QCD mesonic screening masses and restoration of chiral symmetry at high T

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Outline

1. High temperature QCD on the lattice

· Problems and Solutions

2. Screening masses spectrum

- Pseudoscalar-Vector splitting
- Spectrum degeneracy
- Ward Identities and chiral multiplets

The problem of simulating very high temperatures

The problem

- · Renormalize the theory with a hadronic scheme
- Hadronic scale M_{had} and temperature T may differ by orders of magnitude

$$a \ll \frac{1}{T} \ll \frac{1}{M_{\rm had}} \ll L$$

Computationally demanding

A possible solution: finite-volume couplings

Massless scheme: $ar{g}_{
m SF}^2(L_{
m SF})$ at $\mu=1/L_{
m SF}$ [ALPHA Collaboration, 2016-18]

- Set the scale on each lattice with $ar{g}_{ ext{SF}}^2(L_{ ext{SF}})$
- · Explore different scales with step-scaling technique
- Fix the overall scale at low energy with a dimensionful hadronic quantity $M_{
 m had}$ $(r_0,f_\pi,f_{\pi K})$ [M. Bruno et al.,Phys. Rev. D 95 (2017) 074504]



The strategy

Finite volume scheme with SF b.c.

Step-scaling function



Matching with hadronic scheme $(M_{
m had})$



Known relation between

$$M_{\mathrm{had}} \Longleftrightarrow L_{\mathrm{SF}} \Longleftrightarrow \bar{g}_{\mathrm{SF}}(L_{\mathrm{SF}})$$

Finite Temperature with periodic b.c.

Temperature fixed by imposing the relation

$$1/T = L_0 = L_{\rm SF} = 1/\mu$$

Known relation between

$$T \iff \bar{g}_{\rm SF}(L_0)$$

- ➤ No large volumes needed
- ightharpoonup FV effects exponentially small for large LT

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Mesonic screening masses

· Large-distance behaviour of fermionic bilinears

$$C_{\mathcal{O}}(x_3) = \int dx_0 dx_1 dx_2 \langle \mathcal{O}(x)\mathcal{O}(0) \rangle \stackrel{x_3 \to \infty}{=} Ae^{-mx_3} + \dots$$

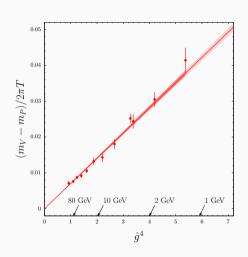
- Here $\mathcal{O}=\{P^a,S^a,V_2^a,A_2^a\}$ from 1 GeV to 160 GeV
- EFT + 1-loop perturbative result ightarrow degeneracy [Laine et al., JHEP 02 (2004) 004]

- \blacktriangleright For asymptotically high temperatures $m \to 2\pi T$ in every channel
- ightharpoonup Probes of chiral symmetry restoration ightharpoonup degeneracy of various channels

Vector-Pseudoscalar mass splitting

- Spin-dependent term of $O(\hat{g}^4)$ in the entire range of temperature
- Mass-splitting visible at the highest temperature

$$\frac{1}{\hat{g}^2(T)} \, \equiv \, \frac{9}{8\pi^2} \ln \frac{2\pi T}{\Lambda_{\overline{\rm MS}}} + \frac{4}{9\pi^2} \ln \left(2 \ln \frac{2\pi T}{\Lambda_{\overline{\rm MS}}} \right)$$

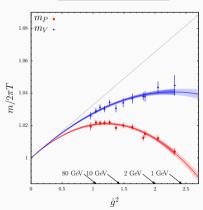


Vector-Pseudoscalar spectrum

$$\begin{cases} \frac{m_P}{2\pi T}(\hat{g}) = p_0 + p_2 \hat{g}^2 + p_3 \hat{g}^3 + p_4 \hat{g}^4 \\ \frac{m_V}{2\pi T}(\hat{g}) = \frac{m_P}{2\pi T} + s_4 \hat{g}^4 \end{cases}$$

- · Results compatible with EFT
- 1-loop matching not reliable:
 - ightharpoonup Pseudoscalar: \hat{g}^4 term cancels \hat{g}^2 term at $T\sim 1~{\rm GeV}$
 - ightharpoonup Vector: \hat{g}^4 term dominant at 1 GeV
 - Mass-splitting at high temperature

| p_3 | 0.0038(22) |
|-------|-------------|
| p_4 | -0.0161(17) |
| s_4 | 0.00704(14) |



Continuum Chiral Ward Identities

Non-singlet axial Ward Identities

$$\mathcal{O} = A^b_{\mu}(z)V^c_{\nu}(y) \longrightarrow \langle V^a_k(z)V^a_k(y)\rangle = \langle A^a_k(z)A^a_k(y)\rangle$$

$$\mathcal{O} = P^b(z)S^0(y) \longrightarrow 2\langle P^a(z)P^a(y)\rangle = -\frac{1}{2}\langle S^0(z)S^0(y)\rangle$$

$$\mathcal{O} = S^b(z)P^0(y) \longrightarrow 2\langle S^a(z)S^a(y)\rangle = -\frac{1}{2}\langle P^0(z)P^0(y)\rangle$$

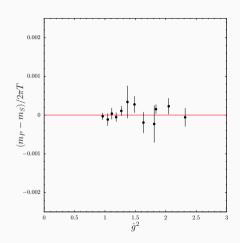
Singlet axial Ward Identities

$$\mathcal{O} = P^{0}(z)S^{0}(y) \longrightarrow \langle S^{0}(z)S^{0}(y)\rangle + \langle P^{0}(z)P^{0}(y)\rangle = N_{f}\langle QP^{0}(z)S^{0}(y)\rangle$$

$$\mathcal{O} = P^{a}(z)S^{a}(y) \longrightarrow \langle S^{a}(z)S^{a}(y)\rangle + \langle P^{a}(z)P^{a}(y)\rangle = N_{f}\langle QP^{a}(z)S^{a}(y)\rangle$$

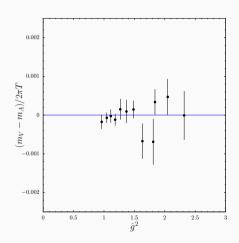
Pseudoscalar-Scalar spectrum

- Complete degeneracy between the Pseudoscalar and the Scalar channels
 - ightharpoonup Only Q=0 sector contributes to the path integral

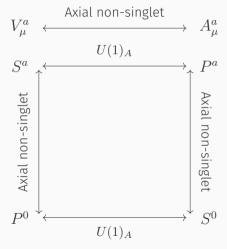


Vector-Axial spectrum

- Complete degeneracy between the Vector and the Axial channels
 - ➤ Non-singlet chiral symmetry restoration



Chiral Multiplets



Conclusion

→ Step-scaling technique provides a solid strategy to study QCD at very high temperature

 $ightharpoonup O(\hat{g}^4)$ terms needed to explain the Vector-Pseudoscalar spectrum in the entire range of temperature

- ➤ No signal of chiral symmetry breaking
- Degeneracy pattern consistent with Ward Identities