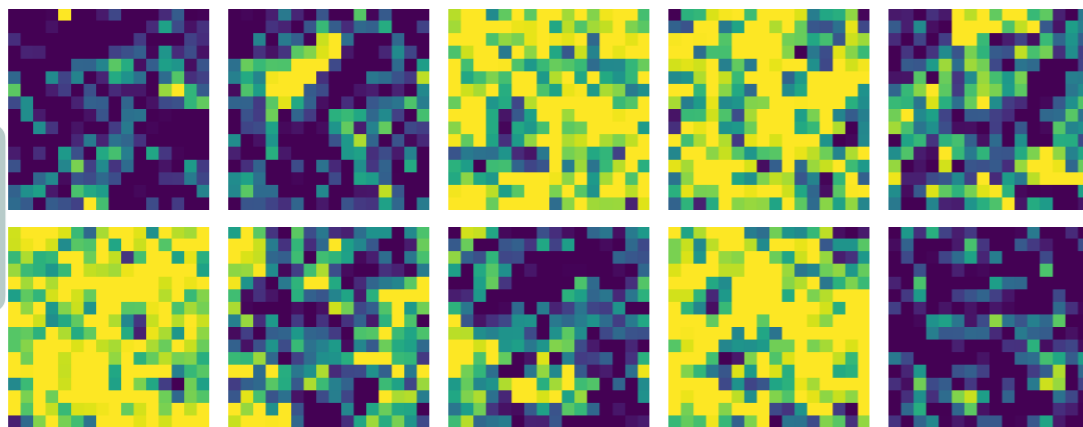
Choose energy interpolation in $m^2(t)$, $\lambda(t)$ for the action given by

$$U_t(\varphi) = \sum_x \left[-2 \sum_{\mu} \varphi_x \varphi_{x+\mu} \right] + (2D + m_t^2) \varphi_x^2 + \lambda_t \varphi_x^4$$

at phase transition



passed phase transition



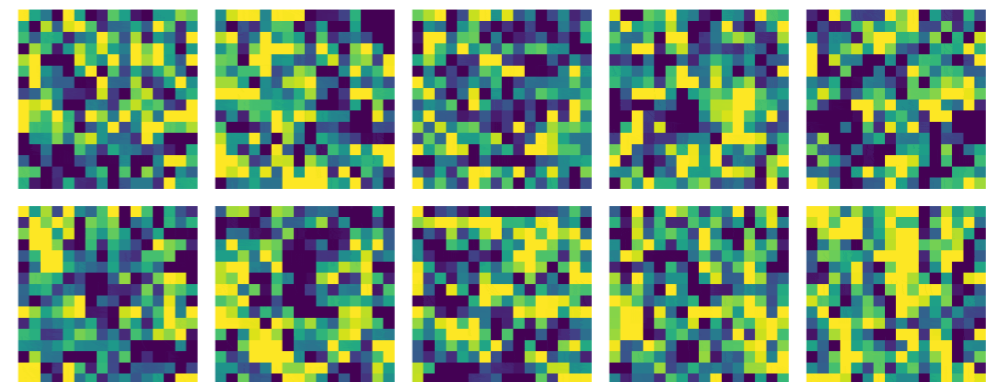
Conclusion

Dynamical formulation of unbiased sampling with transport based on Jarzynski equality

Loss functions do not require backpropagating through SDE

PINN loss is an off-policy loss! (See also Lorenz' previous talk)

Here's one more fun gif! Thanks!



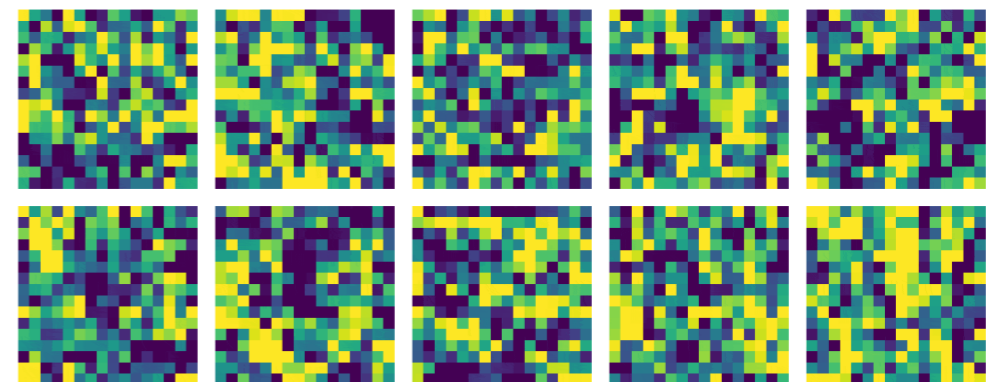
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Backup Slides

Computationally cheaper weights!

Note that the weights do not need a divergence if you use $b_t = \nabla \phi_t$

Proposition

Let (X_t, A_t) be the solution to the coupled SDE/ODE

$$dX_t = -\epsilon_t \nabla U_t(X_t) dt + \nabla \phi_t(X_t) dt + \sqrt{2\epsilon_t} dW_t, \quad X_0 \sim \rho_0$$

$$dB_t = \partial_t U_t(X_t) dt + \frac{1}{\epsilon_t} \partial_t \hat{\phi}_t(X_t) + \frac{1}{\epsilon_t} \left| \nabla \hat{\phi}_t(X_t) \right|^2 dt + \sqrt{\frac{2}{\epsilon_t}} \nabla \hat{\phi}_t(X_t) \cdot dW_t \quad A_0 = 0$$

then for all test functions $h(x)$, we have

$$\int_{\mathbb{R}^d} h(x) \rho_t(x) dx = \frac{\mathbb{E}[e^{A_t} h(x)]}{\mathbb{E}[e^{A_t}]} \quad Z_t / Z_0 = e^{-F_t + F_0} = \mathbb{E}[e^{A_t}]$$

$$\text{where } A_t = \frac{1}{\epsilon_t} [\hat{\phi}_t(X_t) - \hat{\phi}_0(X_0)] - B_t$$

works by using expanding $d\phi_t$ with Ito formula.

Proof:

Definition of the SDE/ODE for X_t, A_t with $b_t = \nabla \phi_t$

$$dX_t = -\varepsilon_t \nabla U(X_t) dt + \hat{\nabla} \phi_t(X_t) dt + \sqrt{2\varepsilon_t} dW_t, \quad \hat{X}_0 \sim \rho_0,$$

$$dA_t = \Delta \hat{\phi}_t(X_t) dt - \nabla U_t(X_t) \cdot \nabla \hat{\phi}_t(X_t) dt - \partial_t U_t(X_t) dt, \quad A_0 = 0,$$

Ito formula says

$$\begin{aligned} d\hat{\phi}_t(X_t) &= \partial_t \hat{\phi}_t(X_t) dt - \varepsilon_t \nabla \hat{\phi}_t(X_t) \cdot \nabla U(X_t) dt + \left| \nabla \hat{\phi}_t(X_t) \right|^2 dt \\ &\quad + \sqrt{2\varepsilon_t} \nabla \hat{\phi}_t(X_t) \cdot dW_t + \varepsilon_t \Delta \hat{\phi}_t(X_t) dt, \end{aligned}$$

Solving for $\Delta \phi_t$ allows us to write the relation

$$dA_t = \frac{1}{\varepsilon_t} d\hat{\phi}_t(X_t) dt + dB_t$$

where
$$dB_t = \partial_t U_t(X_t) dt + \frac{1}{\varepsilon_t} \partial_t \hat{\phi}_t(X_t) dt + \frac{1}{\varepsilon_t} \left| \nabla \hat{\phi}_t(X_t) \right|^2 dt + \sqrt{\frac{2}{\varepsilon_t}} \nabla \hat{\phi}_t(X_t) \cdot dW_t$$

Proposition 5 (KL control). *Let $\hat{\rho}_t$ be the solution to the transport equation*

$$\partial_t \hat{\rho}_t = -\nabla \cdot (\hat{b}_t \rho_t), \quad \hat{\rho}_{t=0} = \rho_0 \quad (27)$$

where $\hat{b}_t(x)$ is some predefined velocity field. Then, given any estimate \hat{F}_t of the exact free energy F_t , we have

$$D_{KL}(\hat{\rho}_{t=1} || \rho_1) \leq \sqrt{L_{PINN}^{T=1}(\hat{b}, \hat{F})}. \quad (28)$$

This proposition is proven in Appendix [5.1](#).