



Determination of Lee-Yang edge singularities in QCD by rational approximations

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Introduction

QCD phase diagram

- ▶ One of our main goals is the study of the **critical points** in the QCD phase diagram
- ▶ We can think of at least three critical regions:
 - (1) The **Roberge-Weiss transition** region
 - (2) The **chiral transition** region
 - (3) The region around the **critical endpoint** of QCD

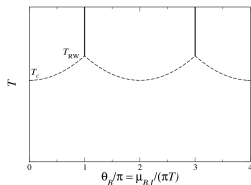


Fig from Phys.Rev. D 93, 74504 (2016)

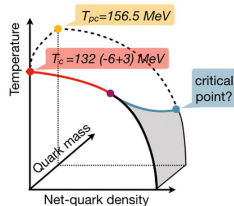


Fig taken from C. Schmidt

Introduction

Yang-Lee edge singularities

In this work we study the **singularities** of QCD in the **complex** μ -plane.

- ▶ There is a deep connection between phase transitions and the singularities in the complex plane.

Consider for instance the **Ising model**

(see also previous talk by F. Di Renzo).

The free energy has a branch cut on the imaginary axis for the symmetry breaking field h . The branch point is known as **Yang-Lee edge singularity**.

[Yang, Lee, 1952]

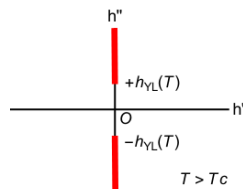


Fig from New J. Phys. 19, 083009 (2017)

- ▶ Why are YLE singularities useful?
 - (1) If for $T \rightarrow T_c$ the YLE singularities end up on real axis this signals the presence of a **physical phase transition**
 - (2) At $T \neq T_c$ the presence of a YLE singularities determines a finite **radius of convergence**

Introduction

The multi-point Padé method

- ▶ How do we get the information that we need from lattice data?
Because of the **sign problem** we cannot directly explore regions at real (or even complex) chemical potentials.
- ▶ In our approach:
 - (1) We perform simulations at **imaginary μ**
 - (2) We calculate the **Taylor coefficients** of a given observable
 - (3) We merge the results using **multi-point Padé** approximants
- ▶ This is a combination of the **Taylor expansion** and **analytic continuation** methods.

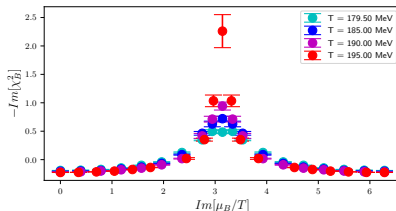
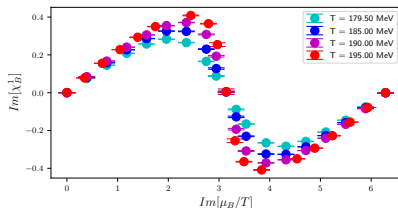
We can extract information about the singularities of the observable by studying the uncanceled poles of the rational approximants and we can analytically continue the rational approximants to real μ .

RW transition region

Numerical set-up

Let's first consider the **high-temperature** regime. The numerical set-up that we adopted is the following.

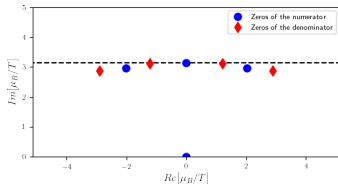
- ▶ We simulate highly-improved staggered quarks (HISQ) in $N_f = 2 + 1$ QCD using $36^3 \times 6$ lattices
- ▶ Simulations performed at $T = 195.0, 190.0, 185.0$ and 179.5 MeV
- ▶ We have measured the cumulants $\chi_{B,Q}^1 \equiv \frac{1}{Z} \frac{\partial Z}{\partial \mu_{B,Q}}$ and their first derivative w.r.t $\hat{\mu}_B$ at $O(10)$ **imaginary** chemical potentials



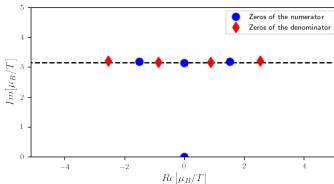
RW transition region

Singularity structure from Padé analysis (I)

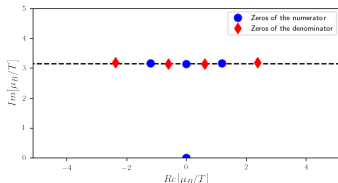
$T = 179.5 \text{ MeV}$



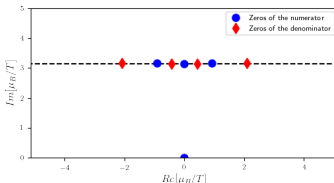
$T = 185.0 \text{ MeV}$



$T = 190.0 \text{ MeV}$



$T = 195.0 \text{ MeV}$

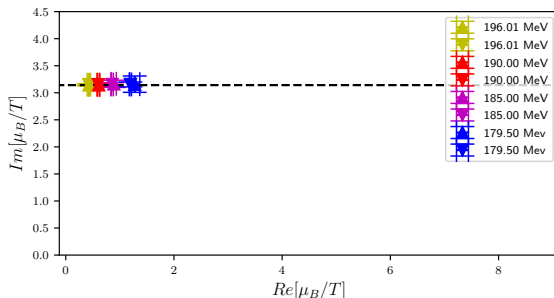


⇒ signature of branch cut singularities pinching the imag axis

RW transition region

Singularity structure from Padé analysis (II)

Summary:



- ▶ For all 4 temperatures we find singularities lying on $\hat{\mu}_B^I = \pi$. We find consistent results from the Padé approximants for χ_B^1 and χ_Q^1 .
- ▶ Are these the **YLE singularities** associated to the **RW critical point**?

RW transition region

Scaling analysis (I)

Theoretical expectations:

- ▶ Magnetic EoS: $M(h, t) = h^{\frac{1}{\delta}} f_G(z) + M_{reg}(h, t)$, where
 - ▶ $t = t_0^{-1} \frac{T - T_{RW}}{T_{RW}}$ and $h = h_0^{-1} \frac{\hat{\mu}_B - i\pi}{i\pi}$ are scaling fields
 - ▶ f_G is a scaling function depending only on the scaling variable $z = t/h^{\frac{1}{\beta\delta}}$
 - ▶ β and δ are the critical exponents [3d, Z(2) universality class]
- ▶ We can solve $t/h^{\frac{1}{\beta\delta}} = z_c \equiv |z_c| e^{i\frac{\pi}{2\beta\delta}}$
 $\rightarrow \hat{\mu}_{YLE}^R = \pm \pi \left(\frac{z_0}{|z_c|} \right)^{\frac{1}{\beta\delta}} \left(\frac{T - T_{RW}}{T_{RW}} \right)^{\frac{1}{\beta\delta}}$, $\hat{\mu}_{YLE}^I = \pi$

Scaling analysis:

- ▶ We can fit the real part of the singularities using the ansatz:

$$\hat{\mu}_{YLE}^R = a \left(\frac{T - T_{RW}}{T_{RW}} \right)^{\frac{1}{\beta\delta}} + b$$

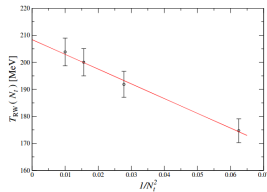
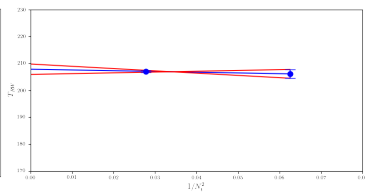
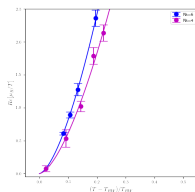
\rightarrow estimate non-universal parameters T_{RW} , z_0 from fit parameters

RW transition region

Scaling analysis (II)

- ▶ By fitting our data using $\beta\delta = 1.5635$ we obtain the estimate $T_{RW} = 207.08(35) \text{ MeV}$.

A rudimentary **continuum extrapolation** using our previous data for $N_\tau = 4$ yields $T_{RW}^{\text{cont}} = 207.1(2.4) \text{ MeV}$.



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- ▶ Using $|z_c| = 2.452$ from [Connely et al., 2020] we also find $z_0 = 10.2-10.5$

Chiral transition and CEP regions

Scaling analysis

▶ Chiral transition region

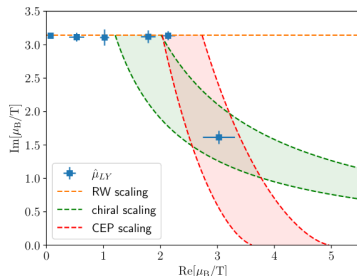
- ▶ $t = t_0^{-1} \left[\frac{T - T_c}{T_c} + k_2^B \left(\frac{\mu_B}{2} \right)^2 \right]$,
 $h = h_0^{-1} \frac{m_l}{m_s^{phys}}$ [Mukherjee and Skokov, 2021]
- ▶ $\hat{\mu}_{YLE} = \left[\frac{1}{k_2^B} \frac{z_c}{z_0} \left(\frac{m_l}{m_s^{phys}} \right)^{\frac{1}{\beta\delta}} - \frac{T - T_c}{T_c} \right]^{\frac{1}{2}}$
- ▶ non-universal parameters T_c, k_2^B, z_0
from [HotQCD data], $|z_c|$ from [Connelly]

▶ CEP region

- ▶ linear ansatz [Basar, 2021]

$$t = \alpha_t (T - T_{CEP}) + \beta_t (\mu_B - \mu_{CEP}) ,$$
$$h = \alpha_h (T - T_{CEP}) + \beta_h (\mu_B - \mu_{CEP})$$

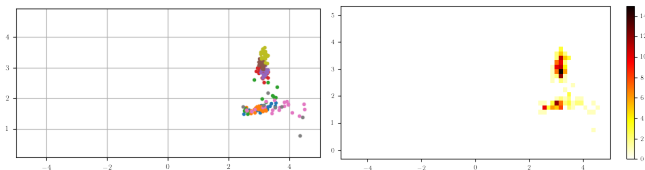
- ▶ $\mu_{YLE} \sim \mu_{CEP} - c_1 (T - T_{CEP}) + ic_2 |z_c|^{-\beta\delta} (T - T_{CEP})^{\beta\delta}$
- ▶ **parameters unknown:** to visualize one can guessestimate
 $\mu_{CEP} = 500 - 630 \text{ MeV}$, $T_{CEP} = T_{pc} \left(1 - k_2^B \left(\frac{\mu_{CEP}}{T_{pc}} \right)^2 \right)$



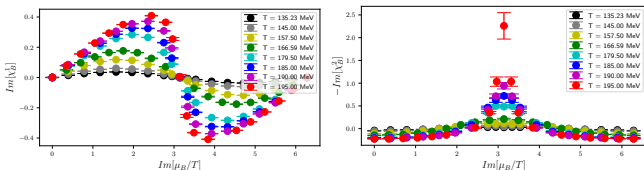
Chiral transition and CEP regions

Numerical set-up and interval dependence

- Previously at 145 MeV we had found a candidate for a chiral singularity by restricting the fit interval from $[0, 2\pi]$ to $[0, \pi]$



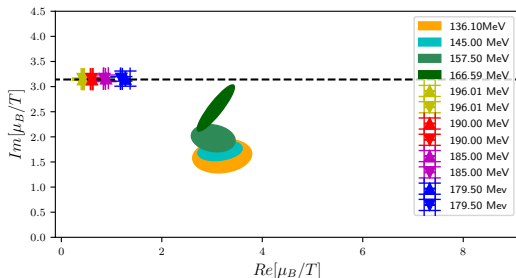
- We performed an additional set of simulations in the **low-temperature** regime ($T = 135.23, 157.50, 166.59$ MeV)



Chiral transition and CEP regions

Singularity structure from Padé analysis (I)

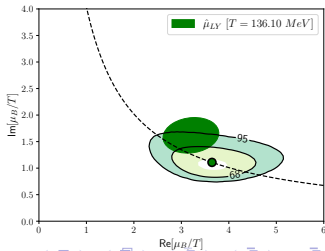
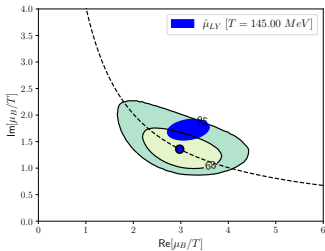
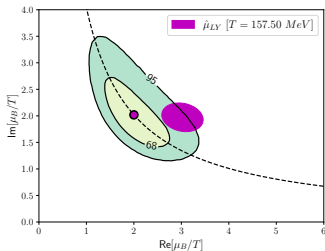
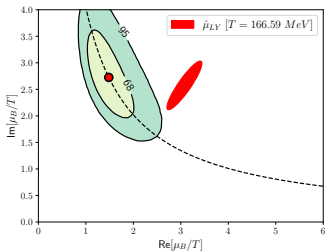
Summary:



- ▶ We observe a **singularity** moving towards the real axis as T is decreased
- ▶ We can infer a **radius of convergence** $\mathcal{R} = |\hat{\mu}_B - 0| \approx 3 \div 4$ for $T = 136 \div 166$ MeV
- ▶ Are these the **YLE singularities** associated to the **chiral transition**?

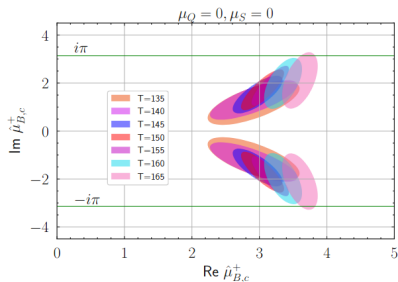
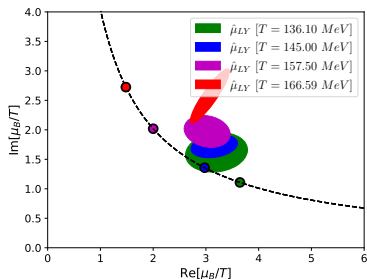
Chiral transition and CEP regions

Singularity structure from Padé analysis (II)



Chiral transition and CEP regions

Singularity structure from Padé analysis (III)



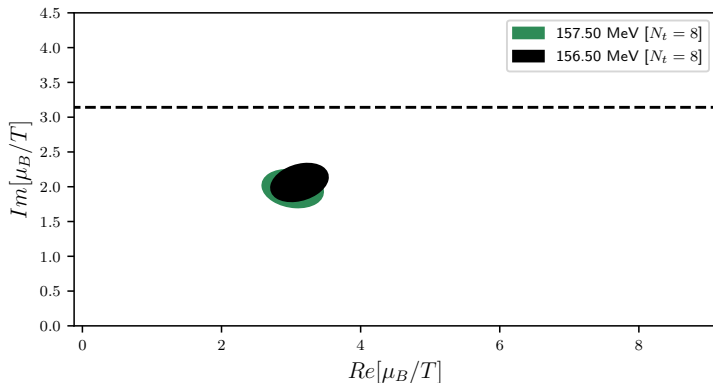
Phys. Rev. D 105, 074511 (2022)

- ▶ Individually the results are in agreement with theoretical expectations within errors, but looking at the big picture the trajectory has a steeper slope than expected...

Chiral transition and CEP regions

$N_\tau = 8$ data

- ▶ We also have data from $N_\tau = 8$ lattices for one temperature. We find a singularity consistent with the singularity that we had found for a very similar temperature on the $N_\tau = 6$ lattice.



Conclusions

- ▶ We have used the [multi-point Padé method](#) to study the singularities of QCD in the complex plane.
- ▶ We have identified the YLE singularities associated to the [Roberge-Weiss critical point](#).

The singularities follow the expected scaling behaviour for a transition belonging to the [3d, Z\(2\) universality class](#).

We also have a preliminary [continuum extrapolation](#) for T_{RW} .

[Future prospects](#): for a proper extrapolation we would need $N_\tau = 8$ data.

- ▶ We have found singularities that *might* be identified as the [YLE singularities](#) associated to the [chiral transition](#). These imply a [radius of convergence](#) $\mathcal{R} \approx 3 \div 4$ for $T = 136 \div 166$ MeV.

[Future prospects](#): simulations with smaller m_l/m_s ratio are currently ongoing.

Thank you for listening!

Backup slide

