# Adjoint fermions at large- $N_c$ on the lattice



Ken-Ichi Ishikawa, Masanori Okawa

Instituto de Física Teorica UAM-CSIC Hiroshima University

pietro.butti@uam.es



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Pietro Butti (IFT)

Instituto de

niversidad Autónoma

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# Large-N<sub>c</sub> on the lattice



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**1** Twisted reduction at large-*N<sub>c</sub>* 

2 Scale setting w/ Wilson flow

**3** Some results for adjoint fermions

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Apply tBC + volume reduction [González-Arroyo, Okawa 1983,2010] ( $b = \frac{1}{a^2 N_c}$ )

$$\mathcal{S}_{w} = b N \sum_{n,\mu 
eq 
u} \operatorname{tr} ig [ \mathbbm{1} - U_{\mu}(n) U_{
u}(n+\mu) U_{\mu}^{\dagger}(n+
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Apply tBC + volume reduction [González-Arroyo, Okawa 1983,2010] ( $b = \frac{1}{\sigma^2 N_c}$ )

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 (volume reduction  $L=1$ )

Apply tBC + volume reduction [González-Arroyo, Okawa 1983,2010] ( $b = \frac{1}{\sigma^2 N_c}$ )

$$S_{\mathbf{w}} = bN \sum_{\substack{n,\mu \neq \nu}} \operatorname{tr} \left[ \mathbb{1} - U_{\mu}(n) U_{\nu}(n+\mu) U_{\mu}^{\dagger}(n+\nu) U_{\nu}^{\dagger}(n) \right]$$
$$U_{\mu}(n) \rightarrow U_{\mu} \qquad (\text{volume reduction } L = 1)$$
$$U_{\mu}(n+\hat{\nu}) \rightarrow \Gamma_{\nu} U_{\mu} \Gamma_{\nu}^{\dagger} \qquad (\text{twisted BC})$$

• Twist-"eaters":  $\Gamma_{\mu}\Gamma_{\nu} = z_{\nu\mu}\Gamma_{\nu}\Gamma_{\mu}$  with  $z_{\mu\nu} = e^{\frac{2\pi ik}{\sqrt{N_c}}\epsilon_{\mu\nu}}$  [Talk by J. Da Silva yesterday]

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$$\begin{split} \mathcal{S}_{\mathbf{w}} &= bN \sum_{n,\mu \neq \nu} \operatorname{tr} \left[ \mathbb{1} - U_{\mu}(n) U_{\nu}(n+\mu) U_{\mu}^{\dagger}(n+\nu) U_{\nu}^{\dagger}(n) \right] \\ & U_{\mu}(n) \rightarrow U_{\mu} \qquad (\text{volume reduction } L=1) \\ & U_{\mu}(n+\hat{\nu}) \rightarrow \Gamma_{\nu} U_{\mu} \Gamma_{\nu}^{\dagger} \qquad (\text{twisted BC}) \\ & V_{\mu} &= U_{\mu} \Gamma_{\mu} \qquad (\text{change of variable}) \end{split}$$

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• "Effective"-lattice 
$$V = (\sqrt{N_c})^4$$
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• "Effective"-lattice 
$$V = (\sqrt{N_c})^4$$
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• Volume independence:  $z(\mathcal{C}) \langle V(\mathcal{C}) \rangle_{\mathsf{TEK}} \xrightarrow{N \to \infty} \langle W(\mathcal{C}) \rangle_{V \to \infty}$ 

• Apply tBC + volume reduction to both  $U_{\mu}(n)$  and  $\Psi(n)$ 

$$\mathcal{D}_{\mathbf{w}} = \mathbb{1} - \kappa_{\mathsf{adj}} \sum_{\mu} \left[ (\mathbb{1} - \gamma_{\mu}) \mathcal{U}_{\mu}^{\mathsf{adj}} + (\mathbb{1} + \gamma_{\mu}) (\mathcal{U}_{\mu}^{\mathsf{adj}})^{\dagger} 
ight] \quad \mathsf{with} \quad \mathcal{U}_{\mu}^{\mathsf{adj}} \Psi = \mathcal{U}_{\mu} \Psi \mathcal{U}_{\mu}^{\dagger}$$

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• Spectroscopy can be performed using [González-Arroyo, Okawa, 2013]

$$C_{AB}(t) = \frac{1}{2N_c^{\frac{3}{2}}} \sum_{n_0=1}^{\sqrt{N_c}-1} e^{-i(at)\frac{\pi n_0}{a\sqrt{N_c}}} \operatorname{tr}\left(AD_w^{-1}(0)BD_w^{-1}(p_0)\right)$$

• Apply tBC + volume reduction to both  $U_{\mu}(n)$  and  $\Psi(n)$ 

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Several exploratory studies for different number of flavors

 $N_f = 0$ : (fundamental) Meson spectrum at large- $N_c$  [JHEP-04(2021)230 ]  $N_f = \frac{1}{2}$ :  $\mathcal{N} = 1$  SUSY Yang-Mills [Butti et al. JHEP-07(2022)074]  $N_f = 1, 2$ : [Phys.Rev. D 88(2013)], [JHEP-08 (2015)034]

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# Wilson flow on a twisted-reduced lattice

• Define flowed energy density in the twisted-reduced lattice

$$E = -\frac{1}{128} \sum_{\mu,\nu} \operatorname{tr} \left[ z_{\nu\mu} (V_{\nu} V_{\mu} V_{\nu}^{\dagger} V_{\mu}^{\dagger} + V_{\mu} V_{\nu}^{\dagger} V_{\mu}^{\dagger} V_{\nu} + V_{\nu}^{\dagger} V_{\mu}^{\dagger} V_{\nu} V_{\mu} + V_{\mu}^{\dagger} V_{\nu} V_{\mu} V_{\nu}^{\dagger} - \text{h.c.} \right]^{2}$$

• Evolve gauge field according to the evolution equation

$$\frac{\mathrm{d}A_{\mu}(t)}{\mathrm{d}t}=D_{\nu}G_{\mu\nu}(t)$$

• Compute  $\Phi(t) = \left\langle \frac{t^2 E(t)}{N_c} \right\rangle$  and solve

$$\Phi(t)\Big|_{t=t_c} = c$$
 and  $t\frac{\mathrm{d}}{\mathrm{d}t}\Phi(t)\Big|_{t=w_c^2} = c$ 

### Wilson flow on a twisted-reduced lattice



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• At  $V = \infty$ , W. flow defines a renormalized coupling ( $\lambda = g^2 N_c$ ) at  $\mu = 1/\sqrt{8t}$ 

$$\lambda_{
m gf}(\mu)=rac{1}{rac{3}{128\pi^2}rac{N_c^2-1}{N_c^2}}\Phi(t)$$

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• On the reduced torus  $V = 1^4$  with twisted BC [Ibañez, García Pérez, 2019]

$$\hat{\lambda}(\mu) = rac{1}{\mathcal{N}\left(rac{\sqrt{8t}}{\sqrt{N_c}}
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• We define [Butti et al. JHEP-07(2022)074]

$$\hat{\Phi}(t) = \frac{\frac{3}{128\pi^2}}{\mathcal{N}_{\mathsf{lat}}(\sqrt{8t/N_c})} \Phi(t, N_c) \quad \text{for } T \in \left[1.25, \gamma^2 \frac{N_c}{8}\right]$$

• Remnant finite- $N_c$  effects are < 3% in the scaling window.

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# Wilson flow corrected with $\mathcal{N}_{N_c}$



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# The lattice spacing vs N<sub>f</sub>



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• 
$$\bar{m}_q \rightarrow \infty$$
 = Yang-Mills ( $a_{\rm YM}$ )



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- $\bar{m}_q \rightarrow \infty$  = Yang-Mills ( $a_{\rm YM}$ )
- $\bar{m}_a = 0$  = SUSY (asusy computation in [Butti et al., [HEP-07(2022)074])



- $ar{m}_q 
  ightarrow \infty$  = Yang-Mills ( $a_{
  m YM}$ )
- $\bar{m}_q = 0$  = SUSY ( $a_{SUSY}$  computation in [Butti et al., JHEP-07(2022)074])
- Lattice scale depends exponentially from the log of the (subtracted) quark mass

$$a(ar{m}_q) = rac{
ho_1 + 
ho_2 ar{m}_q + 
ho_3 ar{m}_q^2}{rac{
ho_1}{
ho_{ ext{SUSY}}} + 
ho_4 ar{m}_q + rac{
ho_3}{
ho_{ ext{SUSY}}} ar{m}_q^2}$$





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• Parameterize and integrate  $\beta$ -function

$$eta(\lambda)=-rac{b_0\lambda^2}{1-rac{b_1}{b_0^2}\lambda}$$



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Integrate

$$-\lograc{a}{\sqrt{8t_1}}=rac{1}{b_0\lambda}+rac{b_1}{b_0^2}\log\lambda+\log\Lambda$$



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• Define improved couplings

$$\lambda_{I}^{(1)} = rac{1}{P_{\chi}(b)b}$$
 and  $\lambda_{I}^{(2)} = 8(1 - P_{\chi}(b))$ 



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• Fit and extract b<sub>0</sub>

 $24\pi^2 b_0 = 11 - 4N_f = 9 / 11$  vs  $24\pi^2 b_0^{\text{fit}} = 9.76(80) / 10.6(1)$ 



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# Thank you!



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# **Backup material: chiral/SUSY limit**



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#### **Backup material: scales comparison**



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### **Backup material: Critical mass for gluinos**

