

國立陽明交通大學

NATIONAL YANG MING CHIAO TUNG UNIVERSITY

PROGRESS IN CALCULATION OF THE FOURTH MELLIN MOMENT OF THE PION LCDA USING THE HOPE METHOD



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with

Calculation of four-point functions in QCD
Will Detmold, Anthony Grebe, Issaku Kanamori, David Lin, Yong Zhao

TAKEAWAYS

- ▶ Light cone distribution amplitude (LCDA) important in factorization theorem for pion electromagnetic form factor.
- ▶ Heavy-quark Operator Product Expansion (HOPE) → moments of LCDA.
- ▶ Preliminary results for $\langle \xi^4 \rangle$ at range of finite lattice spacings, twist-suppressions.

FACTORIZATION FOR $F_\pi(Q^2)$: 1979

- ▶ $\phi_M(x, \mu^2)$: Light Cone Distribution Amplitude (LCDA): Probability amplitude
- ▶ Factorization theorem:

$$F_\pi(Q^2) \underset{\text{large } Q^2}{=} \int_0^1 dx dy \phi_{\overline{M}}(y, Q^2) T_H(x, y, Q^2) \phi_M(x, Q^2)$$

$$\underset{\text{large } Q^2}{=} \int_0^1 dx dy \left(- \begin{array}{c} \text{---} \\ \textcirclearrowleft \end{array} \times \left[\begin{array}{c} \text{---} \\ \textcirclearrowleft \end{array} + \begin{array}{c} \text{---} \\ \textcirclearrowright \end{array} \right] \times \begin{array}{c} \text{---} \\ \textcirclearrowright \end{array} \right)$$

$$\underset{\text{large } Q^2}{=} \frac{16\pi \alpha_S(Q^2)}{Q^2} f_\pi^2 \omega_\phi^2(Q^2) \rightarrow \frac{16\pi \alpha_S(Q^2)}{Q^2} f_\pi^2$$

$$\omega_\phi(Q^2) = \frac{1}{3} \int_0^1 dx \frac{\phi(x, Q^2)}{x}$$

$$F_\pi(Q^2)$$

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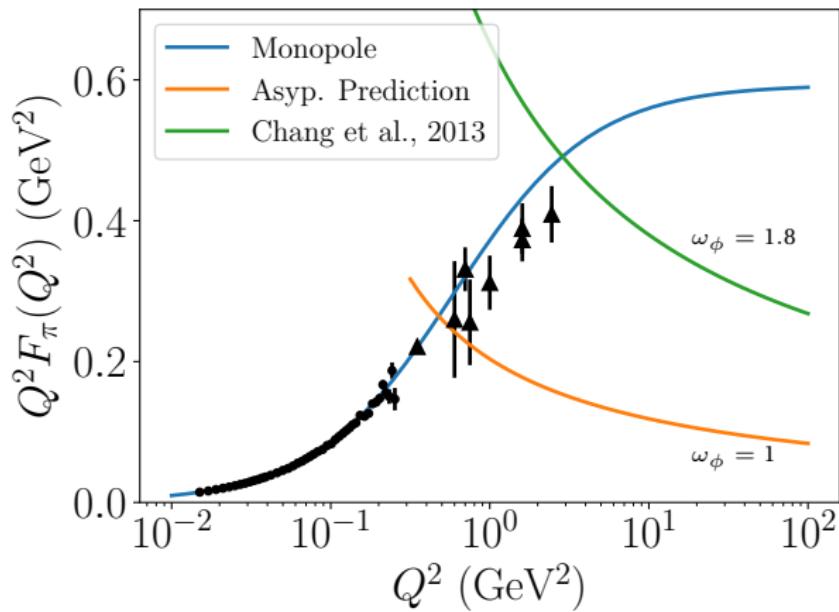
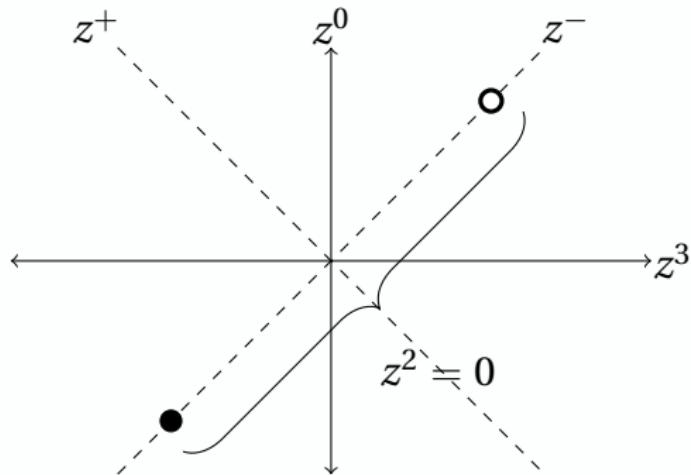


Figure 1: Data from Amendolia et al. (1986), Huber et al. (2008)
Calculation of fourth Mellin moment of pion LCDA: August 12, 2022.

CALCULATING THE PION LCDA

- ▶ Only *ab-initio* method to calculate non-perturbative QCD: Lattice QCD.
- ▶ Problem: LCDA defined as

$$\langle \Omega | \bar{\psi}(z_-) \gamma_\mu \gamma_5 W[z_-, -z_-] \psi(-z_-) | \pi(\mathbf{p}) \rangle = i p_\mu f_\pi \int_{-1}^1 d\xi e^{i\xi p^+ \cdot z_-} \phi_\pi(\xi, \mu^2)$$



HEAVY-QUARK OPE

DETMOLD & LIN, 2006



- ▶ Consider matrix element

$$V^{\mu\nu}(p, q) = \int d^4z e^{iq\cdot z} \langle \Omega | T\{J_\Psi^\mu(z/2)J_\Psi^\nu(-z/2)\} |\pi(\mathbf{p}) \rangle$$
$$J_\Psi^\mu = \bar{\Psi}(x)\Gamma^\mu\Psi(x) + \bar{\psi}(x)\Gamma^\mu\Psi(x)$$

- ▶ Perform operator product expansion:

$$\tilde{Q}^2 = -q^2 + m_\Psi^2 \quad \text{large scale}$$

$$\tilde{\omega} = \frac{1}{\tilde{x}} = \frac{2p \cdot q}{\tilde{Q}^2} \quad \text{expansion parameter}$$

$$V_{\text{HOPE}}^{\mu\nu}(p, q) = K[1 + \tilde{\omega}^2 \langle \xi^2 \rangle + \tilde{\omega}^4 \langle \xi^4 \rangle + \dots] + \underbrace{\mathcal{O}(\alpha_S)}_{\text{Perturbative corrections}} + \underbrace{\mathcal{O}(1/Q^3)}_{\text{Higher twist}}$$

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OVERVIEW OF CALCULATION

- ▶ Pion LCDA $\phi_\pi(\xi, \mu^2)$ important in description of $F_\pi(Q^2)$
- ▶ Long-range sensitive: non-perturbative.
- ▶ HOPE method allows for determination of moments:

$$\langle \xi^n \rangle (\mu^2) = \int_{-1}^1 d\xi \xi^n \phi(\xi, \mu^2)$$

- ▶ HOPE Method:

$$V_{\text{LQCD}}^{\mu\nu}(p, q; a) = \int d^4z e^{iq \cdot z} \langle \Omega | T\{J_\Psi^\mu(z/2) J_\Psi^\nu(-z/2)\} | \pi(\mathbf{p}) \rangle$$

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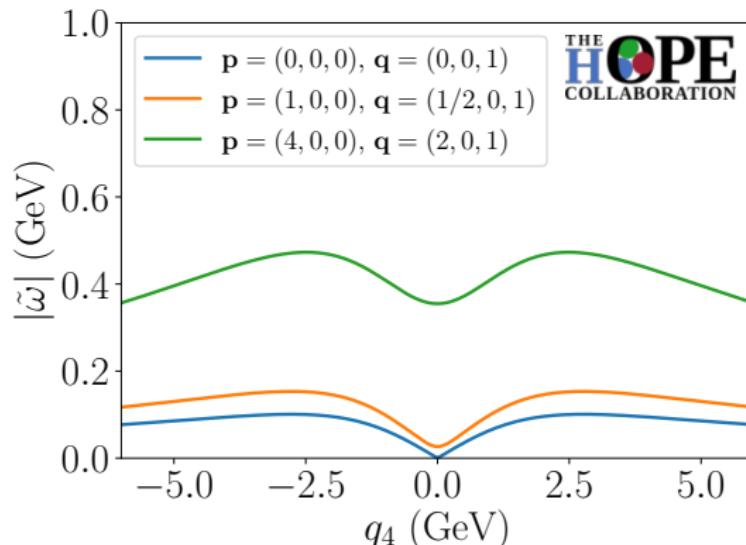
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OPTIMIZING KINEMATICS

$$V^{\mu\nu}(p, q) = K[1 + \tilde{\omega}^2 \langle \xi^2 \rangle + \tilde{\omega}^4 \langle \xi^4 \rangle + \dots], \quad \tilde{\omega} = \frac{2\mathbf{p} \cdot \mathbf{q}}{\tilde{Q}^2}$$

- ▶ Choose $\mathbf{p} = (2, 0, 0) \times 2\pi/L$



LATTICE DETAILS

$L^3 \times T$	a (fm)	N_{cfg}	N_Ψ
$24^3 \times 48$	0.0813	6500	2
$32^3 \times 64$	0.0600	4500	3
$40^3 \times 80$	0.0502	$\mathcal{O}(5000)$	4
$48^3 \times 96$	0.0407	$\mathcal{O}(5000)$	5

- ▶ Quenched approximation with $m_\pi = 550$ MeV
- ▶ Wilson-clover fermions with non-perturbatively tuned c_{SW}
- ▶ With clover term, results fully $O(a)$ improved

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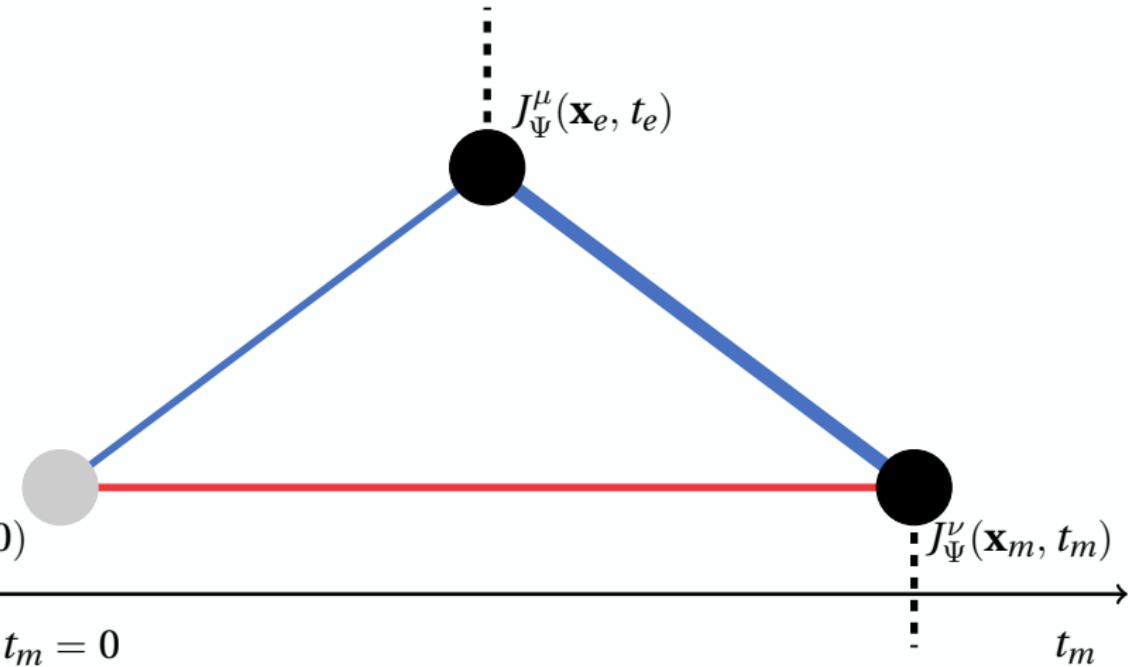
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Still to come...

- ▶ Quenched approximation with $m_\pi = 550$ MeV
- ▶ Wilson-clover fermions with non-perturbatively tuned c_{SW}
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MATRIX ELEMENT CALCULATED

$$C_3^{\mu\nu}(t_e, t_m; \mathbf{p}_e, \mathbf{p}_m) = \int d^3x_e d^3x_m e^{i\mathbf{p}_e \cdot \mathbf{x}_e + i\mathbf{p}_m \cdot \mathbf{x}_m} \langle \Omega | T\{J_\Psi^\mu(x_e) J_\Psi^\nu(x_m) \mathcal{O}_\pi^\dagger(0)\} | \Omega \rangle$$



RATIO METHOD

- ▶ Excited state dependent only on sum $t_e + t_m$.

$$C_3^{\mu\nu}(t_e, t_m; \mathbf{p}_e, \mathbf{p}_m) = R^{\mu\nu}(t_e - t_m; \mathbf{p}, \mathbf{q}) \frac{Z_\pi(\mathbf{p})}{2E_\pi(\mathbf{p})} e^{-E_\pi(\mathbf{p})(t_e + t_m)/2} + \dots,$$

- ▶ Define $t_+ = t_e + t_m$, $t_- = t_e - t_m$.

$$C_3^{\mu\nu}(t_e, t_m; \mathbf{p}_e, \mathbf{p}_m) = R^{\mu\nu}(t_-; \mathbf{p}, \mathbf{q}) \frac{Z_\pi(\mathbf{p})}{2E_\pi(\mathbf{p})} e^{-E_\pi(\mathbf{p})t_+/2} + \dots,$$

- ▶ Consider two sets of time; (t_e, t_m) and $(t'_e, t'_m) = (t_e + \delta, t_m - \delta)$

$$t'_+ = (t_e + \delta) + (t_m - \delta) = t_e + t_m = t_+$$

$$t'_- = (t_e + \delta) - (t_m - \delta) = t_e - t_m + 2\delta \neq t_-$$

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RATIO METHOD

- ▶ Construct ratio with fixed t_+ , varying t_- ($\delta = -1$)

$$\begin{aligned}\mathcal{R} &= \frac{C_3^{\mu\nu}(t_e - 1, t_m + 1; \mathbf{p}_e, \mathbf{p}_m)}{C_3^{\mu\nu}(t_e, t_m; \mathbf{p}_e, \mathbf{p}_m)} \\ &= \frac{R^{\mu\nu}(t_e - t_m - 2; \mathbf{p}, \mathbf{q})}{R^{\mu\nu}(t_e - t_m; \mathbf{p}, \mathbf{q})} \frac{\frac{Z_\pi(\mathbf{p})}{2E_\pi(\mathbf{p})} e^{-E_\pi(\mathbf{p})(t_e + t_m)/2}}{\frac{Z_\pi(\mathbf{p})}{2E_\pi(\mathbf{p})} e^{-E_\pi(\mathbf{p})(t_e + t_m)/2}} \left[1 + \dots \right]\end{aligned}$$

- ▶ Need two t_e , ie t_e and $t_e - 1$
- ▶ No need for 2-point data!
- ▶ No renormalization required.
- ▶ No f_π

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OPERATOR OPTIMIZATION

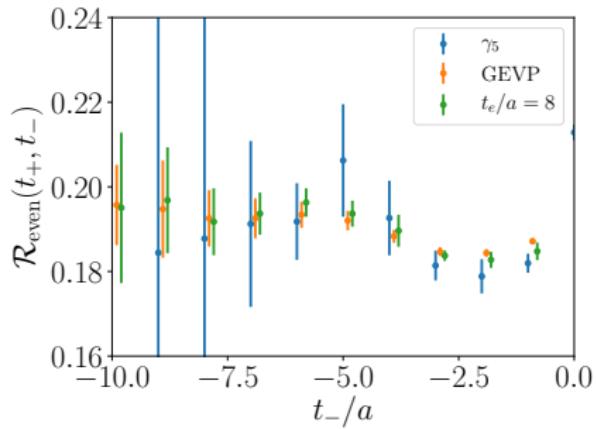
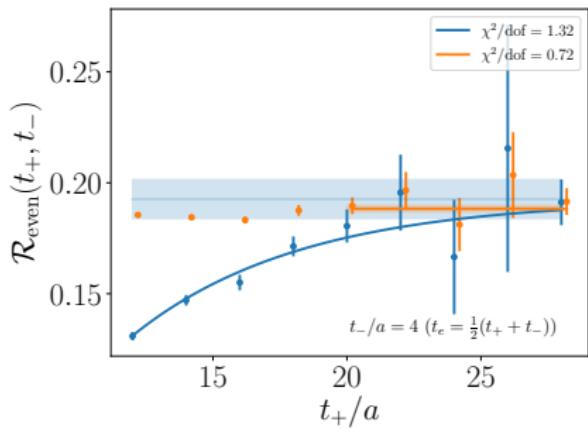
- ▶ Momentum smearing (Bali et al)
- ▶ Variational analysis:

$$\mathcal{O}_\pi(x) = c_1 \mathcal{O}_1(x) + c_2 \mathcal{O}_2(x), \quad \mathcal{O}_1(x) = \bar{\psi} \gamma_5 \psi, \quad \mathcal{O}_2(x) = \bar{\psi} \gamma_4 \gamma_5 \psi$$

$$\begin{aligned} C_{3,\text{GEVP}}^{\mu\nu}(t_e, t_m; \mathbf{p}_e, \mathbf{p}_m) &= c_1 C_{3,\gamma_5}^{\mu\nu}(t_e, t_m; \mathbf{p}_e, \mathbf{p}_m) \\ &\quad + c_2 C_{3,\gamma_4\gamma_5}^{\mu\nu}(t_e, t_m; \mathbf{p}_e, \mathbf{p}_m) \end{aligned}$$

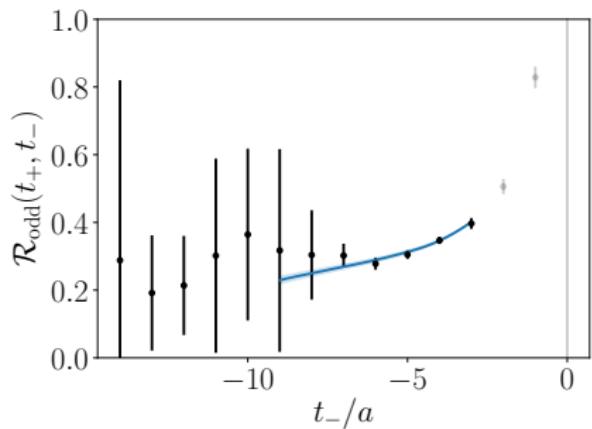
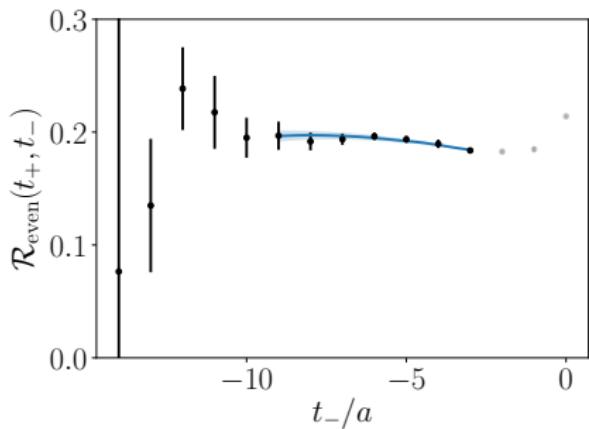
EXCITED STATE CONTAMINATION

$$\mathcal{R} == \frac{R^{\mu\nu}(t_e - t_m - 2; \mathbf{p}, \mathbf{q})}{R^{\mu\nu}(t_e - t_m; \mathbf{p}, \mathbf{q})} \frac{\frac{Z_\pi(\mathbf{p})}{2E_\pi(\mathbf{p})} e^{-E_\pi(\mathbf{p})(t_e + t_m)/2}}{\frac{Z_\pi(\mathbf{p})}{2E_\pi(\mathbf{p})} e^{-E_\pi(\mathbf{p})(t_e + t_m)/2}} \left[1 + \dots \right]$$



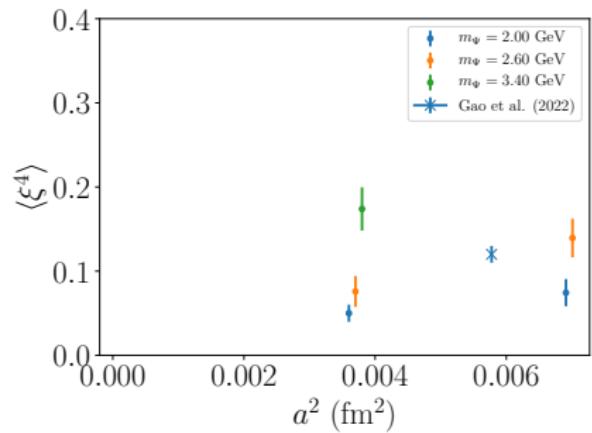
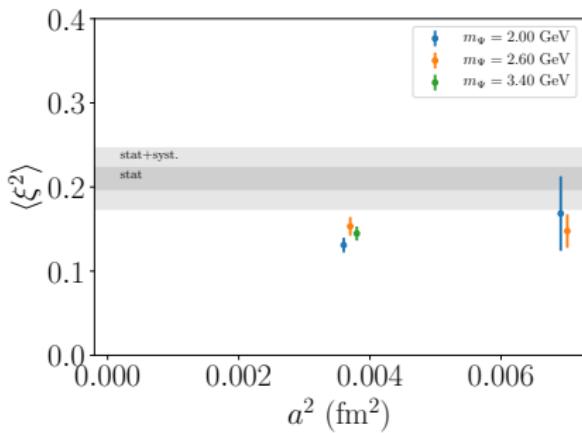
EXAMPLE AT FIXED LATTICE SPACING

$$\mathcal{R}(t_-, \mathbf{p}, \mathbf{q}, m_\psi, \langle \xi^2 \rangle, \langle \xi^2 \rangle; a) = \mathcal{R}_{\text{HOPE}}(t_-, \mathbf{p}, \mathbf{q}, m_\psi, \langle \xi^2 \rangle, \langle \xi^2 \rangle) + \mathcal{O}(a^2)$$



$$L/a = 24, m_\Psi = 2.0 \text{ GeV}, \langle \xi^2 \rangle = 0.17 \pm 0.04, \langle \xi^4 \rangle = 0.07 \pm 0.02$$

STATUS OF CALCULATION



$$\langle \xi^n \rangle (a, m_\Psi) = \langle \xi^n \rangle + \frac{A}{m_\Psi} + Ba^2 + Ca^2 m_\Psi + Da^2 m_\Psi^2$$

FURTHER WORK & CONCLUSIONS

- ▶ HOPE can be used to extract Mellin moments.
- ▶ Introduced ratio method: no renormalization.
- ▶ 2 lattice spacings, 3 heavy quark masses.
- ▶ Aim for 2 more lattice spacings + more heavy quark masses to enable continuum, twist-2 extrapolation.

HIGHER-TWIST/LATTICE ARTIFACTS TRADEOFF

