

# Hadronic vacuum polarization from step scaling

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# Adler function from low to high energies

# Hadronic vacuum polarization

- The hadronic contributions to  $\alpha_{em}(M_Z^2)$  are given by the HVP  $\Pi(M_Z^2)$

$$\Pi(Q^2)(Q^2\delta_{\mu\nu} - Q_\mu Q_\nu) = \Pi_{\mu\nu}(Q^2) = \int d^4x e^{iQ \cdot x} \langle j_\mu(x) j_\nu(0) \rangle$$

- Interested in the difference

$$\Pi(Q_n^2) - \Pi(Q_0^2)$$

$$Q_0 \sim 1 \text{ GeV}, \quad Q_n \sim 100 \text{ GeV}$$

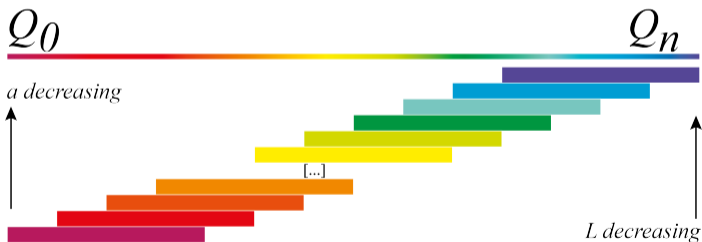
- Computationally unfeasible
- Use a strategy similar to step scaling [Lüscher, Weisz, Wolff, Sommer ...]
- Divide the range into many steps  $Q_0, Q_1 = 2Q_0, \dots, Q_n = 2^n Q_0$

# Strategy

- Discrete Adler function:

$$\Delta(Q^2) = \Pi(4Q^2) - \Pi(Q^2)$$

$$\Pi(Q_n^2) - \Pi(Q_0^2) = \Delta(Q_0^2) + \Delta(Q_1^2) + \dots + \Delta(Q_n^2)$$



- Compute each step in decreasing volumes
- Higher energies are less sensitive to FV effects

# Test the strategy in $QED_2$

- Partition function:

$$Z = \int [dU] \exp(-S_g[U]) \det M(m, U)$$

$$S_g = \beta \sum_{x \in \Lambda} \operatorname{Re} U_{x,0} U_{x+\hat{0},1} U_{x+\hat{1},0}^* U_{x,1}^*$$

$$M = \sum_{\mu=0,1} \eta_{x,\mu} \nabla_{\mu} + m$$

- Line of constant physics (LCP) defined via

$$\beta m^2 = 0.8$$

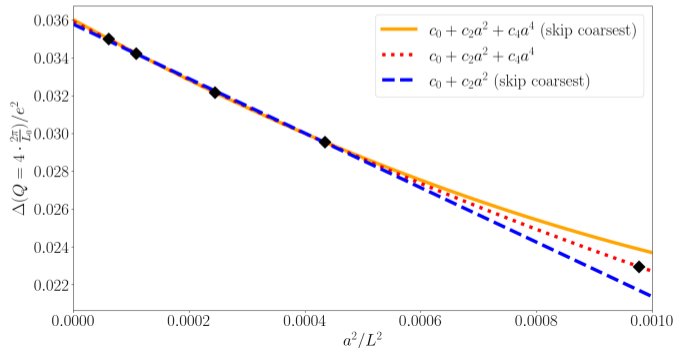
- Volumes given by

$$mL_0 = 16, \quad L_0/L_n = 2^n$$

- Metropolis + top. update, reweighting

# Continuum Extrapolation

- Cont. extrapolation from  $L/a = 32, 48, 64, 96, 128$  ( $\beta \propto a^{-2}, m \propto a$ )
- Sweet spot for lattice artifacts and FV effects,  $Q = 4 \cdot \frac{2\pi}{L}$
- No logarithmically enhanced lattice artefacts  $a^2 \log(a^2)$
- Systematic error from different cont. extrapolations

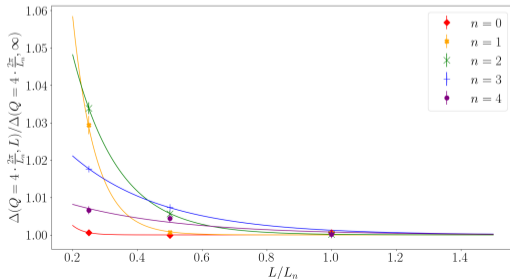


# Finite volume effects

- Same momentum in three different volumes

$$Q_n = 4 \cdot \frac{2\pi}{L_n} = 2 \cdot \frac{2\pi}{L_{n+1}} = 1 \cdot \frac{2\pi}{L_{n+2}}$$

- FV effects  $\propto e^{-m\pi L}$
- $Q = 4 \cdot \frac{2\pi}{L} \Rightarrow$  FV effects below 0.05%



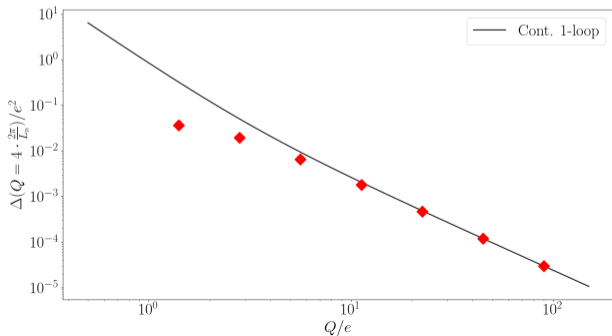


# Discrete Adler function $QED_2$

- 1-loop PT [Adams'98]:

$$\Pi(q^2)_{PT}/e^2 = \frac{1}{\pi} \frac{1}{q^2} \left( 1 + \frac{2m^2}{q^2} \frac{1}{R} \log \frac{1+R}{1-R} \right) \quad R = \sqrt{1 - \frac{4m^2}{Q^2}}$$

- For large  $Q^2$   
 $\Delta(Q^2) \propto \frac{1}{Q^2}$
- Deviation from higher order effects (checked by performing simulations with  $e \rightarrow 0$ )



# Discrete Adler function $QED_2$

$Q/e$	$\Delta(Q)/e^2$	stat.[ $10^{-6}$ ]	cont. extrap.[ $10^{-6}$ ]	FV [ $10^{-6}$ ]	total[ $10^{-6}$ ]
1.405	0.036009	31	16	18	37
2.810	0.019362	17	7	10	13
5.620	0.006515	4	< 1	10	10
11.240	0.001807	< 1	1	4	4
22.480	0.000467	< 1	< 1	< 1	< 1
44.960	0.000118	< 1	< 1	< 1	< 1
89.920	0.000029	< 1	< 1	< 1	< 1

- Adding up gives final result:

$$\Pi(2^{14}Q_0^2) - \Pi(Q_0^2) = 0.064308[41, 0.06\%]$$

- Here we used analytic LCP (known in  $QED_2$ )

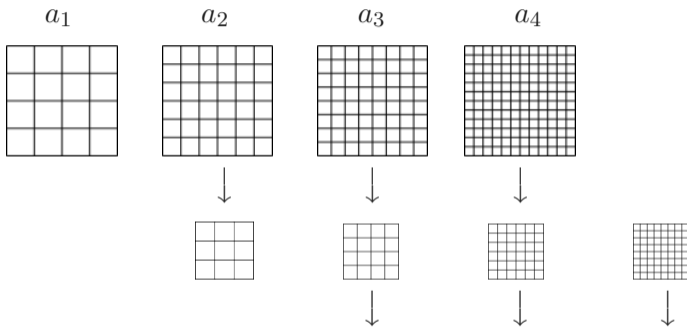


**Lines of constant  
physics  $(\beta(a), m(\beta))$**

# Determination of the LCP

- General strategy to determine  $m(\beta)$  and  $\beta(a)$  twofold:
  - (A) Determine  $\beta(a)$  to ensure that we will have in all of our steps the same physical volume
  - (B) Use  $\Delta$  (Adler function) to fix  $m$  as a function of  $\beta$
  - (C) Please note that this should be done simultaneously

# General strategy



1. Assume we know  $\beta(a_1) = \beta_1$ ,  $\beta(a_2) = \beta_2$ ,  $\beta(a_3) = \beta_3$
2. Determine  $\beta(a_4)$  with cont. extrapolation of an Observable  $\langle O \rangle$  sensitive to  $\beta$
3. Halve phys. size  $L_{2,3,4}/2$  and switch to  $\langle O \rangle$  with smaller FV effects
4. Shift  $a_2 \rightarrow a_1$ ,  $a_3 \rightarrow a_2$ ,  $a_4 \rightarrow a_3$  and go to 1.

# Choosing the observables

- Observable sensitive to  $\beta$ :  $t_0^{(n)}$ -scale from gradient flow

$$\frac{d}{dt} U_t(x, \mu) = - \left[ \partial_{x, \mu} \frac{1}{\beta} S_g(U_t) \right] U_t(x, \mu)$$

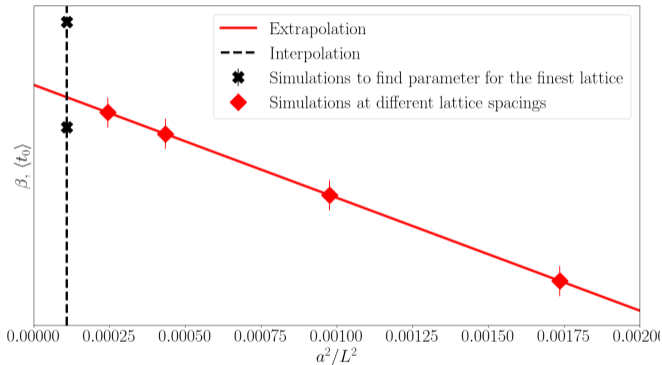
$$t_0^{(n)} \cdot E(U_{t_0^{(n)}}) / (L/a)^2 = 0.375 \cdot \left( \frac{1}{2} \right)^{2n}$$

- Observable sensitive to  $m$ :

$$\Delta^{(n)} = \Delta(Q_n^2, 2^n m)$$

Mass doubling favors mass dependence of the discrete Adler function

# Determination of the LCP



- Intersection point gives the LCP
- Generalization to  $D$  parameters possible ( $D$  observables needed)

# Estimation of uncertainties

- **Systematic uncertainties** from continuum extrapolation
- Many extrapolations weighted with AIC:

$$\exp\left\{-\frac{1}{2}(\chi^2 + 2n_f - n_p)\right\}$$

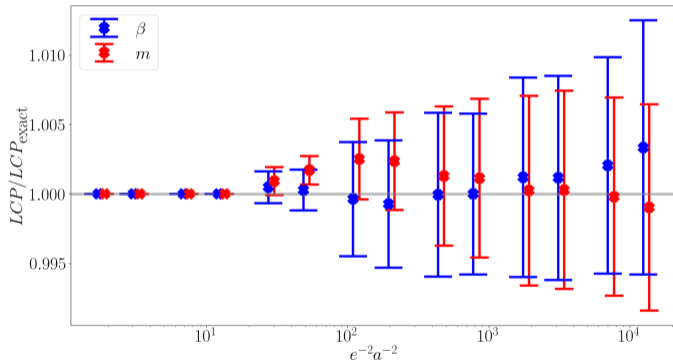
- Recursive structure  $\Rightarrow$  Dependence on the previous estimations
  - 36 different fits  $\Rightarrow 36^6 \approx 10^{10}$  fits after 6 steps
  - We reduce this number by choosing a handful of random representants after every step
- **Statistical uncertainties** from a Jackknife analysis
- Whole procedure repeated on the  $N_J = 48$  Jackknife bins





# Results

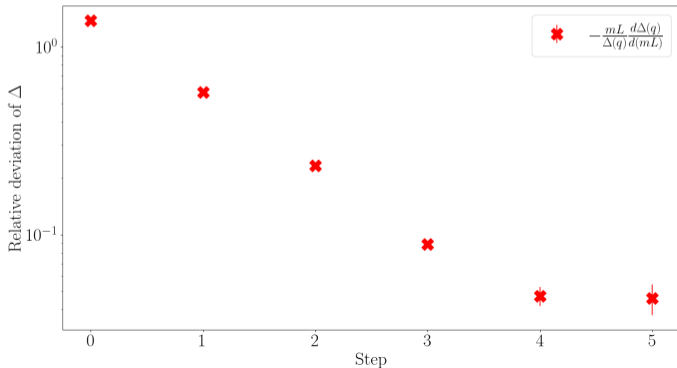
# Results for the LCP in $QED_2$



- Good agreement with exact results
- Uncertainties  $\lesssim 1\%$

# How LCP uncertainty effects Adler function?

- Uncertainty from  $\delta\beta(a)$  estimated from  $\frac{1}{Q^2}$  dependence
- Uncertainty from  $\delta m(a)$  estimated from measuring mass dependence of  $\Delta(Q^2)$



# Results and Conclusion

- The final results of the total study:

$$\Pi(2^{14}Q_0^2) - \Pi(Q_0^2) = 0.064308[186, 0.3\%]$$

- Error budget:

Stat. error:	$32 \cdot 10^{-6}$	0.05%
Finite volume error:	$14 \cdot 10^{-6}$	0.02%
Cont. Extrap. error:	$19 \cdot 10^{-6}$	0.03%
$\beta$ estimation error:	$155 \cdot 10^{-6}$	0.24%
$m$ estimation error:	$95 \cdot 10^{-6}$	0.14%

- Seems to work in  $QED_2$
- Does the method work in QCD to calculate physical observables?

# Thank you for your attention!

I am happy to answer any questions you have!



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