# Unpolarized gluon PDF for the proton using the twisted mass formulation

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The 39th International Symposium on Lattice Field Theory

August 10, 2022



Introduction



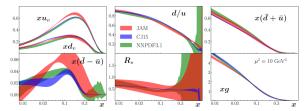
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#### Motivation

- Gluon contributions to physical quantities play a critical role in hadron structure
- Gluon contributions can be large, eq. gluon momentum fraction  $\approx 40\%$
- Dedicated experimental efforts to understand gluonic structure of hadron

[Moffat et al. PRD 104, 016015 (2021)] [Ball et al. EPJC 77, 663 (2017)] [Accardi et al, PRD 93, 114017 (2016)]



Lattice studies of gPDFs can assist in constraining global analysis

## gPDFs on the Lattice

- Several challenges in extracting reliable results
  - purely disconnected diagram
  - at least an order of magnitude more statistics than quark counterparts
  - unavoidable mixing with quark singlet PDFs
- x-dependence of gluon PDFs even more challenging
- Inverse problem in reconstruction of x-dependence due to limited lattice data

- Matrix elements of non-local operators and momentum-boosted proton states
  - Several choices for the form of the gluon operator consisting of two field-strength tensors, separated by spatial distance z, and two straight Wilson lines, connecting points  $0 \rightarrow z$  and  $z \rightarrow 0$

$$M_{\mu i;\nu j}(P,z) = \langle N(P)|F_{\mu i}(z)W(z,0)F_{\nu j}(0)W(0,z)|N(P)\rangle$$

- Choice of indices for  $F_{\mu\nu}$  not unique
- This operator avoids finite mixing under renormalization
  - must subtract vacuum expectation value

$$\mathcal{O} = \frac{1}{2} \sum_{i \neq 3} F_{i3}(x + z\hat{z}) W(x + z\hat{z}, x) F_{i3}(x) - \sum_{i \neq j \neq 3} F_{ij}(x + z\hat{z}) W(x + z\hat{z}, x) F_{ij}(x)$$

- Matrix elements extracted from ratio of 2pt- and 3pt- functions
- Ground state from plateau fit

$$\frac{C^{3pt}(t,\tau,0,\vec{P})}{C^{2pt}(t,0,\vec{P})} \stackrel{0<<\underline{\tau}<< t}{=} \mathcal{M}(\nu,z^2)$$

## Pseudo-PDF Approach (in a nutshell)

 Form the double ratio (reduced ITD) with zero-momentum and local matrix elements to reduce higher twist contributions [Orginos et al., Phys.Rev.D 96 (2017) 9, 094503]

$$\mathfrak{M}(\nu,z^2) \equiv \left(\frac{\textit{M}(\nu,z^2)}{\textit{M}(\nu,0)|_{z=0}}\right) \bigg/ \left(\frac{\textit{M}(0,z^2)|_{p=0}}{\textit{M}(0,0)|_{p=0,z=0}}\right)$$

- Scale evolution and apply matching kernel on ITD
  - neglect mixing with quark singlet
  - $\blacksquare$  normalize with  $\langle x \rangle_g$

$$\mathcal{Q}(\nu, z^2, \mu^2) = \mathfrak{M} + \frac{\alpha_s N_c}{2\pi} \int_0^1 du \, \mathfrak{M}(u\nu, z^2) \left\{ \ln\left(\frac{z^2 \mu^2 e^{2\gamma_E}}{4}\right) B(u) + 4\left[\frac{u + \ln(\overline{u})}{\overline{u}}\right]_+ + \frac{2}{3} \left[1 - u^3\right]_+ \right\}$$

■ Reconstruct x-dependence (Backus-Gilbert, Fourier tansform, etc.)

$$Q(\nu, z^2, \mu^2) = \int_0^1 dx \cos(x\nu) x g(x, \mu^2)$$

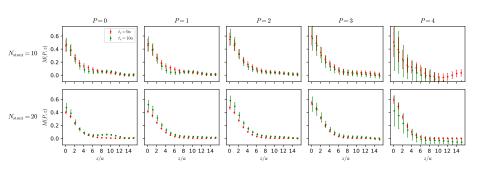
#### Lattice Parameters and Statistics

- $N_f = 2+1+1$  ensemble of twisted-mass clover fermions and Iwasaki improved gluons
  - lacksquare  $m_\pi=$  260 MeV
  - a = 0.09471(39) fm
  - $L^3 \times T = 32^3 \times 64$
  - $\blacksquare$   $Lm_{\pi}=4$
- Stout smearing ( $\omega = 0.129$ )
  - field-strength tensor: 10, 20 steps
  - Wilson line: 0, 10 steps
- Momentum smearing (optimized value  $\xi = 0.6$ ) used for P = 2, 3, 4 [Bali et al, PRD 93, 094515 (2016)]
- Excited states:
  - Numerical results feasible up to  $t_s = 10a$
- Statistics
  - Average over all 6 spatial directions of Wilson line / momentum  $(\pm x, \pm y, \pm z)$
  - Statistics much higher than quark PDFs

$ \mathbf{P_3} \left[\frac{2\pi}{L}\right]$	<b>P</b> <sub>3</sub>   [GeV]	$N_{\rm confs}$	N <sub>src</sub>	N <sub>dir</sub>	Total statistics
0 to 4	0 - 1.67	1,134	200	6	1,360,800

## Matrix Elements: Excited States Contamination and Effect of N<sub>stout</sub>

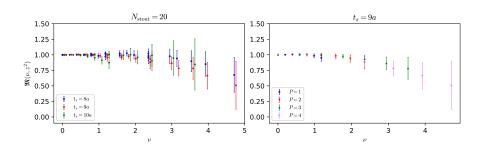
Results



- $\blacksquare$  Various values of  $t_s$  and two stout steps for the gluon operator
- Statistical errors increase with momentum boost and ts
- MEs have expected behavior (higher boosts decay faster to 0)
- Final results use N<sub>stout</sub> = 20

## Double Ratio (Reduced ITD)

■ Double-ratio analysis to narrow down  $t_s$ 

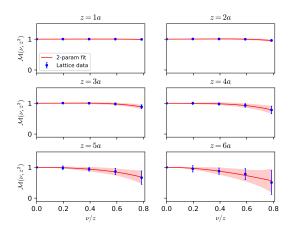


- $\blacksquare$  Values at different  $t_s$  are compatible within uncertainties
- Final results use  $t_s = 9a$ 
  - $z_{max} = 6a = 0.568 \text{ fm}$
- Lattice data form a smooth function
  - Must interpolate for evolution and matching

Results

## Interpolation of Double Ratio

- We interpolate the double-ratio at each z to get a continuous function for the integration
  - interpolation done with linear and 2nd-order polynomial fits

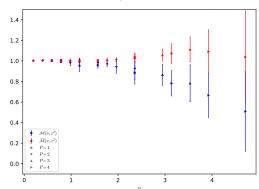


- 2nd-order polynomial fits prove to be the most suitable for evolution and matching
- Choice of fit mostly irrelevant below
  z = 4a

## ITD Development

**a** Apply the evolution kernel to the reduced matrix elements to the scale  $\mu = 2$  GeV ahead of final conversion to  $\overline{\mathrm{MS}}$  scheme

$$\tilde{\mathfrak{M}}(\nu,z^2,\mu^2)=\mathfrak{M}+\frac{\alpha_{s}N_{c}}{2\pi}\int_{0}^{1}du\ln(\frac{z^2\mu^2\mathrm{e}^{2\gamma_{E}}}{4})B(u)\mathfrak{M}(u\nu,z^2)$$

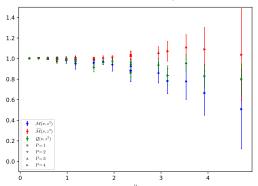


- Data from different (P, z) pairs fall on a universal curve
- We find good agreement up to z = 6a

## ITD Development

 $\blacksquare$  Apply the matching kernel to convert to  $\overline{\rm MS}$  scheme

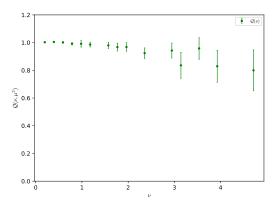
$$Q(\nu, z^2, \mu^2) = \tilde{\mathfrak{M}}(\nu, z^2, \mu^2) + \frac{\alpha_s N_c}{2\pi} \int_0^1 du \ L(u) \mathfrak{M}(u\nu, z^2)$$



- We continue to find good agreement between common values of loffe time from different combinations of momenta and Wilson line lengths
- Matching effects in opposite direction of evolution

# **ITD Development**

 $\blacksquare$  Average over common  $\nu$  for final pseudo-ITD



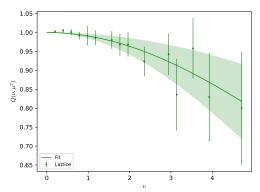
 $\blacksquare$  No information remains regarding initial (P, z) pairs

#### PDF Reconstruction

■ The fit is chosen by the minimization of

$$\chi^2 = \sum_{\nu=0}^{\nu_{max}} \frac{\left(Q(\nu,\mu^2) - Q_f(\nu,\mu^2)\right)^2}{\sigma_Q^2(\nu,\mu^2)}$$

 $Q_f$ : ITDs from fitting ansatz



#### Pseudo-PDF

We fit the pseudo-ITD according to

$$Q_f(\nu,\mu^2) = \int_0^1 dx \cos(\nu x) x g(x)$$

where

$$x g(x) = N x^{a} (1 - x)^{b}$$
Reconstructed

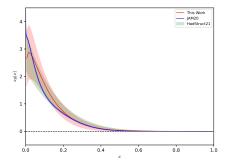
$$x = \frac{3}{3}$$

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- PDF is normalized using gluon momentum fraction  $\langle x \rangle_g^{\overline{\rm MS}}(\mu=2\,{\rm GeV})=0.427(92)$  [Alexandrou et al, PRD 101, 094513 (2020)]
- Other reconstruction method (naive Fourier-transform, Backus-Gilbert method) have proven less suitable [Bhat et al, PRD 103, 034510 (2021)]

# Comparison with Other Works

- Comparison with lattice results from HadStruc collaboration [Khan et al, PRD 104, 094516 (2021)]
  - $m_{\pi} = 358$  MeV, a = 0.094,  $L^3 \times T = 32^3 \times 64$
- JAM20 global analysis [Moffat et al, PRD 104, 016015 (2021)],  $\langle x \rangle_g = 0.40(1)$

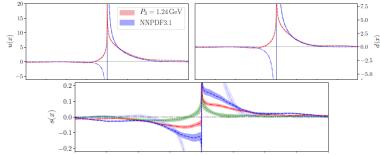


- We find agreement between all results
- This work:  $\nu_{max} = 4.71$ HadStruc:  $\nu_{max} = 7.07$

#### **Future Work**

- Addressing systematic effects
- Increasing the range of accessed loffe times
- Investigation of mixing with quark singlet PDFs

[Alexandrou et al, PRD 104, 054503 (2021)]



Calculations at physical point

## Acknowledgements

- This research is financial support by the U.S. Department of Energy, Office of Nuclear Physics, Early Career Award under Grant No. DE-SC0020405
- The calculations have been partly carried out on HPC resources supported in part by the National Science Foundation through major research instrumentation grant number 1625061 and by the US Army Research Laboratory under contract number W911NF-16-2-0189
- Computations for this work were carried out in part on facilities of the USQCD Collaboration, which are funded by the Office of Science of the U.S. Department of Energy.