

Implications of gradient flow on the static force

The 39th International Symposium on Lattice Field Theory (Lattice 2022) Julian Frederic Mayer-Steudte

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1 Motivation

2 Setup

3 Lattice results

4 Conclusion





Interested in the QCD static energy of a quark-antiquark pair *E*(*r*)
 Given by the Wilson loop

$$E(r) = -\lim_{T \to \infty} \frac{\ln \langle \operatorname{Tr}(W_{r \times T}) \rangle}{T}, \qquad \qquad W_{r \times T} = P\left\{ \exp\left(i \oint_{r \times T} dz_{\mu} g A_{\mu}\right) \right\}$$



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can be used for precise $\alpha_S\text{-}\mathrm{running}$ extraction by comparing PT and lattice





Perturbative form of E(r):

$$E(r) = \Lambda_S - \frac{C_F \alpha_S}{r} \left(1 + \# \alpha_S + \# \alpha_S^2 + \# \alpha_S^3 + \# \alpha_S^3 \ln \alpha_S + \dots \right)$$

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take a derivative of E(r) for the force $F(r) = \partial_r E(r)$

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Direct measurement of F(r):

(A. Vairo Mod. Phys. Lett. A 31 (2016) & EPJ Web Conf. 126 (2016), Brambilla et.al.PRD63 (2001))

$$\begin{split} F(r) &= -\lim_{T \to \infty} \frac{i}{\langle \operatorname{Tr}(W_{r \times T}) \rangle} \left\langle \operatorname{Tr} \left(P \left\{ \exp\left(i \oint_{r \times T} dz_{\mu} g A_{\mu}\right) \hat{\mathbf{r}} \cdot g \mathbf{E}(\mathbf{r}, t^{*}) \right\} \right) \right\rangle \\ &= \frac{\langle \operatorname{Tr}\{PW_{r \times T} g E_{j}(r, t^{*})\} \rangle}{\langle \operatorname{Tr}\{PW_{r \times T}\} \rangle} \end{split}$$



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- Chromoelectric field E inserted into Wilson loop
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• On the lattice: modifying Wilson loop with a discretized *E*-field insertion





Clover discretization of *E*:

$$E_{i} = \frac{1}{2iga^{2}} \left(\Pi_{i0} - \Pi_{i0}^{\dagger} \right) \qquad \qquad \Pi_{\mu\nu} = \frac{1}{4} \left(P_{\mu,\nu} + P_{\nu,-\mu} + P_{-\mu,-\nu} + P_{-\nu,\mu} \right)$$

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- The self energy contribution of E converges slowly to continuum (See e.g. Lepage et.al.PRD48 (1993), G. Bali Phys. Rept. 343 (2001), and many others...) \rightarrow need renormalization Z_E



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We use **Gradient flow** for targeting the renormalization and the signal to noise ratio problems, new scale: flowtime τ_F , flowradius $\sqrt{8\tau_F}$, flowtime ratio τ_F/r^2





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which introduces an additional factor Z_E

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Setup: Continuum results



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- One-loop calculation of the flowed force is known: (Hee Sok Chung et. al. JHEP01(2022)184 (2022))
 - \Box Here: scale $\mu = 1/r$
 - Slight difference to $\mu = 1/\sqrt{r^2 + 8\tau_F}$
 - □ The one-loop behavior should dominate the small τ_F regime





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 Small *\(\tau_F\)* exansion:

$$r^{2}F(r;\tau_{F}) \approx r^{2}F(r;\tau_{F}=0) + \frac{\alpha_{S}^{2}C_{F}}{4\pi} \underbrace{\left[-12\beta_{0} - 6C_{A}c_{L}\right]}_{8n_{f}} \frac{\tau_{F}}{r^{2}} \qquad c_{L} = -\frac{22}{3}$$

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At small flowtime the force is constant in pure gauge $(n_f = 0)$

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Lattice results: setup and parameters



Parameters:

N_S	N_T	eta	a [fm]	$N_{\rm conf}$	Label
20	40	6.284	0.060	6000	L20
26	52	6.481	0.046	6000	L26
30	60	6.594	0.040	6000	L30
40	80	6.816	0.030	2700	L40

Pure gauge configuration produced with overrelaxation and heatbath

Scale setting with $\ln(a/r_0) = -1.6804 - 1.7331(\beta - 6) + 0.7849(\beta - 6)^2 - 0.4428(\beta - 6)^3$ (1S. Necco & R. Sommer. Nucl. Phys. B622 (2002))

Gradient flow with fixed and adaptive solver, with Symanzik action

Lattice results: Crucial steps



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Extraction of the $\lim_{T\to\infty}$ - limit

(William I. Jay, Ethan T. Neil Phys. Rev. D 103, 114502 (2021))

at each fixed separation r and fixed flowtime au_F or flowtime ratio au_F/r^2

Lattice results: Crucial steps

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Working out the continuum limit







Nonperturbative determination of Z_E :

$$Z_E(r) = \frac{\partial_r E(r)}{F(r)}$$



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- $\blacksquare \ Z_E \to 1 \text{ for flowradius } \sqrt{8\tau_F} > a$





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$$Z_E \rightarrow 1$$
 for flowradius $\sqrt{8\tau_F} > a$



Gradient flow reduces discretization effects of field insertions for $\sqrt{8\tau_F} > a$

Lattice results: Continuum results at large *r*



Lattice results: Continuum results at large r







Lattice results: Continuum results at large r





Cornell fit
$$r^2 F(r) = A + \sigma r^2$$

 $\sigma = 5.18 \dots 5.23 \text{ fm}^{-2}$, literature: 5.5 fm⁻²
 $A = 0.2853 \dots 0.2954$

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0.45 0.50

0.55 0.60

0.40

r/[fm]





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Results:

F2I: $\Lambda_0 = 0.2731 \text{ GeV}$ F3ILus: $\Lambda_0 = 0.2666 \text{ GeV}$ F1I: $\Lambda_0 = 0.3033 \text{ GeV}$ A proper error estimation is still

pending...





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Constant fit is too naive for small r





Results (in GeV):



0.44

0 42

--- F2I fit

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 $\tau_{e}/r^{2} = 0.0212$

 $\tau_F/r^2 = 0.0242$

 $\tau_{e}/r^{2} = 0.0288$

пΠ



Results (in GeV):



0.44

0 42

--- F2I fit

--- F2I fit

 $\tau_c/r^2 = 0.0212$

 $\tau_F/r^2 = 0.0242$

 $\tau_c/r^2 = 0.0288$

 Λ_0 varies little, but is within the FLAG error

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- Summary ond observations:
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- Summary ond observations:
 - Gradient flow reduces effectively discretization effects
 - Gradient flow improves qualitatively the signal to noise ratio
 - Good preparation for future applications in NREFTs
- For the future:
 - \Box Proper Λ extraction
 - Go to finer lattices
 - Other operators with field insertions
 - Extend to dynamical fermions



Thank you for your attention!

r_0 scale flow dependence





 $\begin{bmatrix} 13 \\ 12 \\ 10 \\ 9 \\ 8 \\ 0.00 \\ 0.01 \\ 0.02 \\ 0.03 \\ 0.04 \\ 0.04 \\ 0.03 \\ 0.04 \\ 0.04 \\ 0.04 \\ 0.04 \\ 0.04 \\ 0.04 \\ 0.03 \\ 0.04 \\ 0.0$

Figure 1 The r_0 -scale in lattice units for the L20 lattice. The horizontal line corresponds to the expected value.

Figure 2 The r_0 -scale in lattice units for the L30 lattice. The horizontal line corresponds to the expected value.

Tree level behavior





1.00 0.95 eLatt 'F(r, τ_F)^{Tr} 0.90 0.85 F(r, τ_F)^{Tre} 0.80 × $= 0.1028, \tau_F/r^2 = 0.0327$ 0.75 r = 0.2159, $\tau_c/r^2 = 0.0218$ L40 130 L26 L20 lattice spacing a

Figure 3 The ratio of the tree level forces in continuum and on the lattice at zero flowtime.

Figure 4 The ratio of the tree level forces in continuum and on the lattice at finite flowtime.