

Pion pole contribution to HLbL at the physical point

Gurtej Kanwar (University of Bern)

In collaboration with Sebastian Burri,
Marcus Petschlies, Urs Wenger + ETMC

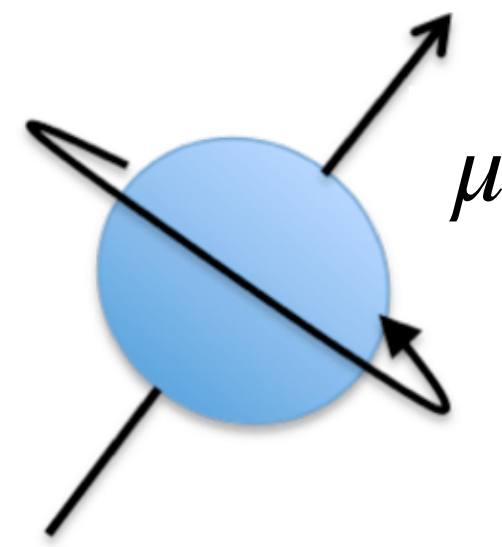


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**UNIVERSITÄT
BERN**

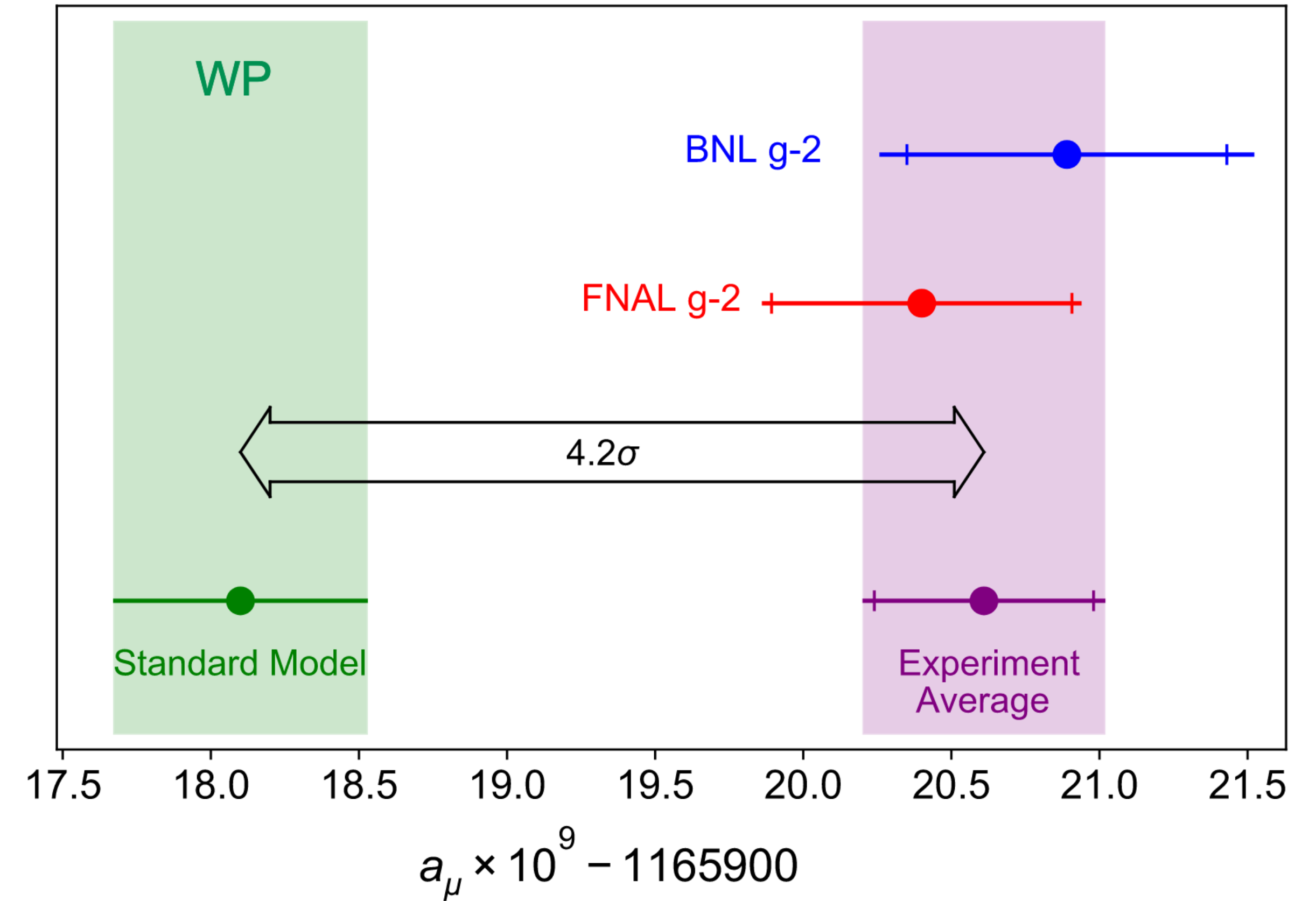
Lattice 2022 (Bonn, Germany)
Aug 11, 2022

Tensions in a_μ

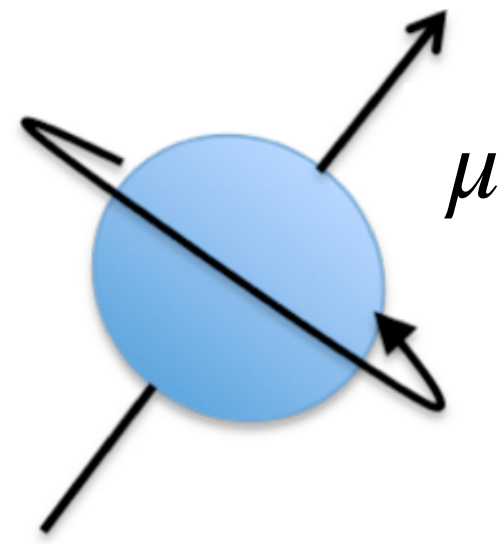


Combined 4.2σ tension

$$\Delta a_\mu = 279(76) \times 10^{-11}$$



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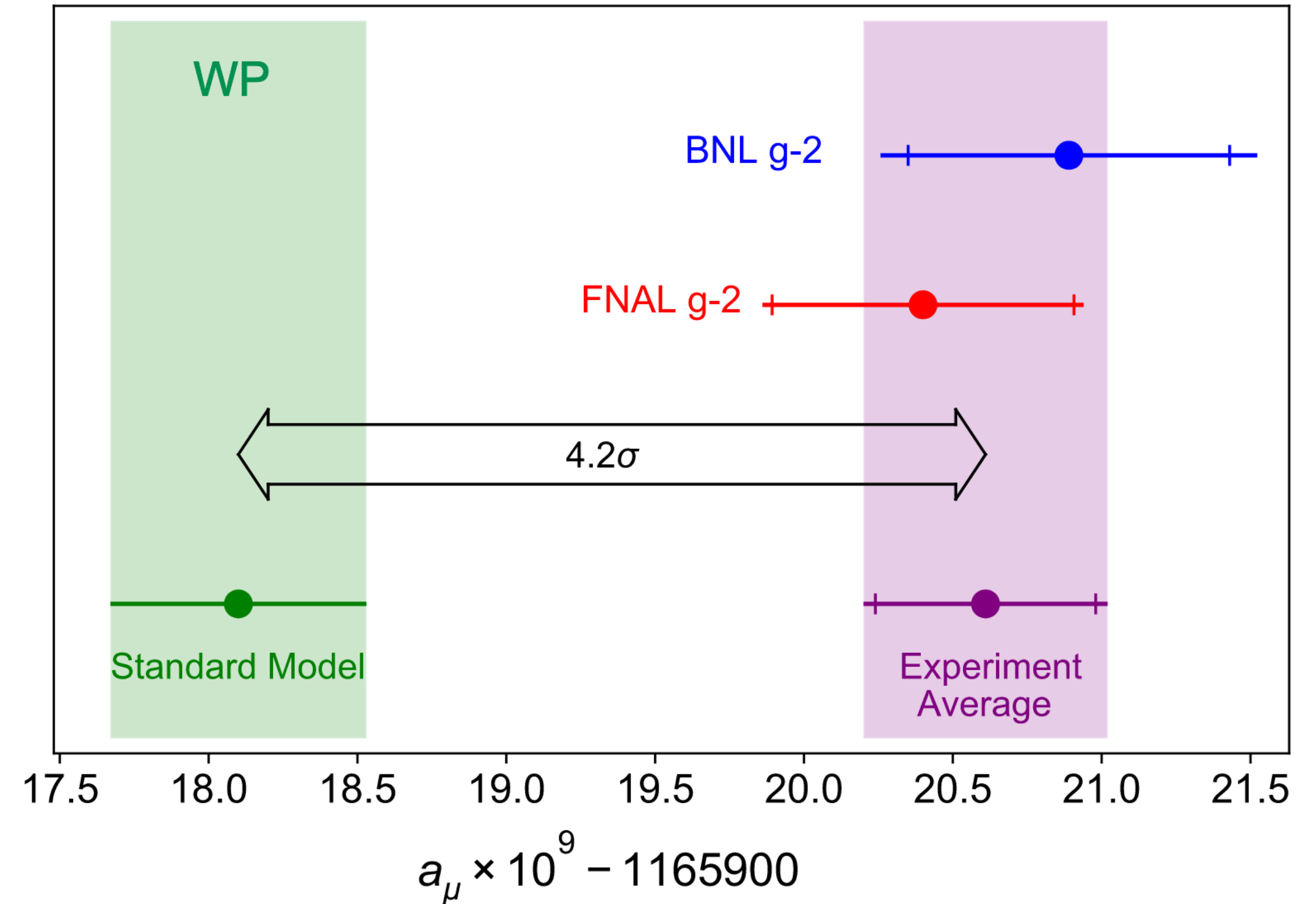


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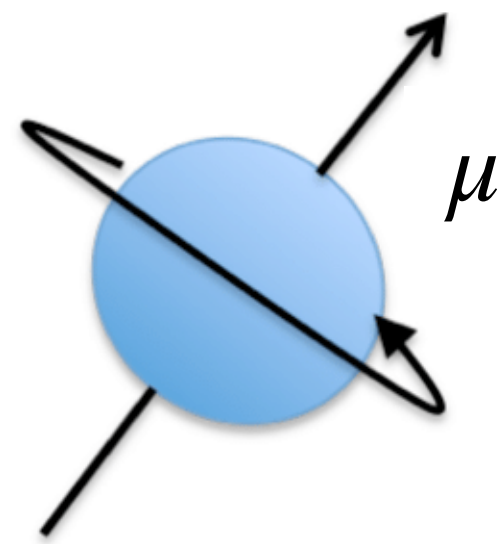
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Expt: BNL E821 + Fermilab E989

- Current err $\sim 60 \times 10^{-11}$
- Expected err $\sim 15 \times 10^{-11}$



Tensions in a_μ



Combined 4.2σ tension

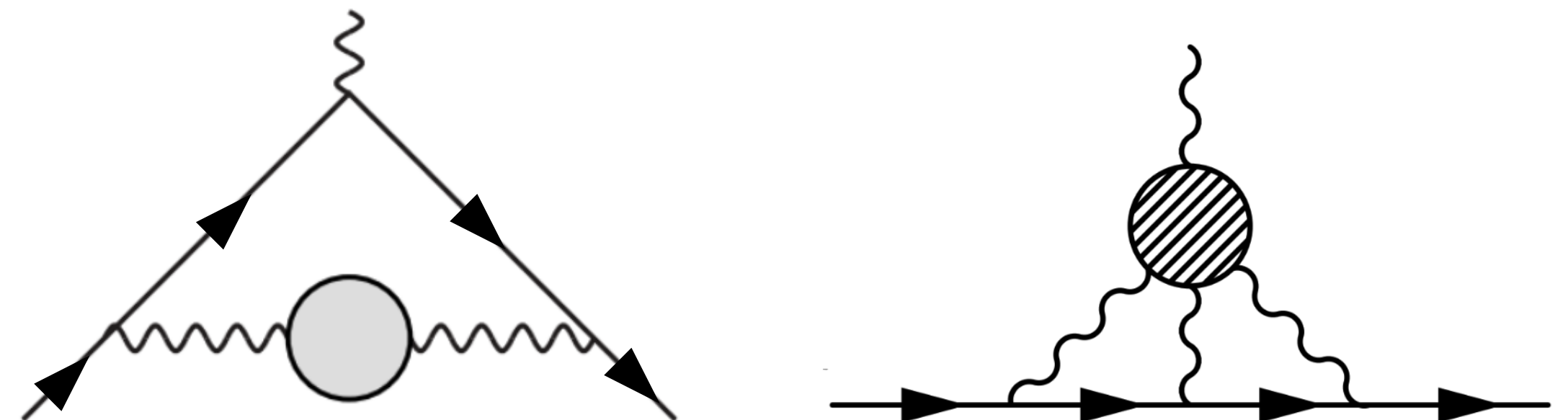
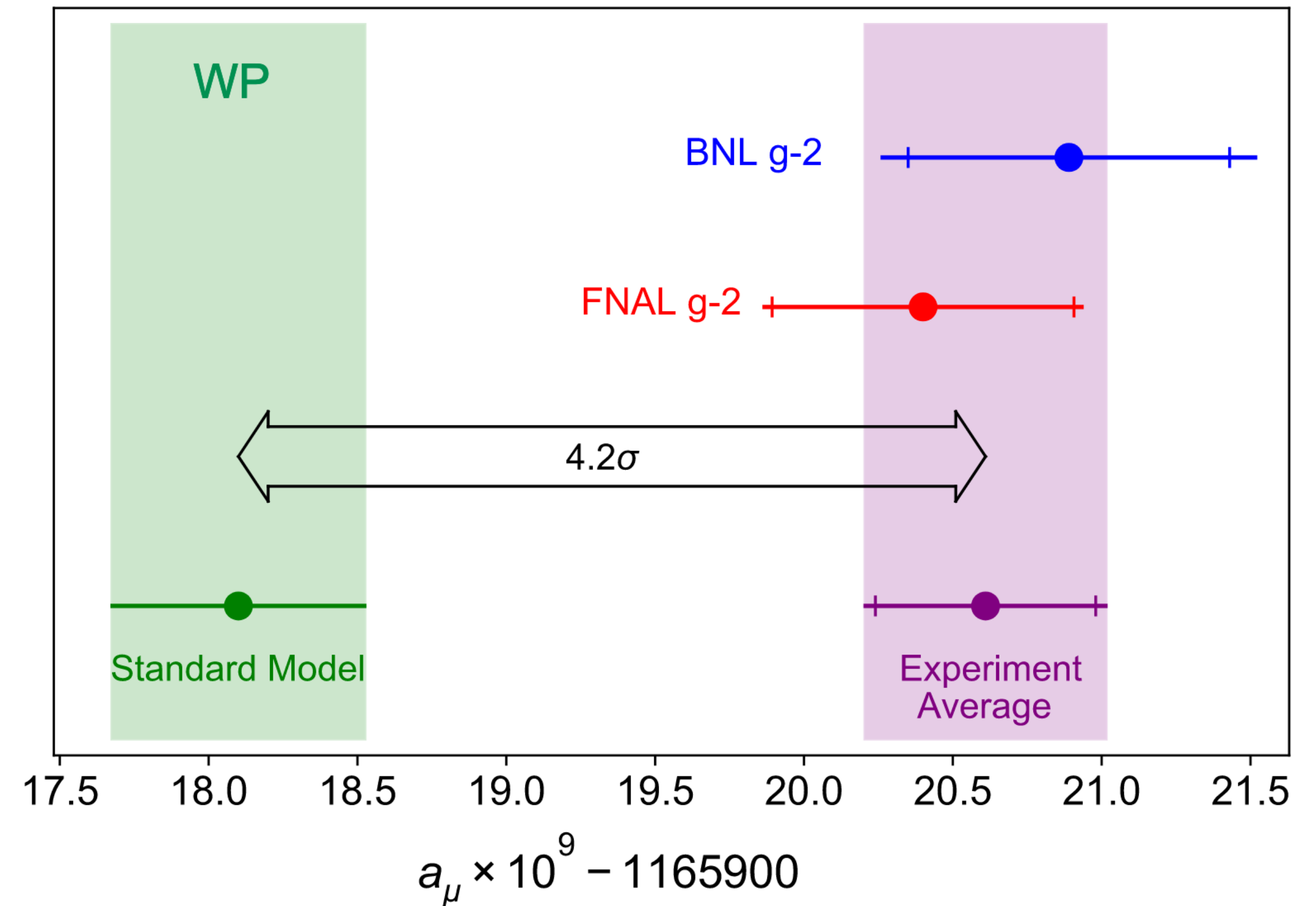
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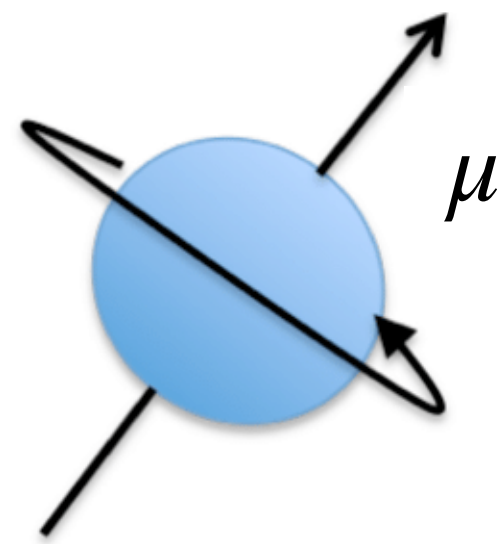
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Theory: Perturbative + pheno + latt

- HVP err $\sim 40 \times 10^{-11}$
HLbL err $\sim 20 \times 10^{-11}$



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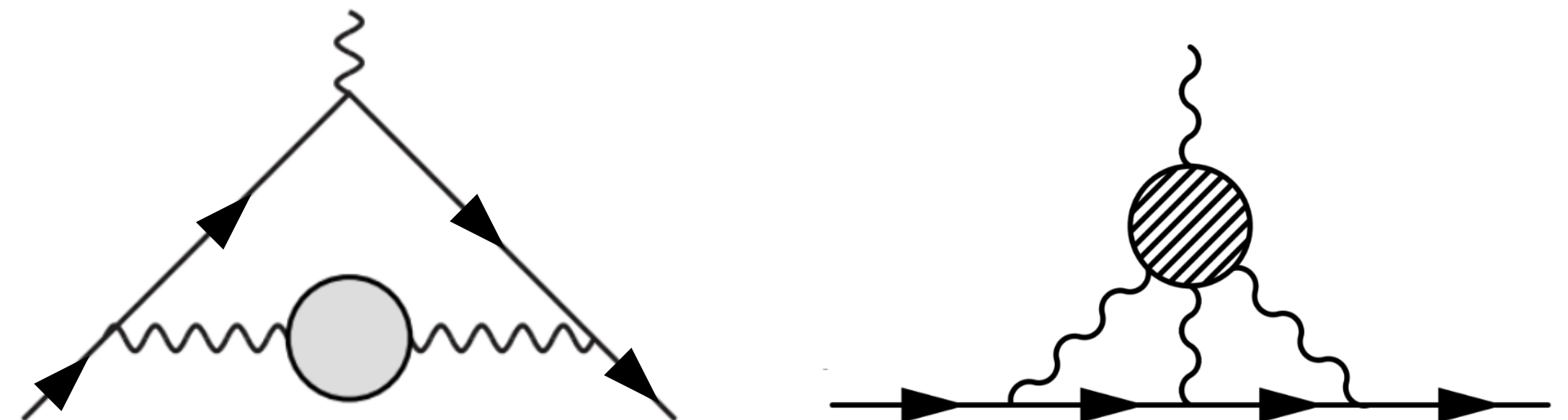
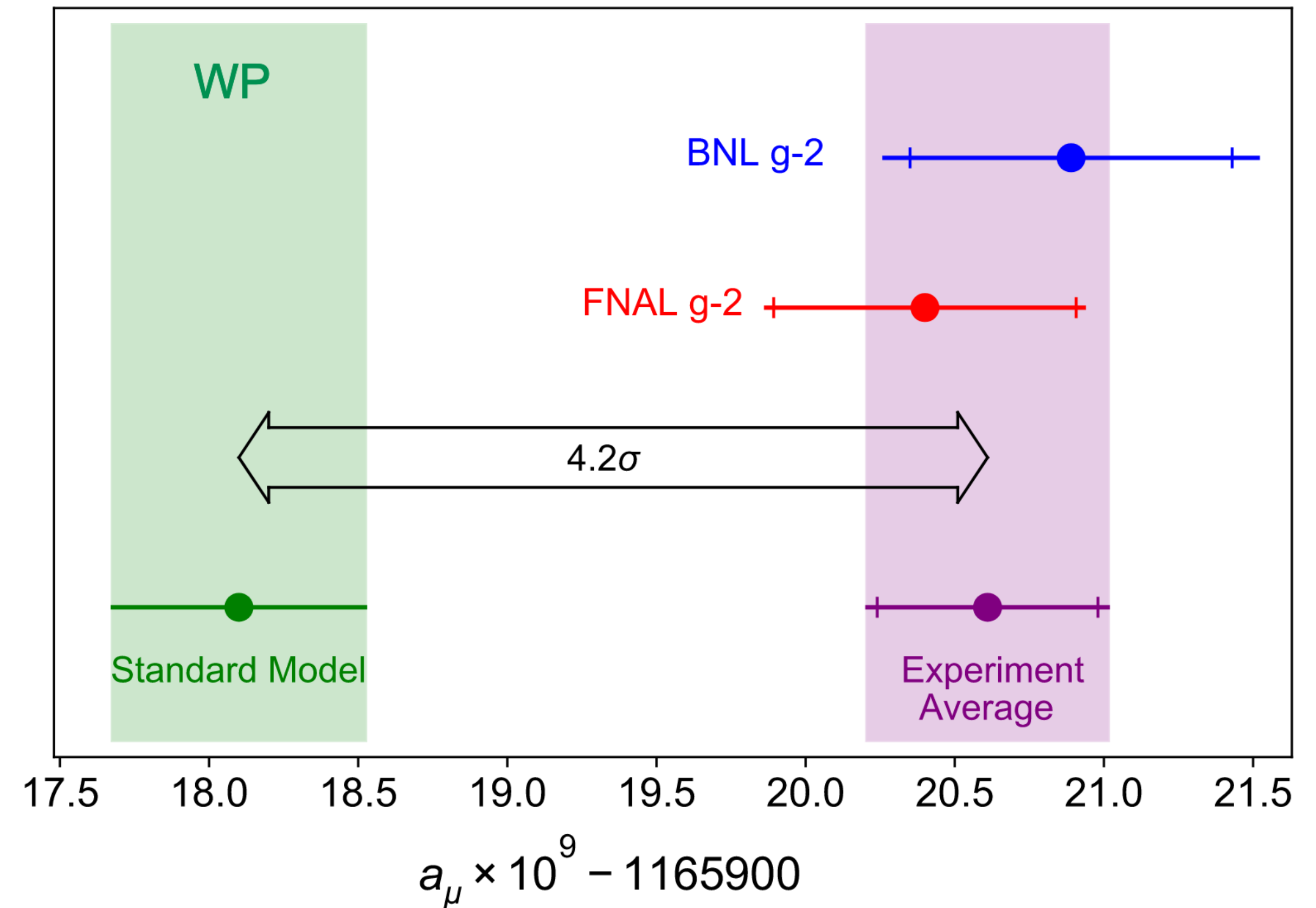
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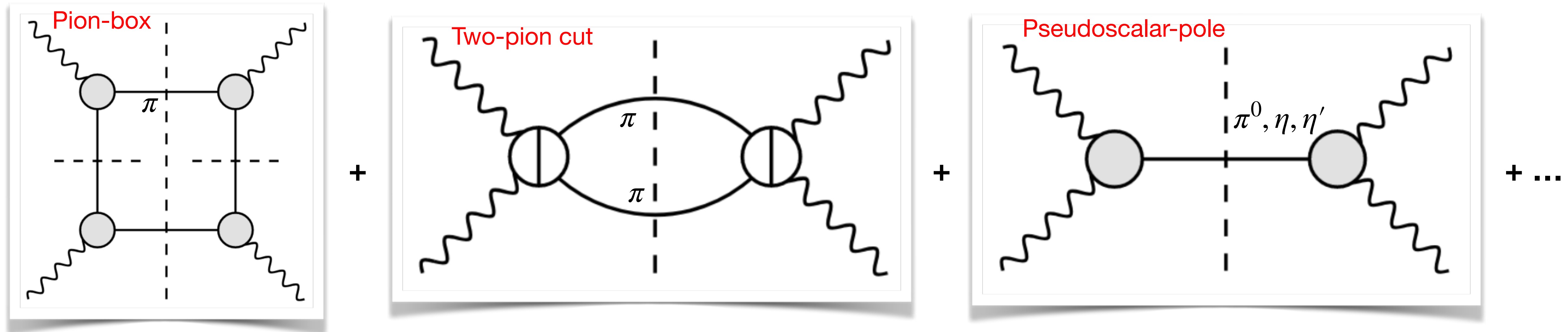
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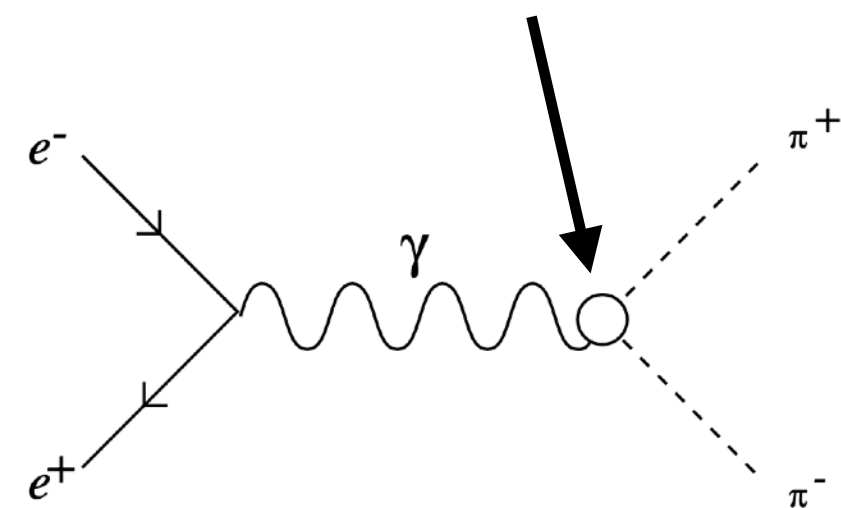
At the order of
E989 projected err



HLbL decomposition



$F_{\gamma^* \rightarrow \pi\pi}$ "VFF"
well-determined
experimentally ✓



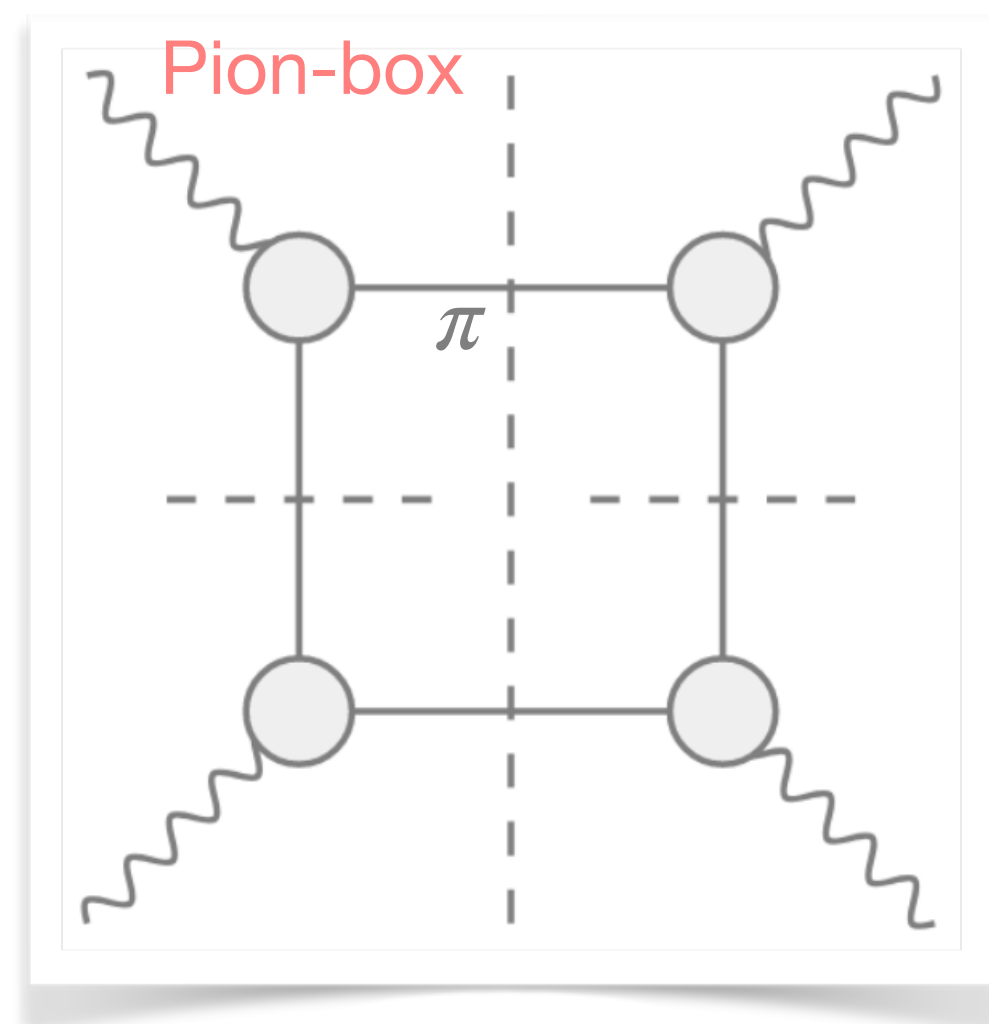
Partial waves for
 $\gamma^* \gamma^* \rightarrow \pi\pi$ required

Result is small and
well-constrained ✓

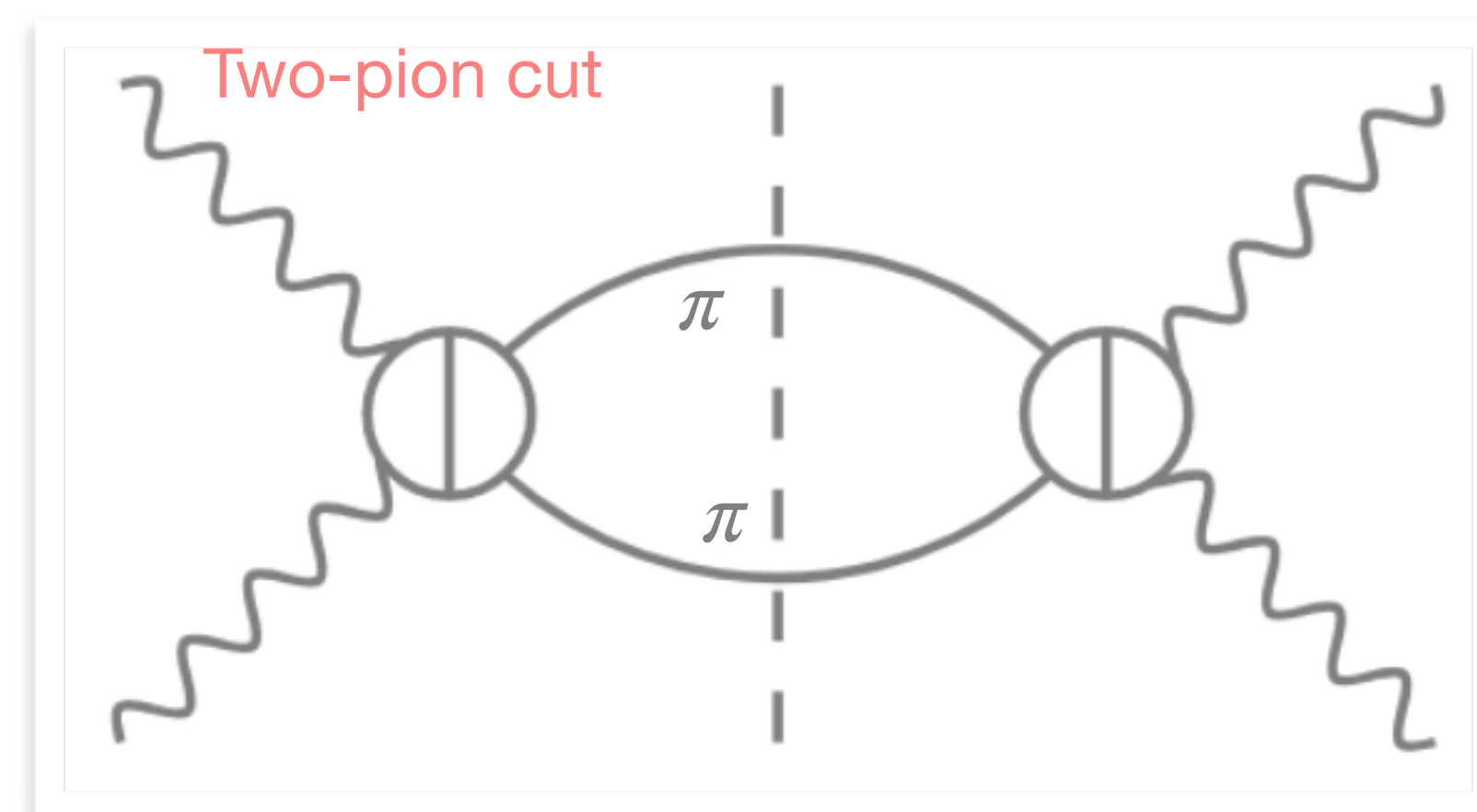
$F_{P \rightarrow \gamma^* \gamma}$ singly-virtual
"TFF" experimentally
accessible

$F_{P \rightarrow \gamma^* \gamma^*}$ doubly-virtual
"TFF" ~**unconstrained**

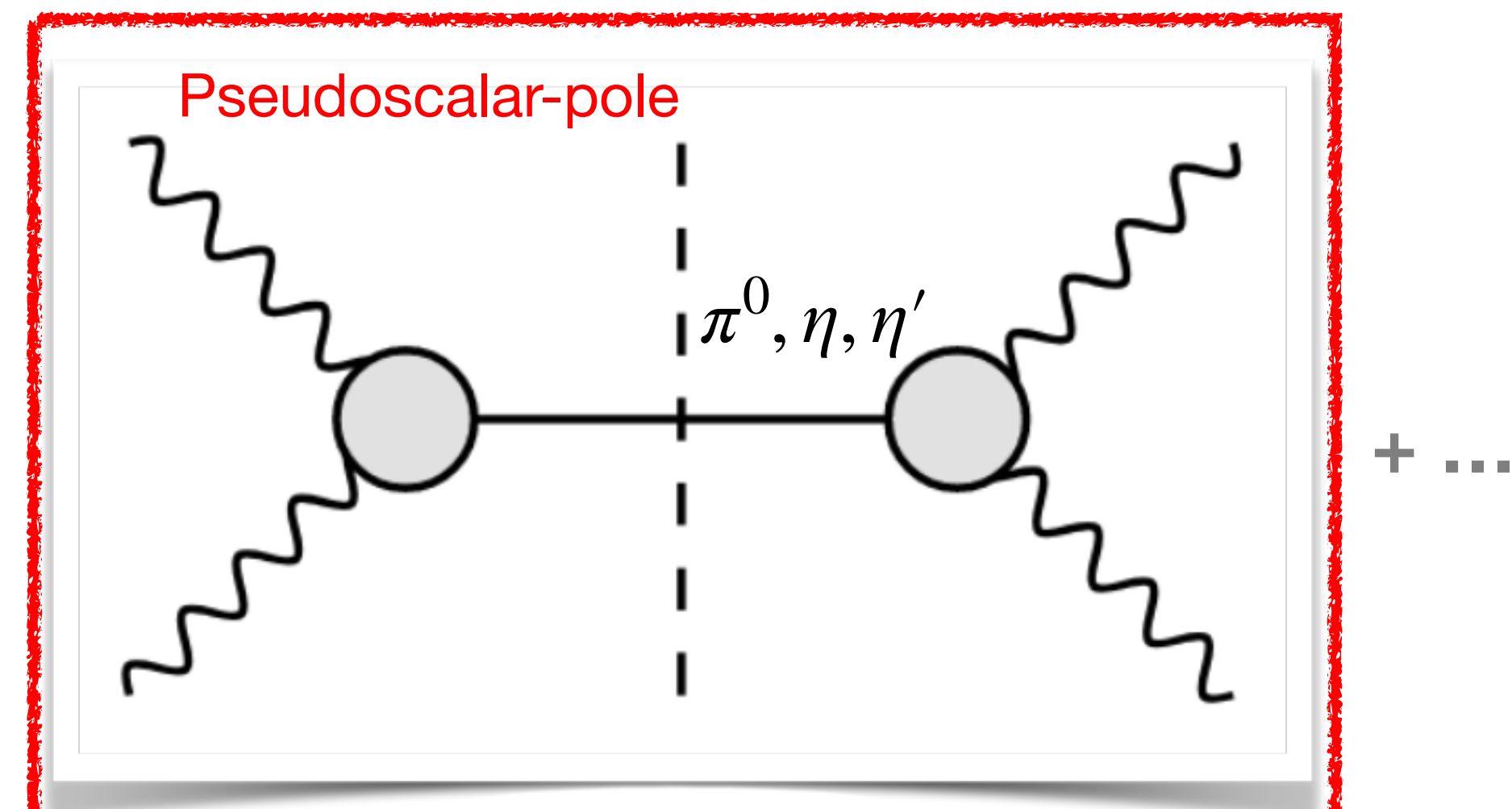
HLbL decomposition



+

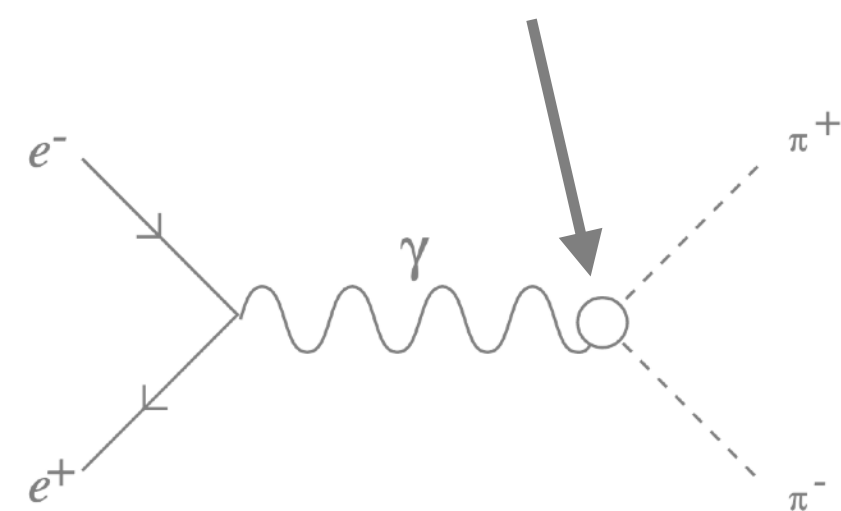


+



+ ...

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THIS TALK

Pseudoscalar TFFs



A Feynman diagram representing the transition form factor $F_{P\gamma\gamma}$. A central grey circle represents the meson P . An incoming horizontal line from the left represents the meson with momentum $p = (\sqrt{m_P^2 + \mathbf{p}^2}, \mathbf{p})$. Two outgoing wavy lines represent photons with momenta $q_1 = (\omega_1, \mathbf{q}_1)$ and $q_2 = (\omega_2, \mathbf{q}_2)$. The photons are labeled $\gamma^{(*)}$.

$$P \in \{\pi^0, \eta, \eta'\} \quad p = (\sqrt{m_P^2 + \mathbf{p}^2}, \mathbf{p}) \quad \gamma^{(*)} \quad q_1 = (\omega_1, \mathbf{q}_1) \quad \gamma^{(*)} \quad q_2 = (\omega_2, \mathbf{q}_2) \quad = F_{P\gamma\gamma}(q_1^2, q_2^2)$$

Singly-virtual, spacelike: $F_{P \rightarrow \gamma^* \gamma} = F_{P\gamma\gamma}(-Q^2, 0)$

Doubly-virtual, spacelike: $F_{P \rightarrow \gamma^* \gamma^*} = F_{P\gamma\gamma}(-Q_1^2, -Q_2^2)$

TFF kinematics

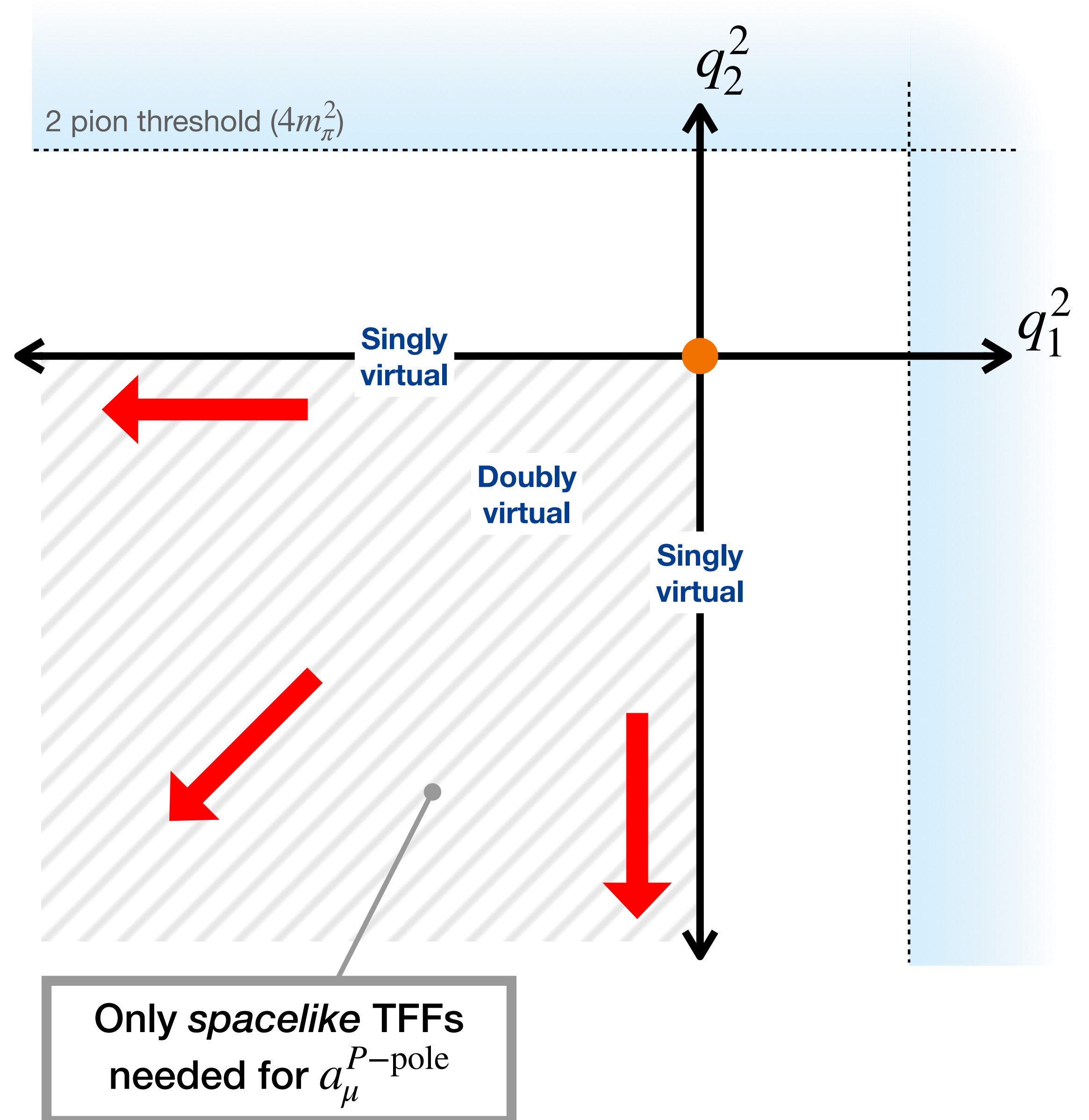
Related to decay width

$$\Gamma(P \rightarrow \gamma\gamma) = \left(\frac{\pi\alpha^2 M_P^3}{4} \right) |F_{P\gamma\gamma}(0,0)|^2$$

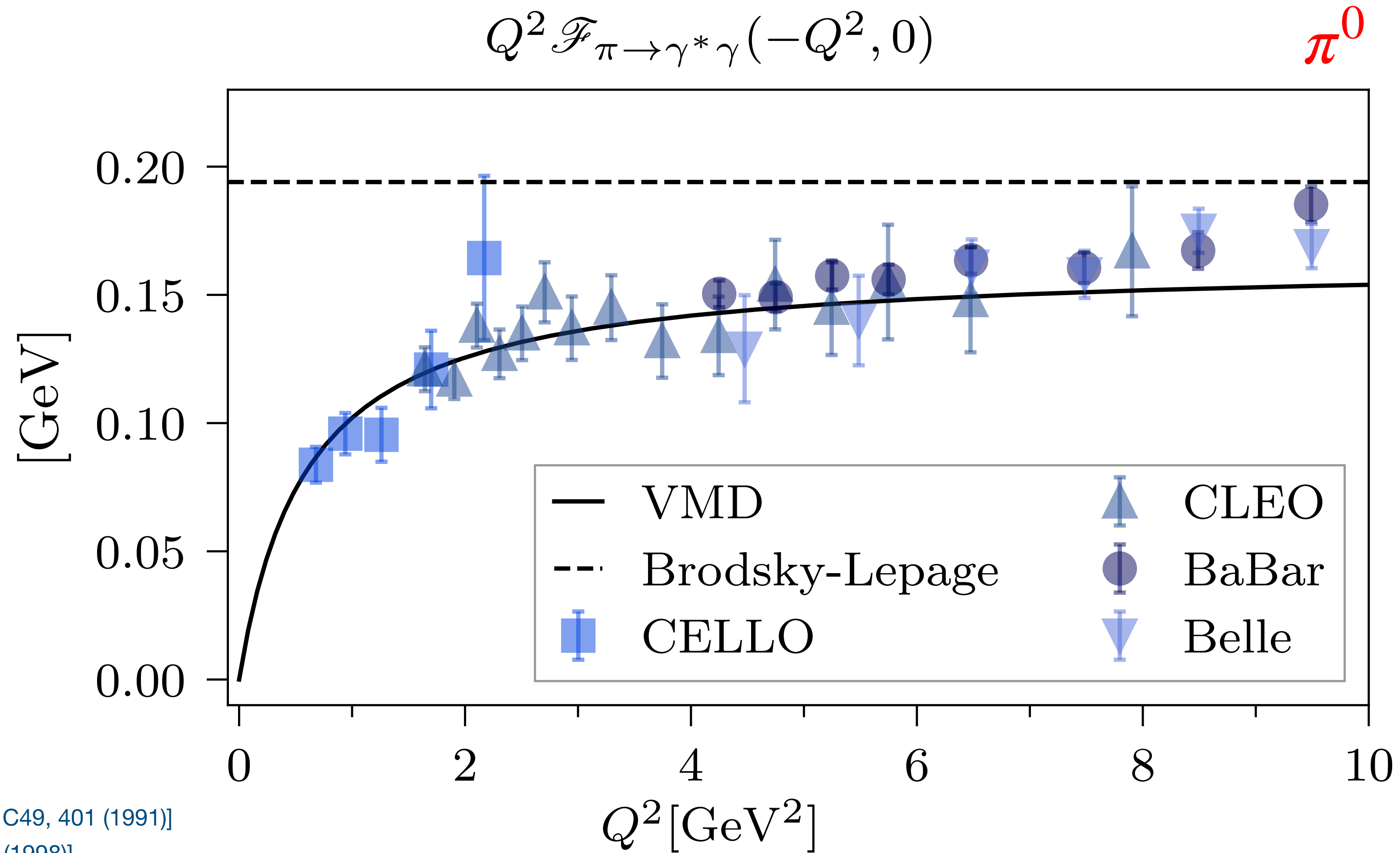
Large- Q^2 limits from pQCD

$$F_{P\gamma\gamma}(-Q^2, 0) \rightarrow \frac{2F_P}{Q^2} \text{ (Brodsky-Lepage)}$$

$$F_{P\gamma\gamma}(-Q^2, -Q^2) \rightarrow \frac{2F_P}{3Q^2} \text{ (OPE)}$$



TFF data: singly virtual



[H. J. Behrend+ (CELLO), Z. Phys. C49, 401 (1991)]

[J. Gronberg+ (CLEO), PRD 57, 33 (1998)]

[B. Aubert+ (BABAR), PRD 80, 052002 (2009)]

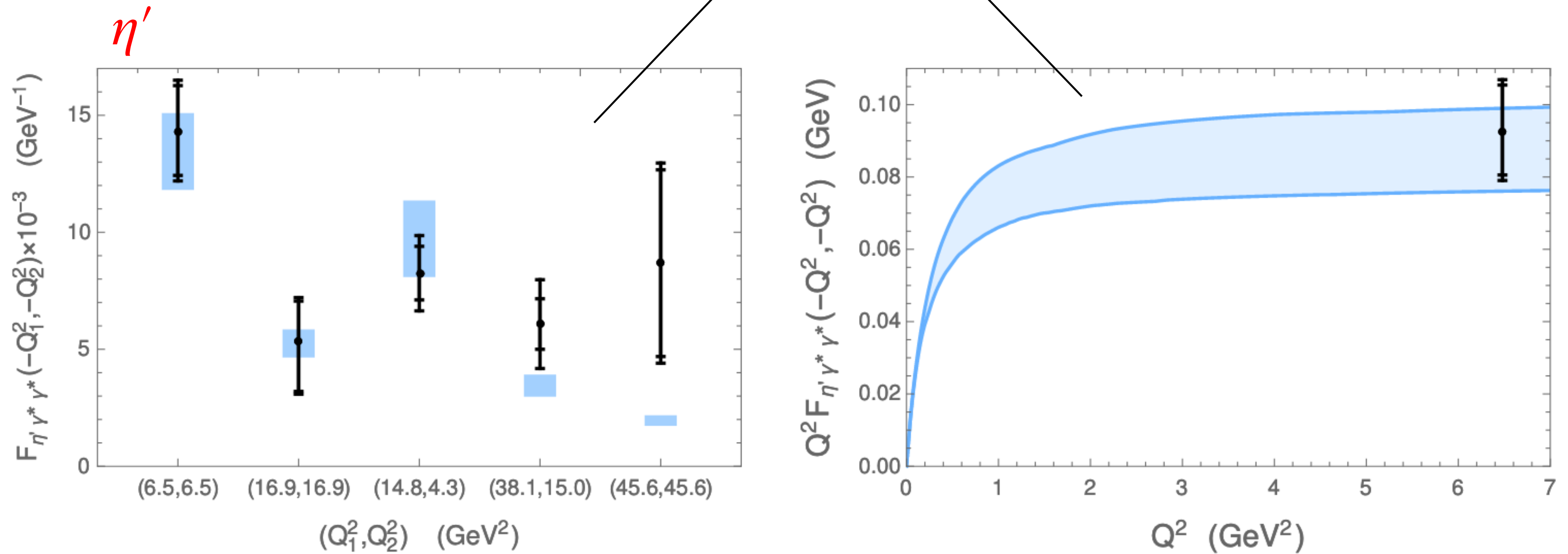
[P. del Amo Sanchez+ (BABAR), PRD 84, 052001 (2011)]

[S. Uehara+ (Belle), PRD 86, 092007 (2012)]

Note: Singly virtual for η, η' are very similar.

TFF data: doubly virtual

Based on ~50 events at BaBar !

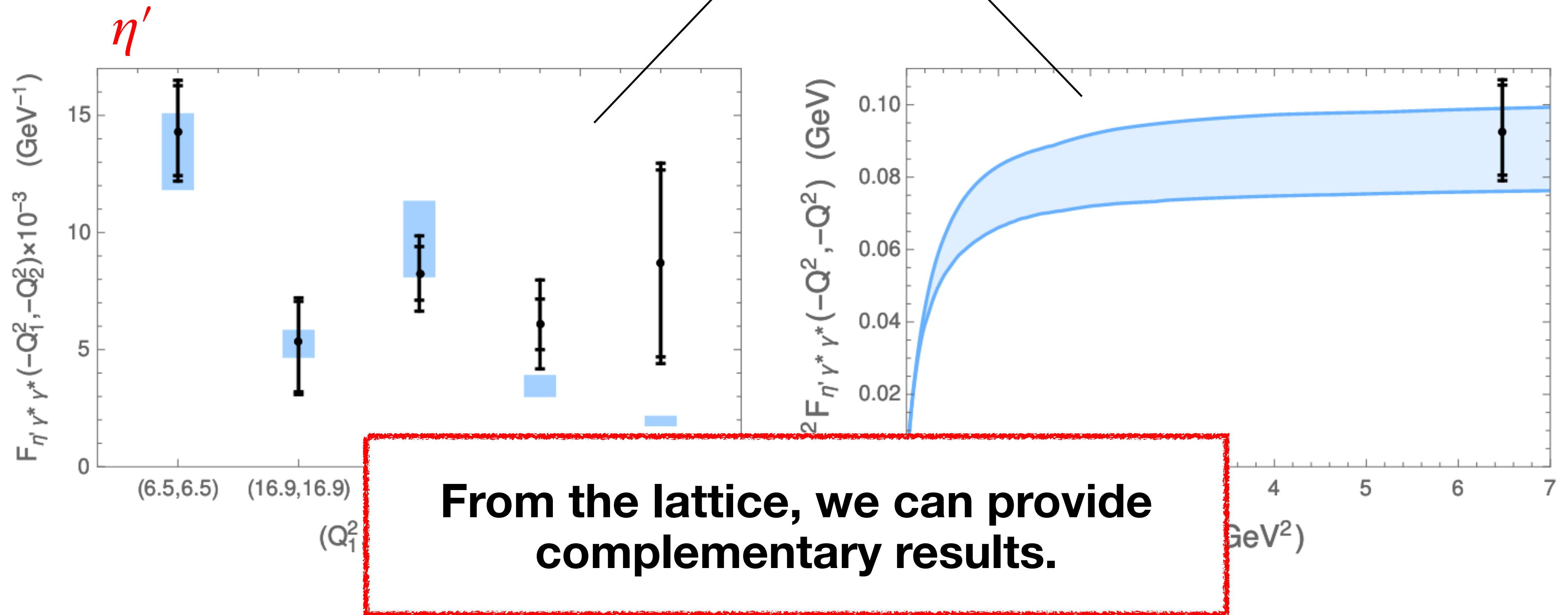


[J. P. Lees+ (BABAR), PRD 98, 112002 (2018)]

Figure from [g-2 WP, Phys. Rep. 887 (2020)]

TFF data: doubly virtual

Based on ~50 events at BaBar !



Lattice details

Action: Nf = 2+1+1 twisted clover, Iwasaki gauge action, *physical point*

Pion analysis:

- Three lattice spacings 
- Preliminary continuum  extrapolated result

ensemble	$L^3 \cdot T/a^4$	m_π [MeV]	a [fm]	L [fm]	$m_\pi \cdot L$
cB072.64	$64^3 \cdot 128$	136.8(6)	0.082	5.22	3.6
cC060.80	$80^3 \cdot 160$	134.2(5)	0.069	5.55	3.8
cD054.96	$96^3 \times 192$	140.8	0.057	5.46	3.9

[ETMC PRD 104, 074520 (2021)]

Eta analysis:

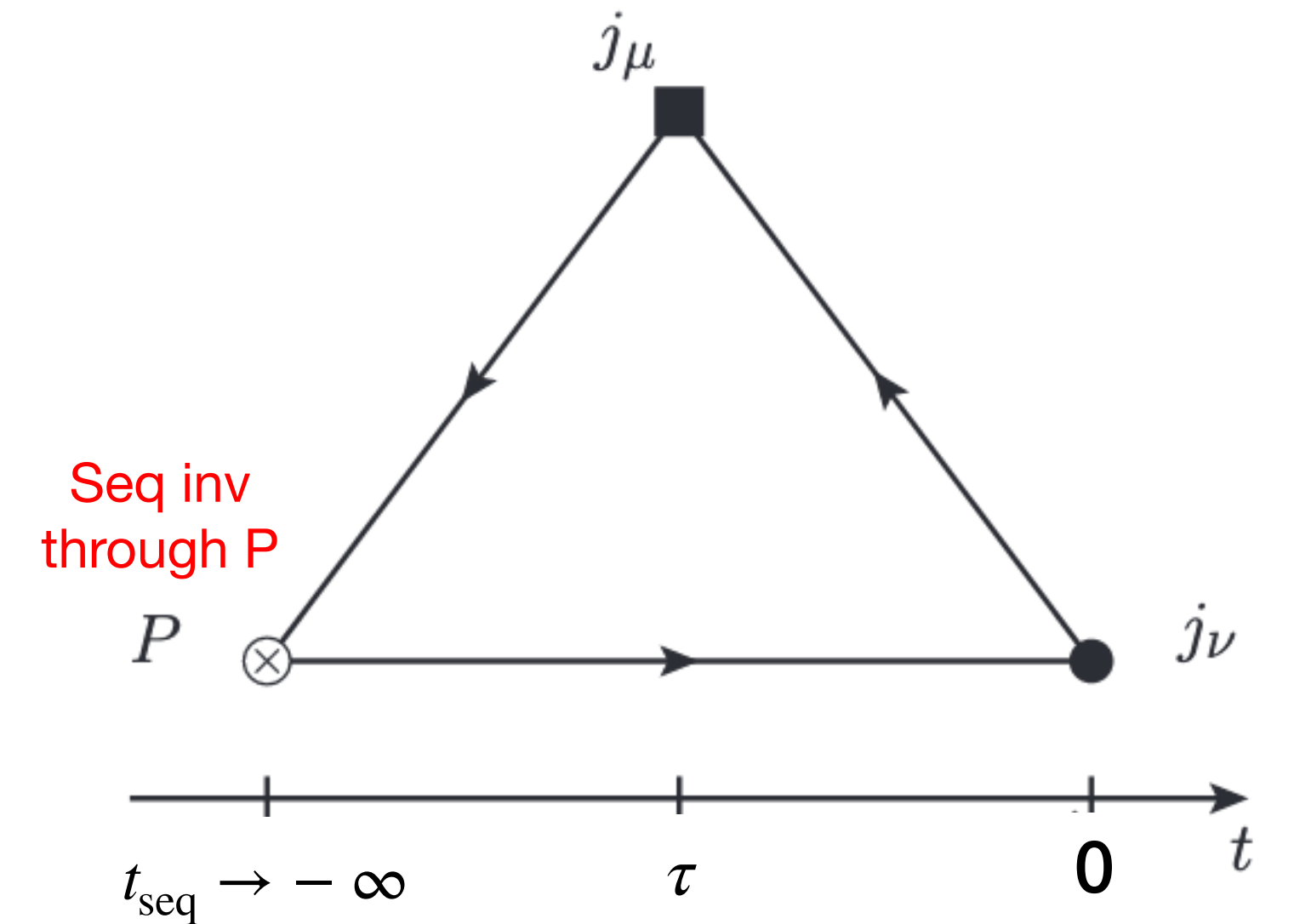
- See talk by Sebastian Burri, today 11:30

Lattice calculation of $F_{P\gamma\gamma}$

1. Euclidean time current-current matrix element

$$\tilde{A}_{\mu\nu}(\tau) = \sum_{\mathbf{x}} e^{-i\mathbf{q}_1 \cdot \mathbf{x}} \left\langle 0 \left| j_{\mu}(\tau; \mathbf{x}) j_{\nu}(0; \mathbf{0}) \right| P(\mathbf{p}) \right\rangle$$

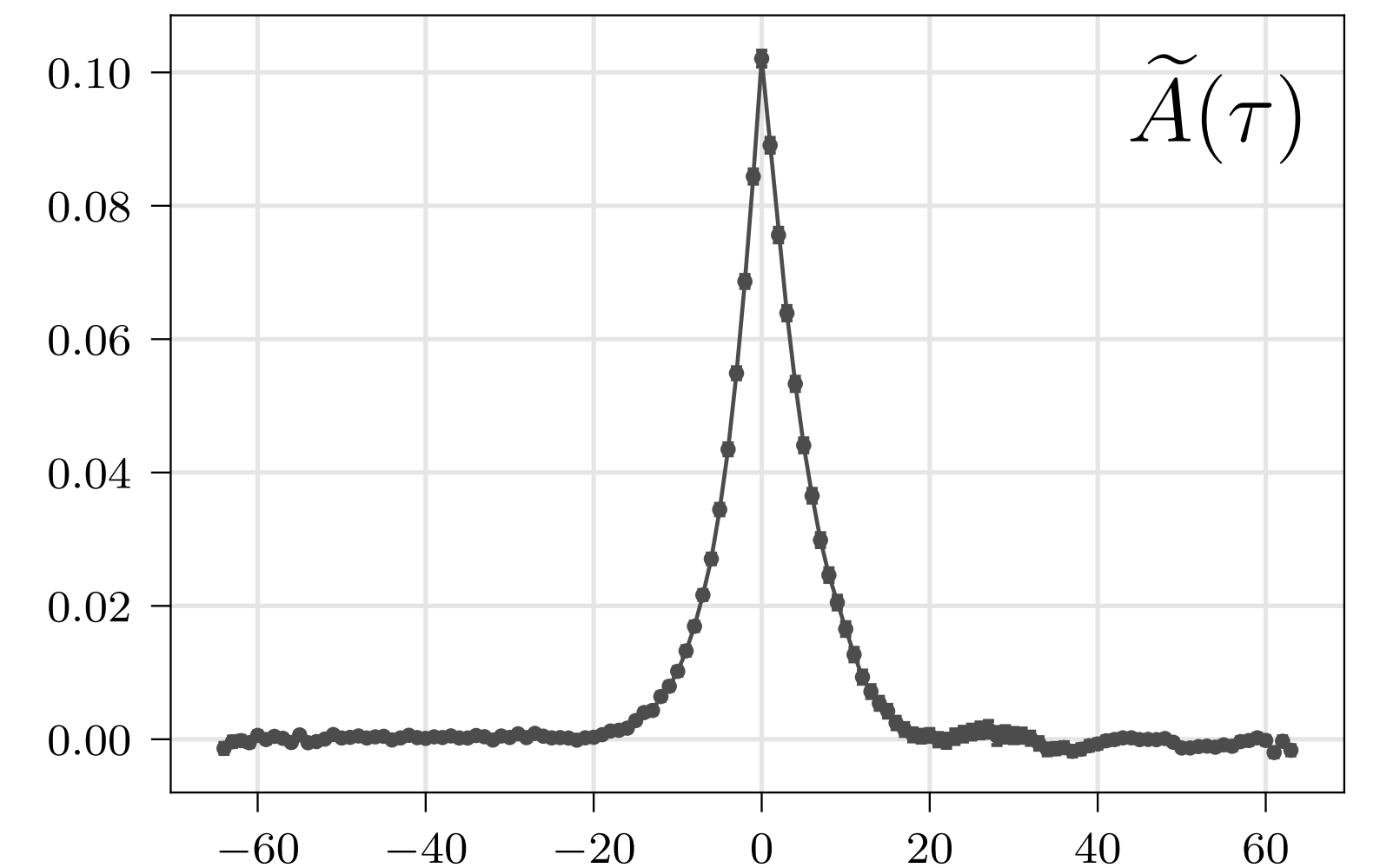
Note: renormalized currents
using precisely known Z_A, Z_V



2. Laplace transform

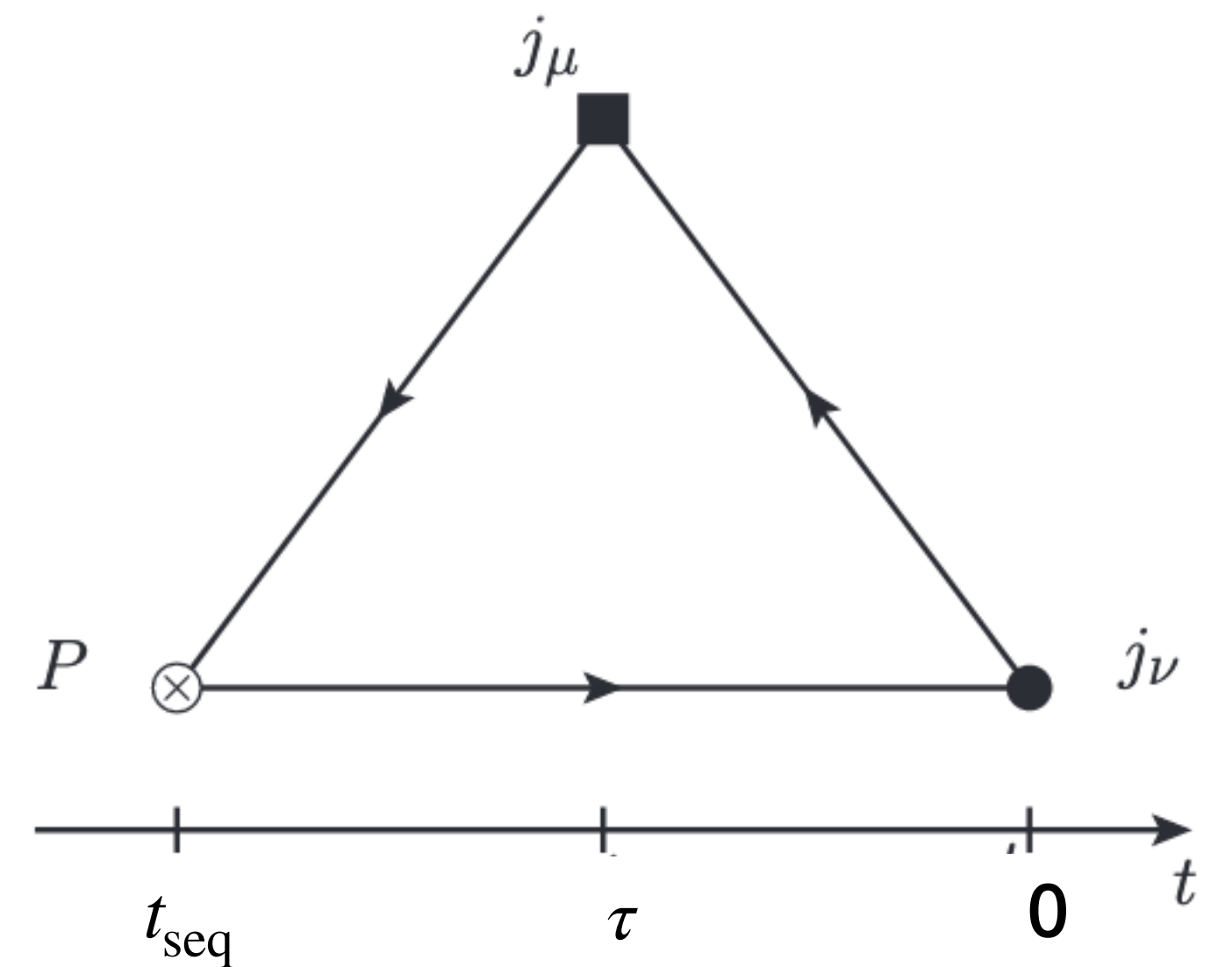
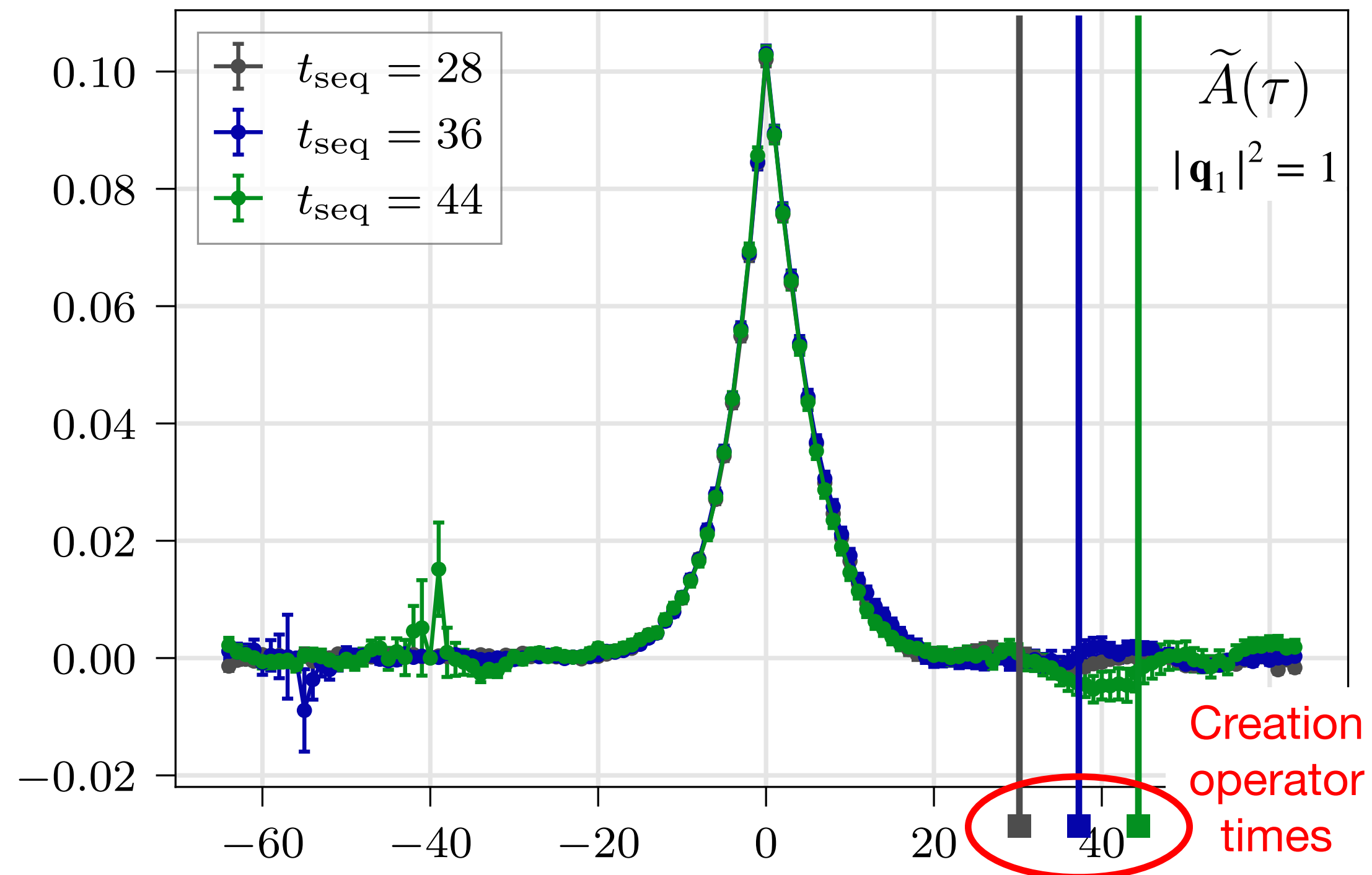
$$\epsilon^{\mu\nu\rho\sigma} q_{1\rho} q_{2\sigma} F_{P\gamma\gamma}(q_1^2, q_2^2) = -i^{n_0} \int_{-\infty}^{\infty} d\tau e^{\omega_1 \tau} \tilde{A}_{\mu\nu}(\tau)$$

where $\mathbf{q}_2 = \mathbf{p} - \mathbf{q}_1$ by momentum conservation



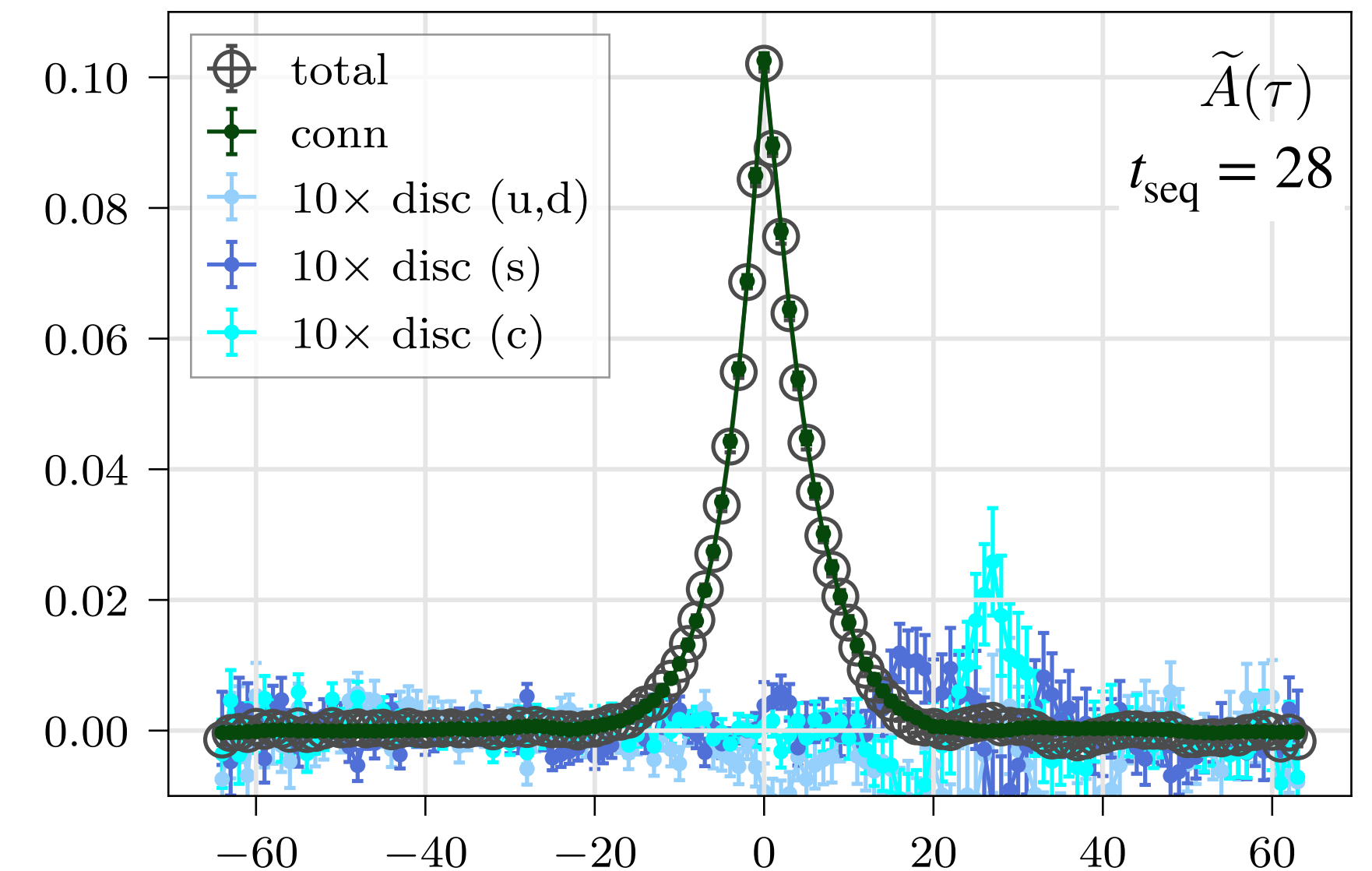
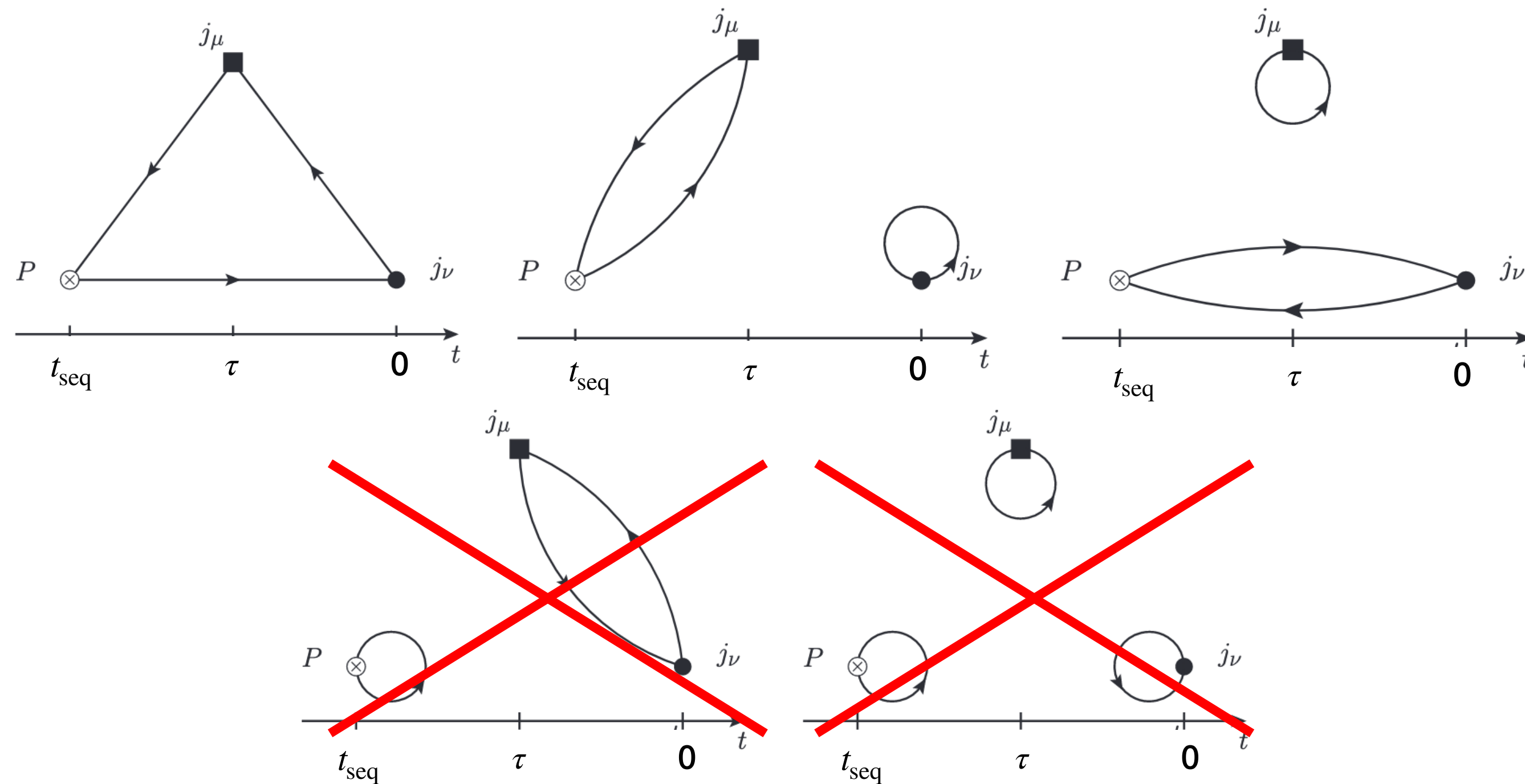
Matrix element from 3-pt function

$$\tilde{A}_{\mu\nu}(\tau) = \lim_{t_{\text{seq}} \rightarrow \infty} \frac{2E_P}{Z_P} e^{E_P t_{\text{seq}}} \sum_{\mathbf{x}, \mathbf{y}} e^{-i\mathbf{q}_1 \cdot \mathbf{x}} e^{i\mathbf{p} \cdot \mathbf{y}} \left\langle j_\mu(\tau, \mathbf{x}) j_\nu(0, \mathbf{0}) P^\dagger(-t_{\text{seq}}, \mathbf{y}) \right\rangle$$



Wick contractions: π^0 isospin rotation

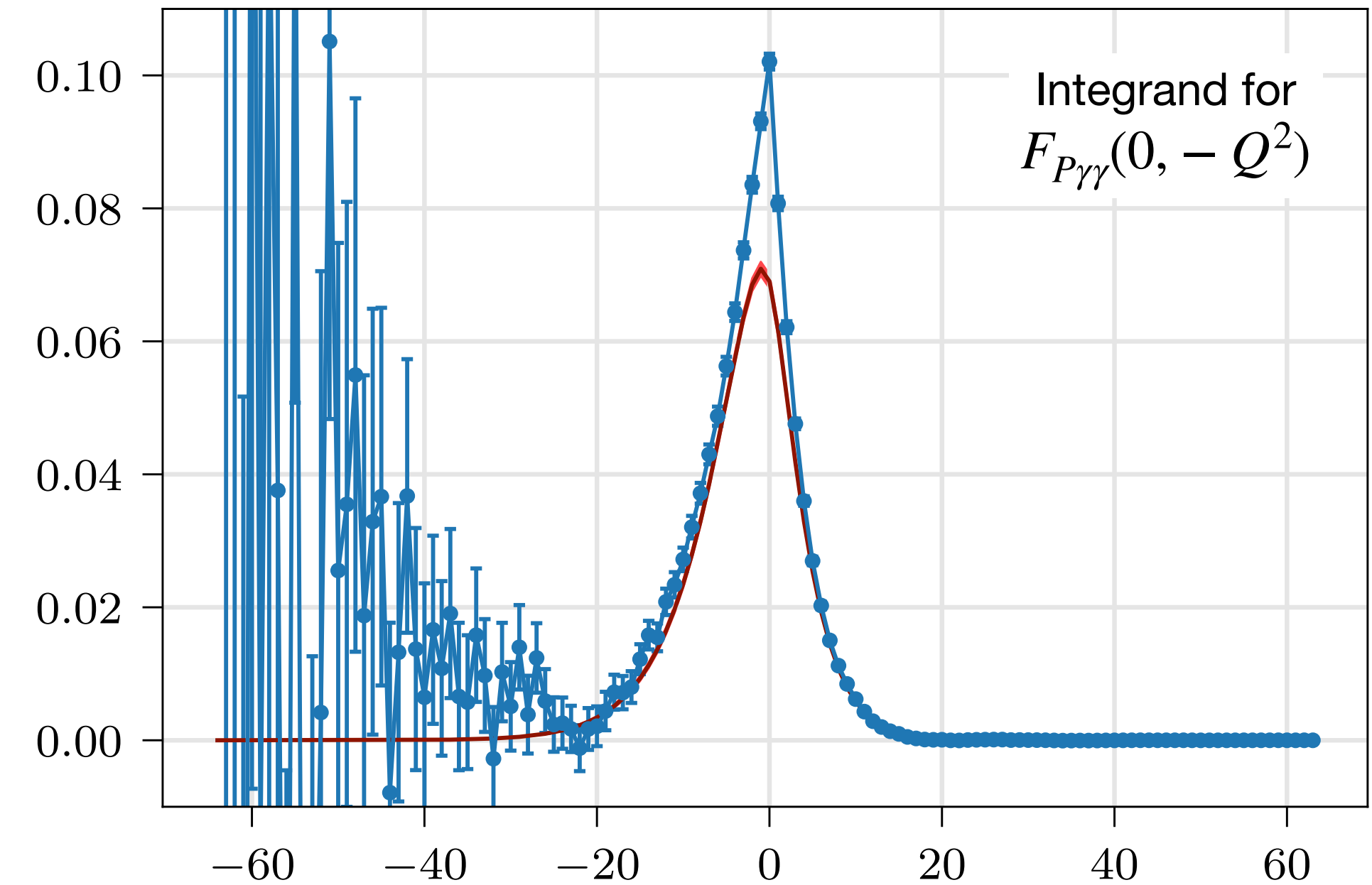
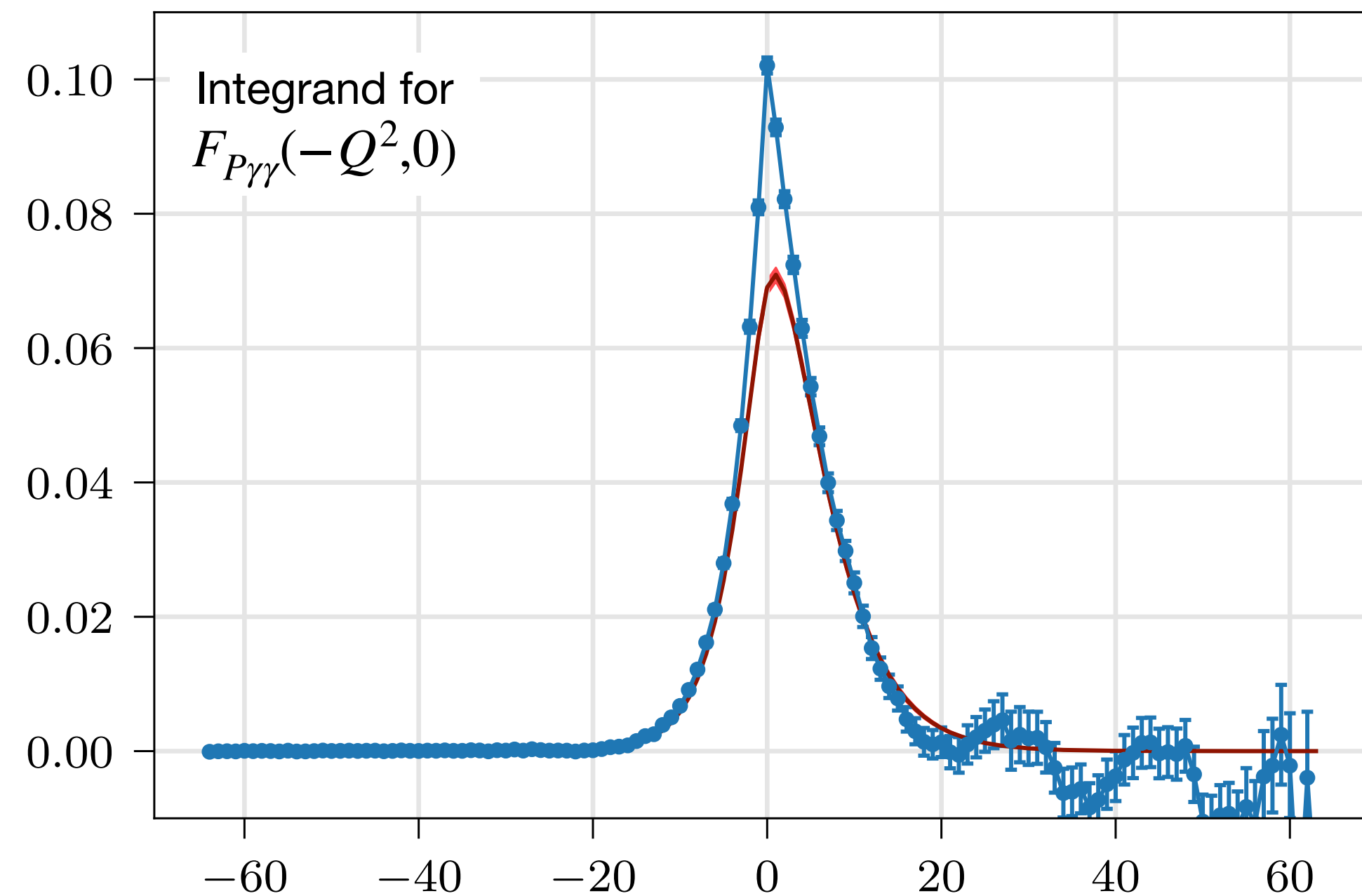
Isospin symmetry allows removing two diagrams:



Note: $O(a^2)$ isospin breaking with twisted mass is part of continuum extrapolation

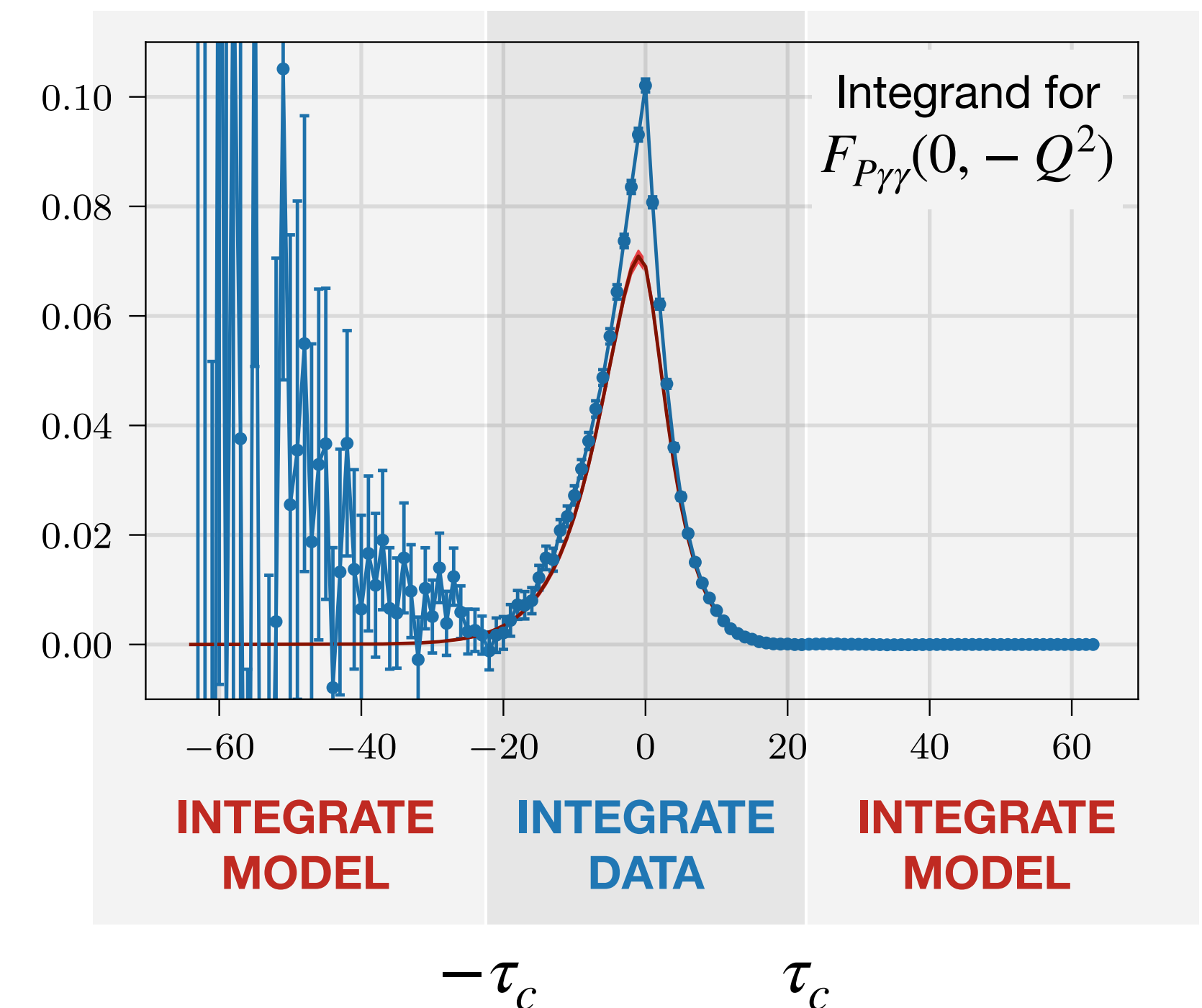
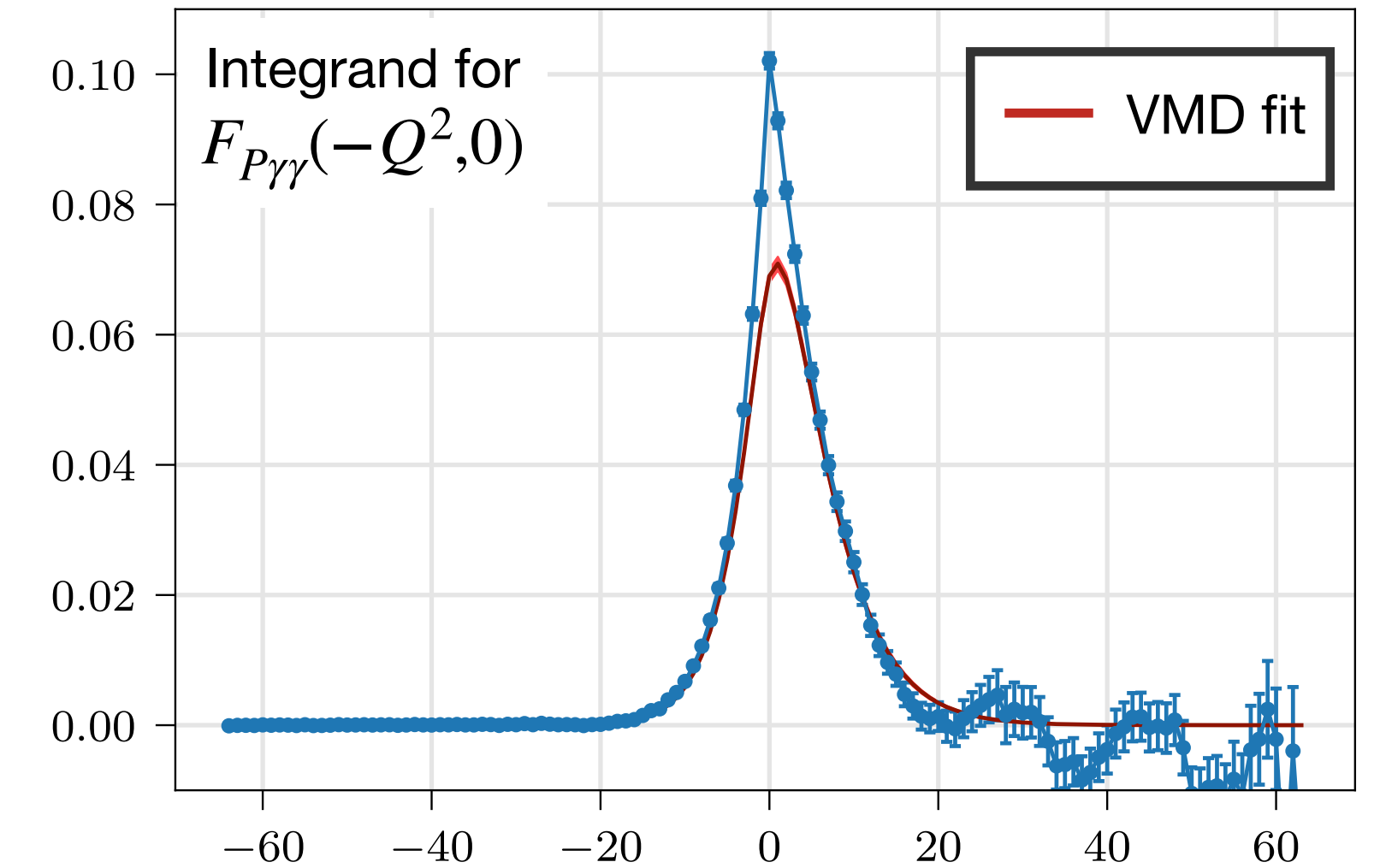
Tail fitting and integration

- Exponentially growing noise in the tails of $\tilde{A}_{\mu\nu}(\tau)$
 - Probed in the Laplace transform $\int_{-\infty}^{\infty} d\tau \boxed{e^{\omega_1 \tau}} \tilde{A}_{\mu\nu}(\tau)$



Tail fitting and integration

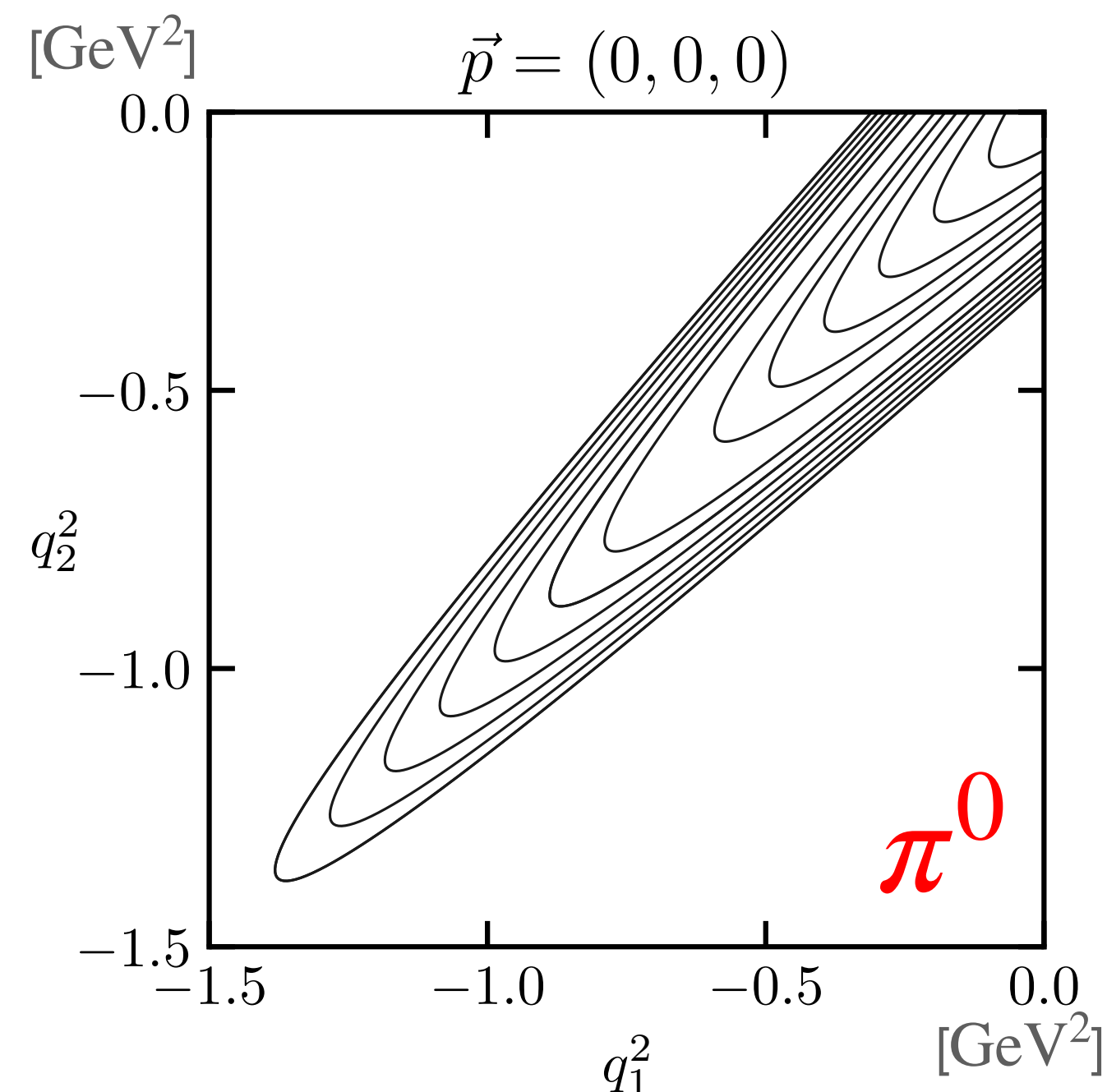
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- To handle this...
 - Fit Vector Meson Dominance (VMD) or Lowest Meson Dominance (LMD) models to the tails
 - Integrate **data** in peak ($|\tau| < \tau_c$)
Integrate **model** in tails ($|\tau| > \tau_c$)



Extrapolating in (q_1^2, q_2^2) plane

Finite volume momenta $\mathbf{q}_1, \mathbf{p}$, arbitrary ω_1

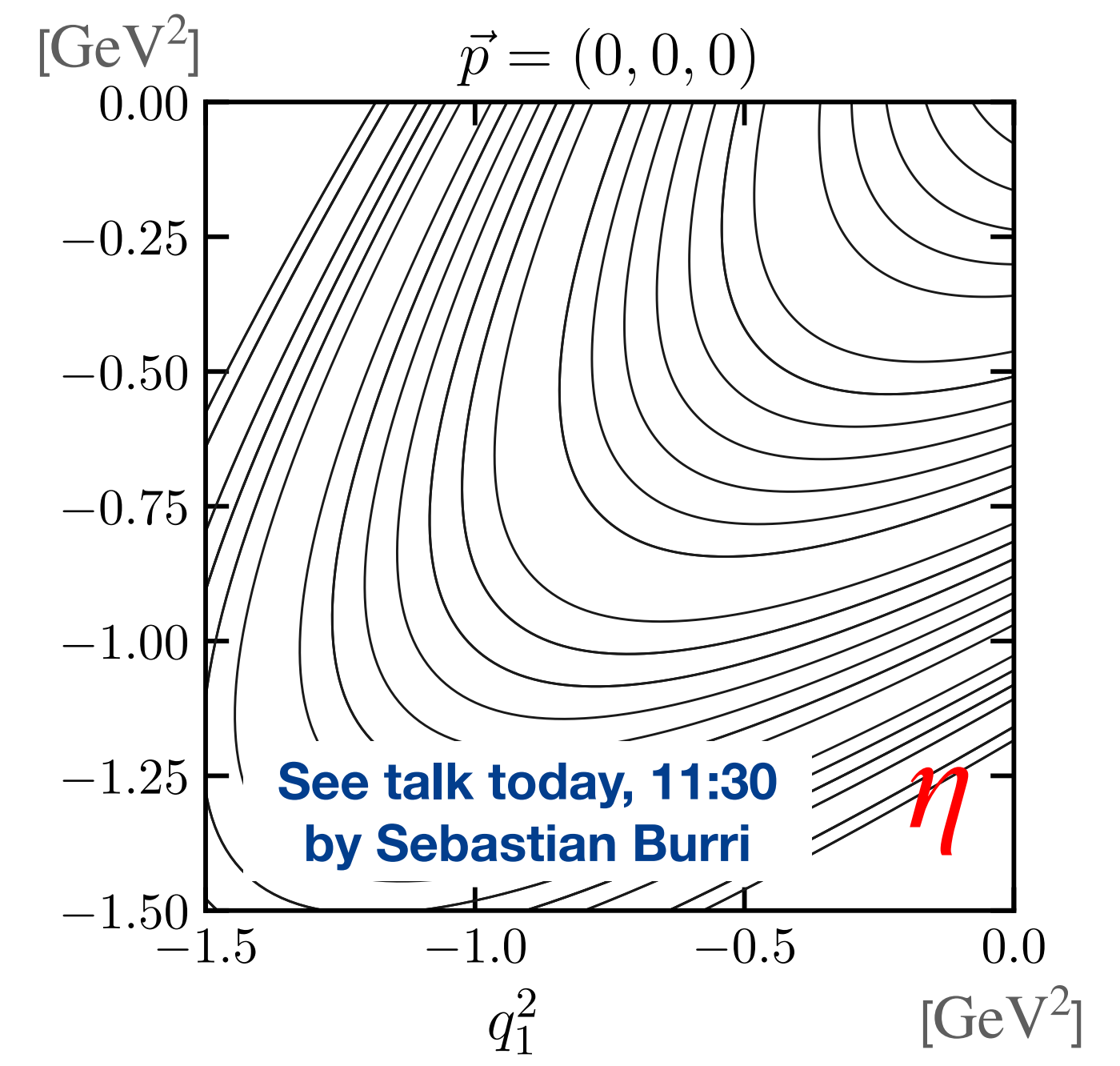
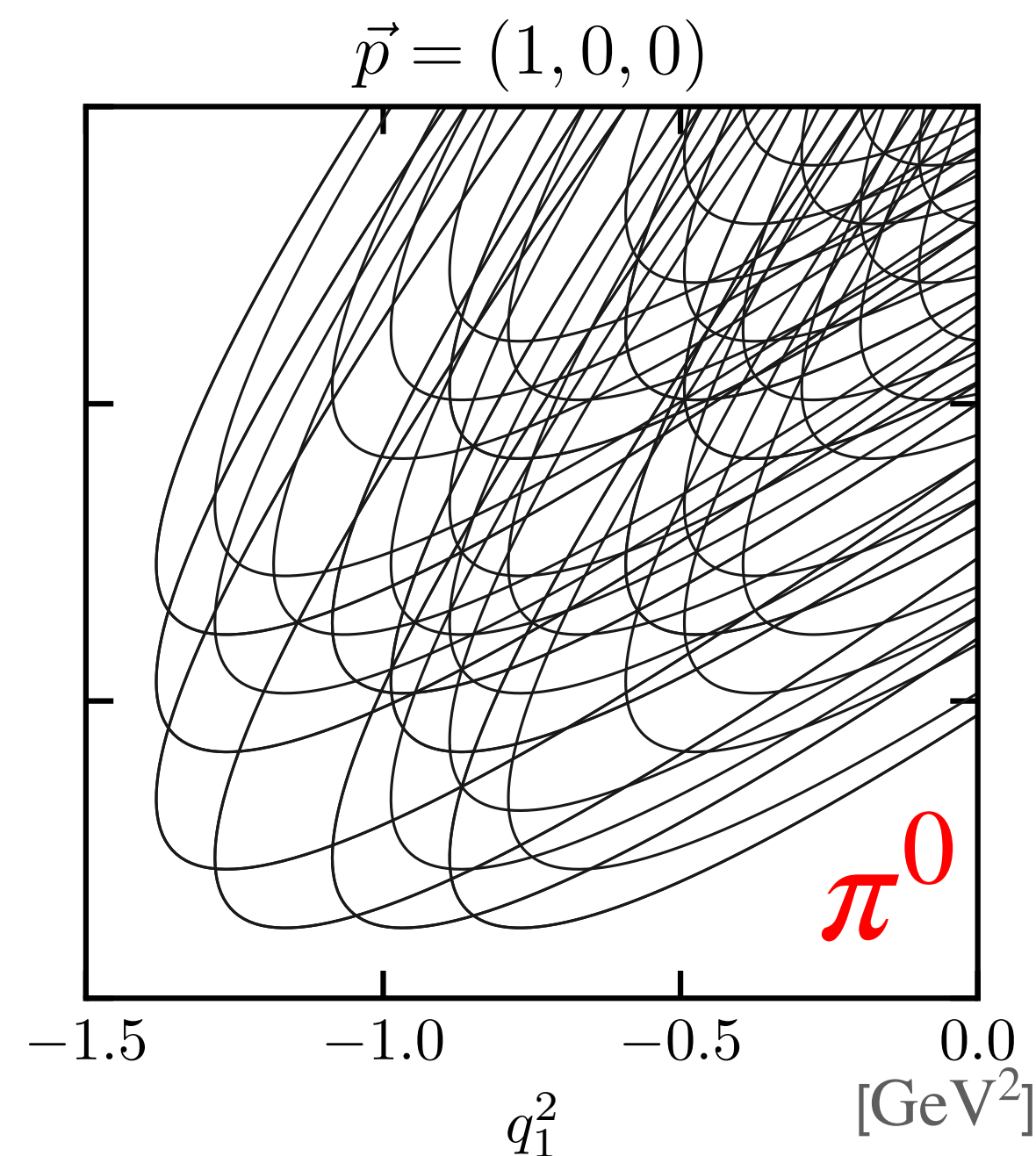
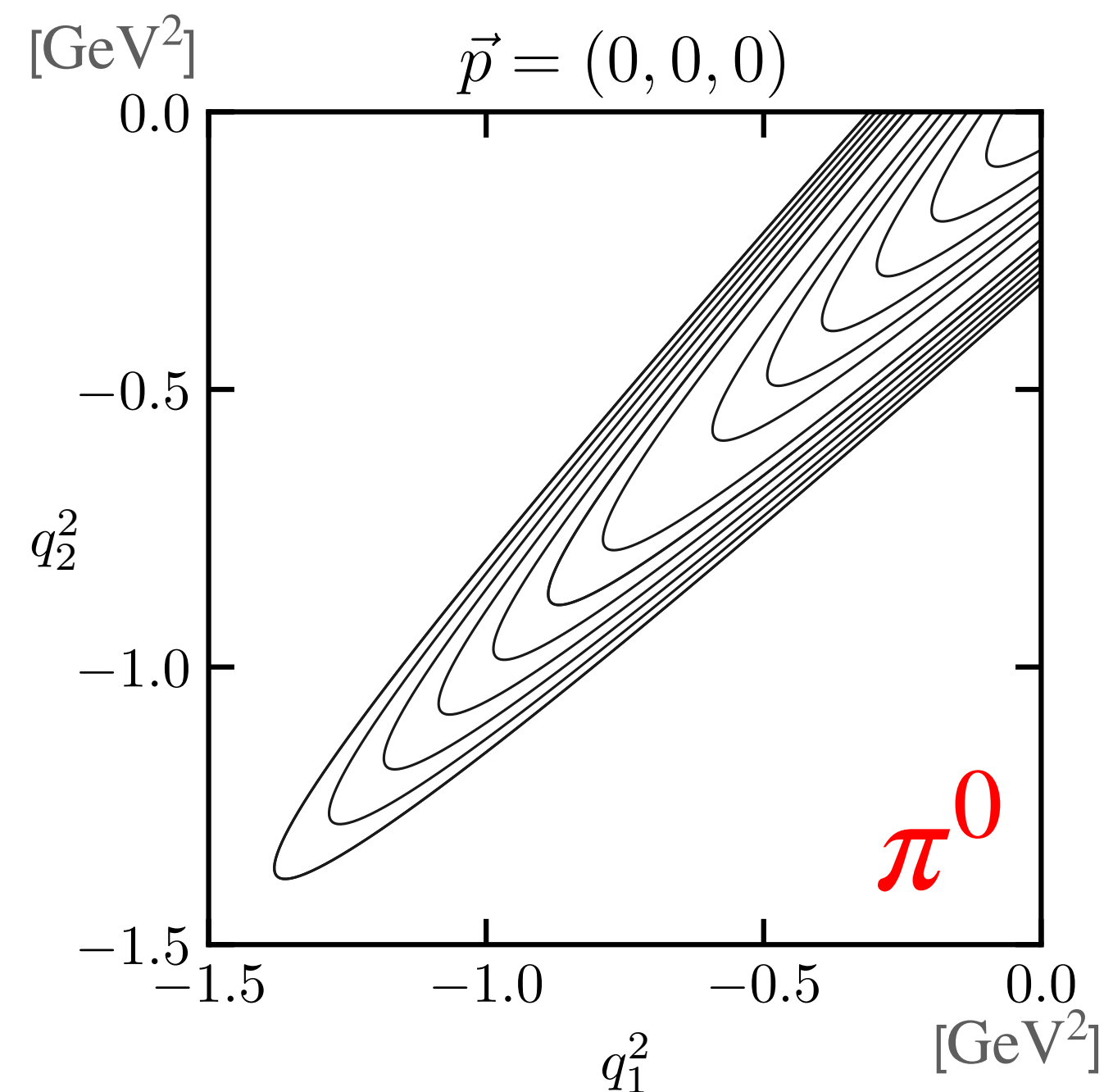
- Parabolic “orbits” of (q_1^2, q_2^2) accessible
- Singly virtual coverage ($\mathbf{p} \neq 0$ & larger m_P help!)



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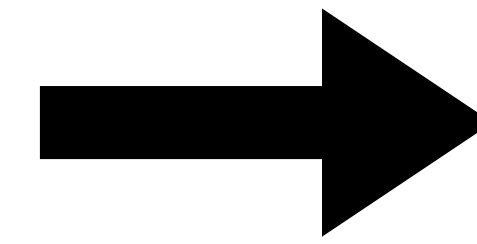
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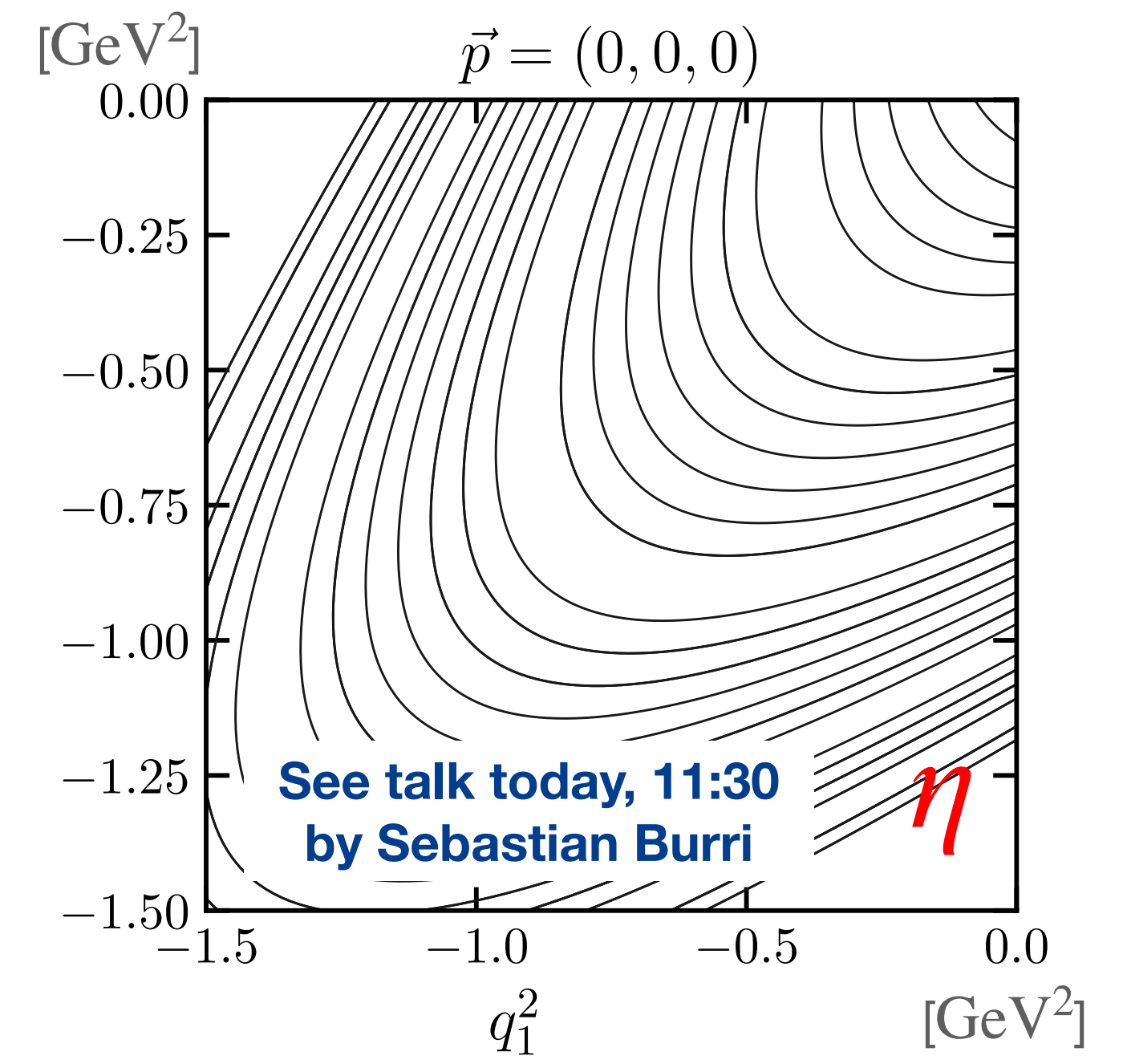
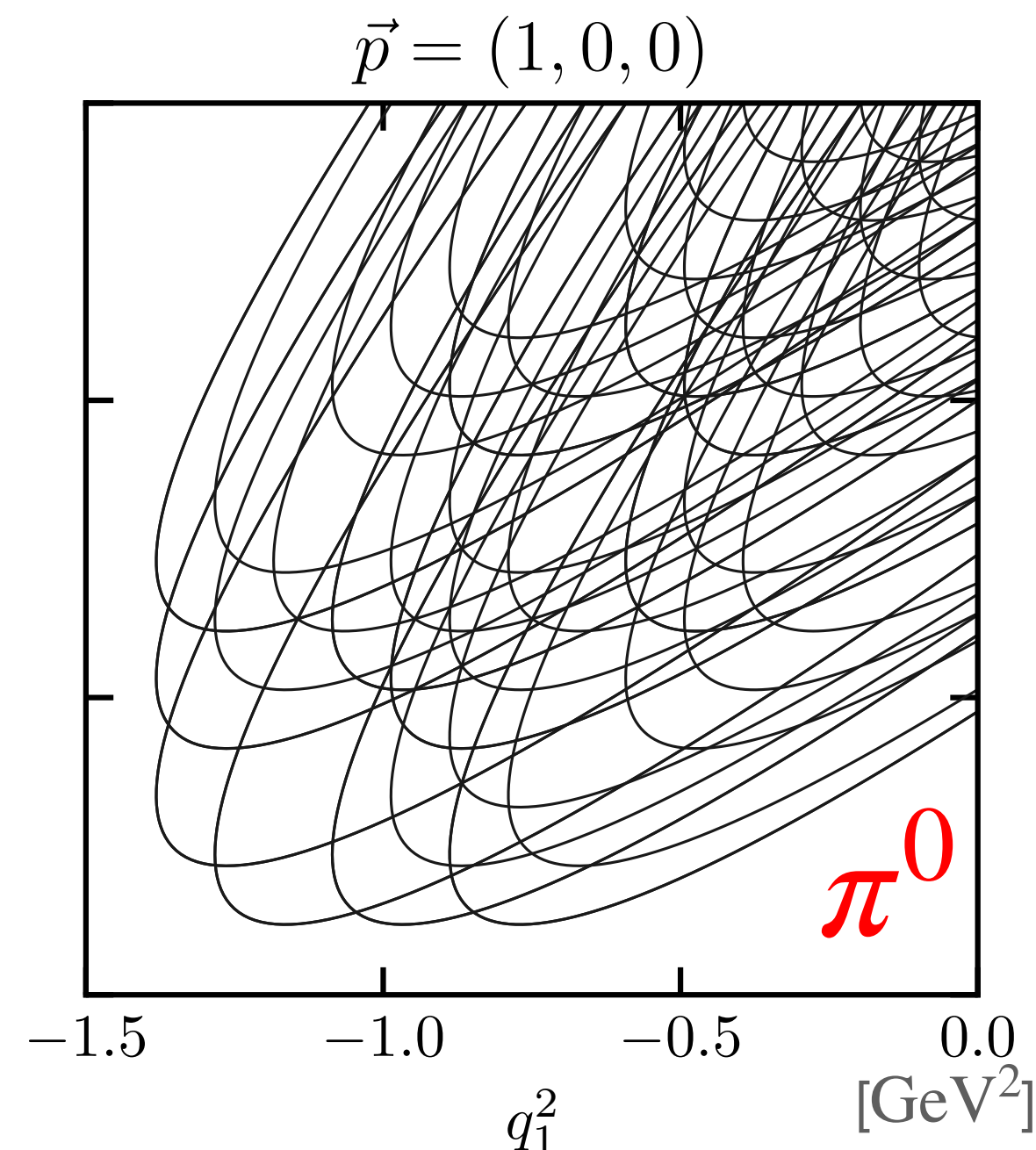
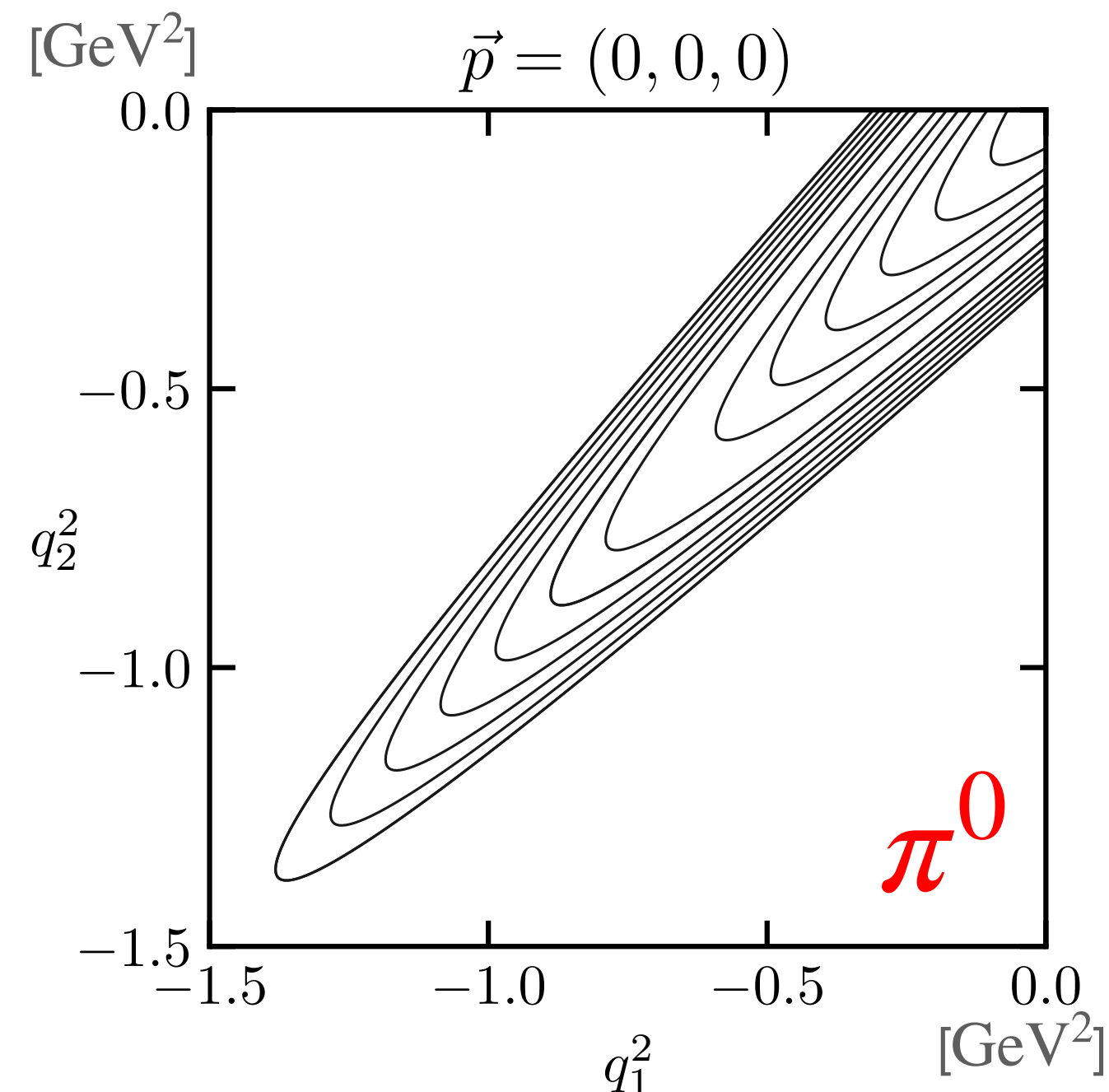
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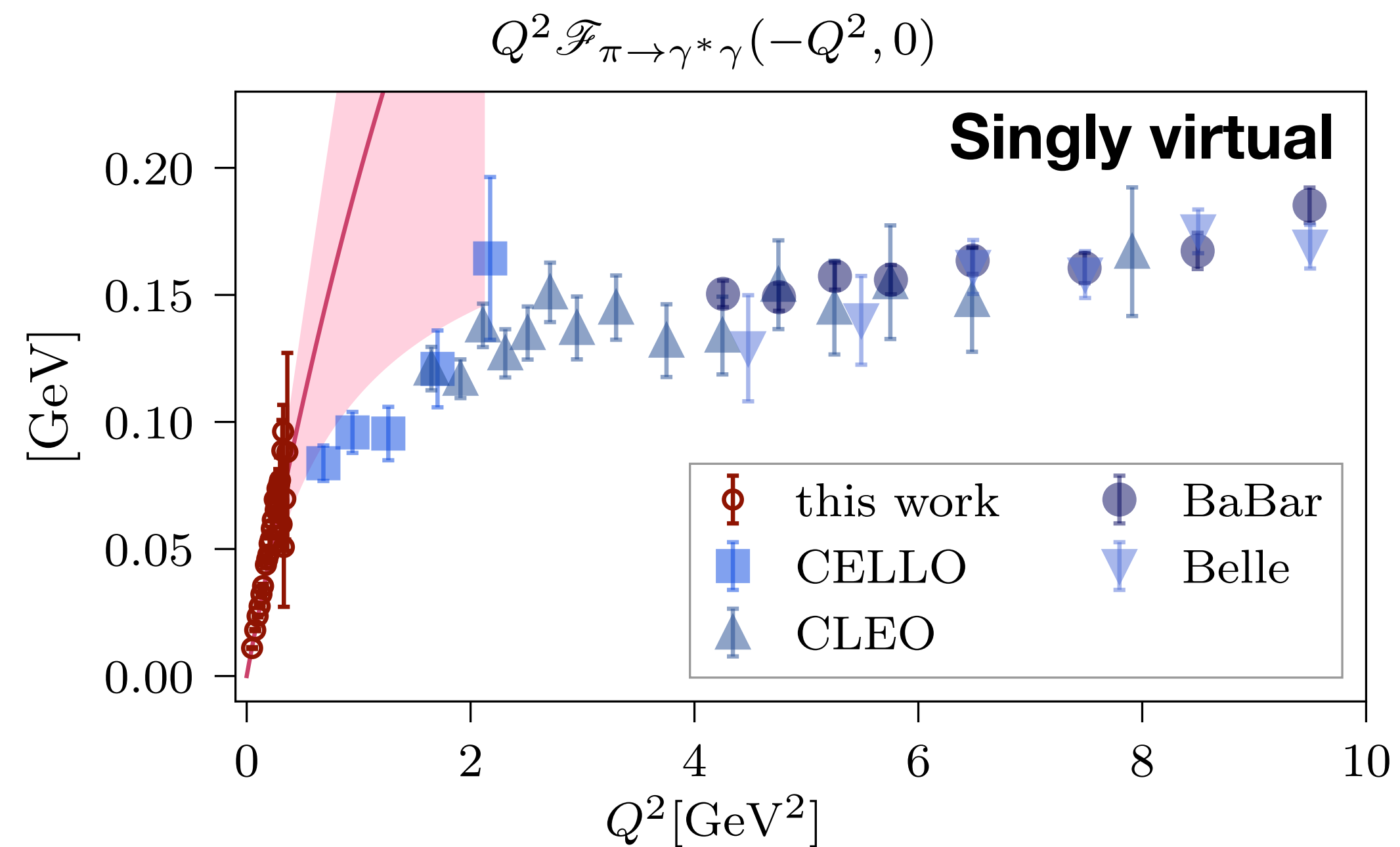
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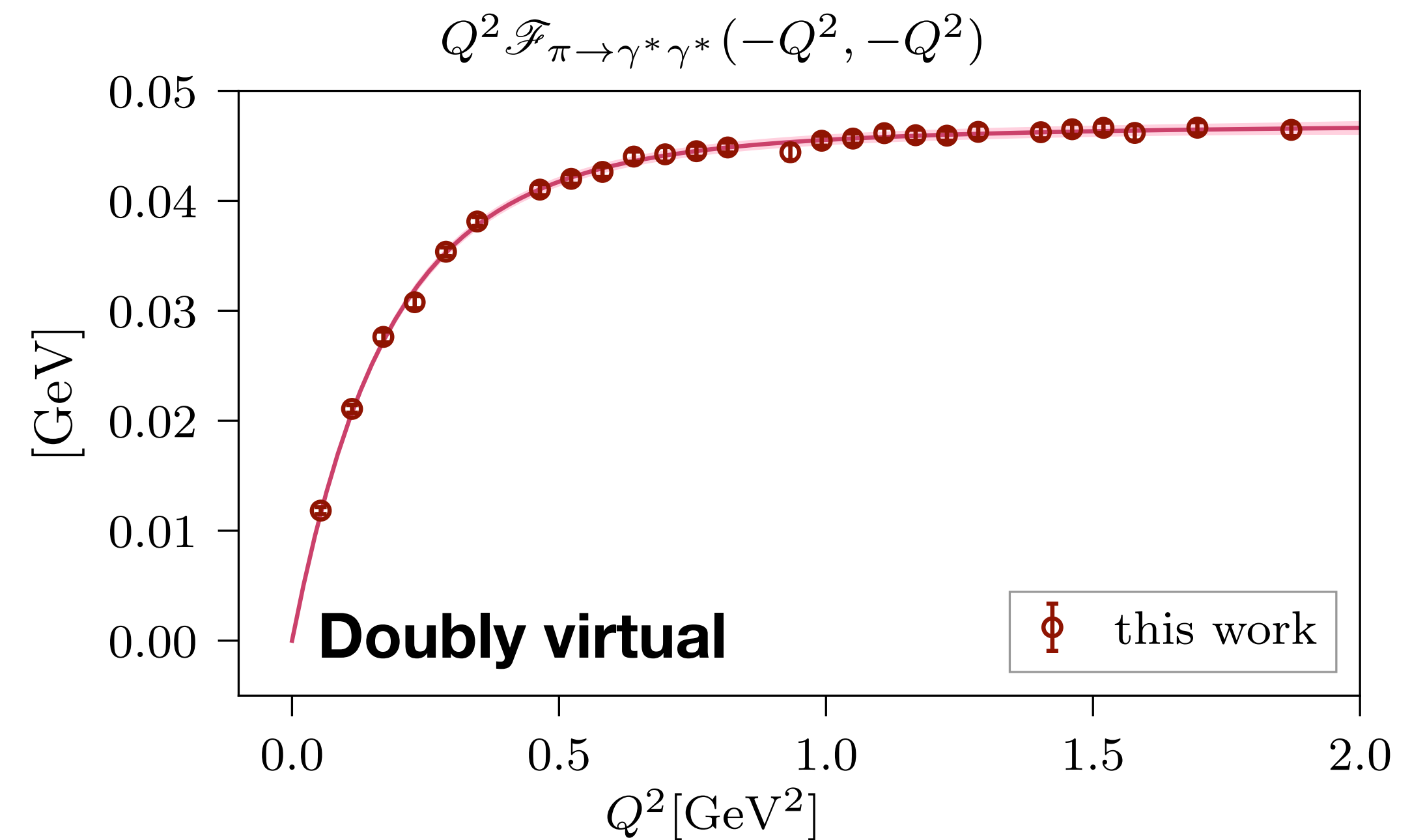
z -expansion for full (q_1^2, q_2^2) dependence



π^0 TFF results



Shown for one choice of ensemble and analysis procedure.

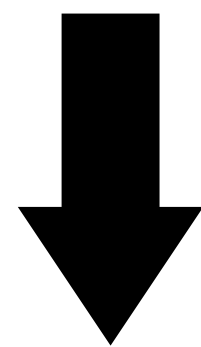


Doubly virtual much more precise than singly virtual due to kinematics.

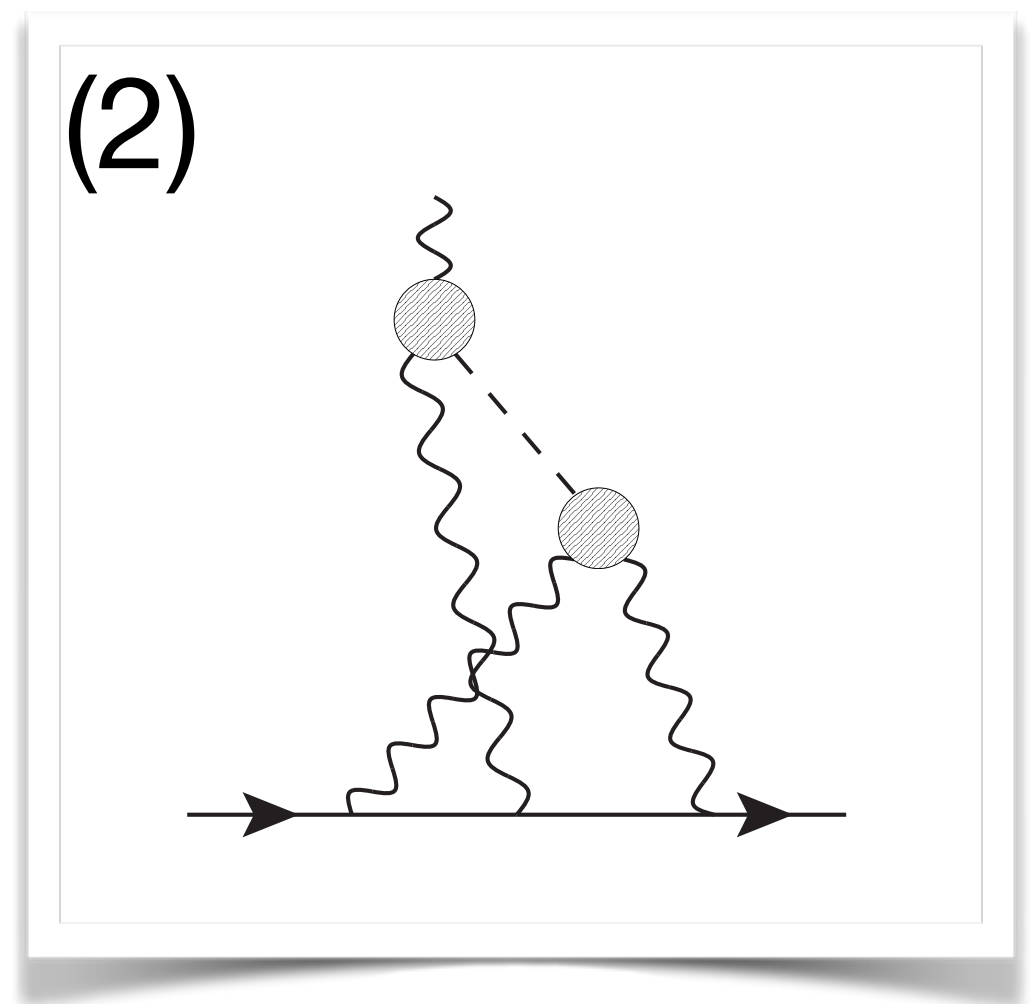
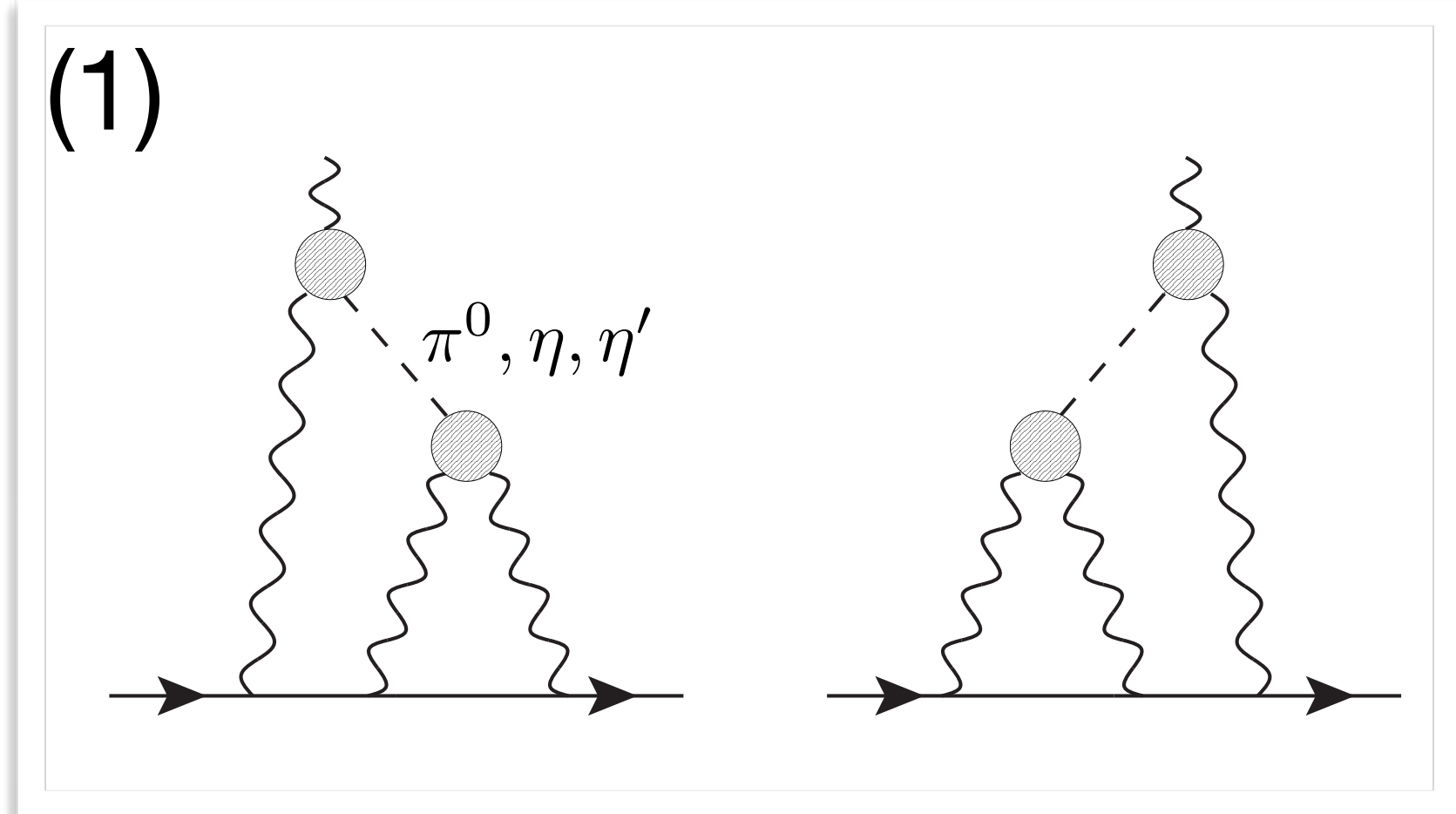
Muon g-2 contribution $a_\mu^{P\text{-pole}}$

$$a_\mu^{P\text{-pole},(1)} = \int_0^\infty dQ_1 \int_0^\infty dQ_2 \int_{-1}^1 d\cos\theta \, w_1(Q_1, Q_2, \cos\theta) \\ \times F_{P\gamma\gamma}(-Q_1^2, -(Q_1 + Q_2)^2) F_{P\gamma\gamma}(-Q_2^2, 0)$$

$$a_\mu^{P\text{-pole},(2)} = \int_0^\infty dQ_1 \int_0^\infty dQ_2 \int_{-1}^1 d\cos\theta \, w_2(Q_1, Q_2, \cos\theta) \\ \times F_{P\gamma\gamma}(-Q_1^2, -Q_2^2) F_{P\gamma\gamma}(-(Q_1 + Q_2)^2, 0)$$



$$a_\mu^{P\text{-pole}} = \left(\frac{\alpha}{\pi}\right)^3 \left[a_\mu^{P\text{-pole},(1)} + a_\mu^{P\text{-pole},(2)} \right]$$

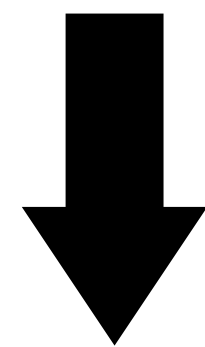


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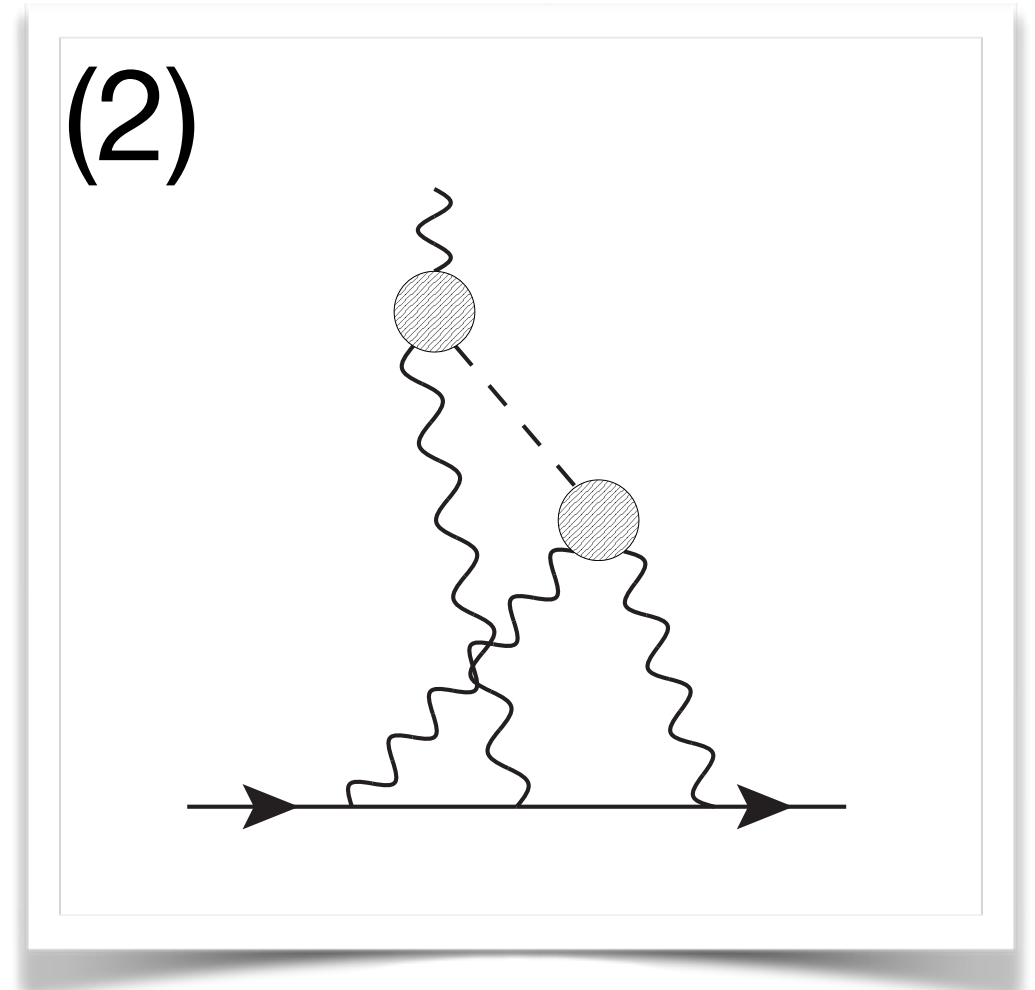
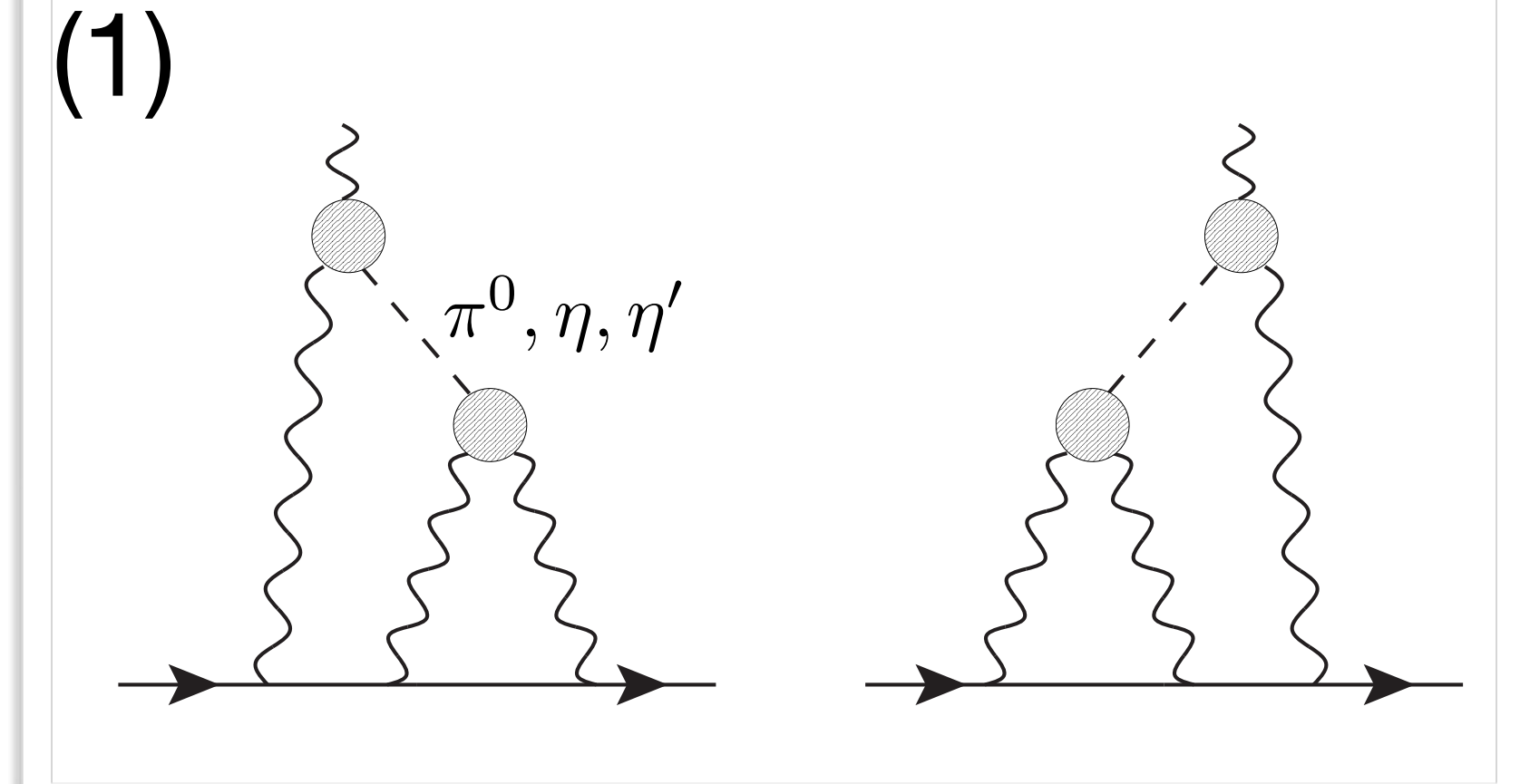
Known kinematic
weight functions

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Muon g-2 contribution $a_\mu^{P\text{-pole}}$

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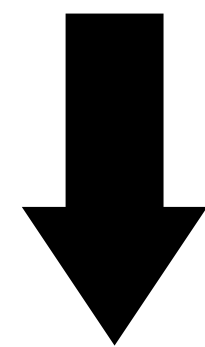
$$a_\mu^{P\text{-pole},(1)} = \int_0^\infty dQ_1 \int_0^\infty dQ_2 \int_{-1}^1 d\cos\theta w_1(Q_1, Q_2, \cos\theta)$$

$$\times F_{P\gamma\gamma}(-Q_1^2, -(Q_1 + Q_2)^2) F_{P\gamma\gamma}(-Q_2^2, 0)$$

$$a_\mu^{P\text{-pole},(2)} = \int_0^\infty dQ_1 \int_0^\infty dQ_2 \int_{-1}^1 d\cos\theta w_2(Q_1, Q_2, \cos\theta)$$

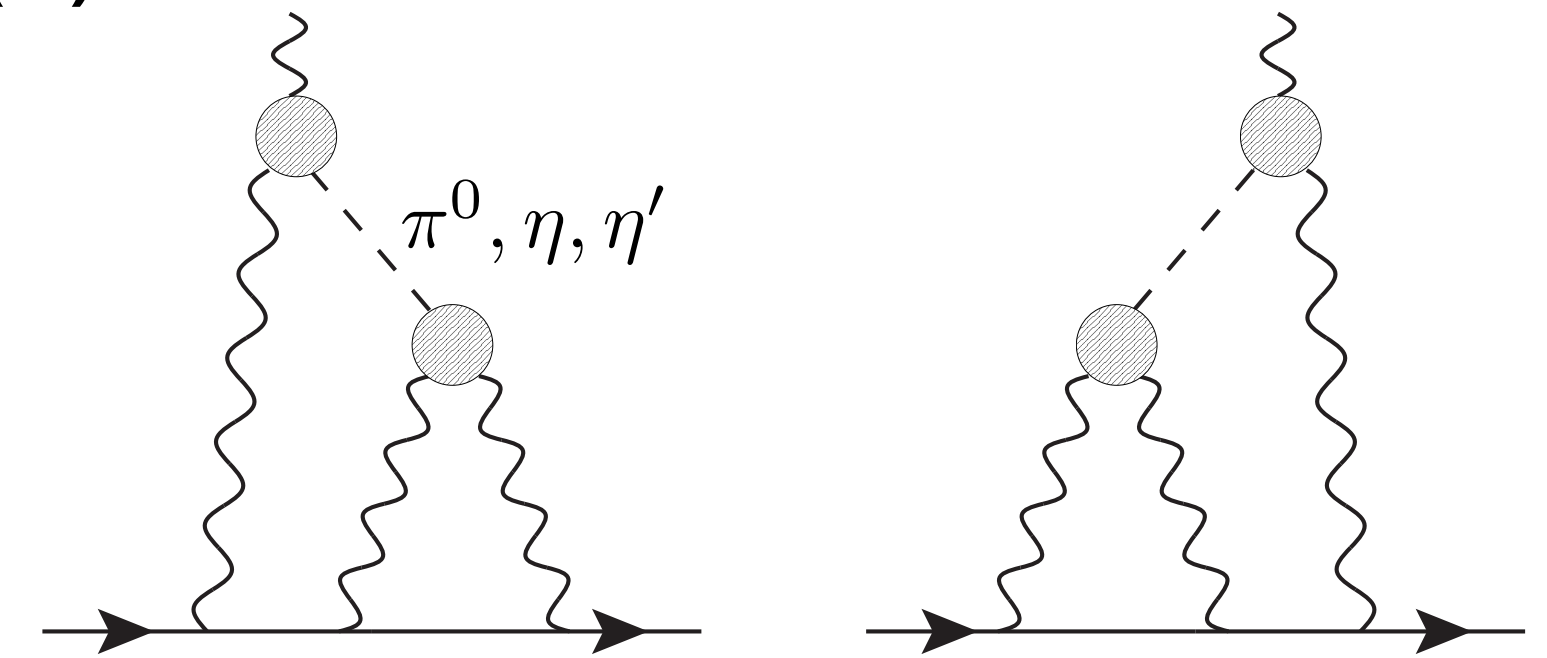
$$\times F_{P\gamma\gamma}(-Q_1^2, -Q_2^2) F_{P\gamma\gamma}(-(Q_1 + Q_2)^2, 0)$$

TFFs with DV x SV,
spacelike structure

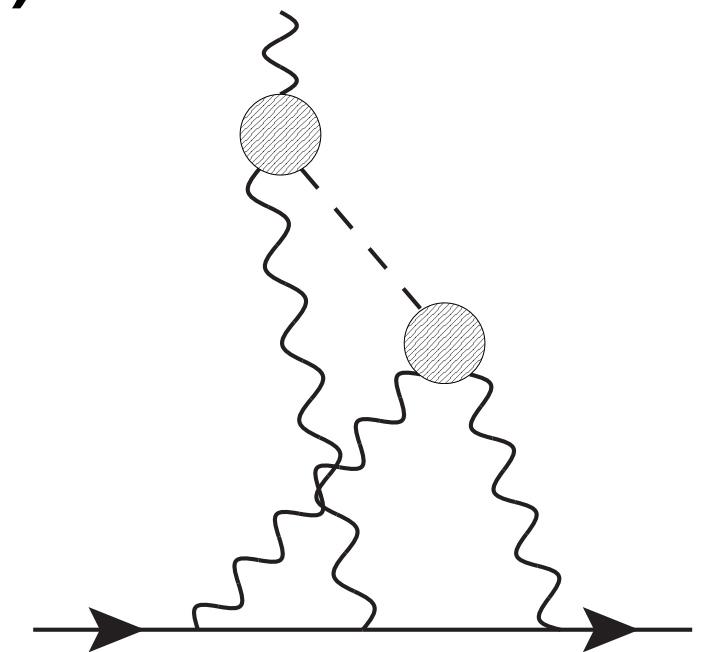


$$a_\mu^{P\text{-pole}} = \left(\frac{\alpha}{\pi}\right)^3 \left[a_\mu^{P\text{-pole},(1)} + a_\mu^{P\text{-pole},(2)} \right]$$

(1)

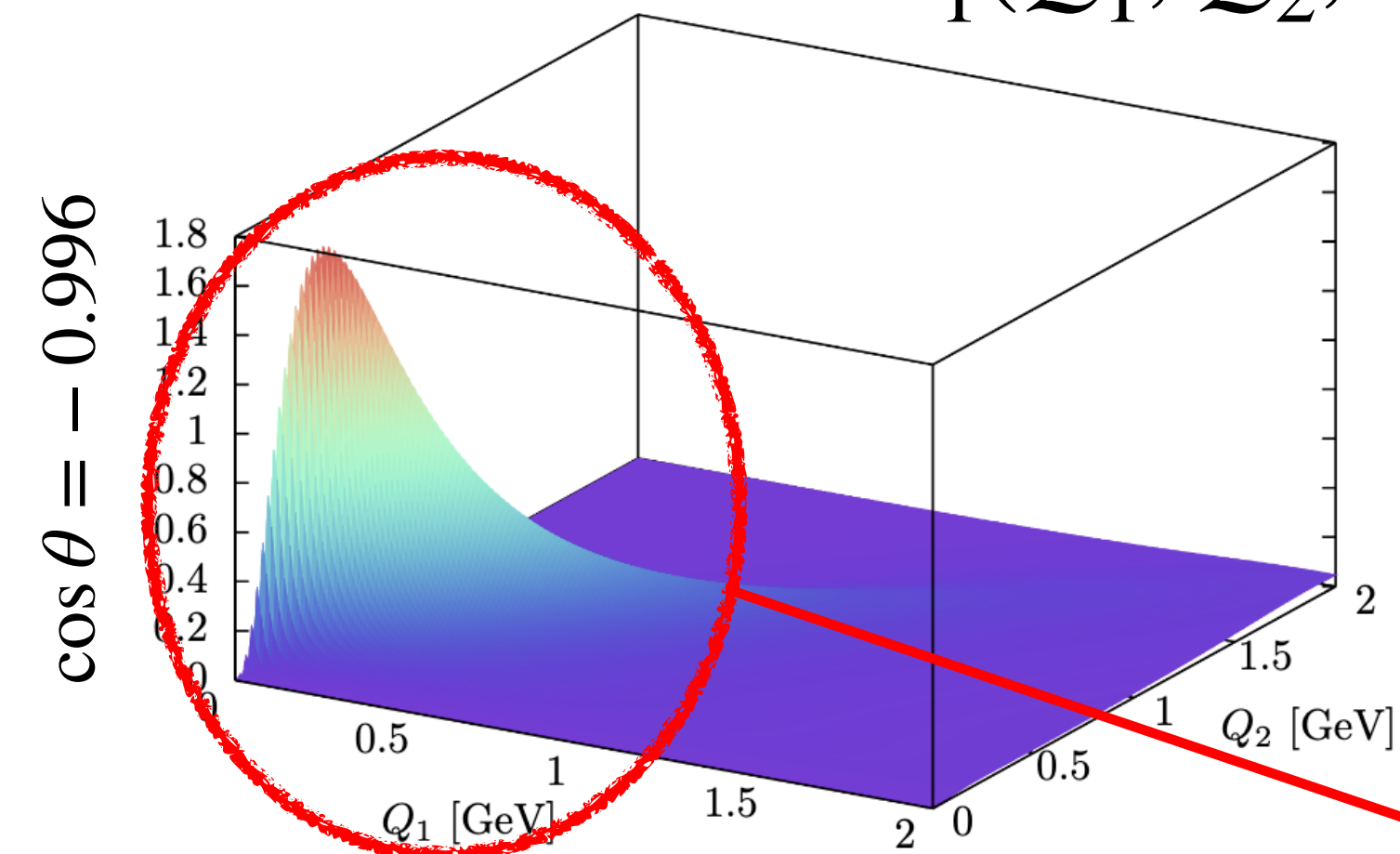


(2)

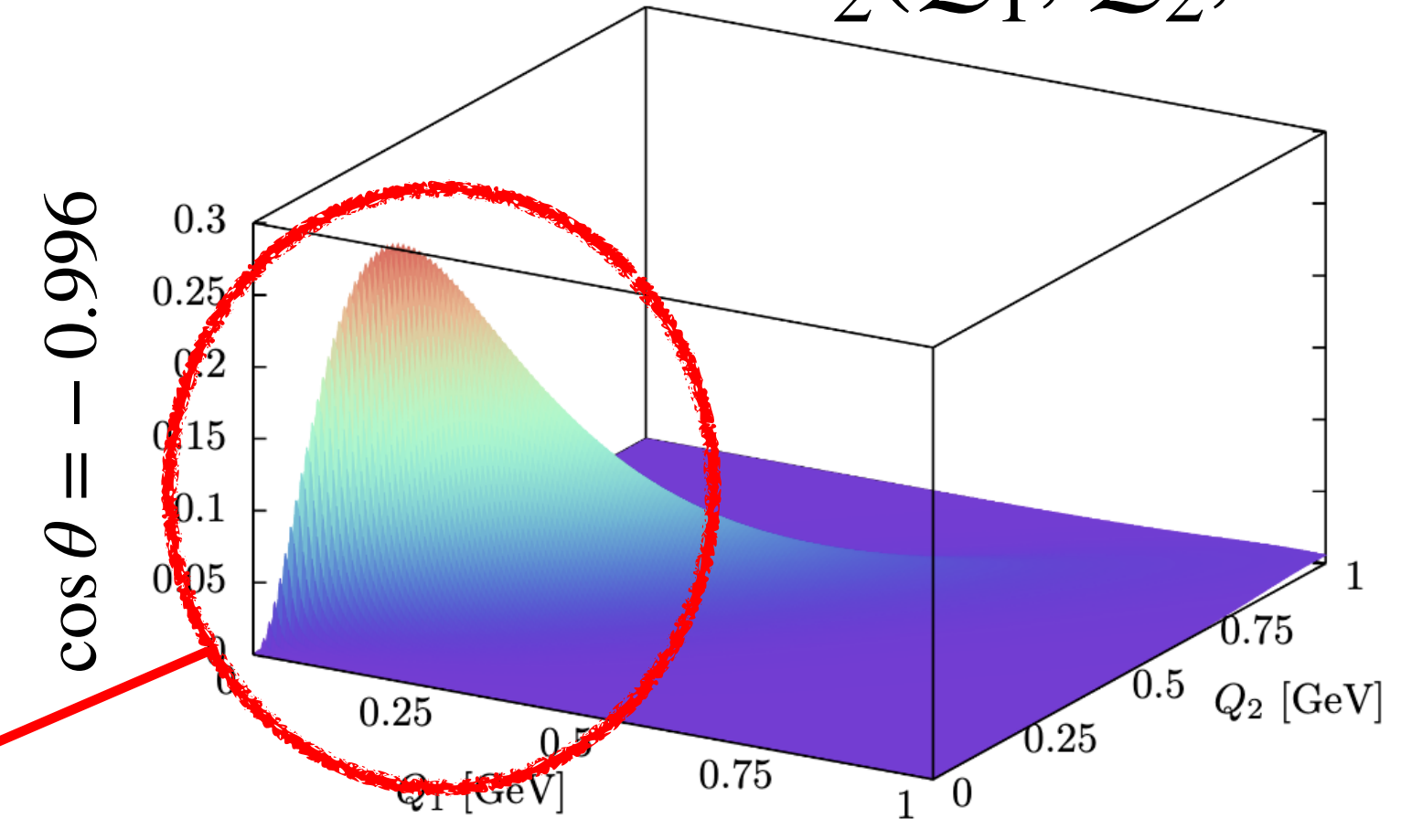


Weight functions for π^0

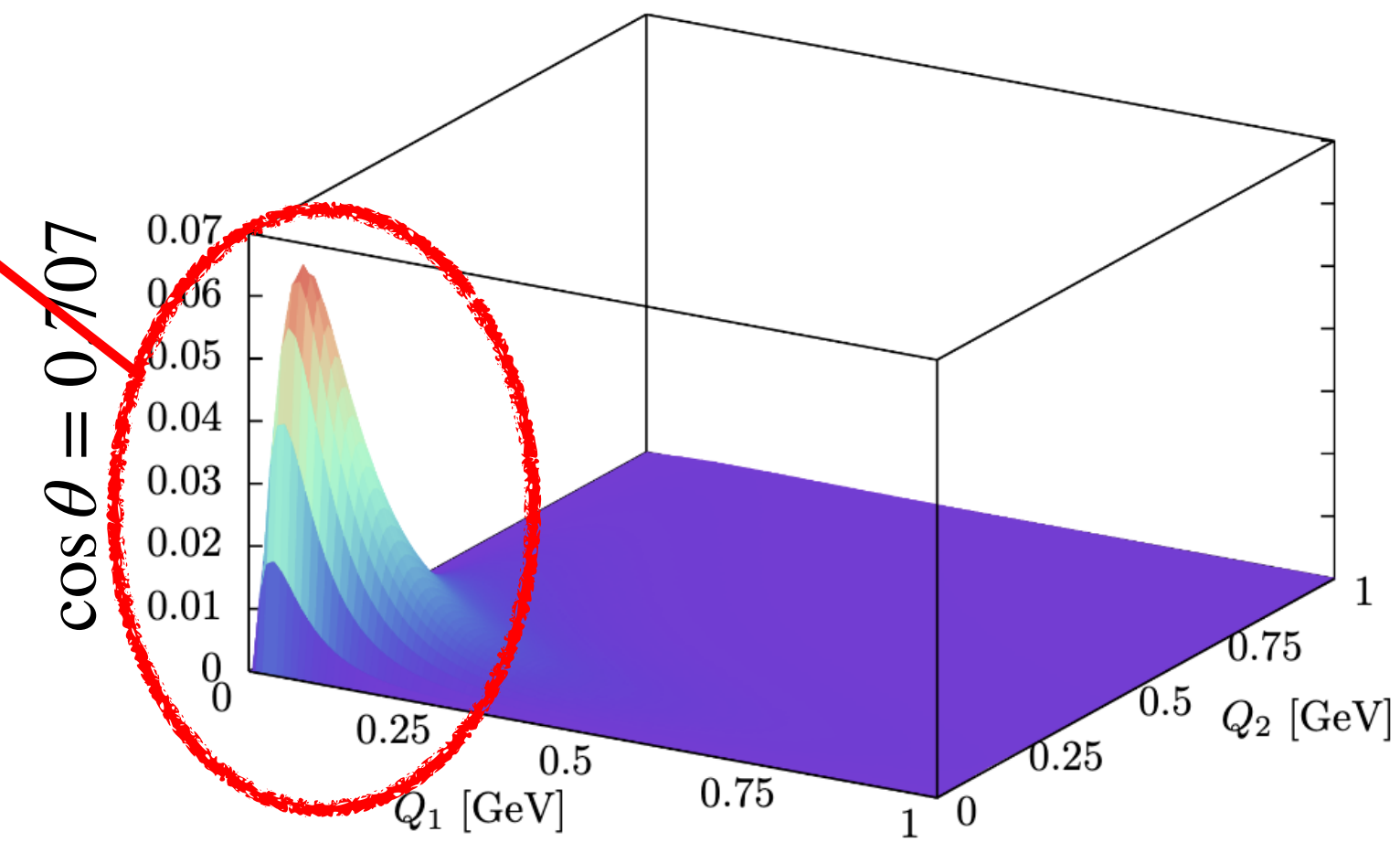
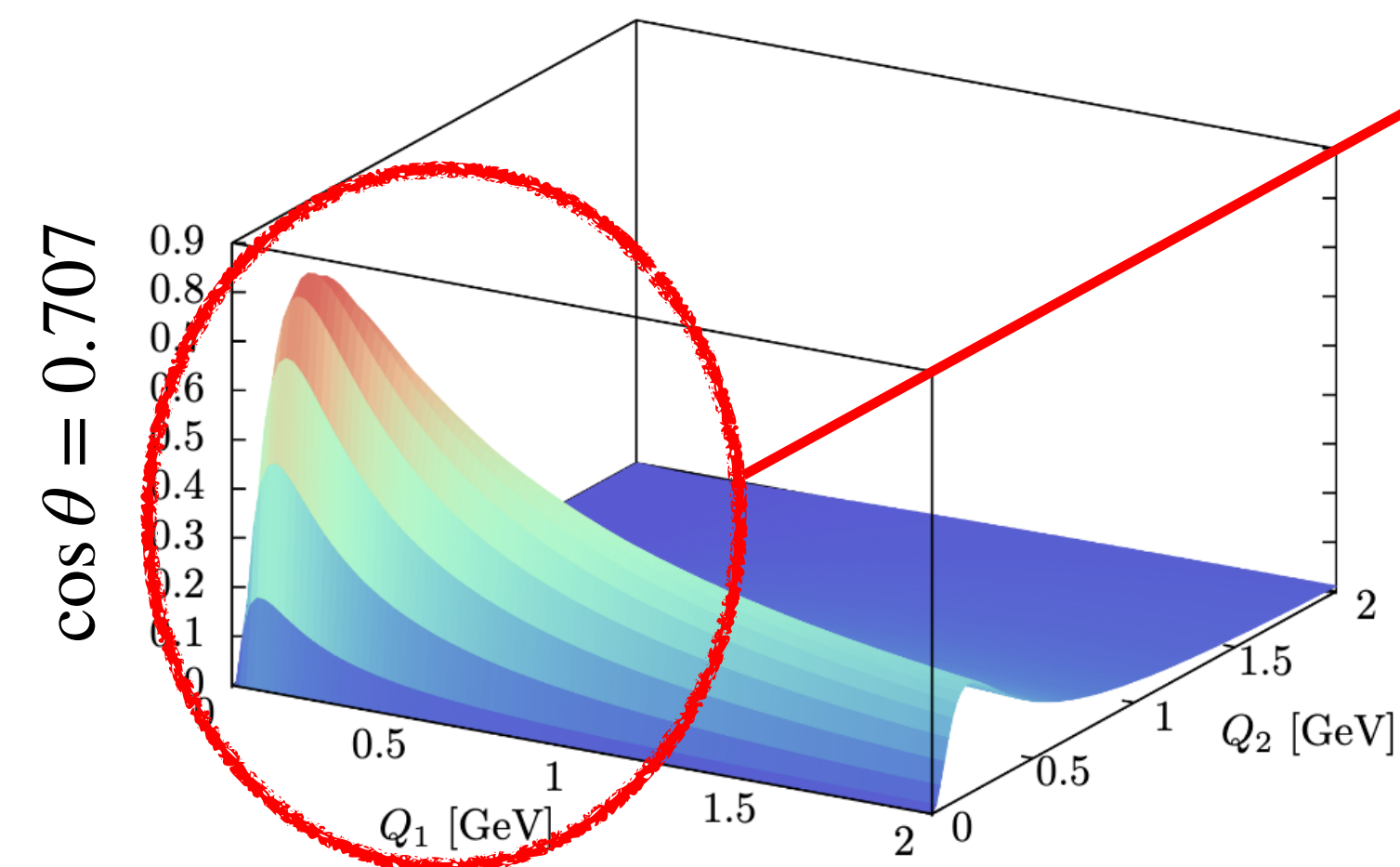
$w_1(Q_1, Q_2, \cos \theta)$



$w_2(Q_1, Q_2, \cos \theta)$



Peaked at
low Q_1^2, Q_2^2



Muon g-2 contribution from the π^0 pole

1. z -expansion fits to FF data per choice of tail fit window, tail cut τ_c

2. Integrated $a_\mu^{\pi^0\text{-pole}}$ per FF

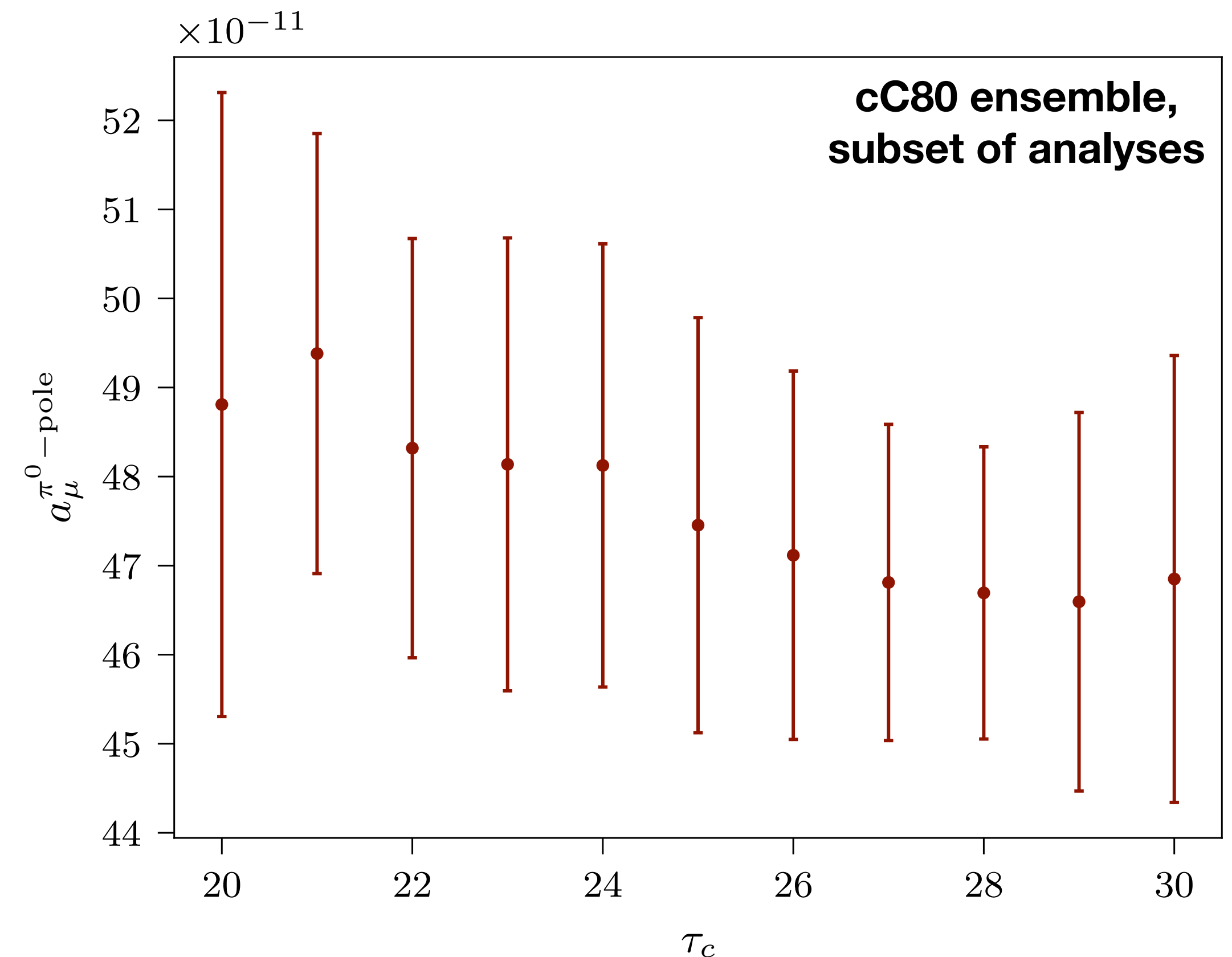


3. Systematic error accounting:

- AIC weighted model averaging

[Borsanyi+ Science 347, 1452 (2015)]

- Checked variation with analysis choices (z -exp order, z -exp fit ranges, t_{seq})



Muon g-2 contribution from the π^0 pole

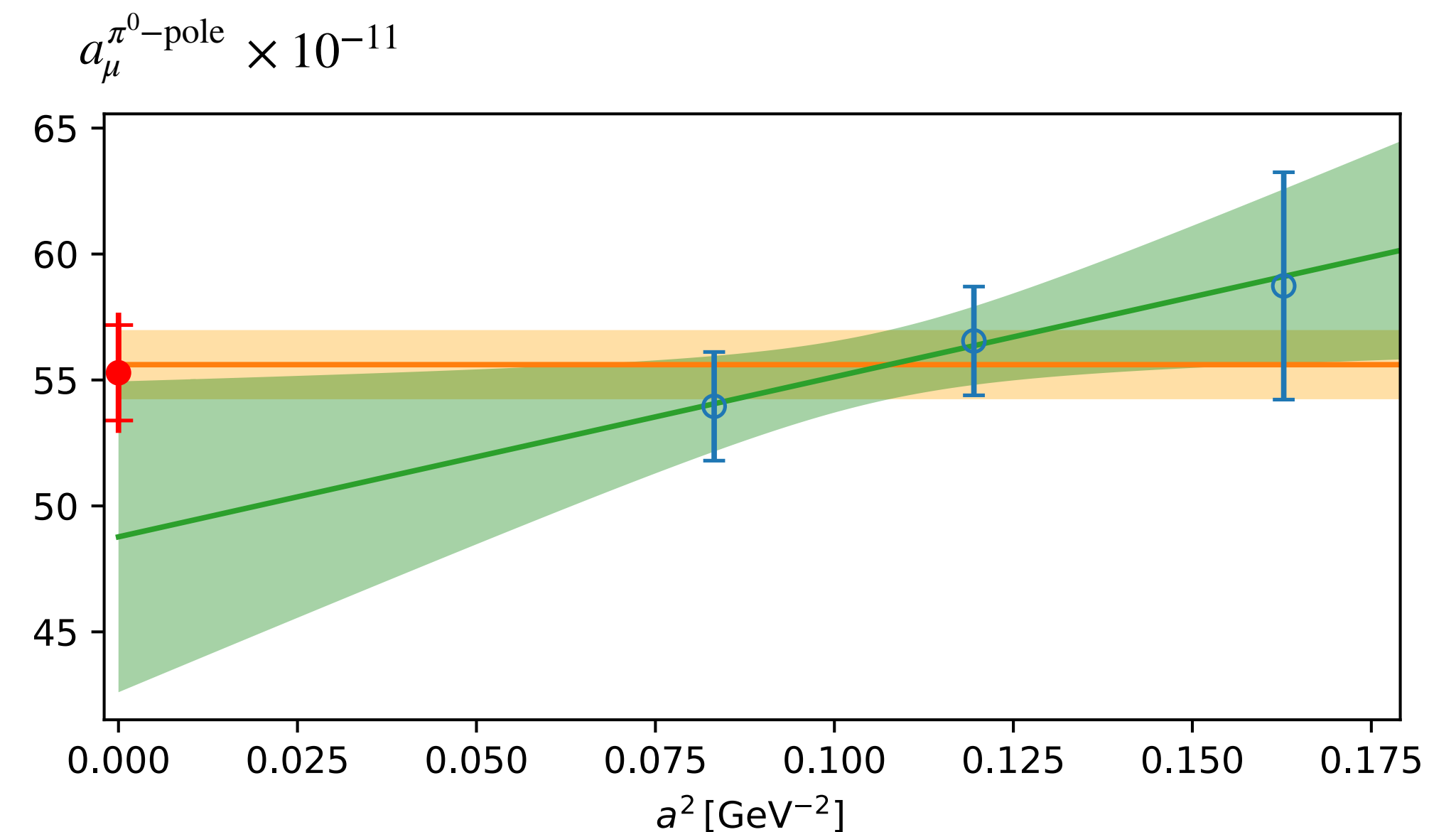
Preliminary continuum limit:

$$a_\mu^{\pi^0\text{-pole}} = 55.3(1.9)_{\text{ctm}}(1.5)_{\text{ctm-syst}} \times 10^{-11}$$

Constant vs linear in a^2 fits.

- Constant: underestimated error
- Linear: overestimated error
- Preliminary result: weighted average + variation between fits as syst. err

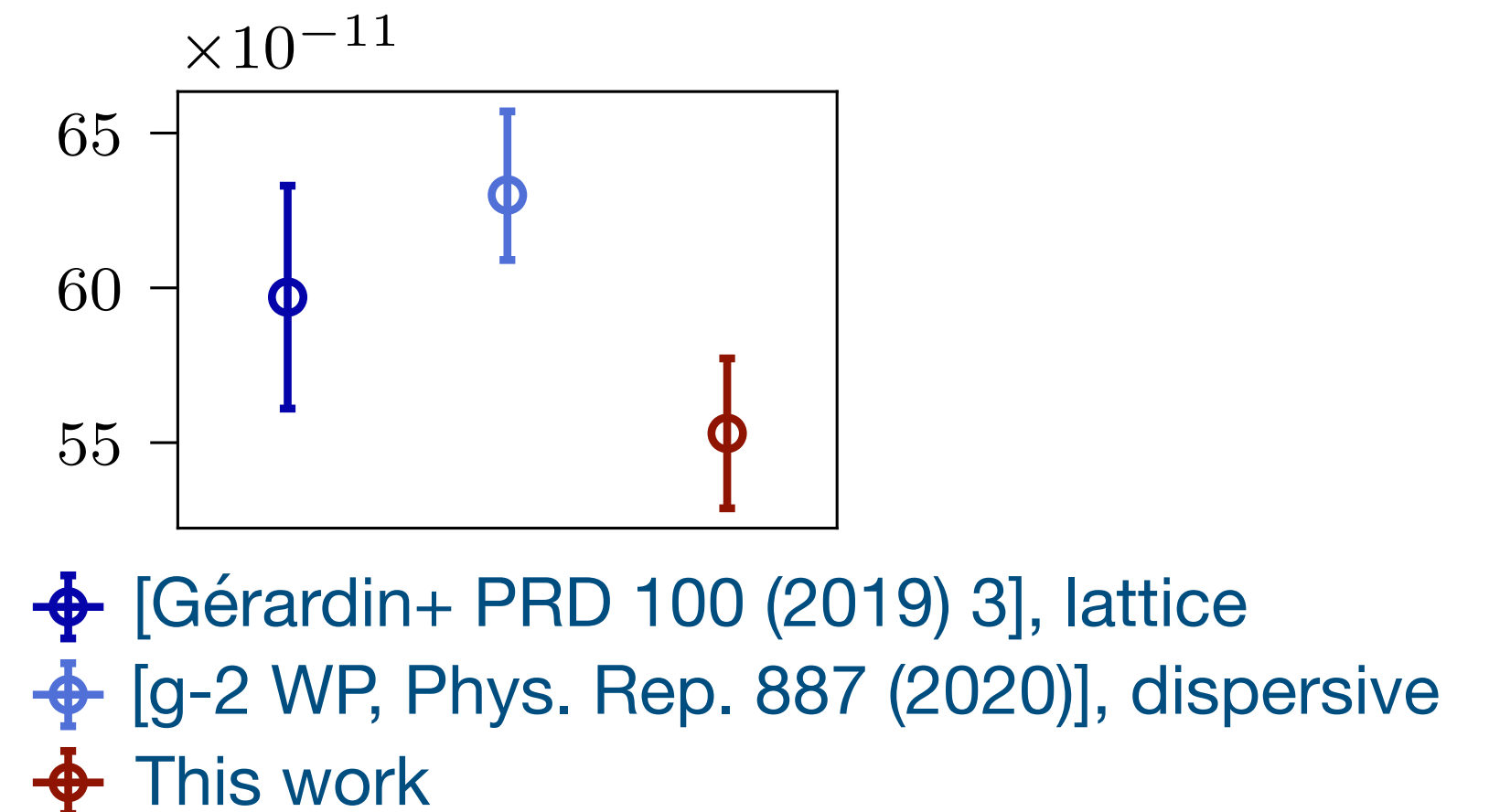
[Alexandrou+ (ETMC) 2104.13408]



Muon g-2 contribution from the π^0 pole

Preliminary continuum limit:

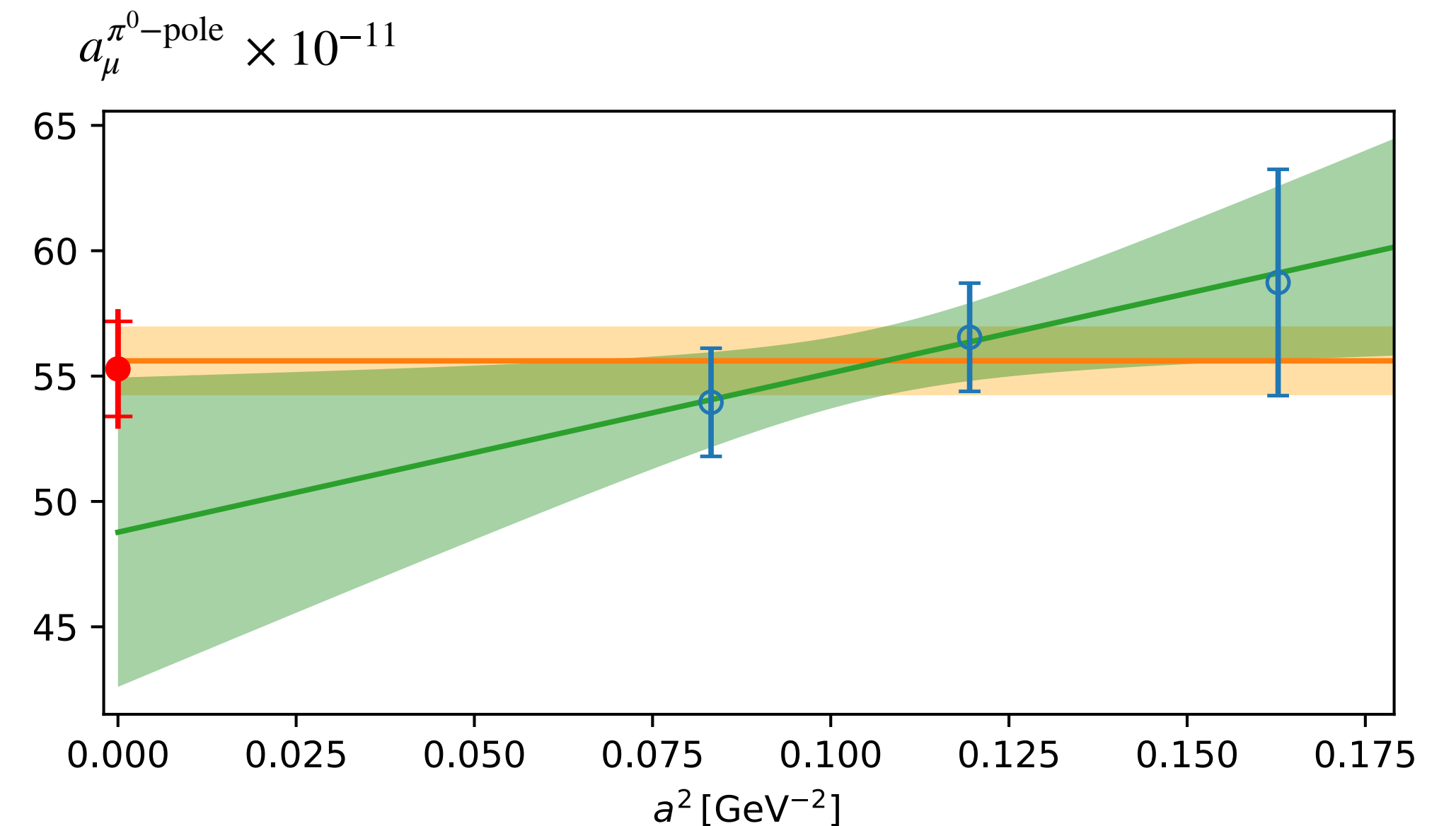
$$a_\mu^{\pi^0\text{-pole}} = 55.3(1.9)_{\text{ctm}}(1.5)_{\text{ctm-syst}} \times 10^{-11}$$



Constant vs linear in a^2 fits.

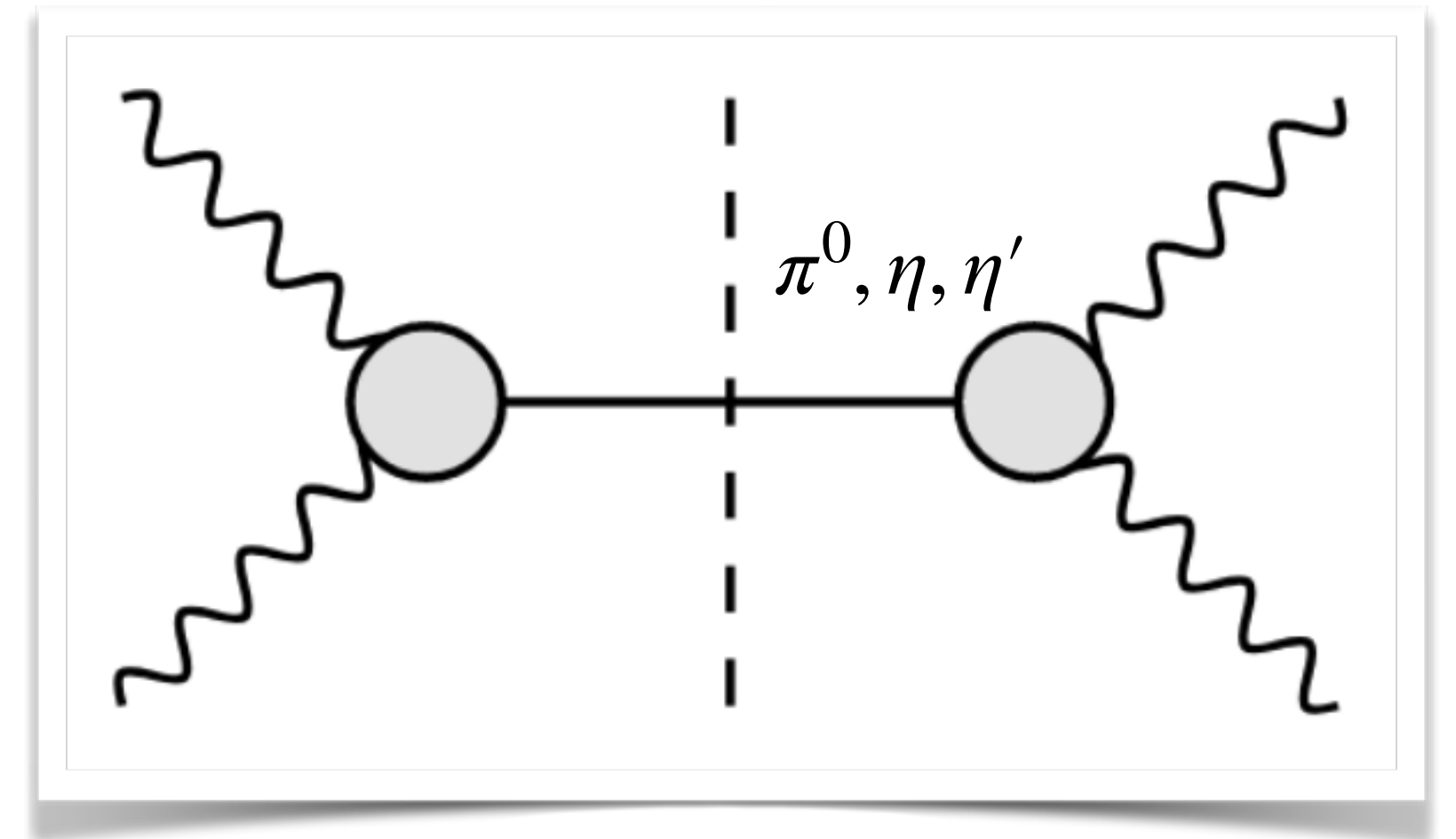
- Constant: underestimated error
- Linear: overestimated error
- Preliminary result: weighted average + variation between fits as syst. err

[Alexandrou+ (ETMC) 2104.13408]



Summary

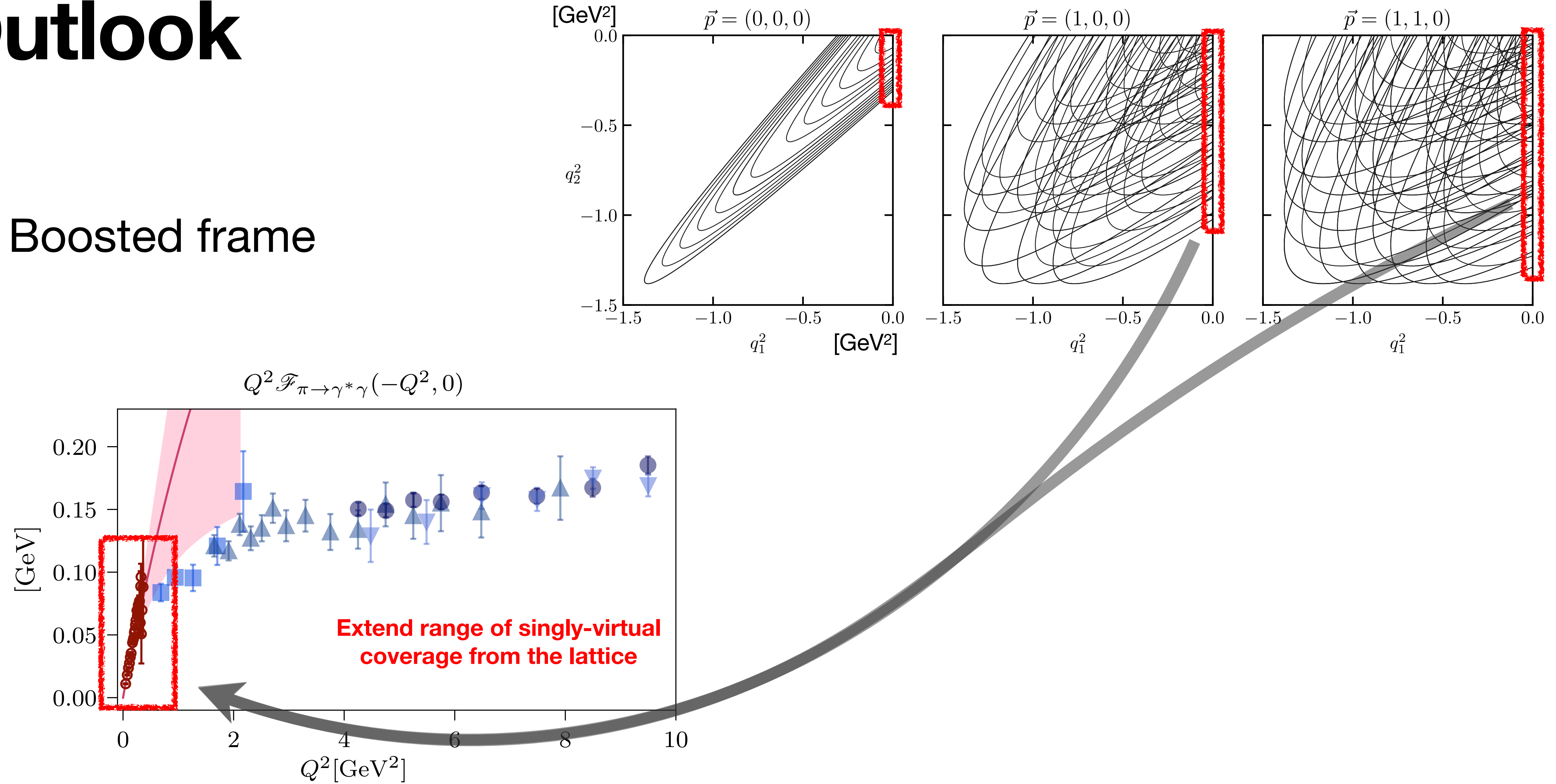
1. Unambiguous decomposition of terms of HLbL
 - Pseudoscalar-pole contribution significant uncertainty
2. Form factors are the key input
 - π^0 dispersively from expt data, η, η' not yet
 - Lattice especially important for **doubly virtual**



3. Comparable determination of $a_\mu^{\pi^0\text{-pole}}$ at **physical point**

Outlook

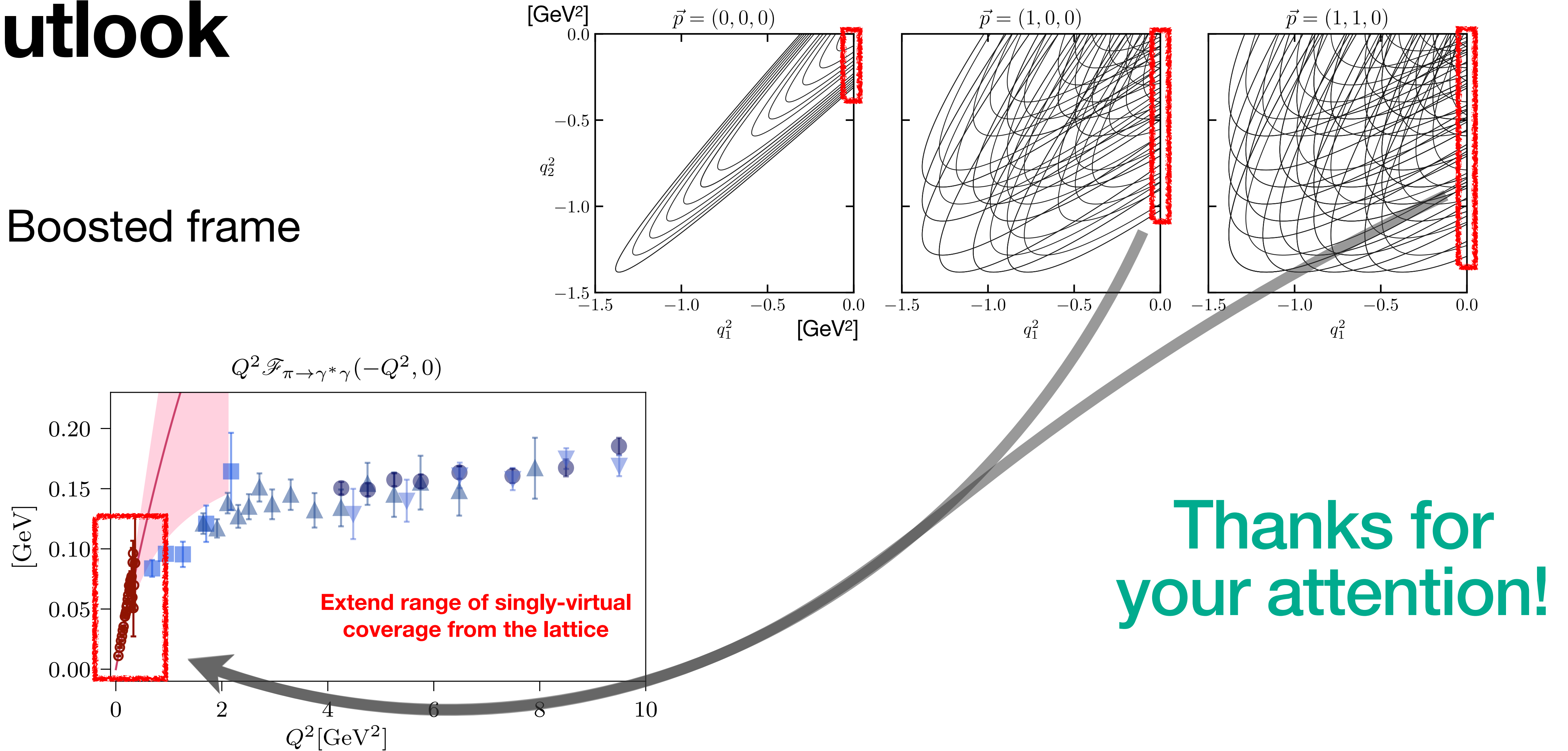
1. Boosted frame



2. Other terms in HLbL decomposition?

Outlook

1. Boosted frame



2. Other terms in HLbL decomposition?

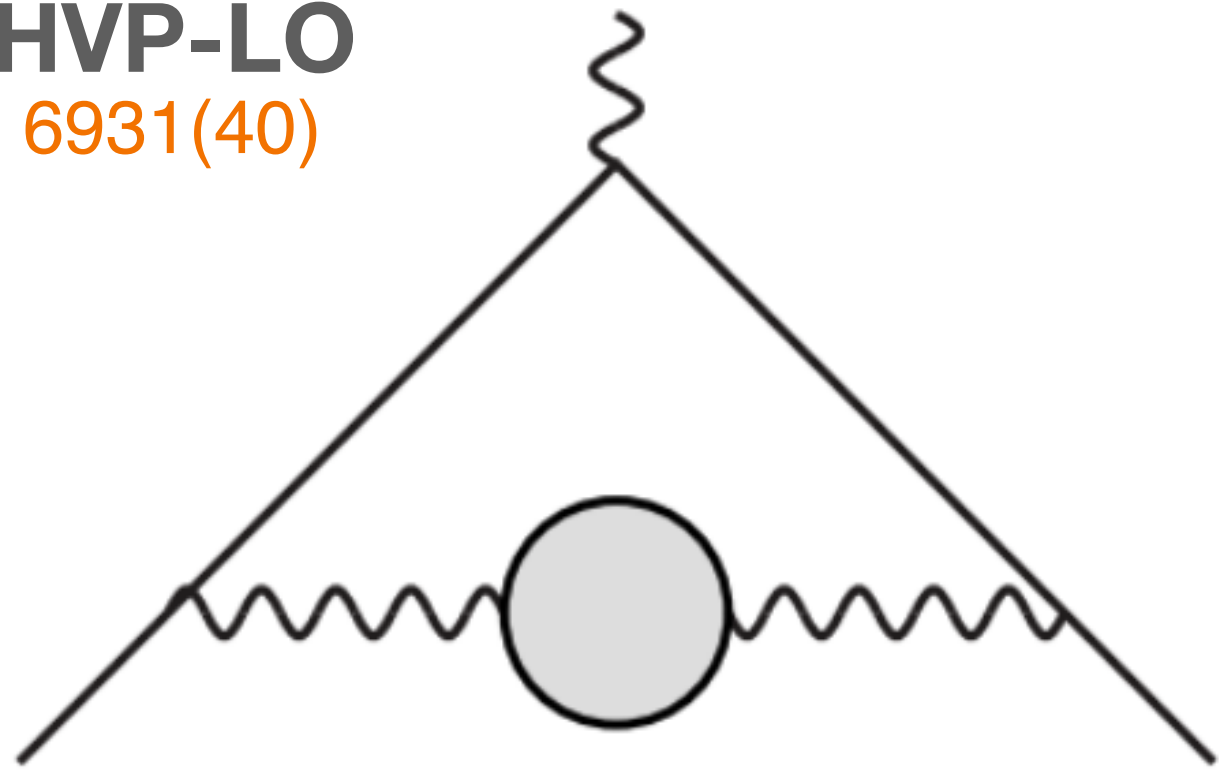
Backup slides

Hadronic contributions to a_μ

All a_μ values $\times 10^{-11}$

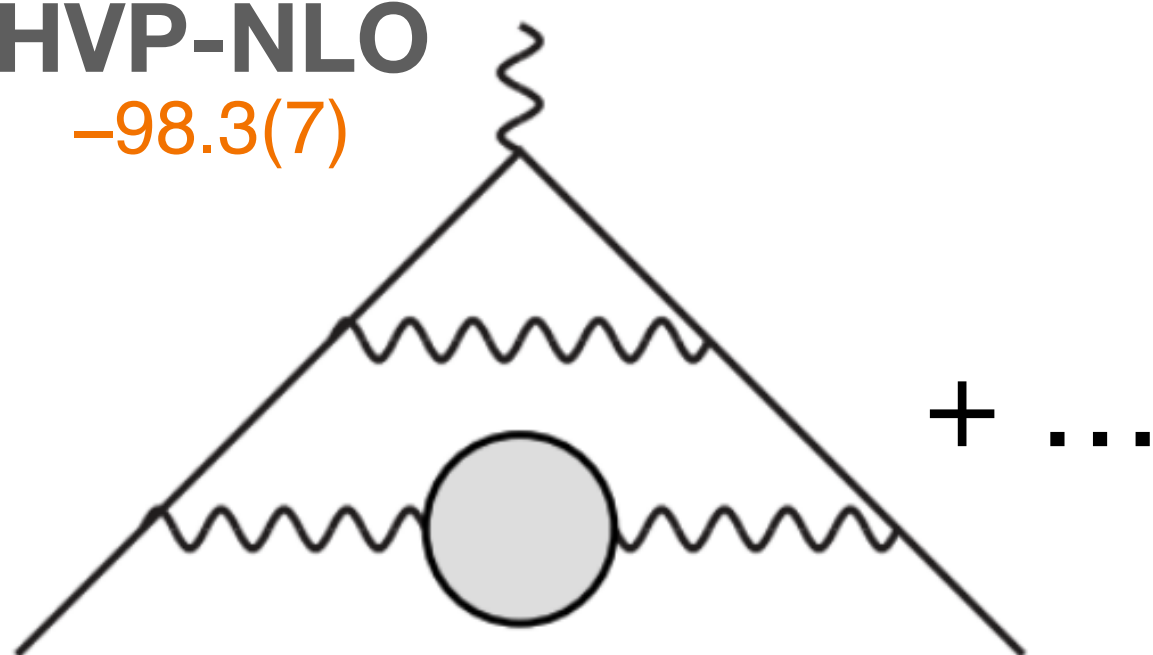
$\mathcal{O}(\alpha^2)$

HVP-LO
6931(40)



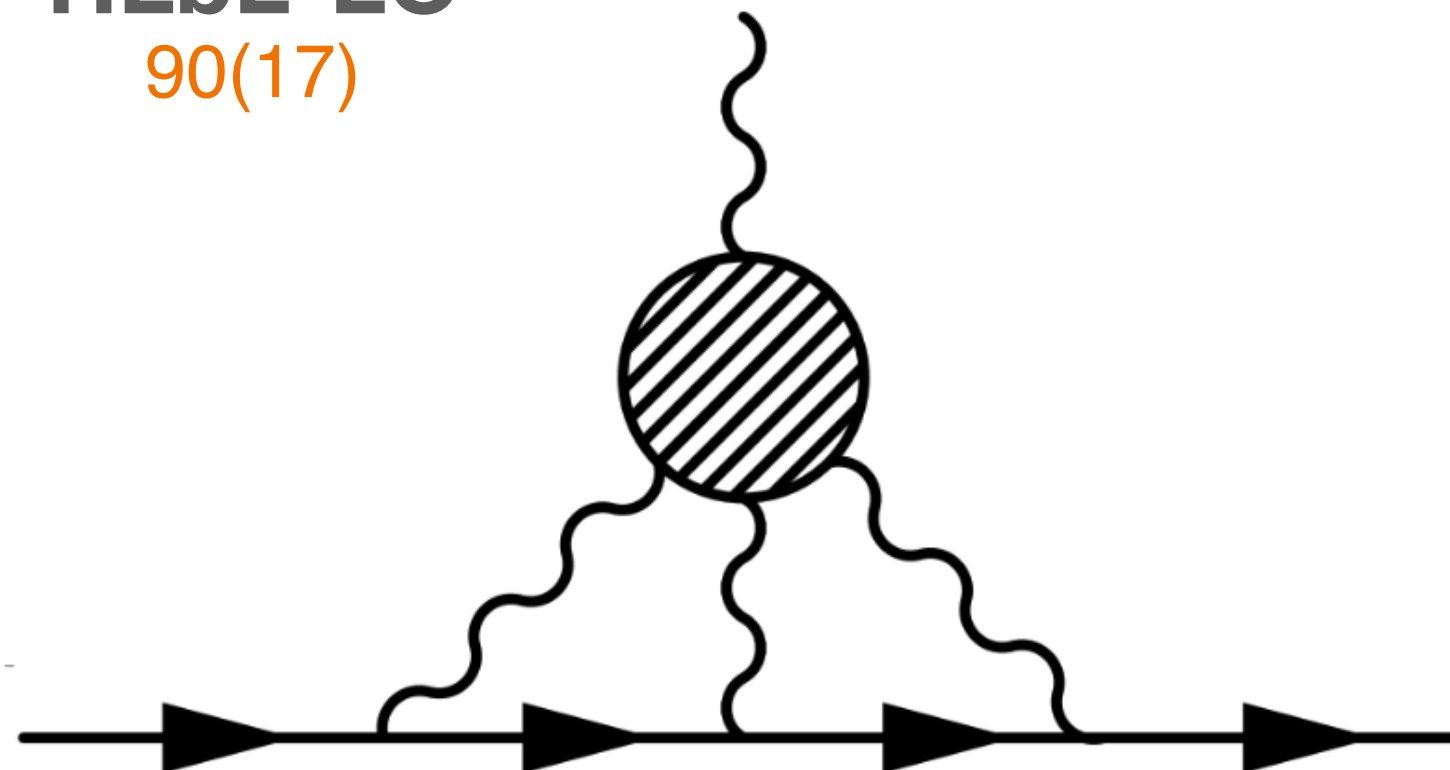
$\mathcal{O}(\alpha^3)$

HVP-NLO
-98.3(7)



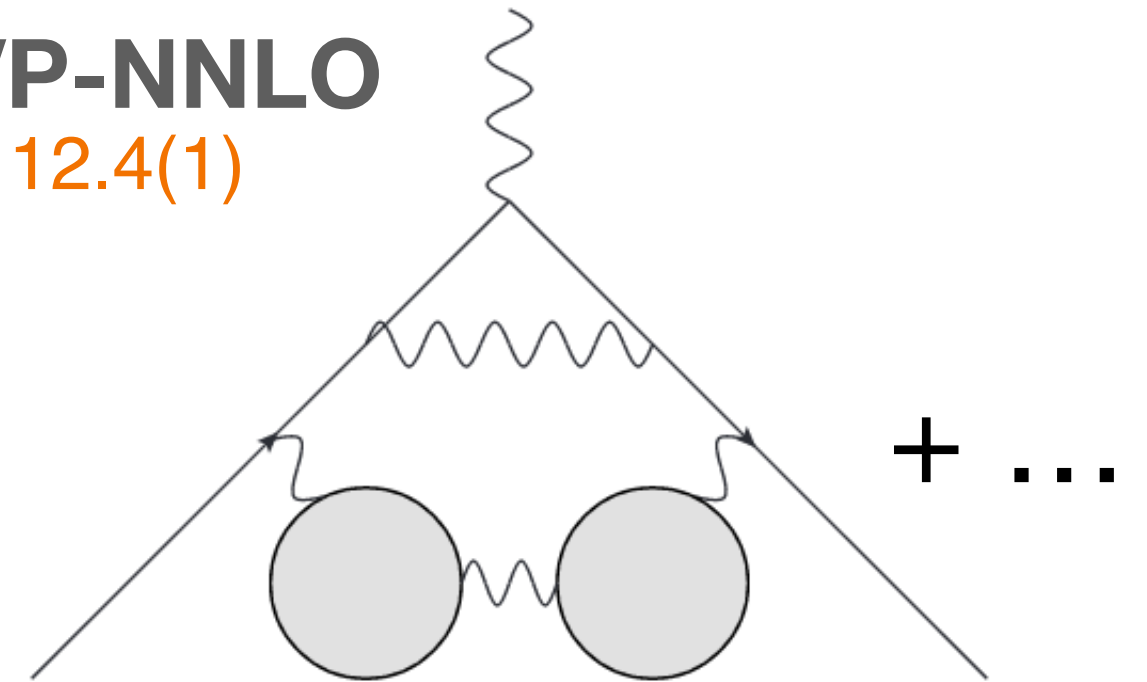
+ ...

HLbL-LO
90(17)



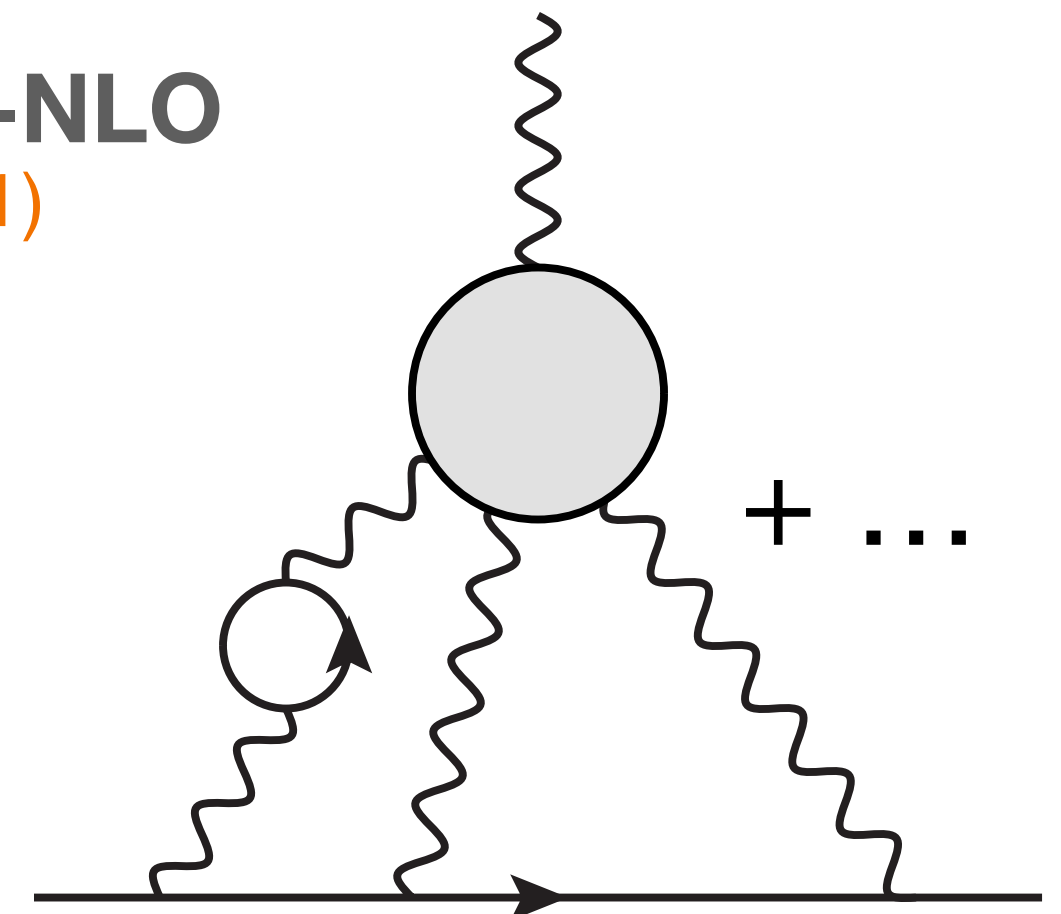
$\mathcal{O}(\alpha^4)$

HVP-NNLO
12.4(1)



+ ...

HLbL-NLO
2(1)



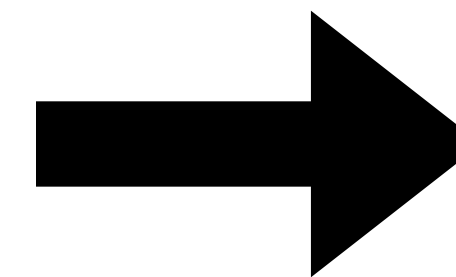
+ ...

VMD & LMD

$$F_{P\gamma\gamma}^{\text{LMD}}(q_1^2, q_2^2) = \frac{\alpha M_V^4 + \beta(q_1^2 + q_2^2)}{(q_1^2 - M_V^2)(q_2^2 - M_V^2)}$$

and

$$F_{P\gamma\gamma}^{\text{VMD}}(q_1^2, q_2^2) = F_{P\gamma\gamma}^{\text{LMD}}(q_1^2, q_2^2) |_{\beta=0}$$



$\tilde{A}^{\text{LMD}}(\tau)$ and $\tilde{A}^{\text{VMD}}(\tau)$ by
inverse Laplace transform

According to VMD, $M_V = M_\rho$,
but left as a free fit parameter
in this work

z -expansion

Conformal transformation:

$$z_k = \frac{\sqrt{t_c - Q_k^2} - \sqrt{t_c - t_0}}{\sqrt{t_c + Q_k^2} + \sqrt{t_c - t_0}}$$

$k \in \{1,2\}$

$t_c = 4m_\pi^2$ given by 2π threshold

t_0 chosen to minimize $|z_k|$ over range of interest

Polynomial fit to cutoff order $N \in \{1,2\}$ used in this work.

Multiplicative factor to match B-L and OPE asymptotics:

$$P(Q_1^2, Q_2^2) = 1 + \frac{Q_1^2 + Q_2^2}{M_V^2}$$

Relevant momenta and relative contribs

Models for $F_{P\gamma\gamma}$

Λ [GeV]	π^0 [LMD+V]	π^0 [VMD]	η [VMD]	η' [VMD]
0.25	14.4 (22.9%)	14.4 (25.2%)	1.8 (12.1%)	1.0 (7.9%)
0.5	36.8 (58.5%)	36.6 (64.2%)	6.9 (47.5%)	4.5 (36.1%)
0.75	48.5 (77.1%)	47.7 (83.8%)	10.7 (73.4%)	7.8 (62.5%)
1.0	54.1 (86.0%)	52.6 (92.3%)	12.6 (86.6%)	9.9 (79.1%)
1.5	58.8 (93.4%)	55.8 (97.8%)	14.0 (96.1%)	11.7 (93.1%)
2.0	60.5 (96.2%)	56.5 (99.2%)	14.3 (98.6%)	12.2 (97.4%)
5.0	62.5 (99.4%)	56.9 (99.9%)	14.5 (100%)	12.5 (99.9%)
20.0	62.9 (100%)	57.0 (100%)	14.5 (100%)	12.5 (100%)

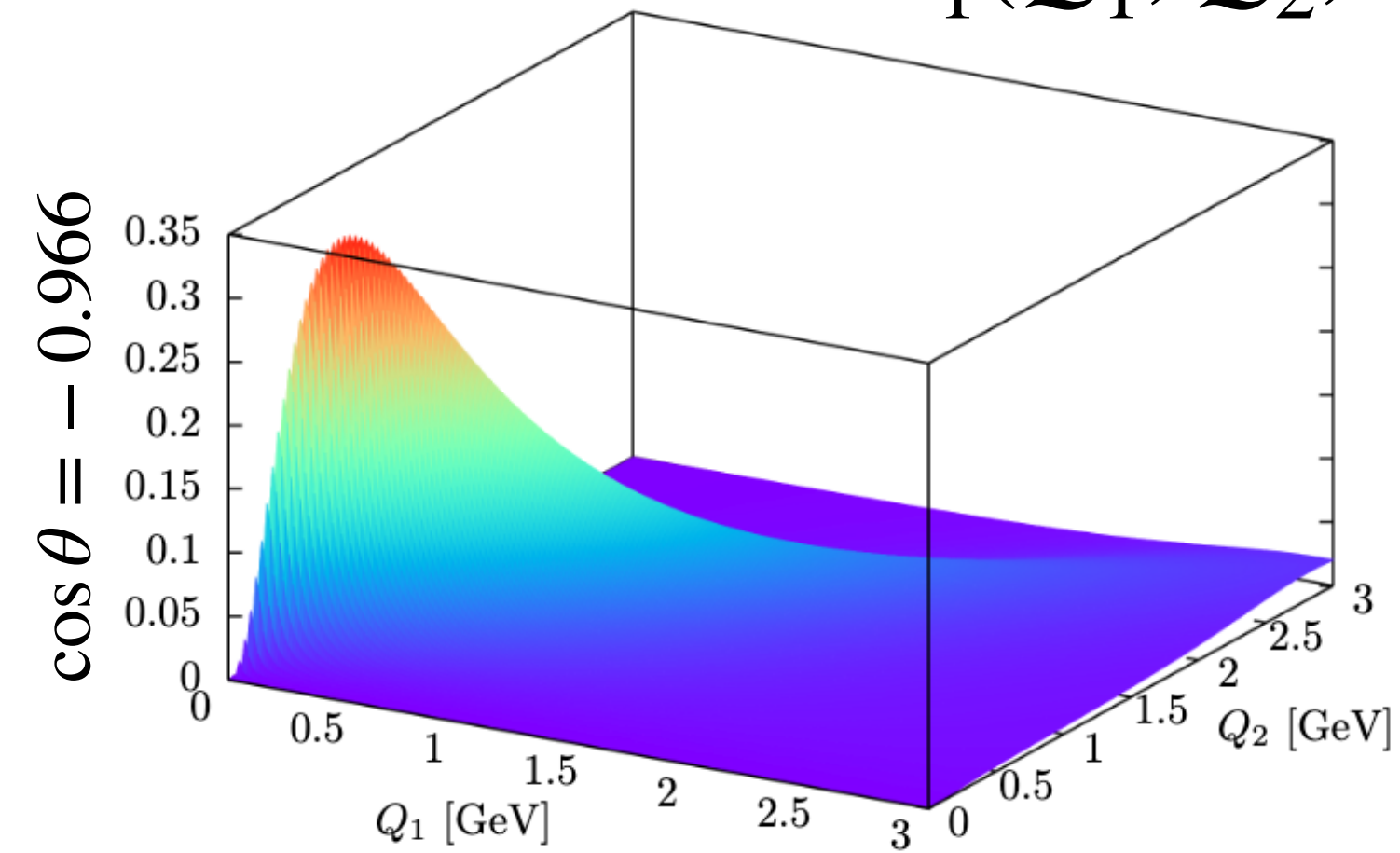
All contribs
saturate
around here

[Table 5 of Nyffeler Phys. Rev. D 94, 053006 (2016)]

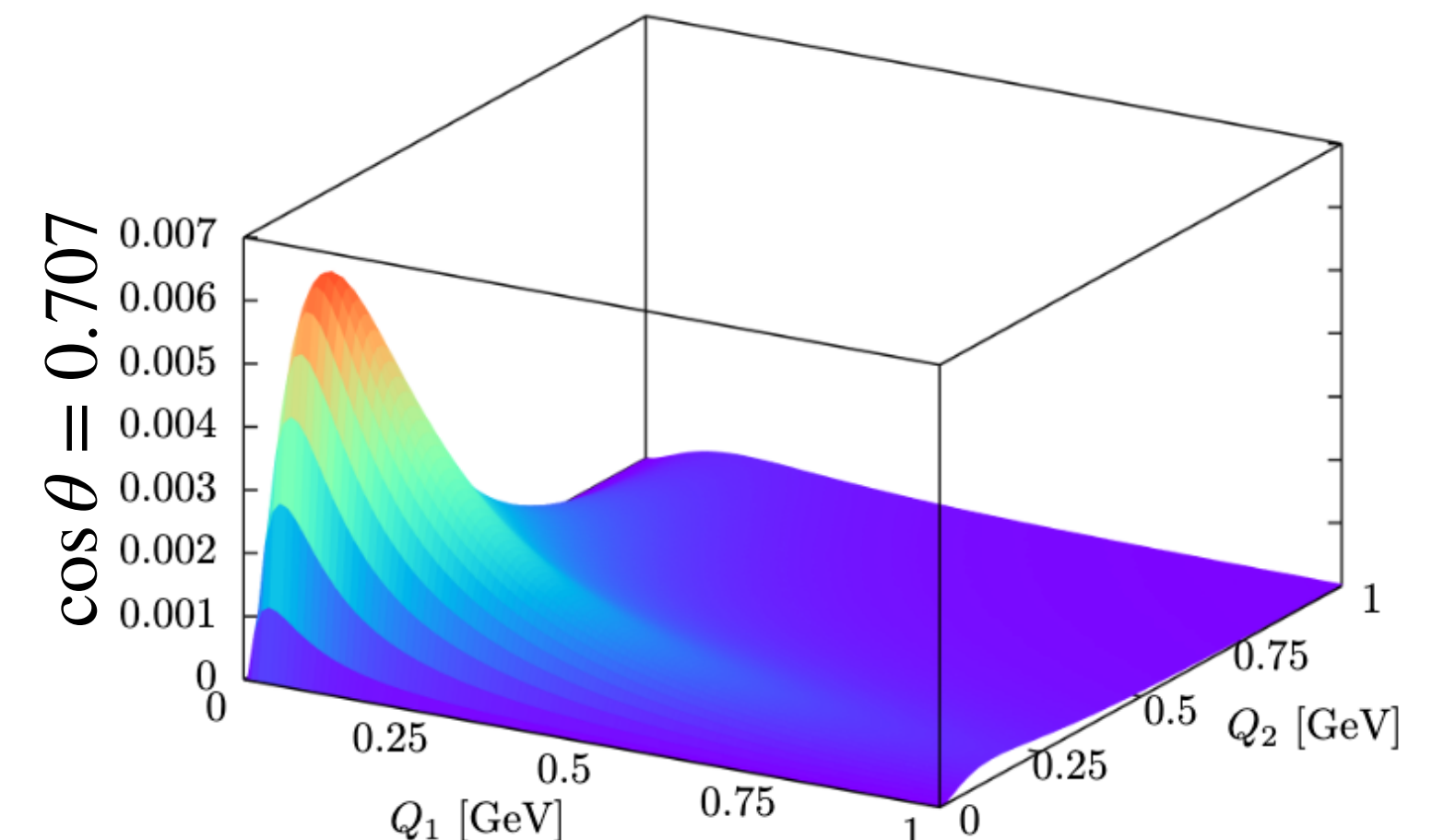
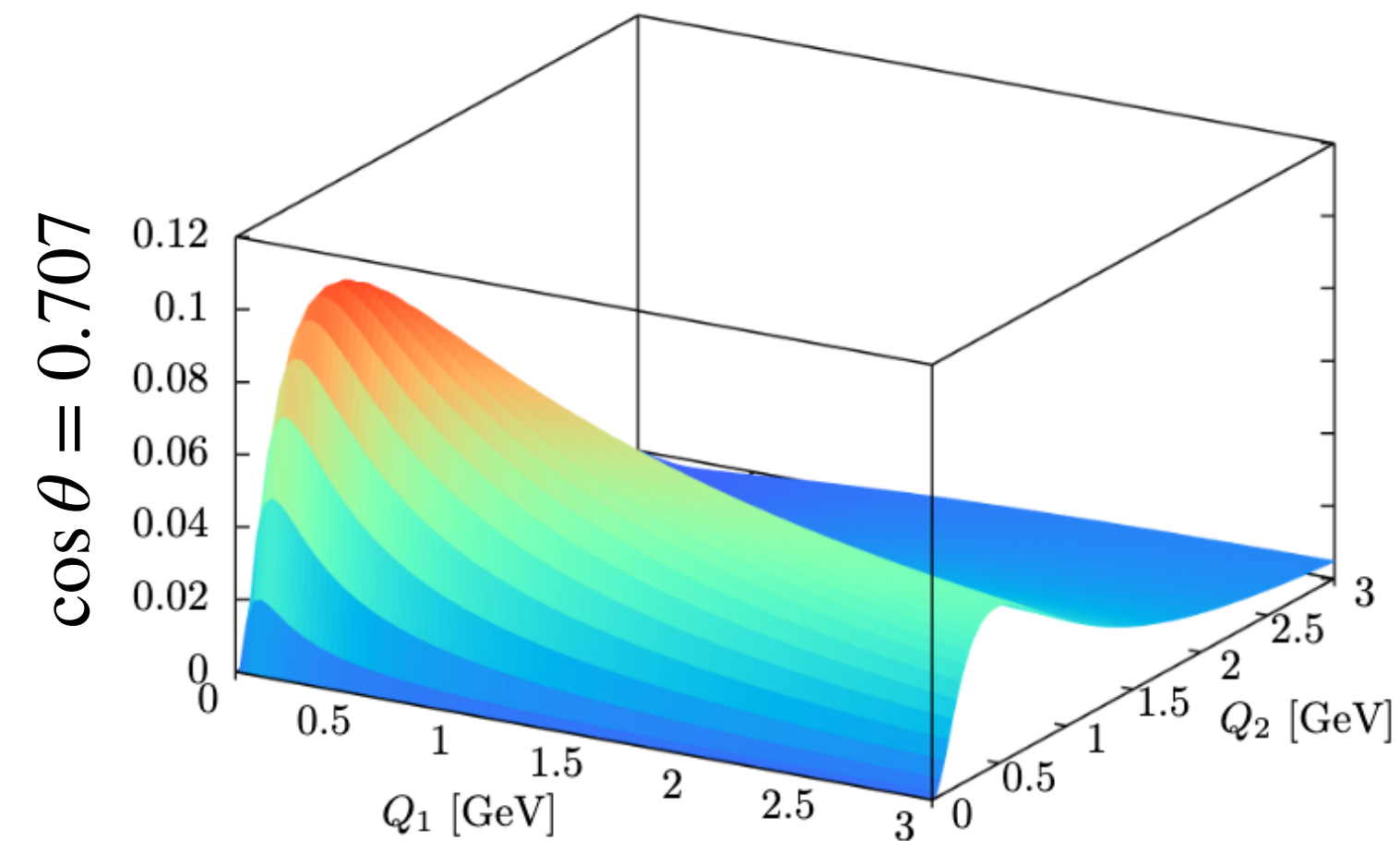
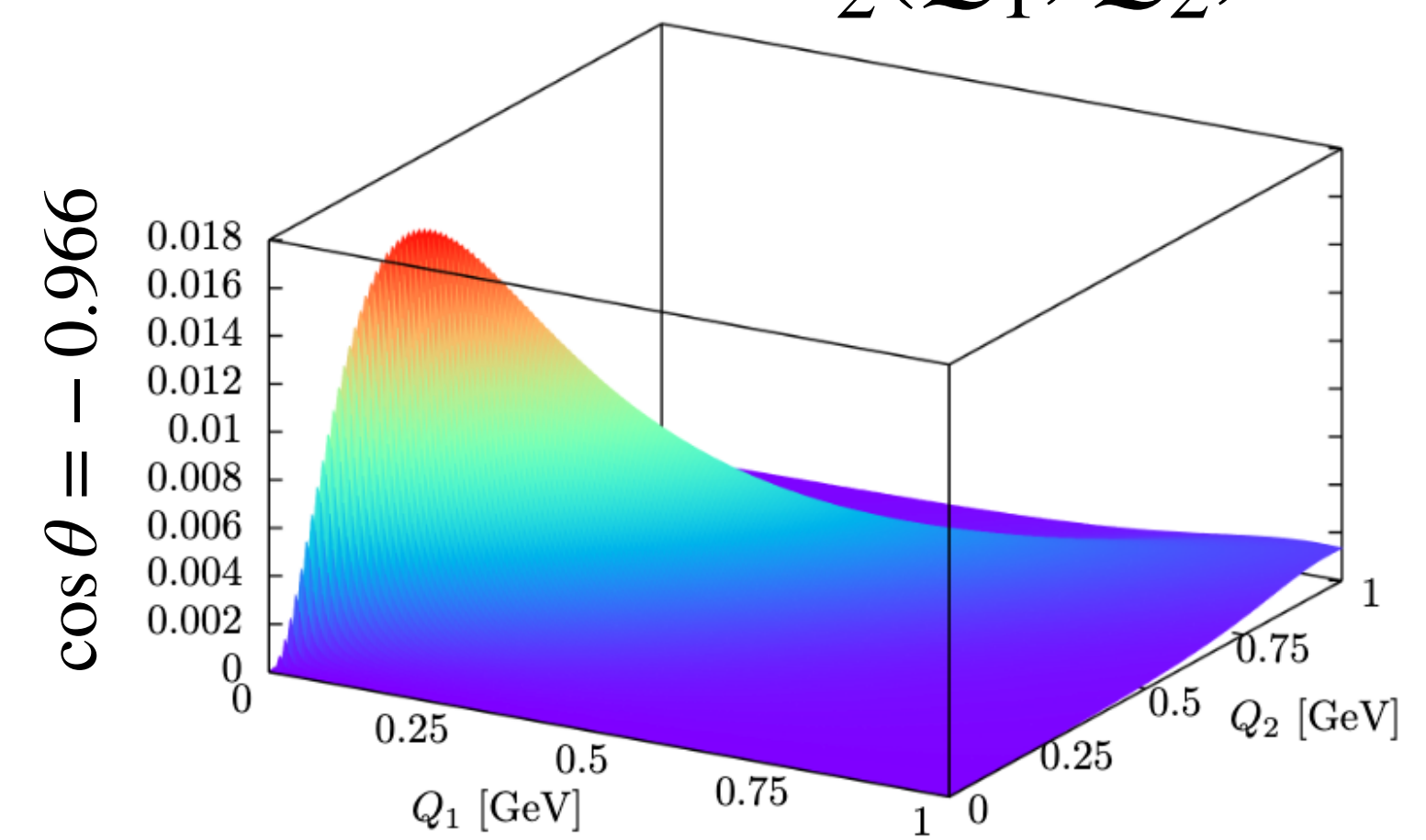
$a_\mu^{P\text{-pole}}$ integrating over
 $0 \leq Q_1, Q_2 \leq \Lambda$

Weight functions: η

$$w_1(Q_1, Q_2, \cos \theta)$$

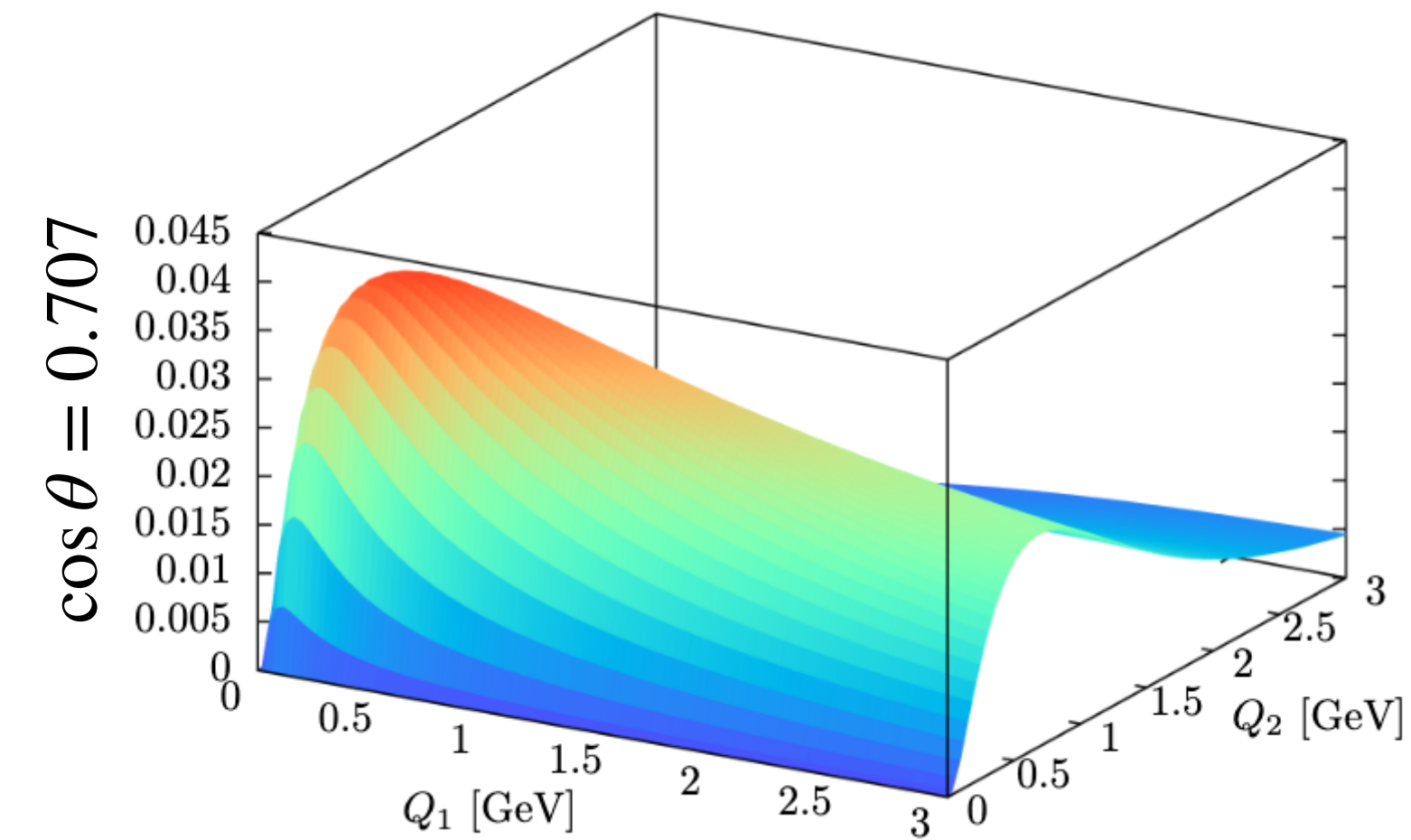
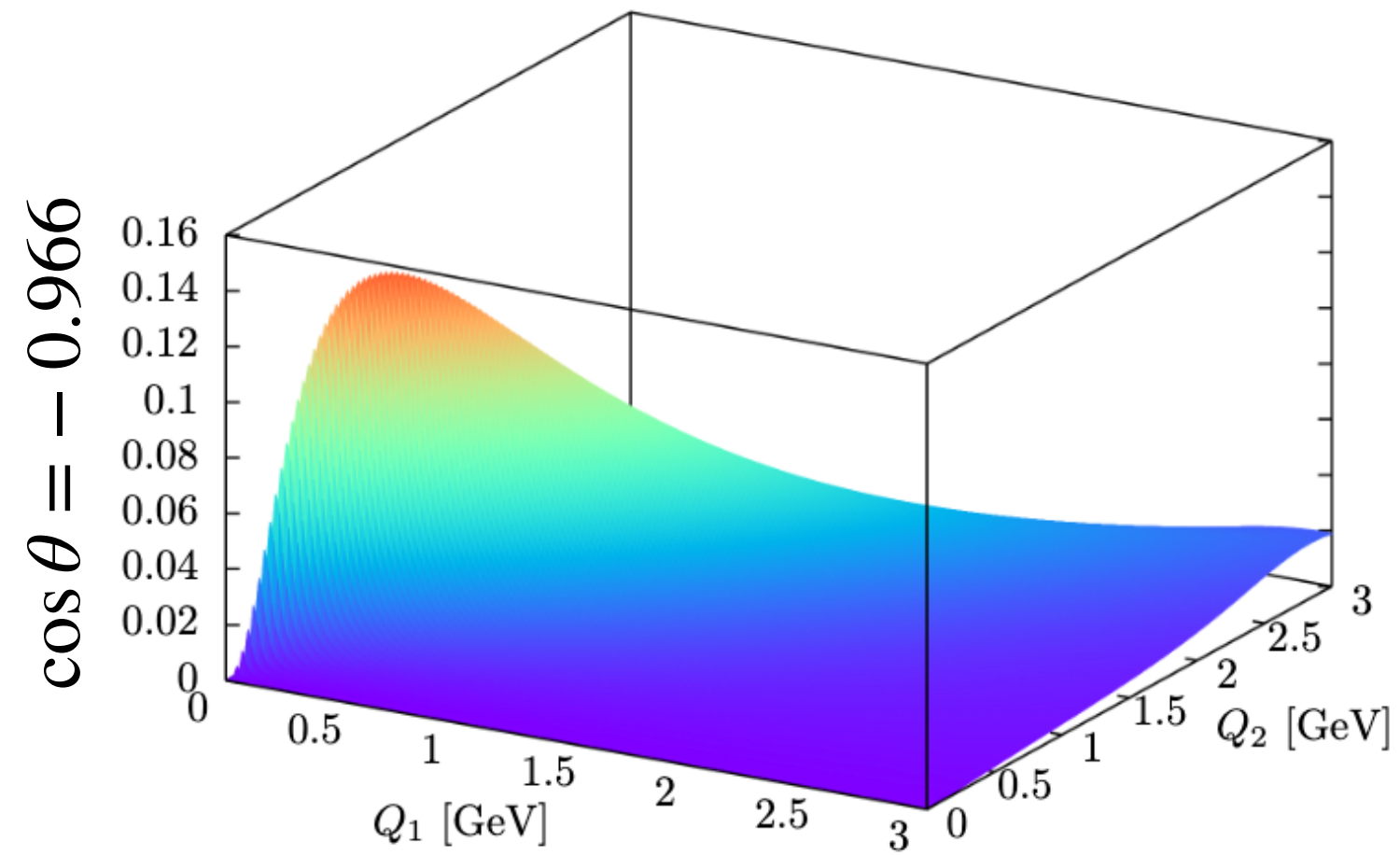


$$w_2(Q_1, Q_2, \cos \theta)$$

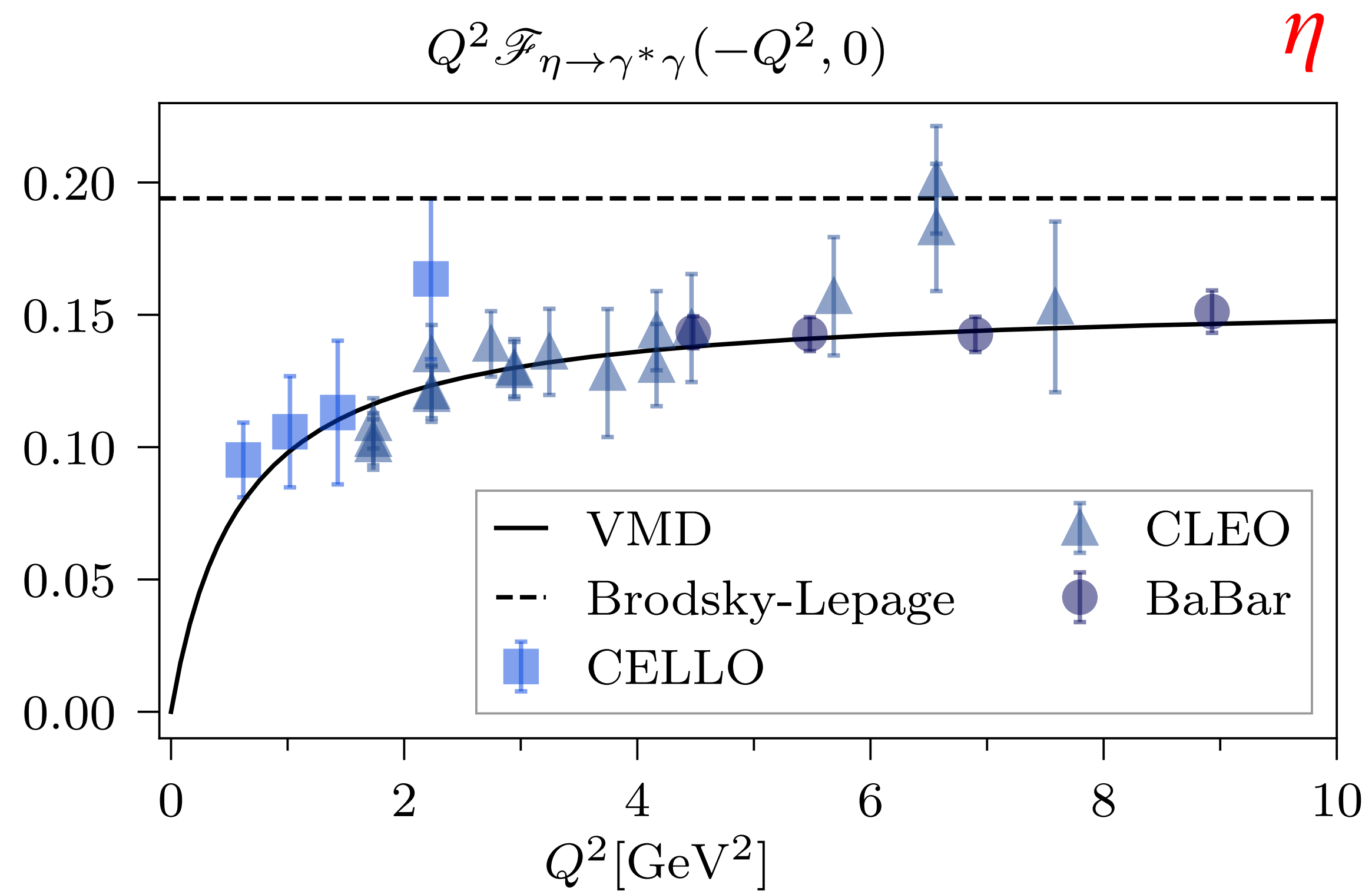


Weight functions: η'

$$w_1(Q_1, Q_2, \cos \theta)$$



TFF data: singly virtual η



Other terms in HLbL decomposition

Contribution	Ref. [6]
π^0, η, η' -poles	93.8(4.0)
π, K -loops/boxes	−16.4(2)
S -wave $\pi\pi$ rescattering	−8(1)
subtotal	69.4(4.1)
scalars	− 1(3)
tensors	
axial vectors	6(6)
u, d, s -loops / short-distance	15(10)
c -loop	3(1)
total	92(19)

Other terms in HLbL decomposition

Contribution	Ref. [6]	
π^0, η, η' -poles	93.8(4.0)	— Addressed by this work
π, K -loops/boxes	−16.4(2)	
S -wave $\pi\pi$ rescattering	−8(1)	
subtotal	69.4(4.1)	
scalars	− 1(3)	
tensors		
axial vectors	6(6)	— AV form factors accessible on the lattice... future work?
u, d, s -loops / short-distance	15(10)	
c -loop	3(1)	Theoretical ambiguities actively being addressed by data-driven community
total	92(19)	