

Constraining Beyond The Standard Model Nucleon Isovector Charges

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Motivation

- Precision studies of neutron decay observables are susceptible to beyond the Standard Model tensor and scalar interactions
- Importantly the Fierz interference term b

$$b^{\text{BSM}} \approx 0.34 g_S \epsilon_S - 5.22 g_T \epsilon_T$$

$$b_V^{\text{BSM}} \approx 0.44 g_S \epsilon_S - 4.85 g_T \epsilon_T$$

- Extracting bounds on new-physics effective couplings ϵ_T and ϵ_S requires knowledge of g_T , g_S at the $\approx 10\%$ level
- To fully utilise the potential of future experimental neutron physics programs, the scalar and tensor charge must be precisely calculated

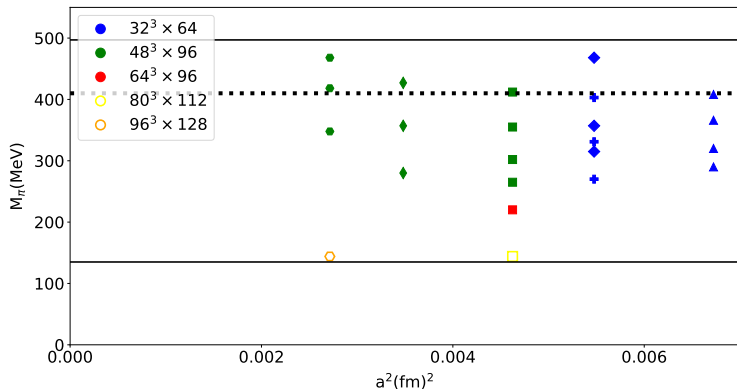
Motivation

- The neutron electric dipole moment (EDM) is sensitive to CP violation, hence is an excellent probe in the search for physics beyond the Standard Model
- The quark EDM contributions to the neutron EDM, d_n , are given by:

$$d_n = d_u g_T^d + d_d g_T^u + d_s g_T^s$$

- g_T^s consistent with 0
- d_u, d_d, d_s contain new CP violating interactions at the TeV scale
- knowing g_T^q put bound on d_n , constrain d_q and BSM theories

Analysis



Clover action comprises the tree-level Symanzik improved gluon action together with a stout smeared fermion action, modified for the use of the FH technique.

Calculating Matrix Elements using The Feynman-Hellmann Theorem

- Relates the derivative of the total energy to the expectation value of the derivative of the action

$$\frac{\partial E_{X,\lambda}(\vec{k})}{\partial \lambda} = \frac{1}{2E_{X,\lambda}(\vec{k})} \langle X, \vec{k} | \frac{\partial S}{\partial \lambda} | X, \vec{k} \rangle_{\lambda},$$

- In lattice calculations, modify the action

$$S \rightarrow S + \lambda \mathcal{O}$$

- Examine the behaviour of hadron energies as λ changes
- Extract the matrix element
- Hadron energies are extracted from two-point functions
- Control of excited state contamination more simple than standard three-point analyses

Calculating Matrix Elements using The Feynman-Hellmann Theorem

- To calculate the tensor charge we have the modified action:

$$S \rightarrow S + \lambda \int d^4x \bar{q}(x) \gamma_5 \sigma_{\mu\nu} q(x),$$

- Applying the Feynman- Hellmann Theorem we get:
Spin-up Spin-down

$$\left. \frac{\partial E_\lambda}{\partial \lambda} \right|_{\lambda=0} = g_T^q \qquad \left. \frac{\partial E_\lambda}{\partial \lambda} \right|_{\lambda=0} = -g_T^q$$

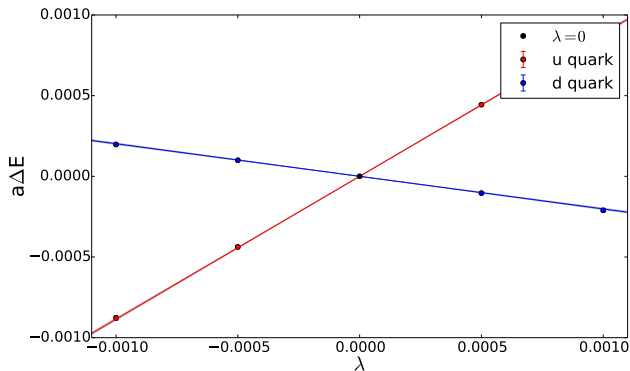
- The energy as a function of λ is therefore given by

$$E(\lambda) = E(0) \pm \lambda g_T^q + \mathcal{O}(\lambda^2).$$

- Nucleon Isovector Charges are then:

$$g_{\mathcal{O}}^{u-d} = g_{\mathcal{O}}^u - g_{\mathcal{O}}^d$$

Results



Calculated at $a = 0.082\text{fm}$, $(\kappa_l, \kappa_s) = (0.119930, 0.119930)$

$$g_T^u = 0.8529(53), \quad g_T^d = -0.1947(42)$$

Weighted Averaging Method

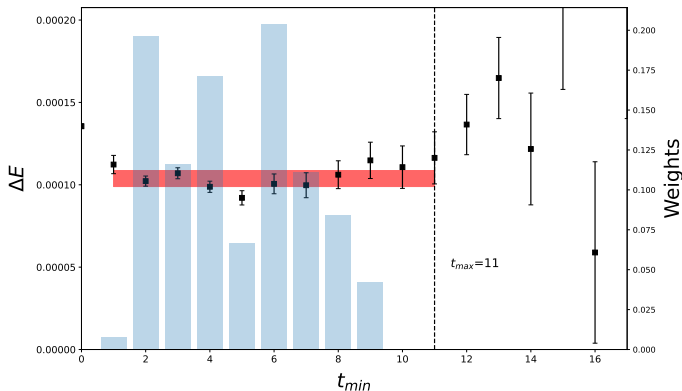
- The dependency of the fits on the time ranges used causes systematic uncertainties
- Time ranges are systematically sampled and each result is assigned a weight:

$$\tilde{w}^f = \frac{p_f (\delta E^f)^{-2}}{\sum_{f'=1}^N p_{f'} (\delta E^{f'})^{-2}},$$

- f labels choice of fit range specified by t_{\min}
- $p_f = \Gamma(N_{\text{dof}}/2, \chi^2/2)/\Gamma(N_{\text{dof}}/2)$ is the p-value of the fit
- δE^f is the uncertainty in the value of the change in energy of the hadron state.
- Taking a weighted average of the N fit findings provides the final unbiased estimate of the energy shift

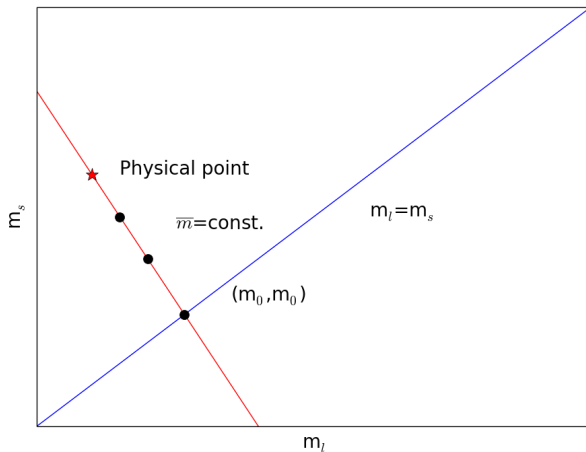
(NPLQCD/QCDSF, Phys. Rev. D, 2021)

Weighted Averaging Method



Energy shift for the down quark at $\lambda = 0.0005$, spin-up. Calculated at $a = 0.082fm$, $(\kappa_l, \kappa_s) = (0.119930, 0.119930)$

Flavour Symmetry Breaking



Flavour Symmetry Breaking

- $SU(3)$ unbroken \Rightarrow operator octet expressed as two couplings (f, d)
- $SU(3)$ broken \Rightarrow construct quantities (D_i, F_i) Example:

$$\begin{aligned}D_1 &= -(A_{\bar{N}\eta N} + A_{\Xi\eta\Xi}) \\ &= -\left(\frac{1}{\sqrt{6}}g_p^u + g_p^d\right) + \frac{1}{\sqrt{6}}(g_{\Xi}^u - 2g_{\Xi}^s)\end{aligned}$$

- Fan plots with slope parameters (r_i, s_i) constraining them
- Calculate each quantity (D_i, F_i) for each quark mass calculated on the lattice

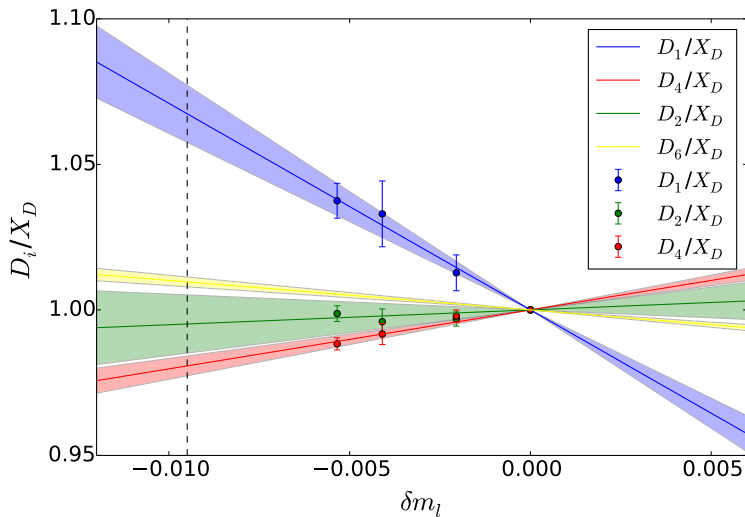
$$D_1 = 2d - 2r_1\delta m_l$$

$$D_2 = 2d + (r_1 + 2\sqrt{3}r_3)\delta m_l$$

$$\text{where } \delta m_l = m_l - \bar{m}$$

- An 'average D' can also be constructed from the diagonal amplitudes $X_D = \frac{1}{6}(D_1 + 2D_2 + 3D_4) = 2d + \mathcal{O}(\delta m_l^2)$

Fan plots



$a = 0.082 \text{ fm}$ using 4 pion masses

Flavour Symmetry Breaking

- Calculate Flavour diagonal matrix elements and hence the nucleon isovector charges through:

$$\langle p | \bar{u} \Gamma u | p \rangle = 2\sqrt{2}f + \left(\sqrt{\frac{3}{2}}r_1 - \sqrt{2}r_3 + \sqrt{2}s_1 - \sqrt{\frac{3}{2}}s_2 \right) \delta m_j^*$$

$$\langle p | \bar{d} \Gamma d | p \rangle = \sqrt{2}(f - \sqrt{3}d) + \left(\sqrt{\frac{3}{2}}r_1 - \sqrt{2}r_3 - \sqrt{2}s_1 - \sqrt{\frac{3}{2}}s_2 \right) \delta m_j^*$$

Where r_1 , r_3 , s_1 and s_2 are the slopes of the fan plots, and:

$$g_O^{u-d} = \langle p | \bar{u} \Gamma u | p \rangle - \langle p | \bar{d} \Gamma d | p \rangle .$$

and using δm_j^* at physical point

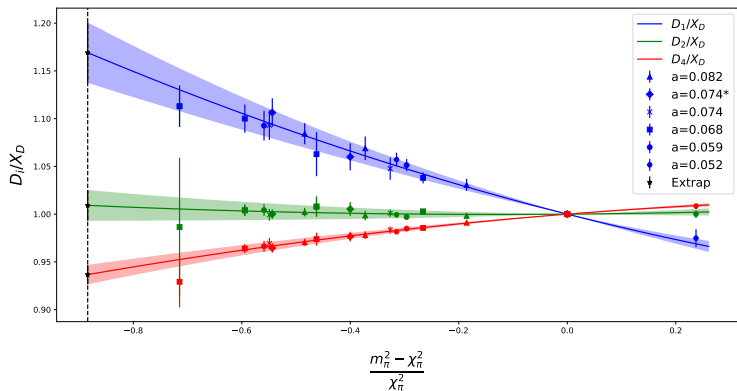
- Flavour breaking expansion method only accounts for the quark mass-dependence of the matrix elements
- In order to quantify systematic uncertainties, we extend this theory to also account for lattice spacing, a , finite volume effects, mL , and second order mass terms, δm_l^2
- Perform a combination of global fits over all ensembles

$$X'_{D,F} = X_{D,F} \left(1 + c_1 \frac{1}{3} [f_L(m_\pi) + 2f_L(m_K)] \right) + c_2 a^2 + c_3 \delta m_l^2,$$

$$\text{where: } f_L(m) = (am)^2 \frac{e^{-mL}}{X_N L},$$

$$D_1 = 2d - 2(r'_1 + b_1 a^2) \delta m_l + d_1 \delta m_l^2.$$

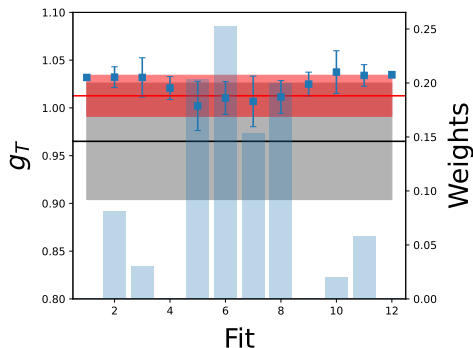
Systematics



a^2 and δm_l^2 correction included. All points shifted in the limit $a \rightarrow 0$

In order to combine all fits we extend our weighted averaging method

Results

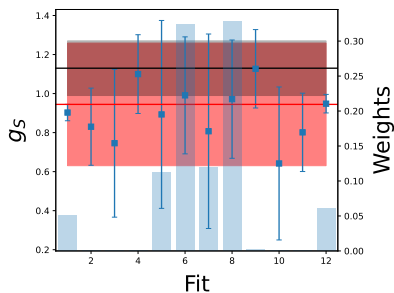


- $g_T = 1.013(21)$
- Black band Flag Review 2021 (arXiv:2111.09849)

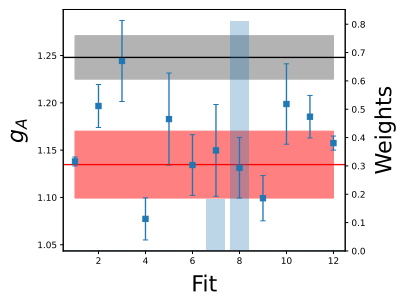
x-axis

1. No Corrections
2. a^2
3. a
4. δm_l^2
5. $a, \delta m_l^2$
6. $a^2, \delta m_l^2$
7. $a, \delta m_l^2, mL$
8. $a^2, \delta m_l^2, mL$
9. $\delta m_l^2, mL$
10. a, mL
11. a^2, mL
12. mL

Results



$$g_s = 0.94(31)$$



$$g_A = 1.135(35)$$

Renormalised at $\mu = 2\text{GeV}$ in the $\overline{\text{MS}}$ scheme

Impact on Phenomenology

- Following the method presented in Bhattacharya et al., PRD, 2012
- Green band is the existing band from $0^+ \rightarrow 0^+$ nuclear beta decay
- $g_T = g_S = 1$
- No uncertainty

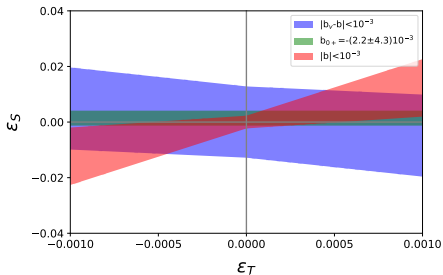
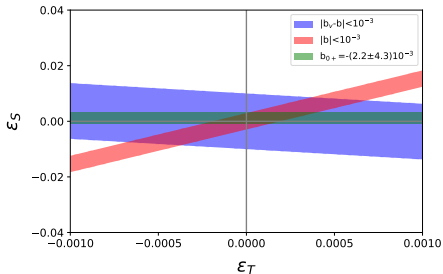
- Tensor and scalar charges using our results:

$$g_T = 1.013(21)$$

$$\delta g_T / g_T \approx 2\%$$

$$g_S = 0.94(31)$$

$$\delta g_S / g_S \approx 30\%$$



- Method presented we have:
 - Control of excited state contamination from FH method
 - Weighted Averaging Method- removing possible systematic uncertainties which arise from bias choices of fit windows
 - Quark mass- Flavour symmetry breaking to extrapolate towards the physical point
 - Extended flavour symmetry breaking method in account for lattice spacing, finite volume effects and second order mass terms, have control over these systematics
- Calculated g_T to $\approx 2\%$ level, successfully reaching the goal of understanding at the 10% level.
- Work is still needed on g_S and g_A to understand it at the same level.