
The hadronic running of the electromagnetic coupling and electroweak mixing angle

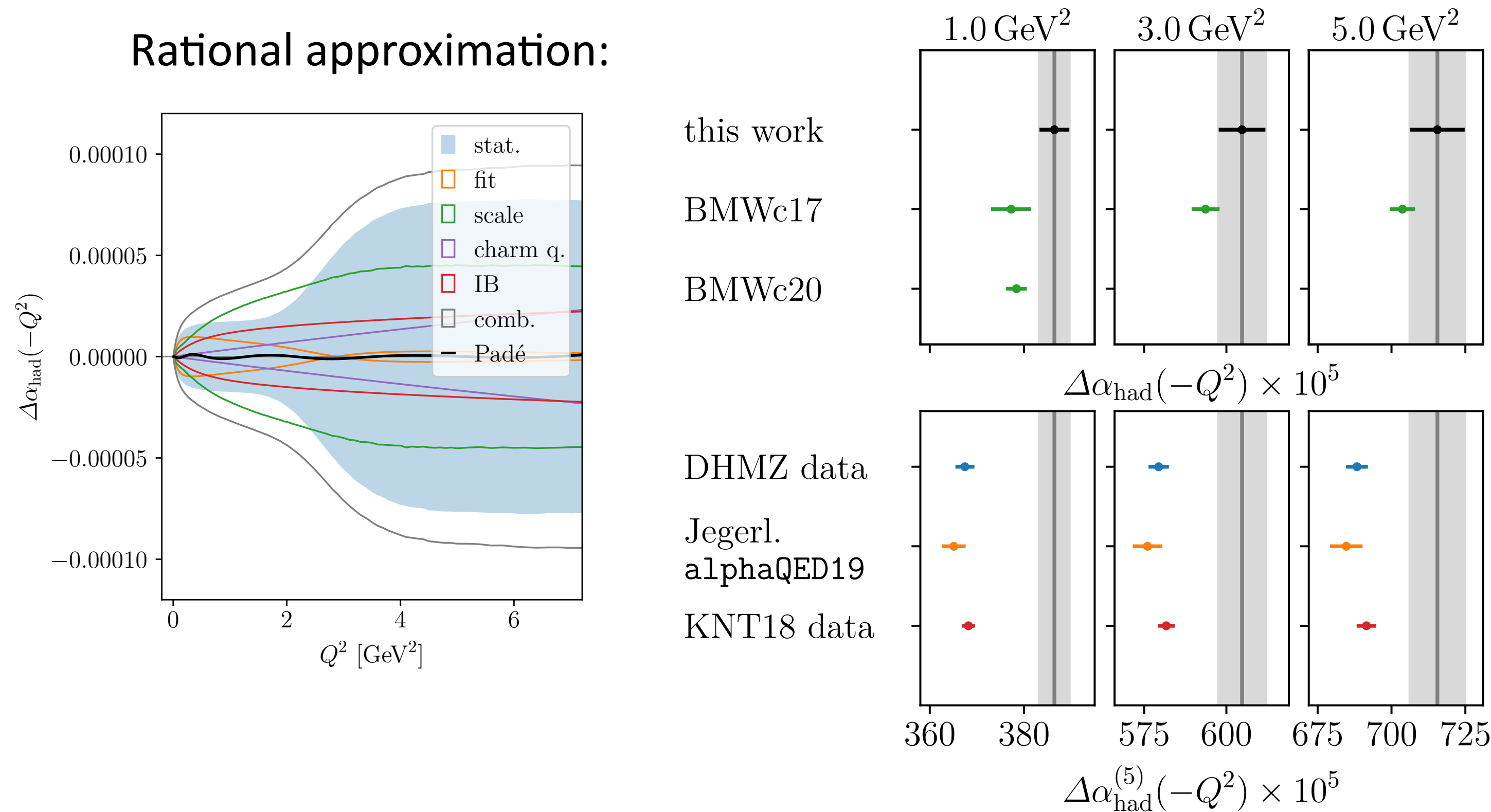
Marco Cè, Antoine Gérardin, Georg von Hippel, Harvey Meyer, Kohtaroh Miura,
Konstantin Ottnad, Andreas Risch, Teseo San José, Jonas Wilhelm, Hartmut Wittig

The 39th International Symposium on Lattice Field Theory — Lattice 2022
Rheinische Friedrich Wilhelms Universität Bonn
8–13 August 2022

Lattice result for the hadronic running of α

[Cè et al., arXiv:2203.08676]

Starting point: Results for $\Delta\alpha_{\text{had}}(-Q^2)$ for **Euclidean** momenta $0 \leq Q^2 \leq 7 \text{ GeV}^2$ [T. San José, TUE 17:10]



- Mainz/CLS and BMWc (2017) differ by $2-3\%$ at the level of $1-2\sigma$
- Tension between Mainz/CLS and phenomenology by $\sim 3\sigma$ for $Q^2 \gtrsim 3 \text{ GeV}^2$
- Tension increases to $\gtrsim 5\sigma$ for $Q^2 \lesssim 2 \text{ GeV}^2$ (smaller statistical error due to ansatz for continuum extrapolation)

Systematic uncertainties from fit ansatz, scale setting, charm quenching, isospin-breaking and missing bottom quark contribution (five flavour theory) included in error budget

Consistency of the Standard Model

Hadronic running at Z -pole: $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) \rightarrow$ key quantity in global electroweak fit

$\Delta\alpha_{\text{had}}$ related to hadronic vacuum polarisation contribution to muon $g - 2$:

$$\Delta\alpha_{\text{had}}^{(5)}(q^2) = -\frac{\alpha q^2}{3\pi} \oint_{m_{\pi^0}^2}^{\infty} ds \frac{R(s)}{s(s - q^2)}, \quad R(s) = \frac{3s}{4\pi\alpha(s)} \sigma(e^+e^- \rightarrow \text{hadrons})$$

$$a_{\mu}^{\text{hvp}} = \left(\frac{\alpha m_{\mu}}{3\pi}\right)^2 \int_{m_{\pi^0}^2}^{\infty} ds \frac{R(s) \hat{K}(s)}{s^2}, \quad 0.63 \lesssim \hat{K}(s) \leq 1$$

Consistency of the Standard Model

Hadronic running at Z -pole: $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) \rightarrow$ key quantity in global electroweak fit

$\Delta\alpha_{\text{had}}$ related to hadronic vacuum polarisation contribution to muon $g - 2$:

$$\Delta\alpha_{\text{had}}^{(5)}(q^2) = -\frac{\alpha q^2}{3\pi} \oint_{m_{\pi^0}^2}^{\infty} ds \frac{R(s)}{s(s - q^2)}, \quad R(s) = \frac{3s}{4\pi\alpha(s)} \sigma(e^+e^- \rightarrow \text{hadrons})$$

$$a_{\mu}^{\text{hvp}} = \left(\frac{\alpha m_{\mu}}{3\pi}\right)^2 \int_{m_{\pi^0}^2}^{\infty} ds \frac{R(s) \hat{K}(s)}{s^2}, \quad 0.63 \lesssim \hat{K}(s) \leq 1$$

Different kernel functions: low-energy region receives smaller weight in $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$

Consistency of the Standard Model

Hadronic running at Z -pole: $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) \rightarrow$ key quantity in global electroweak fit

$\Delta\alpha_{\text{had}}$ related to hadronic vacuum polarisation contribution to muon $g - 2$:

$$\Delta\alpha_{\text{had}}^{(5)}(q^2) = -\frac{\alpha q^2}{3\pi} \oint_{m_{\pi^0}^2}^{\infty} ds \frac{R(s)}{s(s - q^2)}, \quad R(s) = \frac{3s}{4\pi \alpha(s)} \sigma(e^+e^- \rightarrow \text{hadrons})$$

$$a_{\mu}^{\text{hvp}} = \left(\frac{\alpha m_{\mu}}{3\pi}\right)^2 \int_{m_{\pi^0}^2}^{\infty} ds \frac{R(s) \hat{K}(s)}{s^2}, \quad 0.63 \lesssim \hat{K}(s) \leq 1$$

Different kernel functions: low-energy region receives smaller weight in $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$

Task: convert lattice result for $\Delta\alpha_{\text{had}}^{(5)}(-Q^2)$ to an estimate of $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$ and compare to global electroweak fit

How to evaluate $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$

Method 1: Direct dispersion relation (DR)

$$\Delta\alpha_{\text{had}}^{(5)}(q^2) = -\frac{\alpha q^2}{3\pi} \oint_{m_{\pi^0}^2}^{\infty} ds \frac{R(s)}{s(s - q^2)} \quad \text{for } q^2 = M_Z^2$$

→ use combination of perturbation theory and experimental data for R -ratio

How to evaluate $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$

Method 1: Direct dispersion relation (DR)

$$\Delta\alpha_{\text{had}}^{(5)}(q^2) = -\frac{\alpha q^2}{3\pi} \oint_{m_{\pi^0}^2}^{\infty} ds \frac{R(s)}{s(s - q^2)} \quad \text{for } q^2 = M_Z^2$$

→ use combination of perturbation theory and experimental data for R -ratio

Method 2: Adler function approach, aka. “Euclidean split technique”

$$\begin{aligned} \Delta\alpha_{\text{had}}^{(5)}(M_Z^2) &= \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2) \\ &+ [\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2)] \\ &+ [\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-M_Z^2)] \end{aligned}$$

[Chetyrkin et al., Nucl Phys B482 (1996) 213; Eidelman et al., Phys Lett B454 (1999) 369; Jegerlehner, hep-ph/9901386, arXiv:0807.4206]

How to evaluate $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$

Method 1: Direct dispersion relation (DR)

$$\Delta\alpha_{\text{had}}^{(5)}(q^2) = -\frac{\alpha q^2}{3\pi} \oint_{m_{\pi^0}^2}^{\infty} ds \frac{R(s)}{s(s - q^2)} \quad \text{for } q^2 = M_Z^2$$

→ use combination of perturbation theory and experimental data for R -ratio

Method 2: Adler function approach, aka. “Euclidean split technique”

$$\begin{aligned} \Delta\alpha_{\text{had}}^{(5)}(M_Z^2) &= \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2) \quad \leftarrow \text{lattice QCD or DR for } q^2 = -Q_0^2 \\ &+ [\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2)] \\ &+ [\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-M_Z^2)] \end{aligned}$$

[Chetyrkin et al., Nucl Phys B482 (1996) 213; Eidelman et al., Phys Lett B454 (1999) 369; Jegerlehner, hep-ph/9901386, arXiv:0807.4206]

How to evaluate $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$

Method 1: Direct dispersion relation (DR)

$$\Delta\alpha_{\text{had}}^{(5)}(q^2) = -\frac{\alpha q^2}{3\pi} \oint_{m_{\pi^0}^2}^{\infty} ds \frac{R(s)}{s(s - q^2)} \quad \text{for } q^2 = M_Z^2$$

→ use combination of perturbation theory and experimental data for R -ratio

Method 2: Adler function approach, aka. “Euclidean split technique”

$$\begin{aligned} \Delta\alpha_{\text{had}}^{(5)}(M_Z^2) &= \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2) && \leftarrow \text{lattice QCD or DR for } q^2 = -Q_0^2 \\ &+ [\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2)] && \leftarrow \text{Adler function in pQCD or DR} \\ &+ [\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-M_Z^2)] \end{aligned}$$

[Chetyrkin et al., Nucl Phys B482 (1996) 213; Eidelman et al., Phys Lett B454 (1999) 369; Jegerlehner, hep-ph/9901386, arXiv:0807.4206]

How to evaluate $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$

Method 1: Direct dispersion relation (DR)

$$\Delta\alpha_{\text{had}}^{(5)}(q^2) = -\frac{\alpha q^2}{3\pi} \oint_{m_{\pi^0}^2}^{\infty} ds \frac{R(s)}{s(s - q^2)} \quad \text{for } q^2 = M_Z^2$$

→ use combination of perturbation theory and experimental data for R -ratio

Method 2: Adler function approach, aka. “Euclidean split technique”

$$\begin{aligned} \Delta\alpha_{\text{had}}^{(5)}(M_Z^2) &= \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2) && \leftarrow \text{lattice QCD or DR for } q^2 = -Q_0^2 \\ &+ [\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2)] && \leftarrow \text{Adler function in pQCD or DR} \\ &+ [\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-M_Z^2)] && \leftarrow \text{pQCD} \end{aligned}$$

[Chetyrkin et al., Nucl Phys B482 (1996) 213; Eidelman et al., Phys Lett B454 (1999) 369; Jegerlehner, hep-ph/9901386, arXiv:0807.4206]

Euclidean split technique and the Adler function

Adler function:

$$D(-s) = \frac{3\pi}{\alpha} s \frac{d}{ds} \Delta\alpha_{\text{had}}(s)$$

Euclidean split technique and the Adler function

Adler function:

$$D(-s) = \frac{3\pi}{\alpha} s \frac{d}{ds} \Delta\alpha_{\text{had}}(s)$$

$D(Q^2)$ known in massive QCD perturbation theory at three loops

$$\left[\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2) \right]_{\text{pQCD/Adler}} = \frac{\alpha}{3\pi} \int_{Q_0^2}^{M_Z^2} \frac{dQ^2}{Q^2} D(Q^2)$$

Euclidean split technique and the Adler function

Adler function:

$$D(-s) = \frac{3\pi}{\alpha} s \frac{d}{ds} \Delta\alpha_{\text{had}}(s)$$

$D(Q^2)$ known in massive QCD perturbation theory at three loops

$$\left[\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2) \right]_{\text{pQCD/Adler}} = \frac{\alpha}{3\pi} \int_{Q_0^2}^{M_Z^2} \frac{dQ^2}{Q^2} D(Q^2)$$

Relation of $D(Q^2)$ and R -ratio:

$$D(Q^2) = Q^2 \int_{m_{\pi^0}^2}^{\infty} ds \frac{R(s)}{(s + Q^2)^2}$$

Euclidean split technique and the Adler function

Adler function:

$$D(-s) = \frac{3\pi}{\alpha} s \frac{d}{ds} \Delta\alpha_{\text{had}}(s)$$

$D(Q^2)$ known in massive QCD perturbation theory at three loops

$$\left[\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2) \right]_{\text{pQCD/Adler}} = \frac{\alpha}{3\pi} \int_{Q_0^2}^{M_Z^2} \frac{dQ^2}{Q^2} D(Q^2)$$

Relation of $D(Q^2)$ and R -ratio:

$$D(Q^2) = Q^2 \int_{m_{\pi^0}^2}^{\infty} ds \frac{R(s)}{(s + Q^2)^2}$$

Direct DR:

$$\left[\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2) \right]_{\text{DR}} = \frac{\alpha(M_Z^2 - Q_0^2)}{3\pi} \int_{m_{\pi^0}^2}^{\infty} ds \frac{R(s)}{(s + Q_0^2)(s + M_Z^2)}$$

Euclidean split technique and the Adler function

Adler function:

$$D(-s) = \frac{3\pi}{\alpha} s \frac{d}{ds} \Delta\alpha_{\text{had}}(s)$$

$D(Q^2)$ known in massive QCD perturbation theory at three loops

$$\left[\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2) \right]_{\text{pQCD/Adler}} = \frac{\alpha}{3\pi} \int_{Q_0^2}^{M_Z^2} \frac{dQ^2}{Q^2} D(Q^2)$$

Relation of $D(Q^2)$ and R -ratio:

$$D(Q^2) = Q^2 \int_{m_{\pi^0}^2}^{\infty} ds \frac{R(s)}{(s + Q^2)^2}$$

Direct DR:

$$\left[\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2) \right]_{\text{DR}} = \frac{\alpha(M_Z^2 - Q_0^2)}{3\pi} \int_{m_{\pi^0}^2}^{\infty} ds \frac{R(s)}{(s + Q_0^2)(s + M_Z^2)}$$

Perturbation theory:

$$\left[\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) \right] = 0.000\,045(2) \quad [\text{Jegerlehner, CERN Yellow Report, 2020}]$$

Euclidean split technique and the Adler function

Technical advantages of Euclidean split technique:

- Integration over Euclidean squared momentum: resonances and physical thresholds are absent

Euclidean split technique and the Adler function

Technical advantages of Euclidean split technique:

- Integration over Euclidean squared momentum: resonances and physical thresholds are absent
- Only $\Delta\alpha_{\text{had}}(-Q_0^2)$ depends on experimental data in low-energy regime;
direct DR evaluation of $\Delta\alpha_{\text{had}}(M_Z^2)$ requires precise experimental data up to much higher energies

Euclidean split technique and the Adler function

Technical advantages of Euclidean split technique:

- Integration over Euclidean squared momentum: resonances and physical thresholds are absent
- Only $\Delta\alpha_{\text{had}}(-Q_0^2)$ depends on experimental data in low-energy regime;
direct DR evaluation of $\Delta\alpha_{\text{had}}(M_Z^2)$ requires precise experimental data up to much higher energies
- $\Delta\alpha_{\text{had}}(-Q_0^2)$ accessible in lattice QCD or by result from MUonE experiment

Euclidean split technique and the Adler function

Technical advantages of Euclidean split technique:

- Integration over Euclidean squared momentum: resonances and physical thresholds are absent
- Only $\Delta\alpha_{\text{had}}(-Q_0^2)$ depends on experimental data in low-energy regime;
direct DR evaluation of $\Delta\alpha_{\text{had}}(M_Z^2)$ requires precise experimental data up to much higher energies
- $\Delta\alpha_{\text{had}}(-Q_0^2)$ accessible in lattice QCD or by result from MUonE experiment
- Can check validity of perturbative QCD down to $Q_0^2 \sim 5 \text{ GeV}^2$:

$$\left[\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2) \right]_{\text{pQCD/Adler}} \quad \text{versus} \quad \left[\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2) \right]_{\text{DR}}$$

Evaluation of $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$

[Cè et al., arXiv:2203.08676]

Input: Lattice result for $\Delta\alpha_{\text{had}}(-Q_0^2)$ for $Q_0^2 = 3 - 7 \text{ GeV}^2$

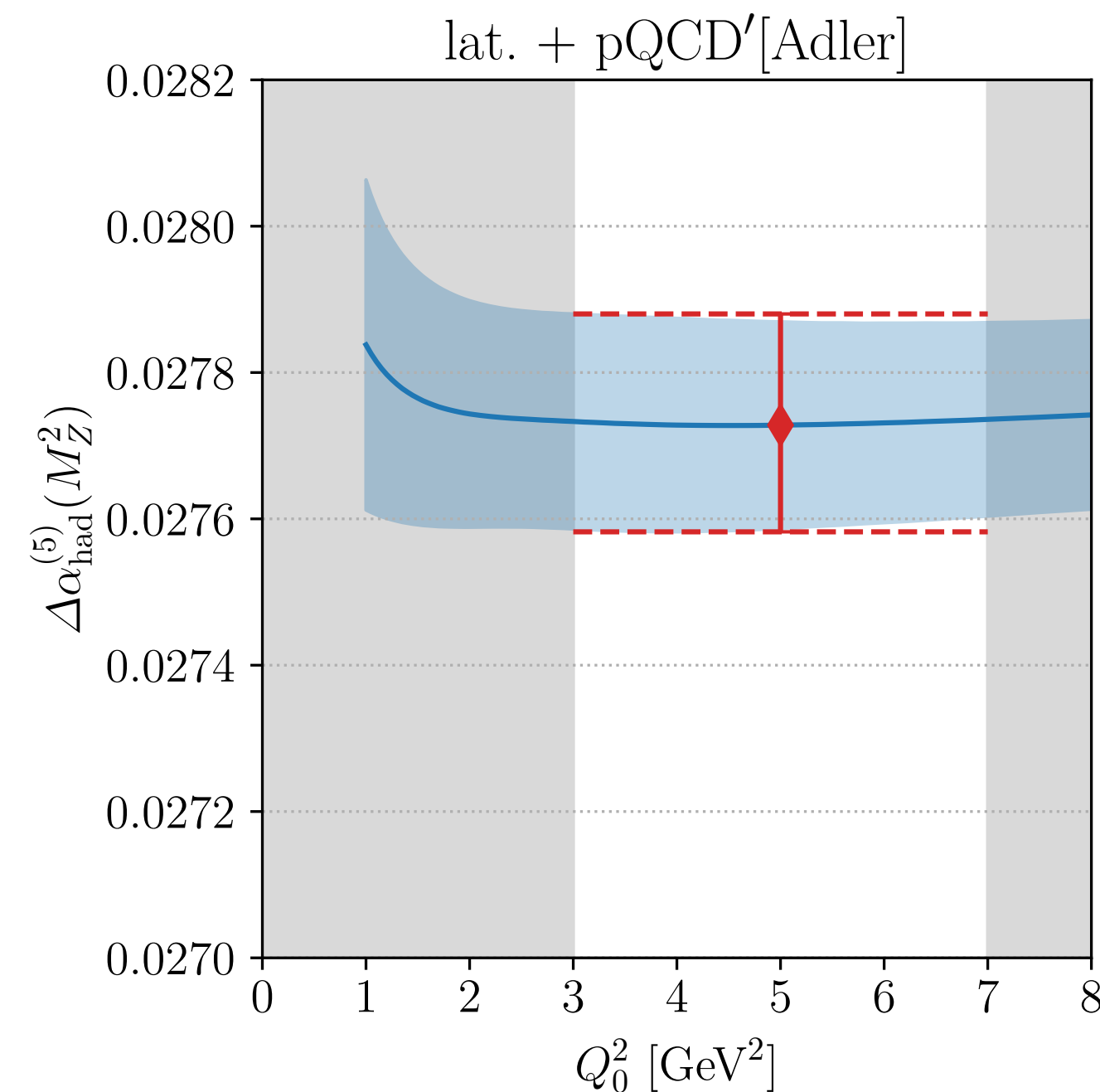
Evaluate $\left[\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2)\right]_{\text{pQCD/Adler}}$ using Jegerlehner's software package **pQCDAdler**

Evaluation of $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$

[Cè et al., arXiv:2203.08676]

Input: Lattice result for $\Delta\alpha_{\text{had}}(-Q_0^2)$ for $Q_0^2 = 3 - 7 \text{ GeV}^2$

Evaluate $\left[\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2)\right]_{\text{pQCD/Adler}}$ using Jegerlehner's software package **pQCDAdler**



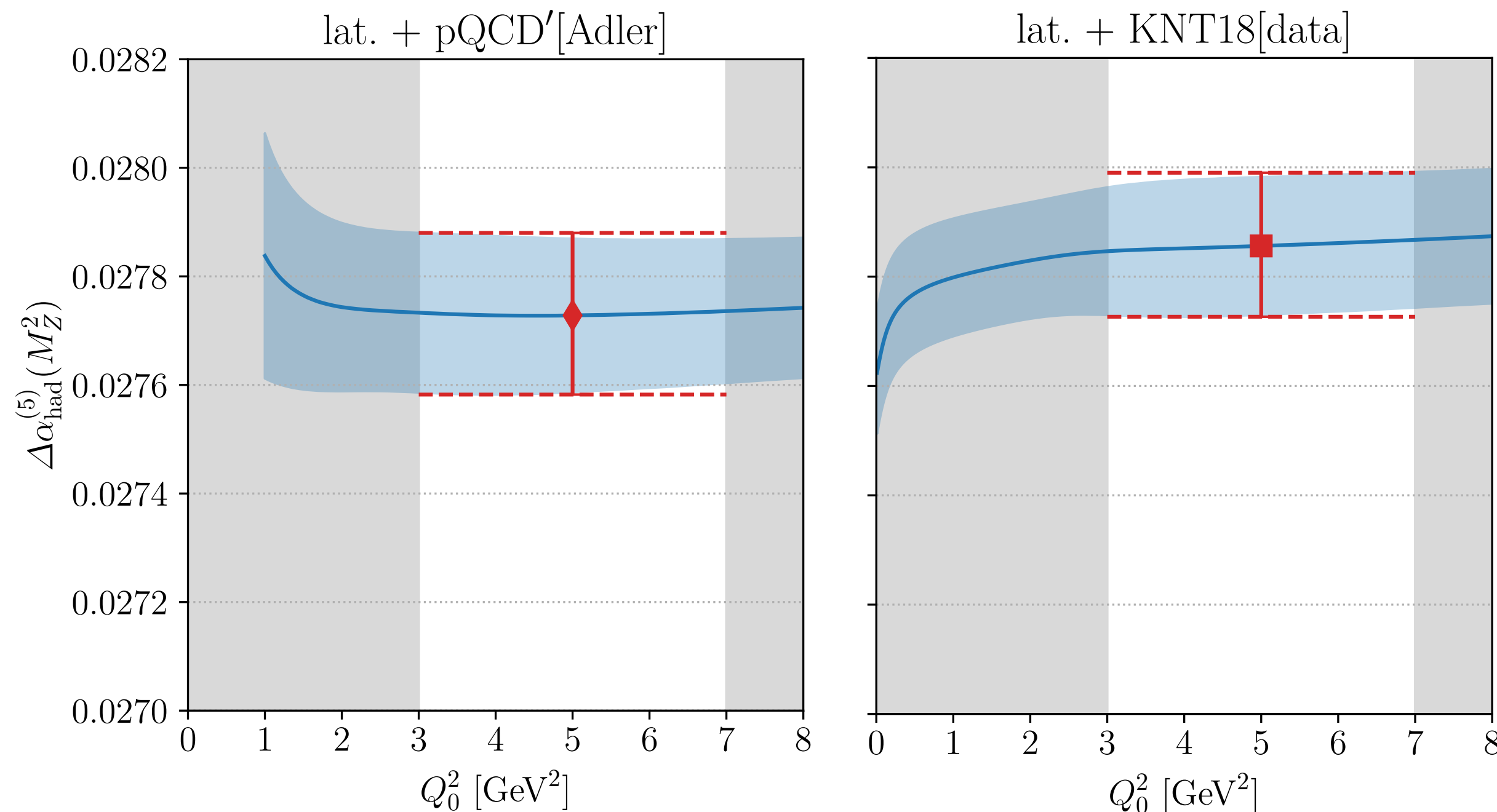
Evaluation of $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$

[Cè et al., arXiv:2203.08676]

Input: Lattice result for $\Delta\alpha_{\text{had}}(-Q_0^2)$ for $Q_0^2 = 3 - 7 \text{ GeV}^2$

Evaluate $\left[\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2)\right]_{\text{pQCD/Adler}}$ using Jegerlehner's software package **pQCDAdler**

Alternatively evaluate $\left[\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2)\right]_{\text{DR}}$ using experimental R -ratio



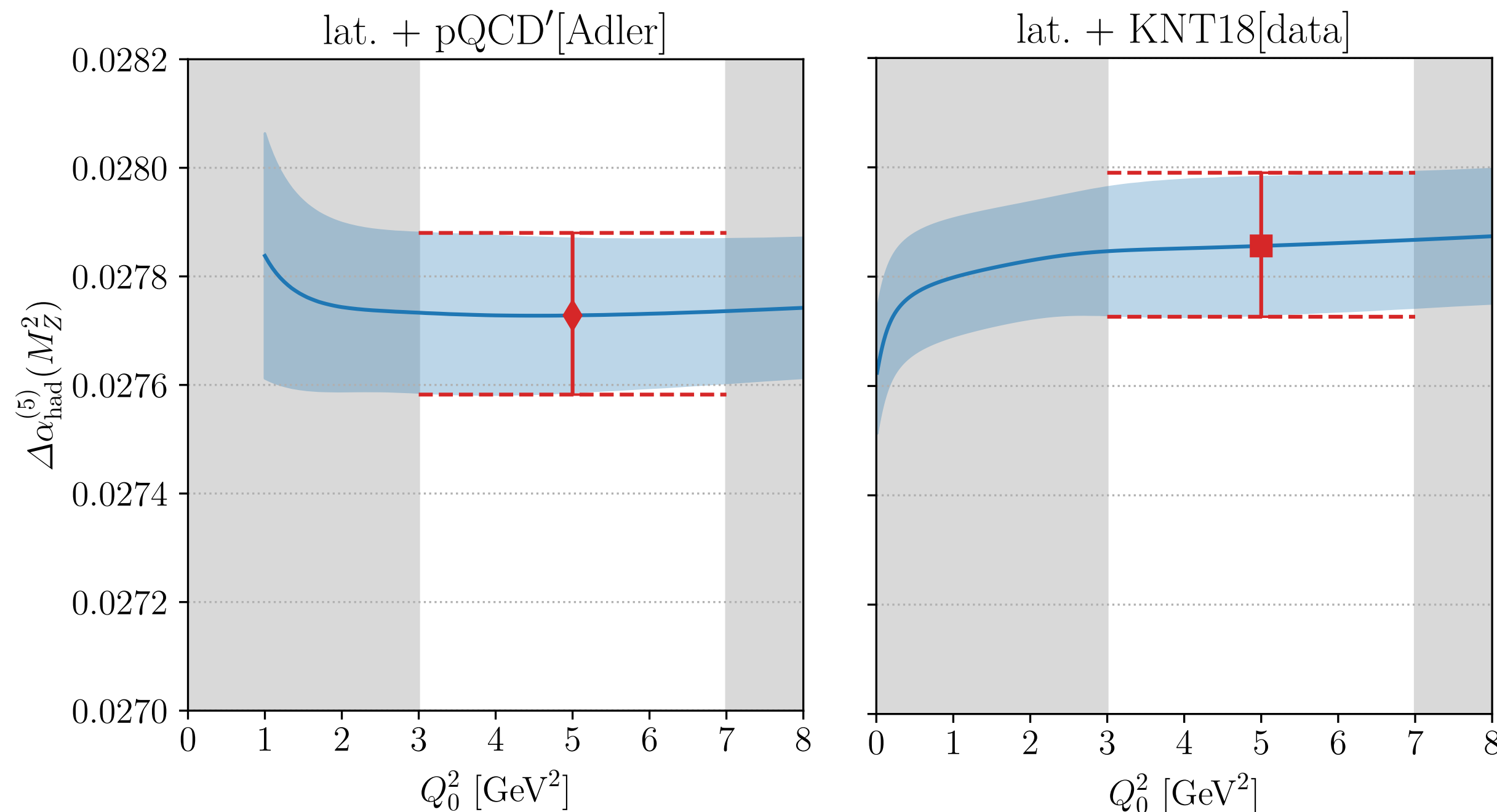
Evaluation of $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$

[Cè et al., arXiv:2203.08676]

Input: Lattice result for $\Delta\alpha_{\text{had}}(-Q_0^2)$ for $Q_0^2 = 3 - 7 \text{ GeV}^2$

Evaluate $\left[\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2)\right]_{\text{pQCD/Adler}}$ using Jegerlehner's software package **pQCDAdler**

Alternatively evaluate $\left[\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-Q_0^2)\right]_{\text{DR}}$ using experimental R -ratio

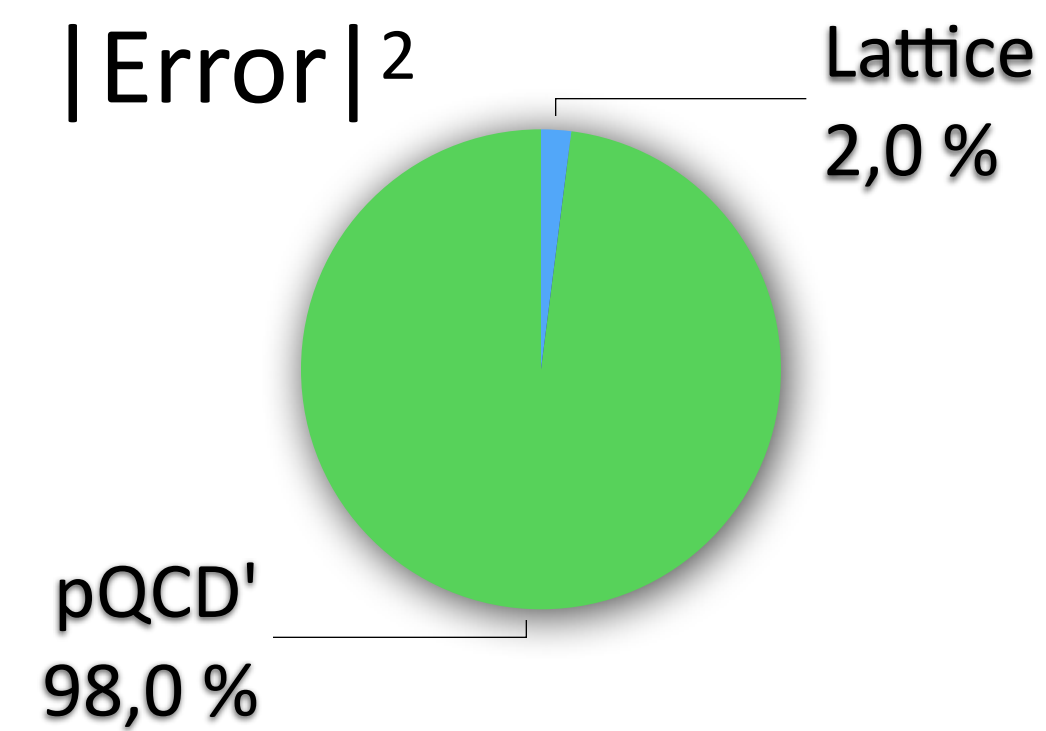
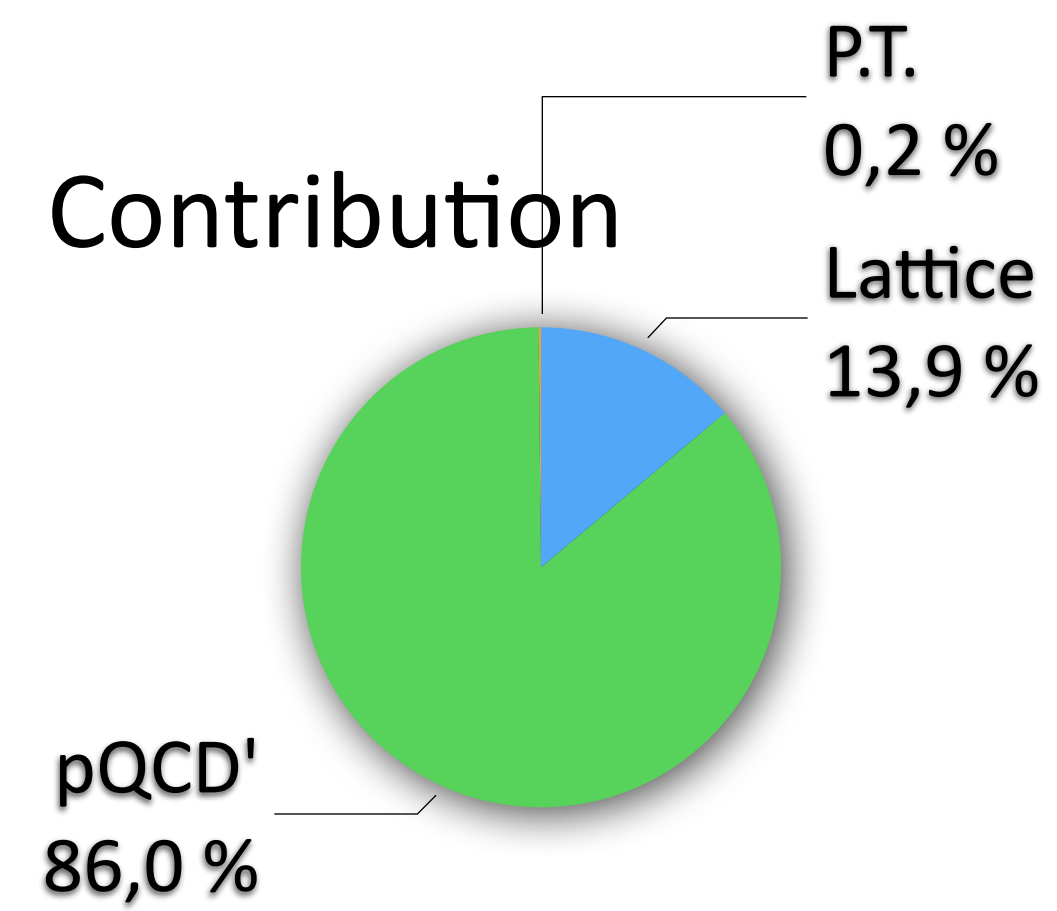


Final estimate: lattice + pQCD/Adler

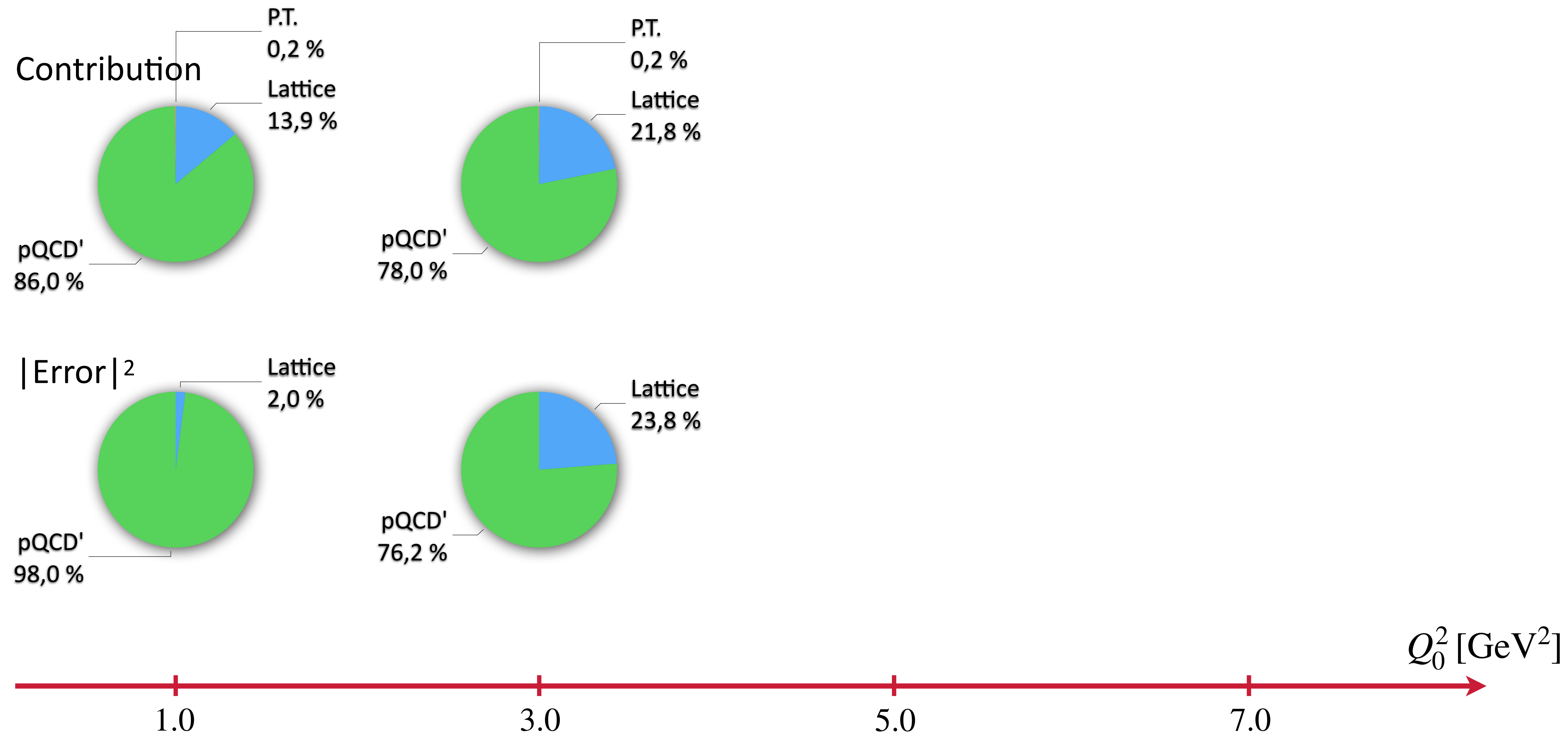
$$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) = 0.027\,73(9)_{\text{lat}}(2)_{\text{btm}}(12)_{\text{pQCD}} \\ = 0.027\,73(15)$$

(error contains ambiguity in the choice of Q_0^2)

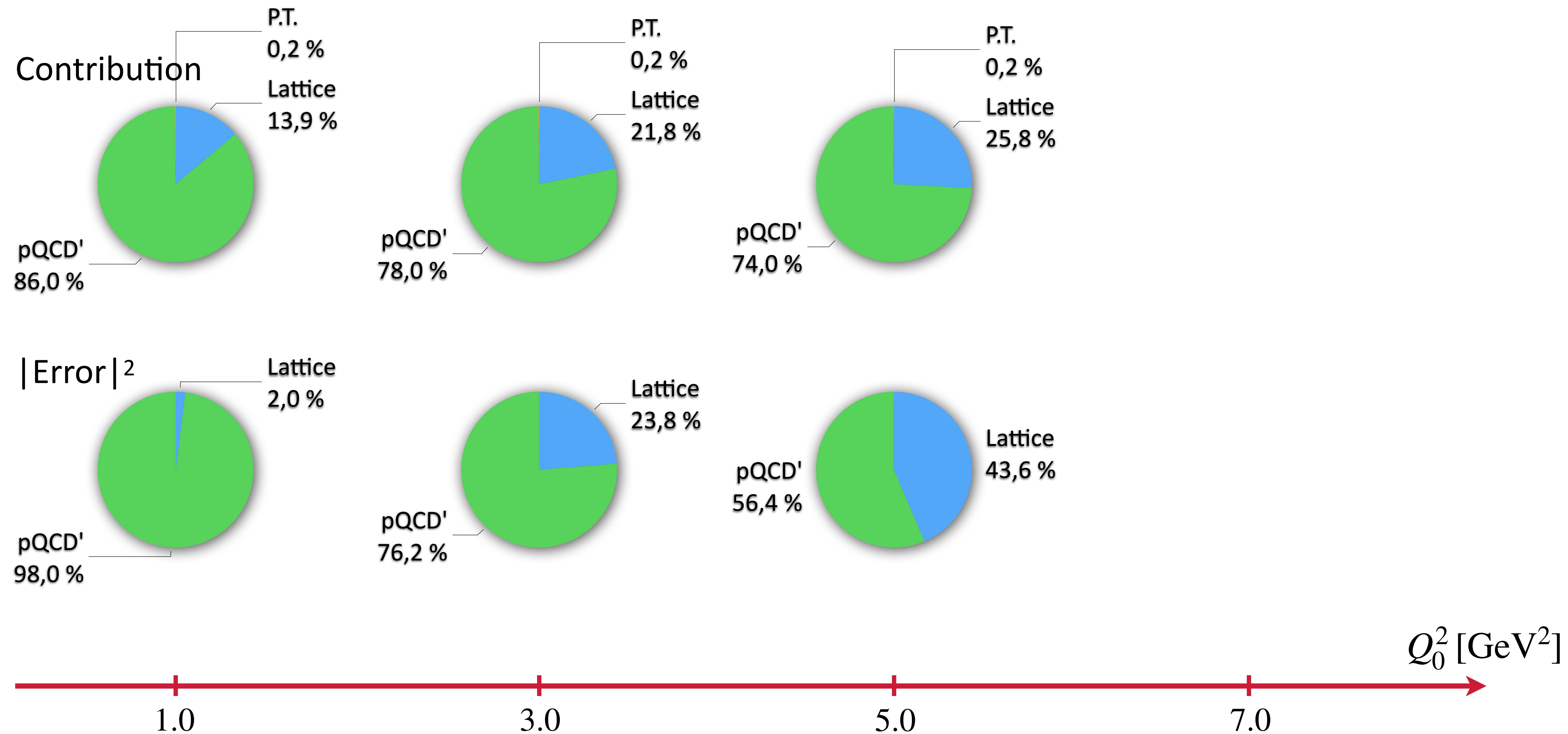
Euclidean split technique: relative contributions to $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$



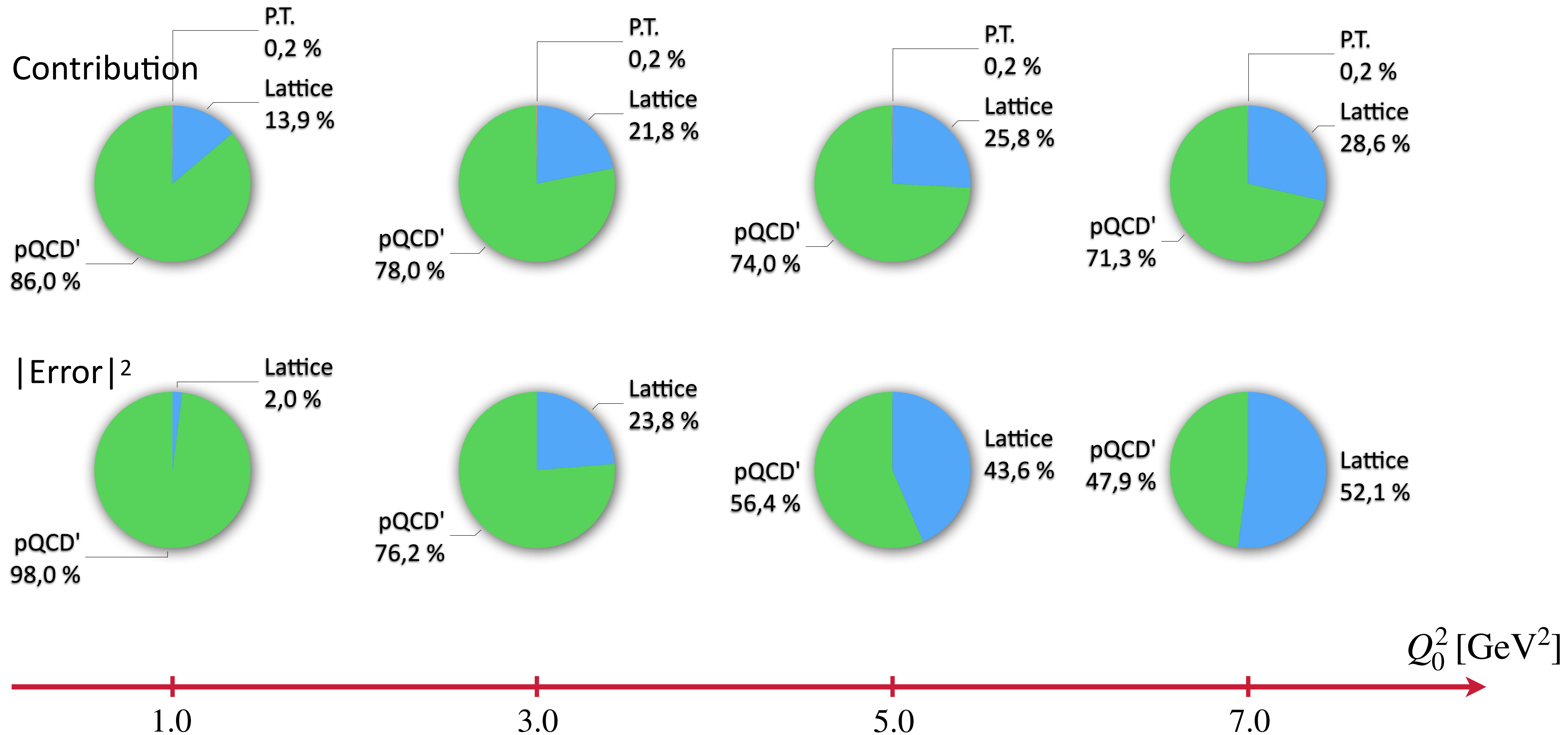
Euclidean split technique: relative contributions to $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$



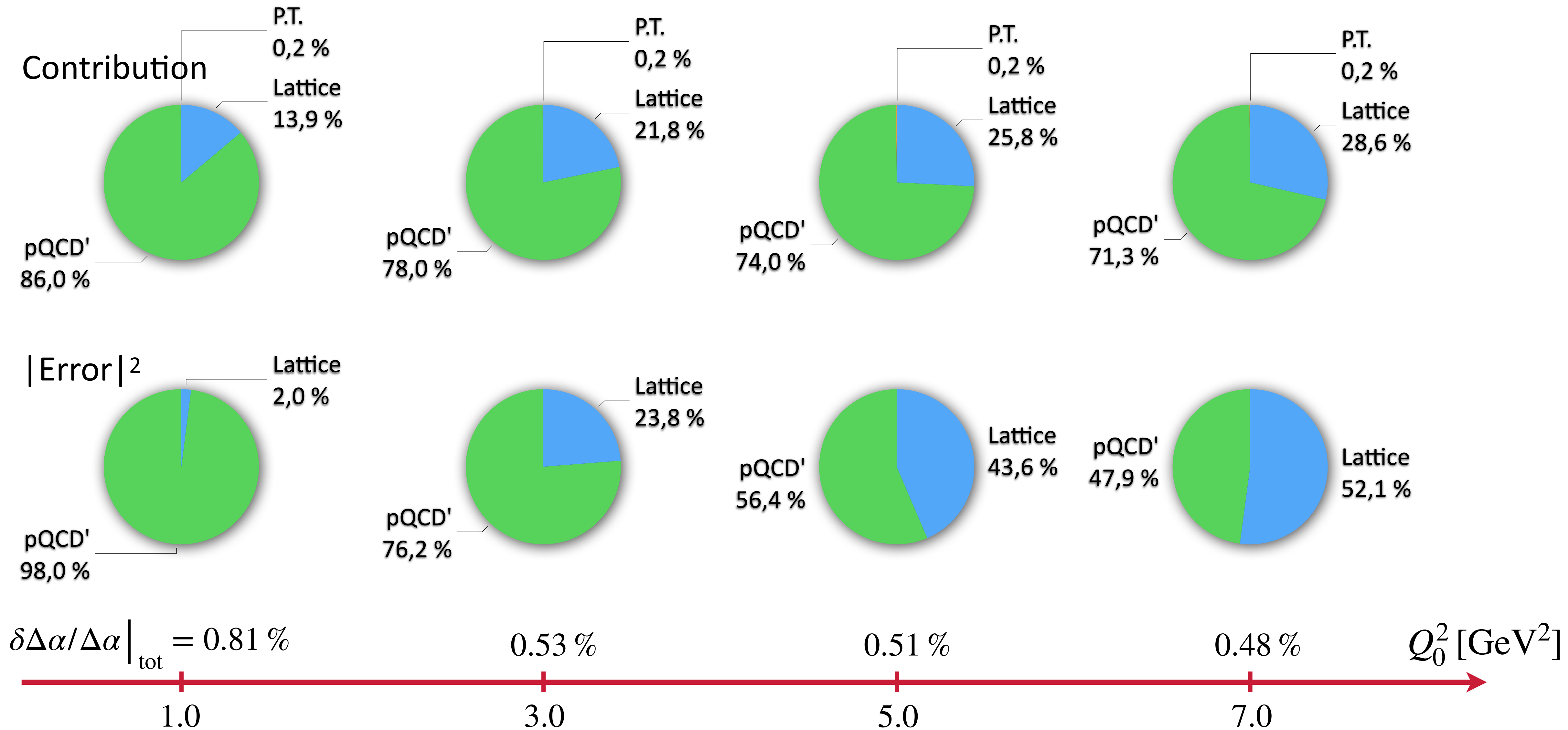
Euclidean split technique: relative contributions to $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$



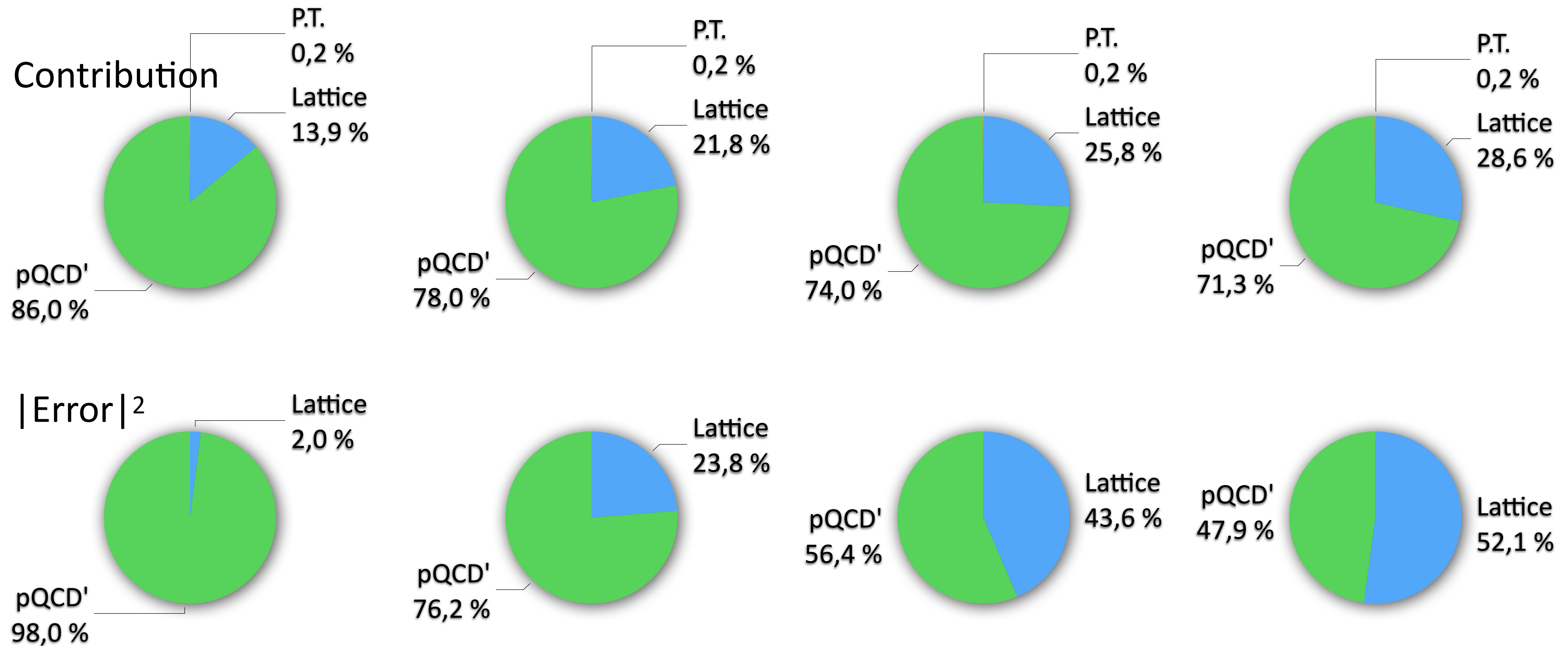
Euclidean split technique: relative contributions to $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$



Euclidean split technique: relative contributions to $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$



Euclidean split technique: relative contributions to $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$



Can tune Q_0^2 to optimise the reliability

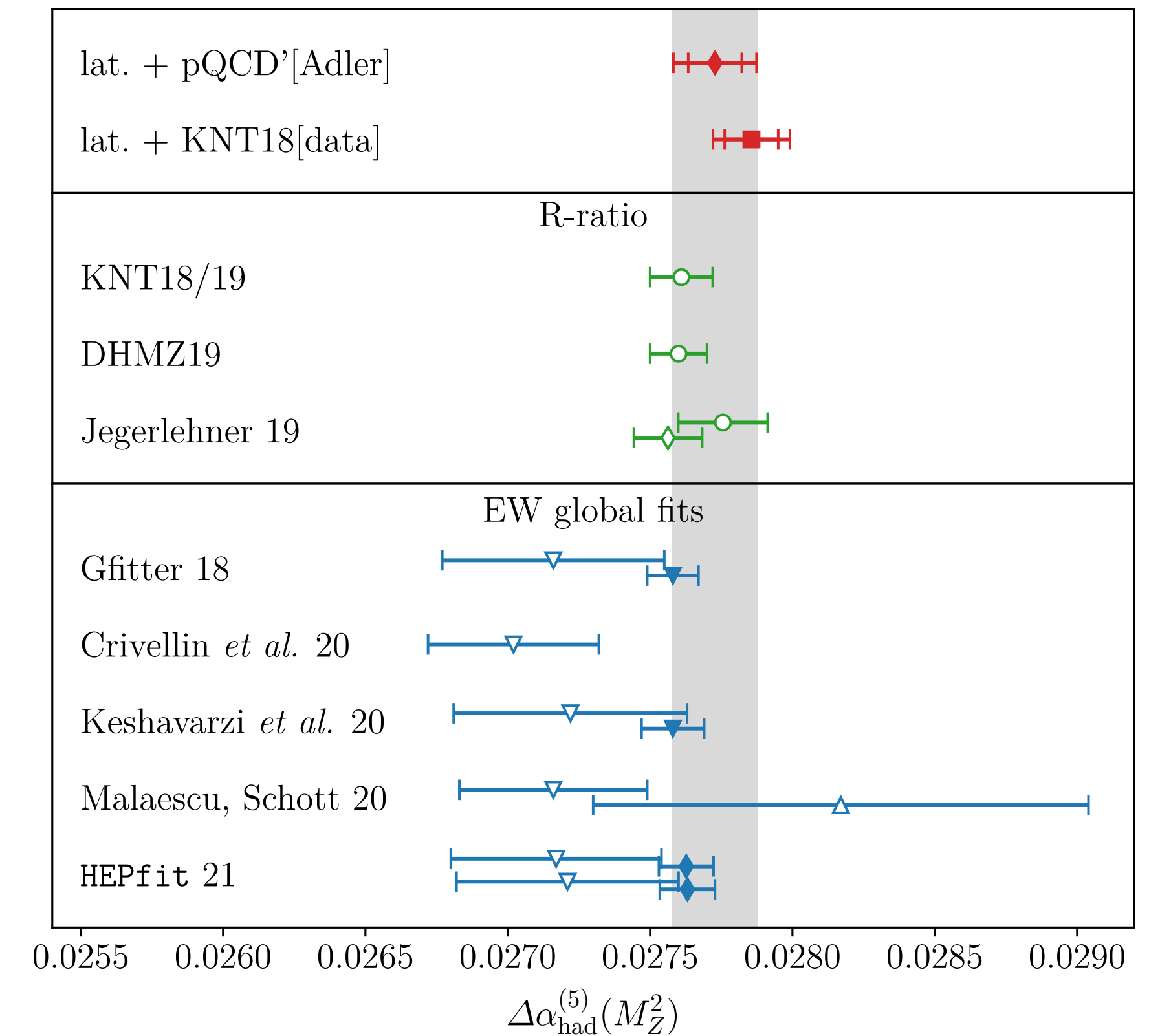
- Lattice QCD accounts for $\sim 25 \%$ of the value of $\Delta\alpha_{\text{had}}(M_Z^2)$ and for $\sim (25 - 50) \%$ of the variance

Comparison with phenomenology and electroweak fit

Mainz/CLS:

$$\Delta\alpha_{\text{had}}(M_Z^2) = 0.027\,73(15)$$

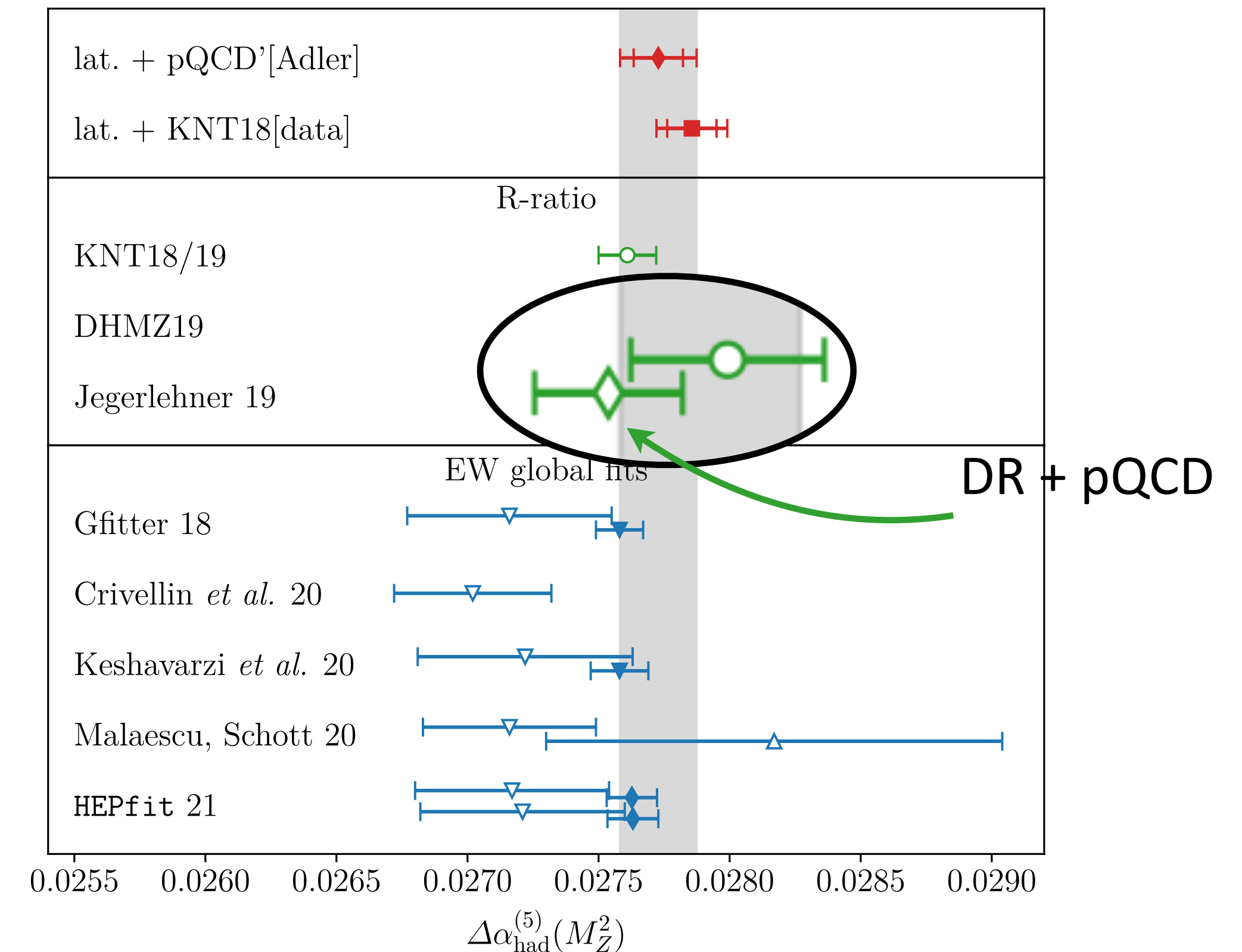
(pQCD/Adler + lattice input)



Comparison with phenomenology and electroweak fit

Mainz/CLS: $\Delta\alpha_{\text{had}}(M_Z^2) = 0.027\,73(15)$
(pQCD/Adler + lattice input)

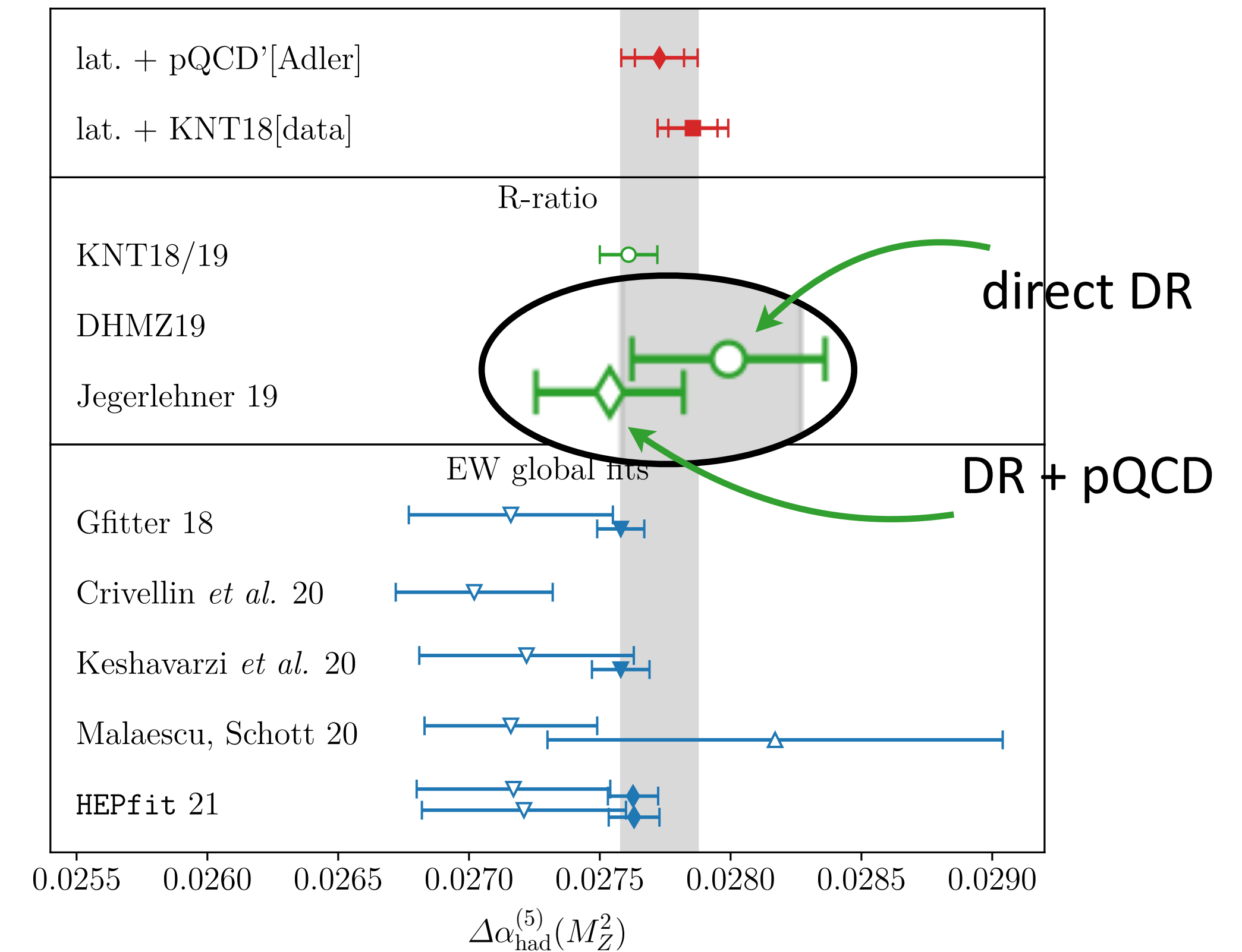
Jegerlehner 19: $\Delta\alpha_{\text{had}}(M_Z^2) = 0.027\,53(12)$
(pQCD/Adler + R -ratio input)



Comparison with phenomenology and electroweak fit

Mainz/CLS: $\Delta\alpha_{\text{had}}(M_Z^2) = 0.027\,73(15)$
(pQCD/Adler + lattice input)

Jegerlehner 19: $\Delta\alpha_{\text{had}}(M_Z^2) = 0.027\,53(12)$
(pQCD/Adler + R -ratio input)

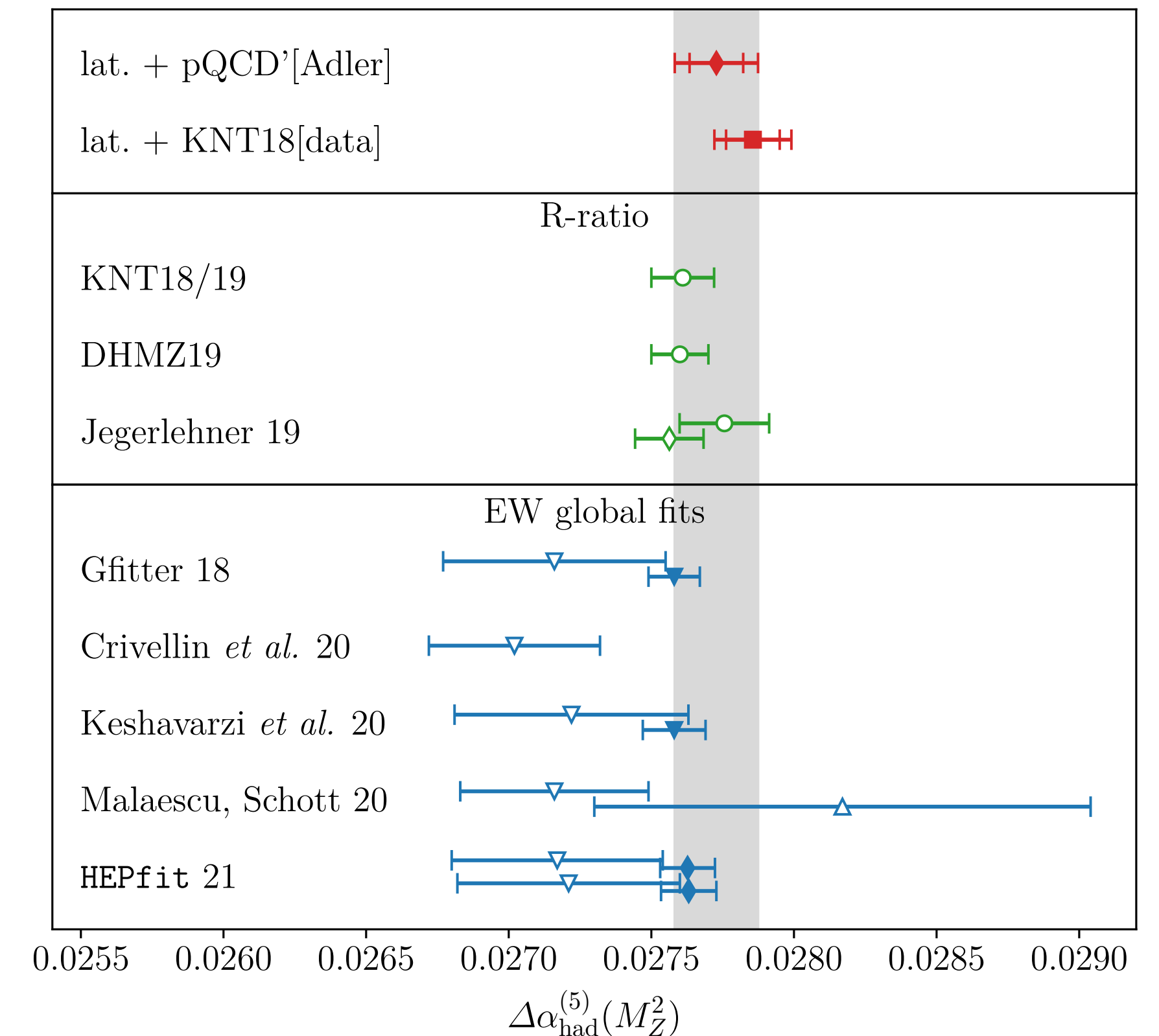


Comparison with phenomenology and electroweak fit

Mainz/CLS: $\Delta\alpha_{\text{had}}(M_Z^2) = 0.027\,73(15)$
(pQCD/Adler + lattice input)

Jegerlehner 19: $\Delta\alpha_{\text{had}}(M_Z^2) = 0.027\,53(12)$
(pQCD/Adler + R -ratio input)

- Agreement within errors at Z -pole obscures the fact that there is a tension of $\sim 3\sigma$ for $Q_0^2 \sim (3 - 7) \text{ GeV}^2$

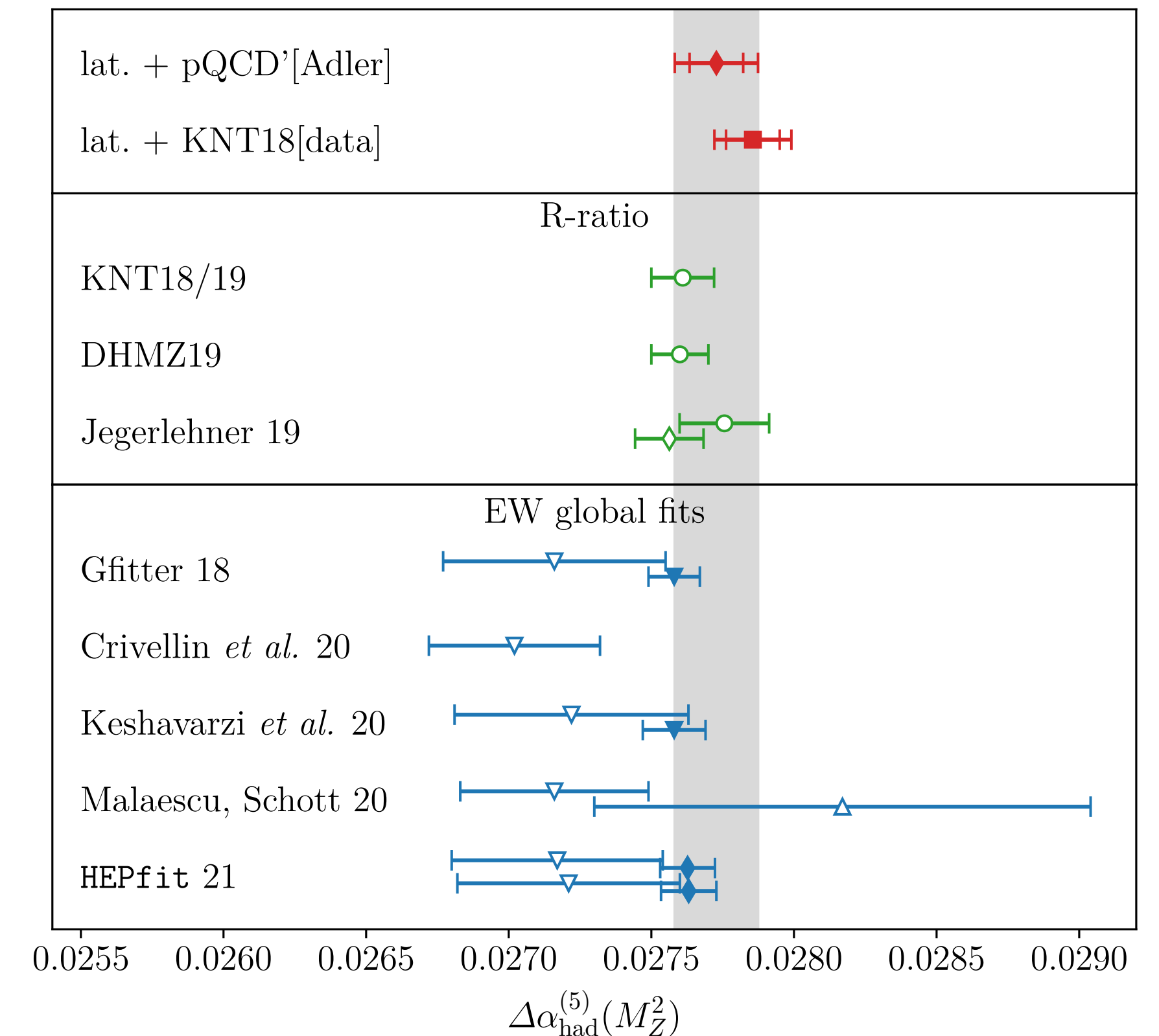


Comparison with phenomenology and electroweak fit

Mainz/CLS: $\Delta\alpha_{\text{had}}(M_Z^2) = 0.027\,73(15)$
(pQCD/Adler + lattice input)

Jegerlehner 19: $\Delta\alpha_{\text{had}}(M_Z^2) = 0.027\,53(12)$
(pQCD/Adler + R -ratio input)

- Agreement within errors at Z -pole obscures the fact that there is a tension of $\sim 3\sigma$ for $Q_0^2 \sim (3 - 7) \text{ GeV}^2$
- Running from $-Q_0^2$ to $-M_Z^2$ is correlated



Comparison with phenomenology and electroweak fit

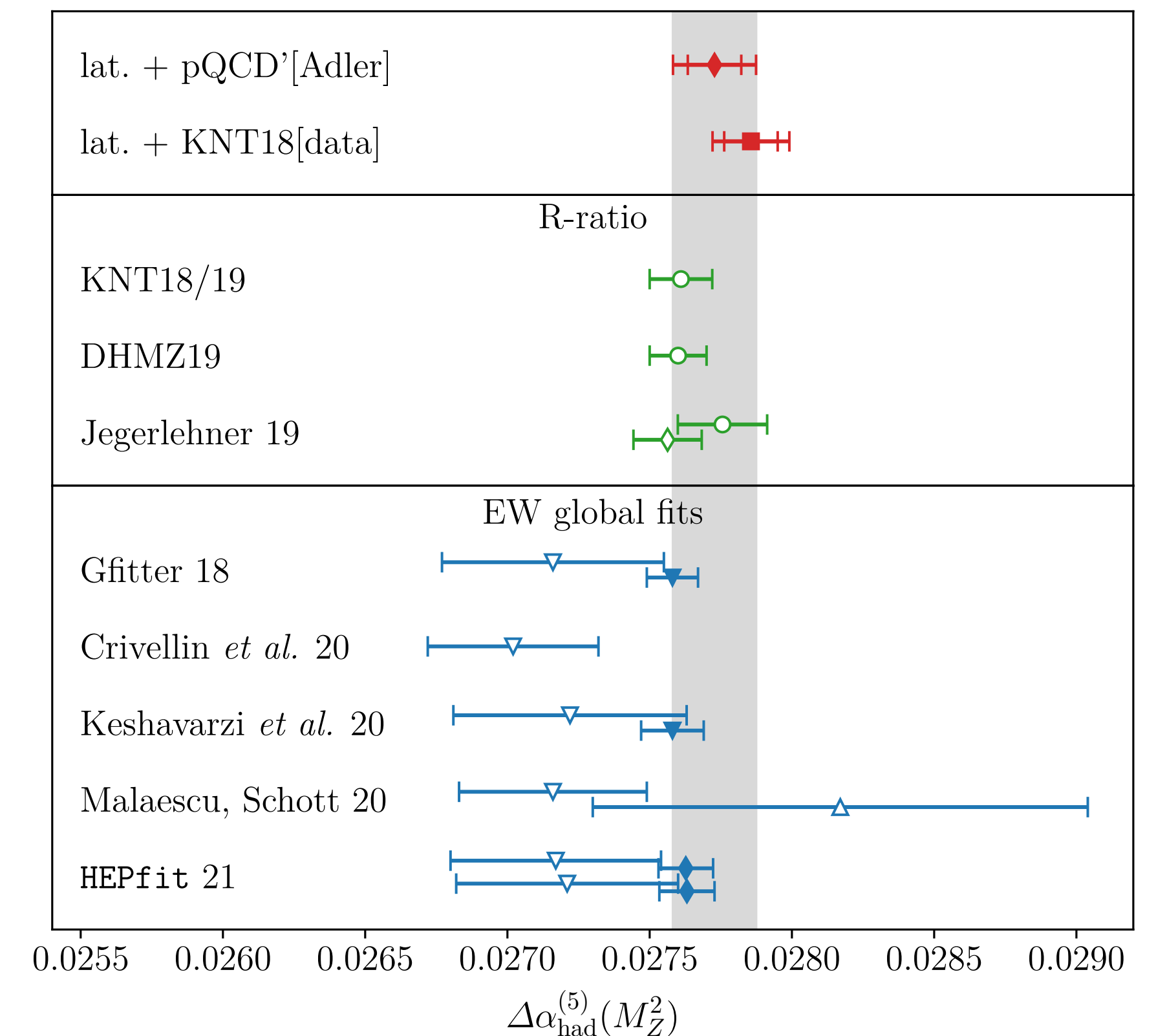
Mainz/CLS: $\Delta\alpha_{\text{had}}(M_Z^2) = 0.027\,73(15)$
(pQCD/Adler + lattice input)

Jegerlehner 19: $\Delta\alpha_{\text{had}}(M_Z^2) = 0.027\,53(12)$
(pQCD/Adler + R -ratio input)

- Agreement within errors at Z -pole obscures the fact that there is a tension of $\sim 3\sigma$ for $Q_0^2 \sim (3 - 7) \text{ GeV}^2$
- Running from $-Q_0^2$ to $-M_Z^2$ is correlated

Electroweak global fit:

∇ : $\Delta\alpha_{\text{had}}(M_Z^2)$ is free fit parameter, determined exclusively from EW precision data



Comparison with phenomenology and electroweak fit

Mainz/CLS: $\Delta\alpha_{\text{had}}(M_Z^2) = 0.027\,73(15)$
(pQCD/Adler + lattice input)

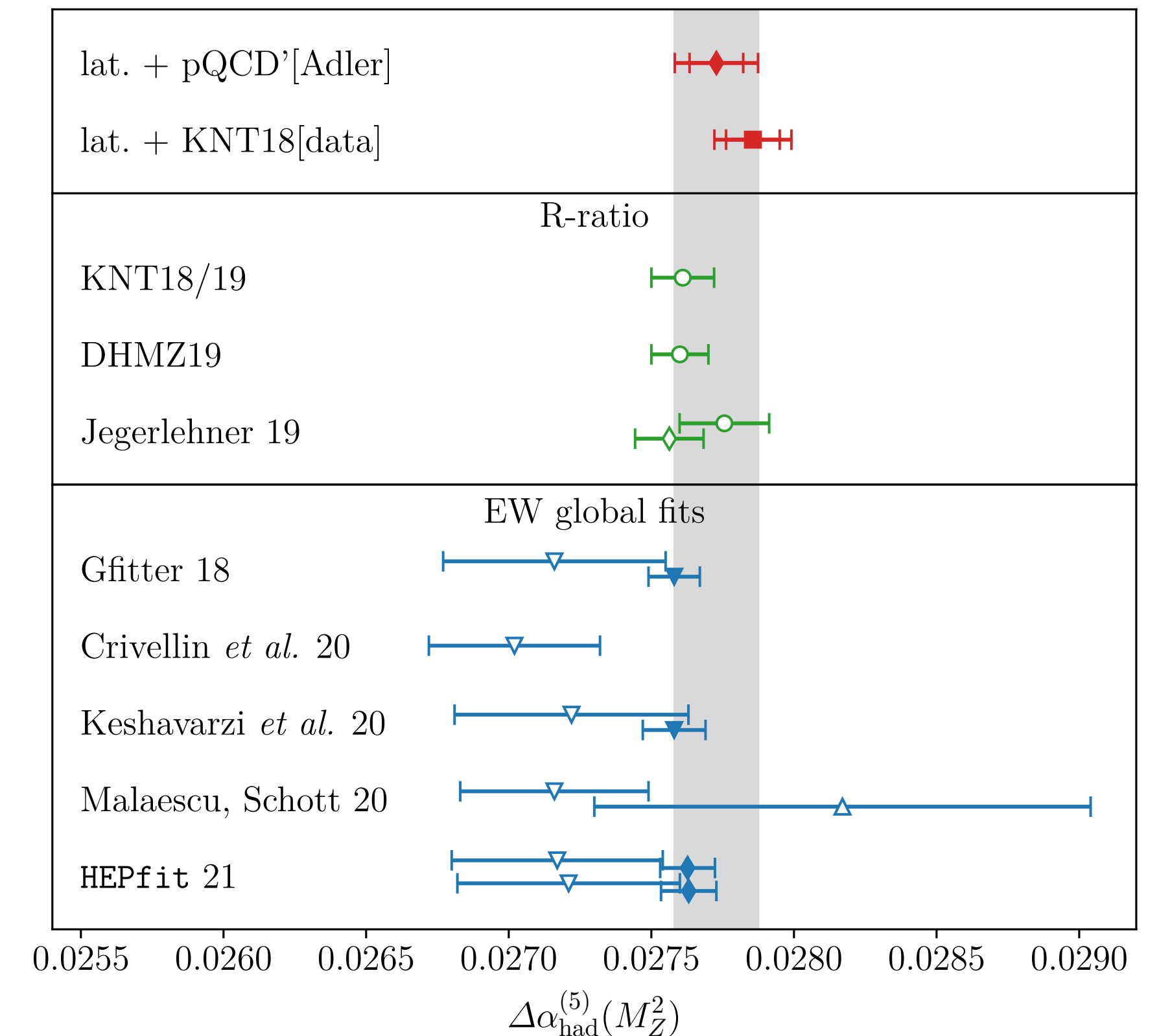
Jegerlehner 19: $\Delta\alpha_{\text{had}}(M_Z^2) = 0.027\,53(12)$
(pQCD/Adler + R -ratio input)

- Agreement within errors at Z -pole obscures the fact that there is a tension of $\sim 3\sigma$ for $Q_0^2 \sim (3 - 7) \text{ GeV}^2$
- Running from $-Q_0^2$ to $-M_Z^2$ is correlated

Electroweak global fit:

∇ : $\Delta\alpha_{\text{had}}(M_Z^2)$ is free fit parameter, determined exclusively from EW precision data

\triangle : $\Delta\alpha_{\text{had}}(M_Z^2)$ and Higgs mass M_H free fit parameters



Comparison with phenomenology and electroweak fit

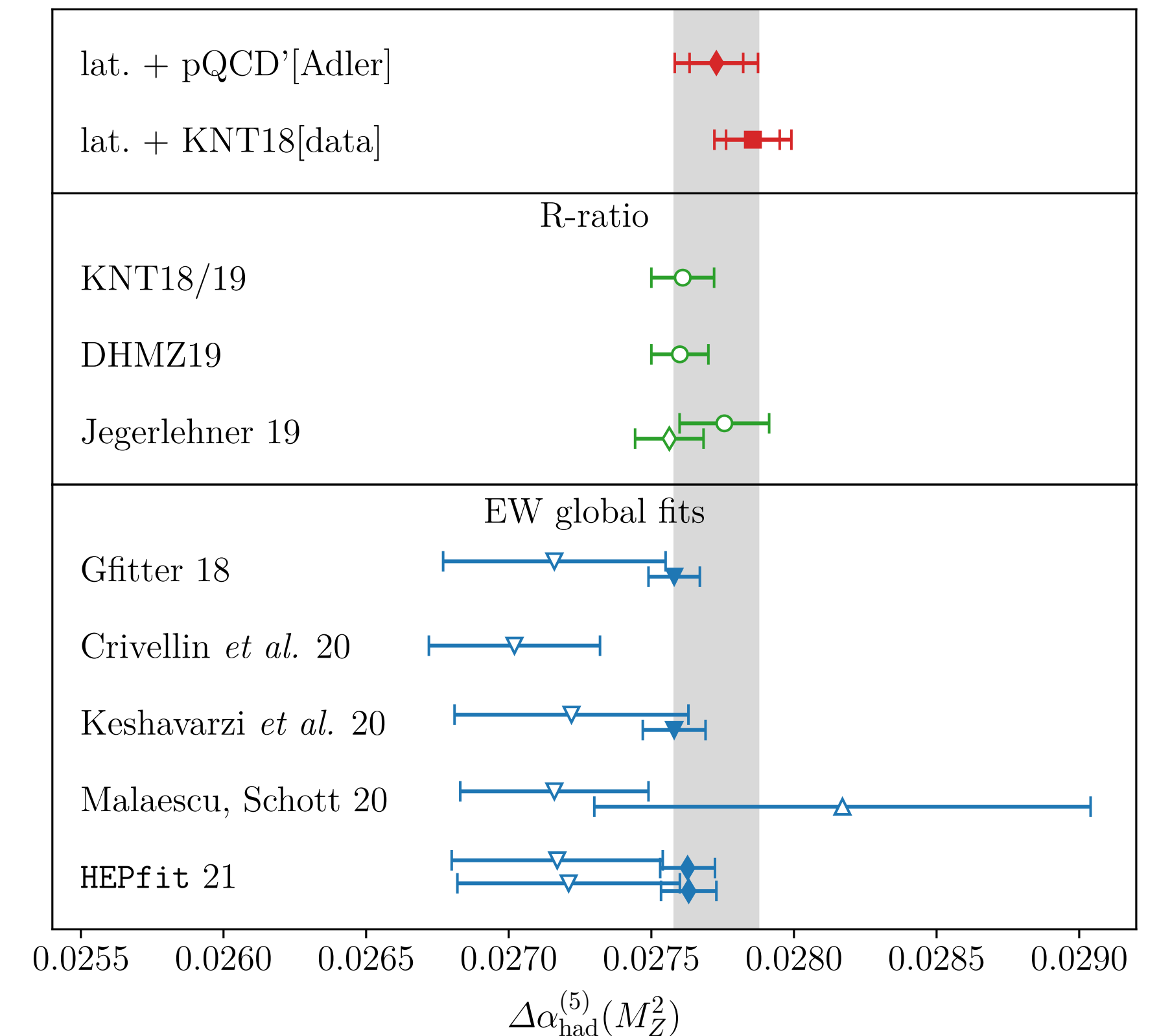
Mainz/CLS: $\Delta\alpha_{\text{had}}(M_Z^2) = 0.027\,73(15)$
(pQCD/Adler + lattice input)

Jegerlehner 19: $\Delta\alpha_{\text{had}}(M_Z^2) = 0.027\,53(12)$
(pQCD/Adler + R -ratio input)

- Agreement within errors at Z -pole obscures the fact that there is a tension of $\sim 3\sigma$ for $Q_0^2 \sim (3 - 7) \text{ GeV}^2$
- Running from $-Q_0^2$ to $-M_Z^2$ is correlated

Electroweak global fit:

- ∇ : $\Delta\alpha_{\text{had}}(M_Z^2)$ is free fit parameter, determined exclusively from EW precision data
- \triangle : $\Delta\alpha_{\text{had}}(M_Z^2)$ and Higgs mass M_H free fit parameters
- \blacktriangledown \blacklozenge : same as ∇ but using priors for $\Delta\alpha_{\text{had}}(M_Z^2)$ centred about R -ratio estimate / BMWc

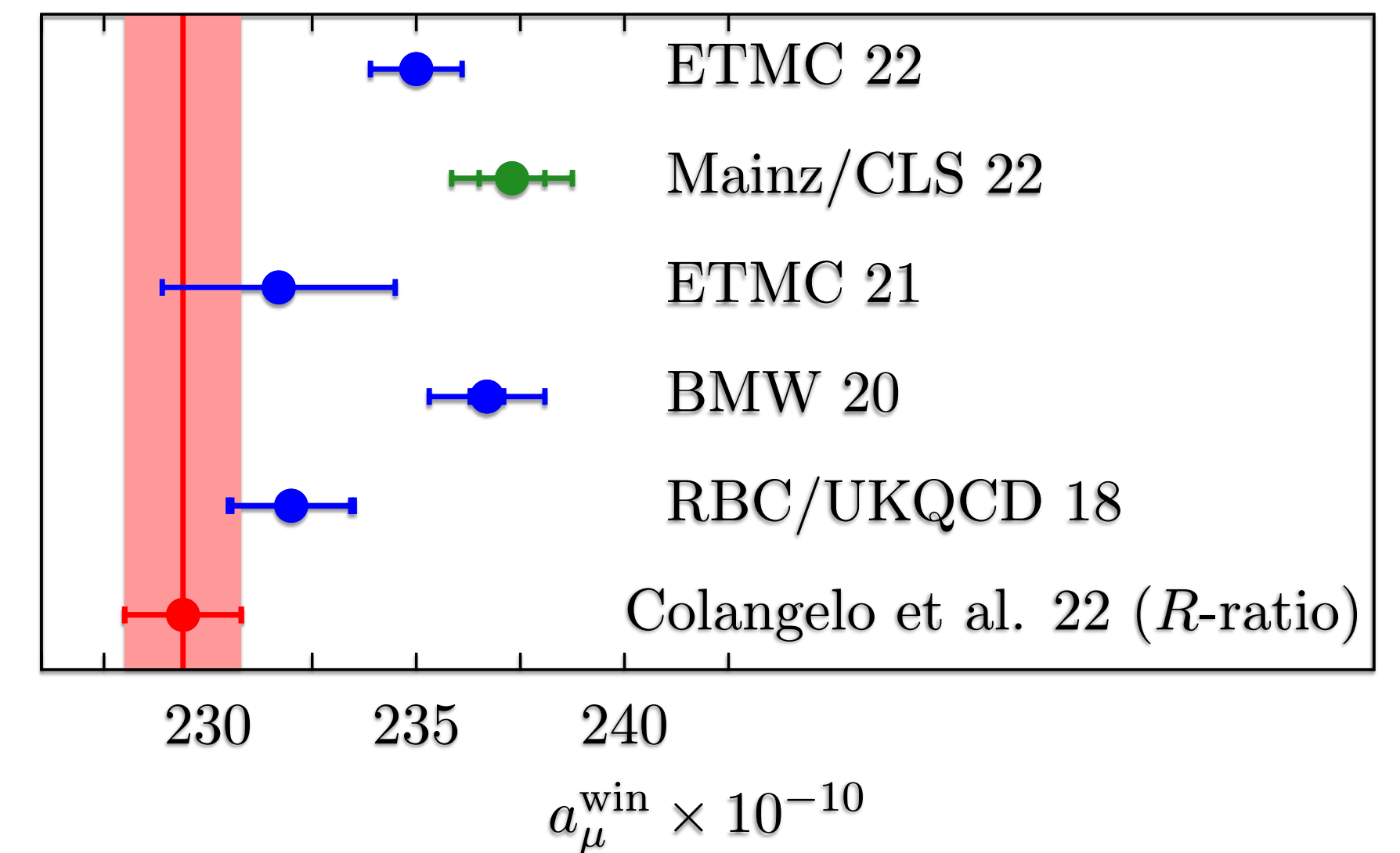


Summary and Discussion

- Lattice+pQCD/Adler estimate for $\Delta\alpha_{\text{had}}(M_Z^2)$ broadly agrees with global electroweak fit
→ no contradiction with the Standard Model
- Lattice estimates for $\Delta\alpha_{\text{had}}(-Q_0^2)$ larger than counterparts derived from data-driven approach
→ tension of $\sim 3\sigma$ for $Q_0^2 \approx 5 \text{ GeV}^2$

Summary and Discussion

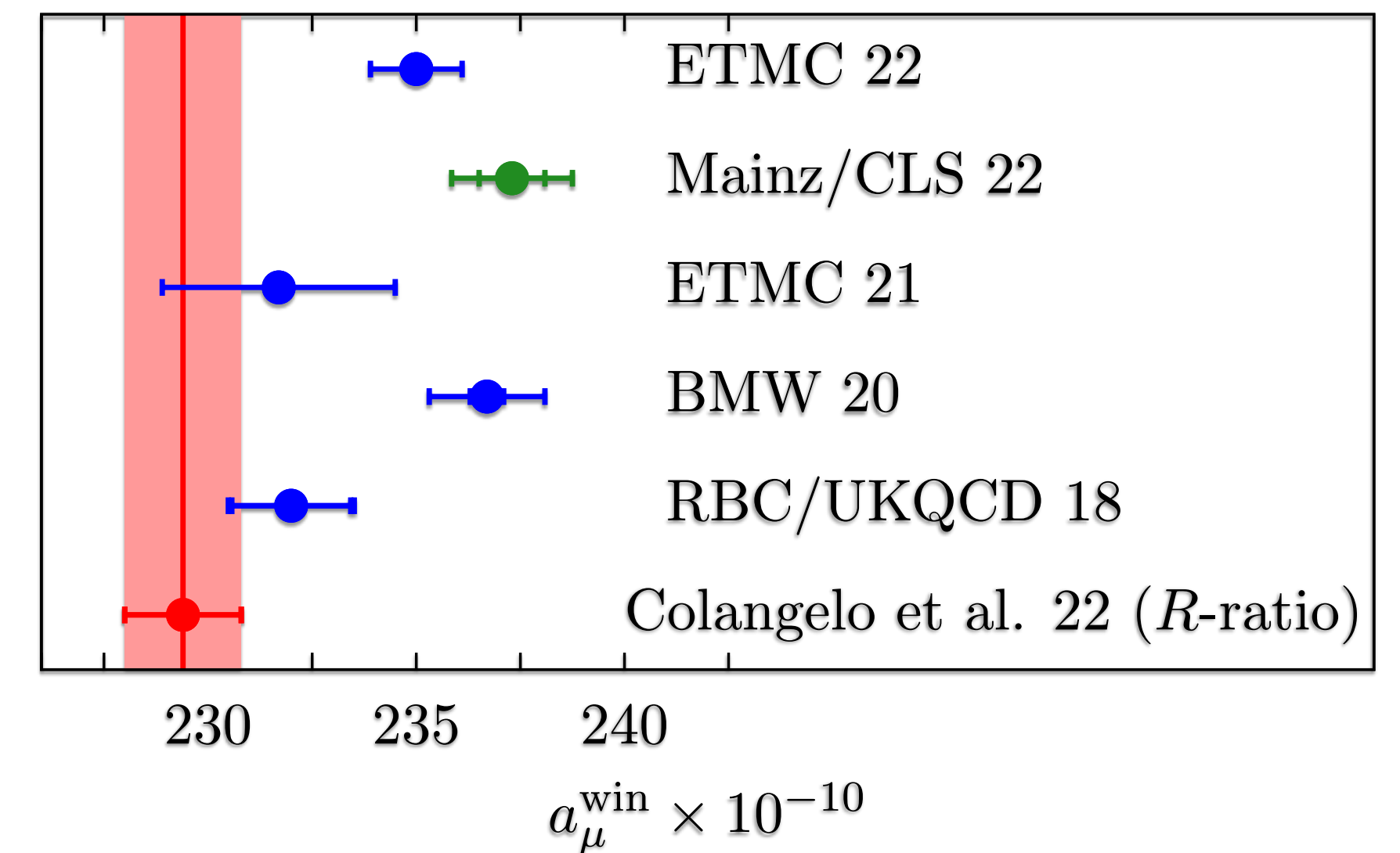
- Lattice+pQCD/Adler estimate for $\Delta\alpha_{\text{had}}(M_Z^2)$ broadly agrees with global electroweak fit
→ no contradiction with the Standard Model
- Lattice estimates for $\Delta\alpha_{\text{had}}(-Q_0^2)$ larger than counterparts derived from data-driven approach
→ tension of $\sim 3\sigma$ for $Q_0^2 \approx 5 \text{ GeV}^2$
- Observation consistent with larger lattice estimates for HVP contribution to a_μ , cf. window observable:



[Cè et al., arXiv:2206.06582]

Summary and Discussion

- Lattice+pQCD/Adler estimate for $\Delta\alpha_{\text{had}}(M_Z^2)$ broadly agrees with global electroweak fit
→ no contradiction with the Standard Model
- Lattice estimates for $\Delta\alpha_{\text{had}}(-Q_0^2)$ larger than counterparts derived from data-driven approach
→ tension of $\sim 3\sigma$ for $Q_0^2 \approx 5 \text{ GeV}^2$
- Observation consistent with larger lattice estimates for HVP contribution to a_μ , cf. window observable:
⇒ Standard Model can accommodate a larger value for a_μ without contradicting electroweak precision data



[Cè et al., arXiv:2206.06582]