# Two- and three-particle scattering in the $(1+1)$-dimensional $O(3)$ non-linear sigma model 

Jorge Baeza-Ballesteros

In collaboration with M. T. Hansen

IFIC, University of Valencia-CSIC
Lattice22-8th August 2022


$$
\begin{aligned}
& \text { VNIVERSITAT } \\
& \text { IEOQVALENCIA }
\end{aligned}
$$



## Finite-volume quantization conditions

Finite-volume spectrum:

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Two-particle QC: Lüscher's formalism [1986]

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Three-particle QC: various approaches

- Relativistic field theory (RFT) [Hansen, Sharpe $(2014,2015)]$
- Non-relativistic effective field theory (NREFT) [Hammer, et al. (2017)]
- Finite-volume unitarity (FVU) [Döring, Mai $(2016,2017)$ ]


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1. Constrain two- and three-body $\mathcal{K}$ matrices

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Goal: Test RFT formalism on a solvable toy model

## $\mathrm{O}(3)$ non-linear sigma model in $(1+1)$ dimensions

We study the $\mathrm{O}(3)$ non-linear sigma model (NLSM) in $1+1$ dimensions

$$
\begin{array}{ll}
S[\sigma]=\frac{\beta}{2} \int \mathrm{~d}^{2} x \partial_{\mu} \sigma(x) \cdot \partial_{\mu} \sigma(x) & \sigma=\left(\sigma^{1}, \sigma^{2}, \sigma^{3}\right) \\
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- Dynamical mass gap $m$
- Global $\mathrm{O}(3)$ isospin symmetry


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Low-energy spectrum: isospin-1 multiplet

- Global O(3) isospin symmetry

Integrable at low energies [Zamolodchikov, Zamolodchikov (1977)]:


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\begin{array}{r}
3 \otimes 3=5 \oplus 3 \oplus 1 \\
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Three particles: 7 irreps $\rightarrow 4$ isospin channels

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3 \otimes 3 \otimes 3 & =7 \oplus(5 \oplus 5) \oplus(3 \oplus 3 \oplus 3) \oplus \\
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1 \\
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$\oplus(5$

$\oplus 3$ $\qquad$ 3) $\oplus 1$
$I=3 \quad I=2 \quad I=1 \quad I=0$
$C_{l=3}, C_{l=2}$ D

$2 \times 2$ matrix for
each momenta combination

## Quantization conditions in $1+1$ dimensions

QCs in $1+1$ dimensions $\rightarrow$ no angular momentum

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$\omega_{k}=E_{\mathrm{CM}} / 2$ :

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\cot \delta(k)=-\frac{1}{2}\left[\cot \left(\frac{L \gamma\left(k+\omega_{k} \beta\right)}{2}\right)+\cot \left(\frac{L \gamma\left(k-\omega_{k} \beta\right)}{2}\right)\right]
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- Three-particle QC: analogous to $(3+1)$ [Work in progress]


## Lattice simulations

We use the standard lattice action

$$
S[\sigma]=-\beta \sum_{x} \sum_{\mu} \sigma(x) \cdot \sigma(x+a \hat{\mu})
$$

We generate the configurations and evaluate $n$-point functions using a cluster algorithm:

* Overcomes critical slowing down
* Improves the signal-to-noise ratio


## Single-cluster algorithm

## Cluster algorithm:

1. Choose a random unit vector $r \in \mathbb{R}^{3}$ and a random "seed" site
2. Grow the cluster, $C$, by adding neighbors

$$
p_{\text {add }}=1-\exp \left[\min \left\{-2 \beta \sigma_{r}(x) \sigma_{r}(x+a \hat{\mu}), 0\right\}\right] \quad \sigma_{r}(x)=\sigma(x) \cdot r
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3. Update the cluster

$$
\sigma(x) \rightarrow \sigma(x)-2 \sigma_{r}(x) r
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4. Measure on the cluster

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Example: Two-point function

$$
\sigma_{r}(\tau, p)=\sum_{x \in C} \mathrm{e}^{i p x} \sigma_{r}(\tau, x) \longrightarrow C_{2 \mathrm{pt}}(\tau, p)=3\left\langle\sigma_{r}(\tau, p) \sigma_{r}^{*}(0, p)\right\rangle
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Need to average over many updates!

## Three-cluster algorithm

To measure six-point functions we need a three-cluster algorithm:

1. Choose three random orthogonal unit vector $r, u, v \in \mathbb{R}^{3}$ and three random "seed" site
2. Grow each cluster, $C_{r}, C_{u}$ and $C_{v}$, separately
3. Update each cluster
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Example: Four- and six-point function

|  | $\propto\left\langle\sigma_{r}\left(\tau, q_{2}\right) \sigma_{u}\left(\tau, q_{1}\right) \sigma_{u}^{*}\left(0, p_{2}\right) \sigma_{r}^{*}\left(0, p_{1}\right)\right\rangle$ |
| :---: | :---: |
| $\begin{aligned} & p_{1} \\ & p_{2} \\ & p_{3} \end{aligned} \bullet \bullet q_{1}$ | $\propto\left\langle\sigma_{r}\left(\tau, q_{3}\right) \sigma_{u}\left(\tau, q_{2}\right) \sigma_{v}\left(\tau, q_{1}\right) \sigma_{r}^{*}\left(0, p_{3}\right) \sigma_{v}^{*}\left(0, p_{2}\right) \sigma_{u}^{*}\left(0, p_{1}\right)\right\rangle$ |

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* Tuned $m L$ and $m T \sim 20.5$
* 256/512/1024 replicas
$\star \sim 10$ million thermalization updates
$\star$ Averaged over $\sim 1$ million measurement updates


## Energy spectrum from lattice simulations

Solve GEVP and extract energy spectrum

$$
C^{-1 / 2}\left(t_{0}\right) C(t) C^{-1 / 2}\left(t_{0}\right) v_{n}=\lambda_{n}(t) v_{n} \longrightarrow \lambda_{n}(t) \xrightarrow{T \gg t \gg t_{0}} A_{n} \mathrm{e}^{-E_{n} t}
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2 particles, $I=2$ channel


3 particles, $I=3$ channel


## Continuum limit

The O(3) NLSM has large discretization effects [Balog, et al (2009, 2010)]:

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Q(m a)=Q(0)+C(m a)^{2} \beta^{3}\left[1+\sum_{k=1}^{\infty} c_{k} \beta^{-k}\right]+\mathcal{O}\left(a^{4}\right)
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## Three-particle scattering

We have determined 3-particle finite-volume energies at maximal isospin

$$
P=\frac{2 \pi}{L} n \quad I=3 \text { channel (Preliminary) }
$$



## Summary and outlook

Goal: Test 3-particle RFT formalism on the $(1+1)$-dimensional $\mathrm{O}(3)$ non-linear sigma model

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We have extrapolated the finite-volume energies to the continuum
We have found very good agreement (sub-percent precision) between lattice and analytical results in the two-particle sector

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## Thank you for your attention!

## Exact two-particle S-matrix

Integrable at low energies [Zamolodchikov, Zamolodchikov (1977)]:

## Unitarity + Crossing symmetry $)+$ Factorization



$$
\sigma_{2}(\theta)=\frac{\theta(i \pi-\theta)}{(i 2 \pi-\theta)(i \pi+\theta)}, \quad \sigma_{1}(\theta)=\frac{-i 2 \pi}{i \pi-\theta} \sigma_{2}(\theta), \quad \sigma_{3}(\theta)=\frac{-i 2 \pi}{\theta} \sigma_{2}(\theta)
$$

$s=2 m^{2}[1+\cosh (\theta)]$

## Two-particle $I=0$ channel

$$
P=\frac{2 \pi}{L} n \quad I=0 \text { channel }
$$



## Three-particle $I=3$ and $I=2$ channels


$C_{l=2}=\left(\begin{array}{c}B_{1}-\frac{1}{2} B_{2}+B_{3}-\frac{1}{2} B_{4}-\frac{1}{2} B_{5}-\frac{1}{2} B_{6} \\ \frac{\sqrt{3}}{2}\left[B_{2}-B_{4}-B_{5}+B_{6}\right]\end{array}\right.$

$$
\left.\begin{array}{c}
\frac{\sqrt{3}}{2}\left[B_{2}-B_{4}+B_{5}-B_{6}\right] \\
B_{1}+\frac{1}{2} B_{2}-B_{3}+\frac{1}{2} B_{4}-\frac{1}{2} B_{5}-\frac{1}{2} B_{6}
\end{array}\right)
$$

$1 \times 1$ matrix if a pair of inital/final momenta are equal
No contribution if all three initial/final momenta are equal

