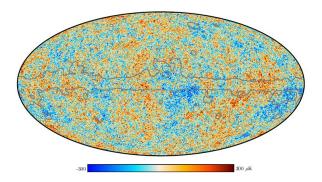
# $\sigma_{\pi N}$ with N<sub>f</sub> = 2 + 1 O(a)improved Wilson fermions

Andria Agadjanov, <u>Dalibor Djukanovic</u>, Georg von Hippel, Konstantin Ottnad, Harvey B. Meyer, Hartmut Wittig

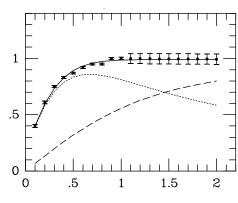
# **Motivation-Dark Matter**

#### Rotation Curves of Galaxies Plot from arXiv:2104.1148



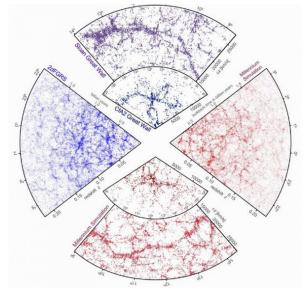
#### Cosmic Microwave Background

Plot from Y. Akrami, et al., Planck 2018 results. I. Overview and the cosmological legacy of Planck, arXiv:1807.06205.



Evidence for Dark Matter





#### Simulation of Galaxy Structures

Plot from V. Springel, C. S. Frenk, S. D. White, The large-scale structure of the Universe, Nature, 440 (2006) 1137. arXiv:astro-ph/0604561

#### **Gravitational Lensing**

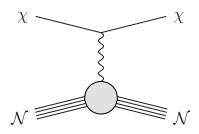
Credit: NASA/CXC/CfA/M. Markevitch et al.; NASA/STScI; Magel- lan/U.Arizona/D.

Clowe et al.; NASA/STScI; ESO WFI



### **DM Direct Searches**

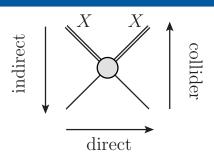
- DM Candidate: WIMP
- WIMP-Nucleus-Scattering



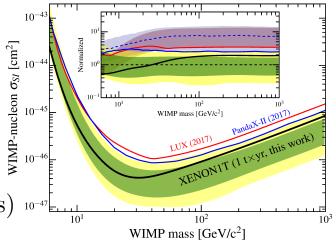
 Rate WIMP-Nucleus Scattering

$$R \sim \frac{\sigma^{\rm SI}}{M_{\chi}\mu_N^2} \times (\text{Nuclear Physics}) \times (\text{Astrophysics})$$

- $\sigma^{\mathrm{SI}}$  depends on Sigma-Term
- Crucial input for interpretation of experiments



Plot from E. Aprile et al,
Dark Matter Search Results from a One Tonne×Year
Exposure of XENON1T
arXiV:1805.12562



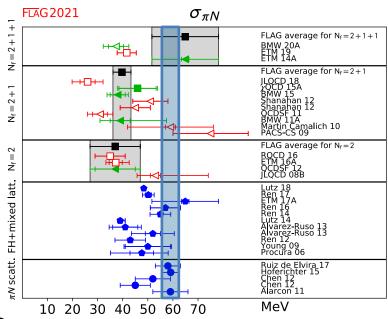
# Sigma-Term

- No scalar probe!
- Phenomenologically via Pion-Nucleon-Scattering (Chang-Dashen-Theorem + extrap.)
- Lattice calculation

$$\sigma_{\pi N} = m_l \langle N | \bar{u}u + \bar{d}d | N \rangle = m_l \frac{\partial m_N}{\partial m_l}$$

Directly or via Mass

 Some tension between Roy-Steiner based estimate and Lattice

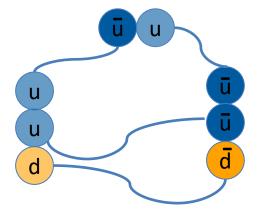


#### **Direct Determination**

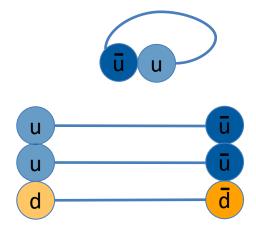
- Connected part
  - Sequential Source
  - Zero Momentum at sink

$$C_{2}(t; \mathbf{p}) = \Gamma_{\alpha\beta} \sum_{\mathbf{x}} e^{-i\mathbf{p}\mathbf{x}} \Big\langle \Psi_{\beta}(\mathbf{x}, t) \overline{\Psi}_{\alpha}(0) \Big\rangle,$$

$$C_{3}(t, t_{s}; \mathbf{q}) = \Gamma'_{\alpha\beta} \sum_{\mathbf{x}, \mathbf{y}} e^{i\mathbf{q}\mathbf{y}} \Big\langle \Psi_{\beta}(\mathbf{x}, t_{s}) \mathcal{O}_{S}(\mathbf{y}, t) \overline{\Psi}_{\alpha}(0) \Big\rangle,$$



- Disconnected part
  - Loops All-to-All: OET+HPE+HP
  - Noisy:
    - → Additional two-point functions

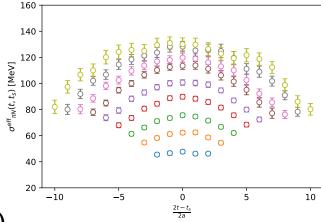


$$C_3^{\text{disc}}(t, t_s; \mathbf{0}) = \left\langle L_S(\mathbf{0}, z_0) \cdot C_2(\mathbf{p}', y_0, x; \Gamma') \right\rangle - \left\langle L_S(\mathbf{0}, z_0) \right\rangle \cdot \left\langle C_2(\mathbf{p}', y_0, x; \Gamma') \right\rangle$$

#### Excited States – Summation

Usual Ratio (forward limit):

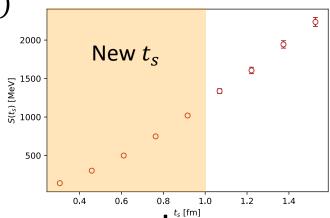
$$R(t,t_s) = \frac{C_3(t,t_s)}{C_2(t_s)} \qquad \operatorname{Re} R(t,t_s) \xrightarrow{t,(t_s-t)\gg 0} G_{S} \xrightarrow{\begin{cases} \frac{1}{2} & 100 \\ \frac{1}{2} & 100$$



Excited states  $\sim e^{-\Delta t}$ ,  $e^{-\Delta(t_s-t)}$ 

• Summed correlator:

$$S(t_s) = \sum_{t=t_c}^{t_s - t_c} \sigma_{\pi N}^{\text{eff}}(t, t_s)$$



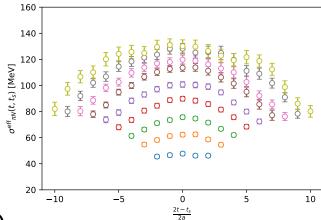
Excited states parametrically suppressed

### Excited States – Summation

Usual Ratio (forward limit):

$$R(t,t_s) = \frac{C_3(t,t_s)}{C_2(t_s)} \qquad \text{Re } R(t,t_s) \xrightarrow{t,(t_s-t)\gg 0} G_S \xrightarrow{\mathbb{Z}}_{0}^{120} G_S$$

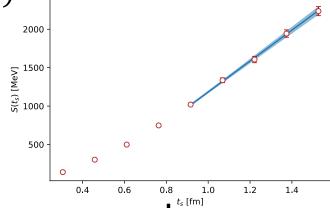
$$G_{\rm S}^{\rm eff}(t,t_s)$$
 = Re  $R(t,t_s)$ 



Excited states  $\sim e^{-\Delta t}$ ,  $e^{-\Delta(t_s-t)}$ 

• Summed correlator:

$$S(t_s) = \sum_{t=t_c}^{t_s - t_c} \sigma_{\pi N}^{\text{eff}}(t, t_s)$$



Excited states parametrically suppressed

$$S(t_s) = (\sigma_{\pi N})$$

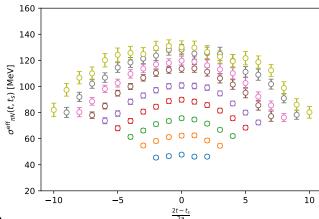
$$(1+t_{s}-2t_{c})$$

### Excited States – Summation

Usual Ratio (forward limit):

$$R(t,t_s) = \frac{C_3(t,t_s)}{C_2(t_s)} \qquad \operatorname{Re} R(t,t_s) \xrightarrow{t,(t_s-t)\gg 0} G_{S} \xrightarrow{\mathbb{S}^{\frac{120-5}{2000}}} G_{S}$$

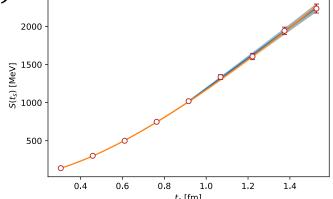
$$G_{S}^{\text{eff}}(t,t_s) = \operatorname{Re} R(t,t_s) \xrightarrow{40-6} G_{S}$$



Excited states  $\sim e^{-\Delta t}$ ,  $e^{-\Delta(t_s-t)}$ 

• Summed correlator:

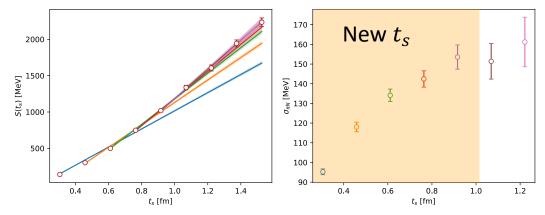
$$S(t_s) = \sum_{t=t_c}^{t_s - t_c} \sigma_{\pi N}^{\text{eff}}(t, t_s)$$



Excited states parametrically suppressed

$$S(t_s) = (\sigma_{\pi N} + m_{11}e^{-\Delta t_s}) (1 + t_s - 2t_c) + e^{-\Delta t_s} \frac{2m_{10} \left(e^{\Delta(1 - t_c + t_s)} - e^{\Delta t_c}\right)}{e^{\Delta} - 1} + \dots$$

#### Summation

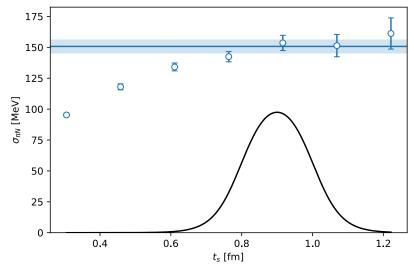


- Excited State Fits need priors for gap  $\Delta$  (like explicit 2-state-Fit)
- Linear Fits:
  - Not trustworthy for small  $t_{s}$
  - Error increases with larger starting  $t_{s}$
  - Several possibilities
    - Choose one, use weights according to AIC, p-values, ...
    - Define a fit-range in physical units (see Talks by J. Koponen,
       M. Salg) → Window

$$w_i = \frac{1}{2} \tanh \frac{t_s - t_{lo}}{\Delta t} - \frac{1}{2} \tanh \frac{t_s - t_{up}}{\Delta t}$$



#### Summation



• Instead of one particular starting  $t_{\mathcal{S}}$  use window function as weights with

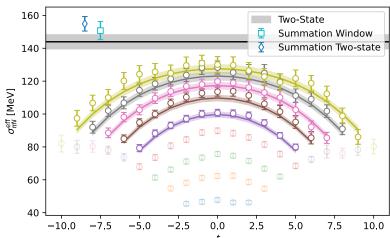
$$t_{\rm lo} = 0.8 \, {\rm fm}, \quad t_{\rm up} = 1.0 \, {\rm fm} \quad {\rm and} \quad \Delta t = 0.08 \, {\rm fm}$$

for all ensembles

- Close to "plateau" average for every ensemble
- Less affected by single-point-estimate-fluctuation

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# Comparison of Methods



Two-State Direct Fit uses:

$$\sigma_{\pi N}^{\text{eff}} = \sigma_{\pi N} + m_{10} \left( e^{-\Delta t} + e^{-\Delta (t_s - t)} \right) + m_{11} e^{-\Delta t_s} e^{(t - \frac{t_s}{2})/a}$$

Needs gaussian priors  $\Delta=2m_{\pi}$  (5% width)

- Summation Two-State:
  - Fits to summed correlator including first excited state (excluding  $m_{11}$ )
  - Needs gaussian priors  $\Delta=2m_\pi$  (5% width)
  - Single starting  $t_s$  (no average)
- Summation Window:
  - Fits to summed correlator no excited states
  - Window average
  - No priors
- For all ensembles **Summation Window** compatible within **2**  $\sigma$



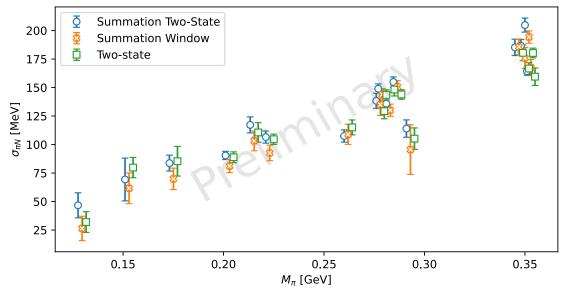
## Lattice Setup

ID	a [fm]	T/a	L/a	$M_{\pi}[\mathrm{MeV}]$	$M_{\pi}L$	$t_{ m sep}[{ m fm}]$	$N_{ m cfg}$
H102	0.086	96	32	354	4.96	0.35, 0.43, 0.52, 0.6, 0.69, 0.78, 0.86, 0.95, 1.04, 1.12, 1.21, 1.3, 1.38, 1.47	2005
H105		96	32	280	3.93		1027
C101		96	48	225	4.73		2000
N101		128	48	281	5.91		1596
S400	0.076	128	32	350	4.33	0.31, 0.46, 0.61, 0.76, 0.92, 1.07, 1.22, 1.37, 1.53	2873
N451		128	48	286	5.31		1011
D450		128	64	216	5.35		500
D452		128	64	153	3.79		1000
N203	0.064	128	48	346	5.41	0.26, 0.39, 0.51, 0.64, 0.77, 0.9, 1.03, 1.16, 1.29, 1.41	1543
N200		128	48	281	4.39		1712
D200		128	64	203	4.22		2000
E250		192	96	129	4.04		400
S201		128	32	293	3.05		2093
N302	0.050	128	48	348	4.22	0.2, 0.3, 0.4, 0.5, 0.6, 0.7, 0.8, 0.9, 1., 1.1, 1.2, 1.3, 1.39	2201
J303		192	64	260	4.19		1073
E300		192	64	174	4.21		570

- Enlarged range in  $t_{sep}$ 
  - → Monitor excited state contribution
- Roughly same statistics at every  $t_{sep}$  (setup see talk K. Ottnad)
  - $\rightarrow$  Number of sources adapted to  $t_{sep}$
- Chiral/Continuum/Finite-Size extrapolation possible



#### Results



- For ensembles < 250 MeV</li>
   Two-State analysis (including priors for Δ) generally slightly above
- Two-State analysis very sensitive to priors



Chiral Expansion

$$\sigma_{\pi N} = (k_1 + k_a a^2) M^2 + k_2 M^3 + 2k_3 M^4 \log \frac{M}{\mu} + k_4 M^4 + k_L M^2 \left(\frac{1}{L} - \frac{M}{2}\right) e^{-ML}$$

- $k_1, k_2, k_3$  depend on (known) LECs
- Challenging fits, leaving all parameters free
  - Fits not stable, especially with cuts in pion mass
  - Only fit up to  $M^3$ 
    - → Coefficients inconsistent with ChPT
  - $-k_3$  and  $k_4$  competing (cancellations)
    - $\rightarrow$  drop  $k_3$
  - Or take ChPT information into account via priors



Chiral Expansion

$$\sigma_{\pi N} = (k_1 + k_a a^2) M^2 + k_2 M^3 + 2k_3 M^4 \log \frac{M}{\mu} + k_4 M^4 + k_L M^2 \left(\frac{1}{L} - \frac{M}{2}\right) e^{-ML}$$

- $k_1, k_2, k_3$  depend on (known) LECs
- Challenging fits, leaving all parameters free
  - Fits not stable, especially with cuts in pion mass
  - Restrict  $k_1$ ,  $k_2$  (,  $k_3$ ) via gaussian priors (still fitted)

$$k_1 = -4c_1 = (4.44 \pm 0.12) \,\text{GeV}^{-1}$$

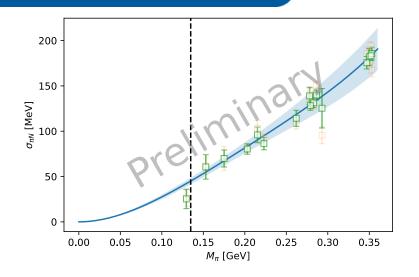
$$k_2 = -\frac{9g_A^2}{64\pi F_\pi^2} = (-8.52 \pm 0.04) \,\text{GeV}^{-2}$$

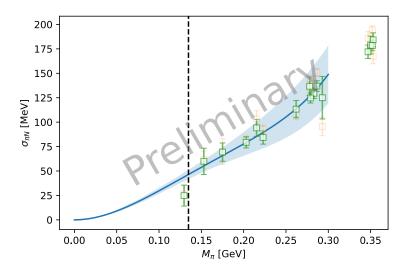
$$k_3 = -\frac{3}{32\pi^2 F_\pi^2} \left(\frac{g_A^2}{m_N} - 8c_1 + c_2 + 4c_3\right) - \frac{c_1}{8\pi^2 F_\pi^2}$$

$$= (-11.38 \pm 0.35) \,\text{GeV}^{-3}$$

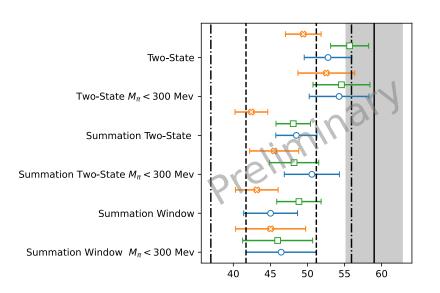
– Width 5x the error for  $k_1$ ,  $k_2$  (,  $k_3$ )







- Data from window average of summation data
- Green points:
   After correction of lattice artefacts (only central value)
   of original data (orange points)
- Left: No Pion mass cut
  - Right: Pion mass cut 300 MeV



- With Restrictions from ChPT: Result stable w.r.t
  - Cuts in pion mass
  - Window vs Two-State
- Two-State result higher but within  $2 \sigma$  of Summation Window result
- Grey Area: Result using Roy-Steiner
- CCF result depends on priors
- Fits: blue, green, orange = (k1,k2),(k1,k2,k3),(k1,k2,no log)
- Further Improvements:
  - Increase statistics for the disconnected part
  - Simultaneous fit with nucleon mass



# Thanks!