

Novel Bottomonium Spectral Results with Backus-Gilbert

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Outline:

- 1 Spectral Reconstruction in Non-Relativistic (NR)QCD
- 2 The Backus-Gilbert Method
- 3 Laplace Shifting
- 4 Conditioning and Studies into the Covariance Matrix

Spectral Reconstruction in NRQCD

The spectral representation of the Euclidean correlation function $G_{\Gamma}(\tau, \mathbf{x}) = \langle O_{\Gamma}(\tau, \mathbf{x}) O_{\Gamma}(0, \mathbf{0}) \rangle$ is

$$G(\tau; T) = \int_0^{\infty} \rho(\omega; T) K(\omega, \tau) d\omega \quad (1)$$

where $K(\omega, \tau; T) = \cosh[\omega(\tau - 1/2T)] / \sinh(\omega/2T) \rightarrow e^{-\omega\tau}$ when $M \gg T$.

Aim: Invert the Laplace transform

$$G(\tau; T) = \int_0^{\infty} \rho(\omega; T) e^{-\omega\tau} \quad (2)$$

\implies This is ill-posed!

One method: Backus-Gilbert¹

¹Backus, Gilbert 1968, *Geophysical Journal International*

Construct sampling functions $A(\omega, \omega_0) = \sum_{\tau} c_{\tau}(\omega_0)K(\omega, \tau)$ which sample $\rho(\omega)$ at a particular value ω_0 :

$$\begin{aligned}\hat{\rho}(\omega_0) &= \sum_{\tau} c_{\tau}(\omega_0)G(\tau) \\ &= \int A(\omega, \omega_0)\rho(\omega)d\omega\end{aligned}\tag{3}$$

$c_{\tau}(\omega_0)$ chosen to minimise the 'width' of $A(\omega, \omega_0)$ (i.e. maximise resolving power).

Our choice of width:²

$$J(\omega_0) = \int [A(\omega, \omega_0) - \delta(\omega - \omega_0)]^2 d\omega\tag{4}$$

²Oldenburg 1984, IEEE GE-22 vol. 6

Minimising $J(\omega_0)$ requires inversion of a near-singular matrix \mathcal{K} :

$$\mathcal{K}_{\tau\tau'} = \int K(\omega, \tau)K(\omega, \tau') d\omega = \int e^{-\omega(\tau+\tau')} d\omega \quad (5)$$

but $0 \leq \tau, \tau' < 128$ for our lattices \Rightarrow This requires conditioning.

Our approach:

$$\mathcal{K}(\alpha) = \mathcal{K} + \alpha I \quad (6)$$

previously used³:

$$\mathcal{K}(\alpha) = \alpha\mathcal{K} + (1 - \alpha)\text{Cov}[G] \quad (7)$$

³FASTSUM 2021, 2112.02075, 2112.04201

- The functions $A(\omega, \omega_0)$ replicate $\delta(\omega - \omega_0)$ *imperfectly* – $A(\omega, \omega_0)$ narrower for small ω_0
- $A(\omega, \omega_0)$ generated in absence of $G(\tau)$

Idea: Use Laplace shift transform to ‘increase’ resolution about a feature⁴:

$$e^{\Delta \cdot \tau} G(\tau) \xrightarrow{\mathcal{L}} \rho(\omega - \Delta) \quad (8)$$

Drawback: Statistical errors amplified

⁴FASTSUM 2021, 2112.02075, 2112.04201

Laplace-Shifting: An example

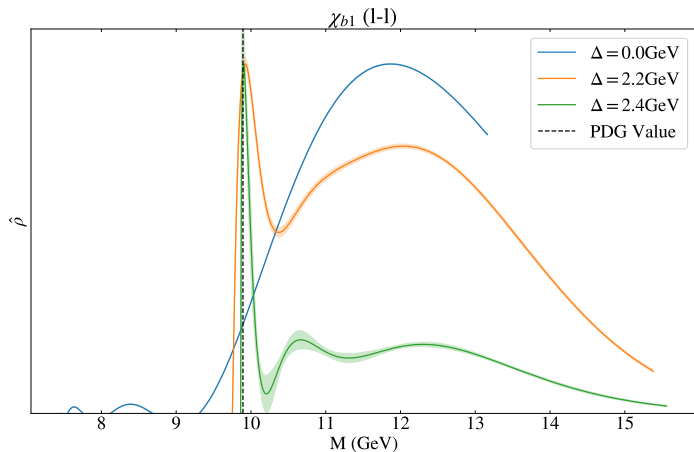
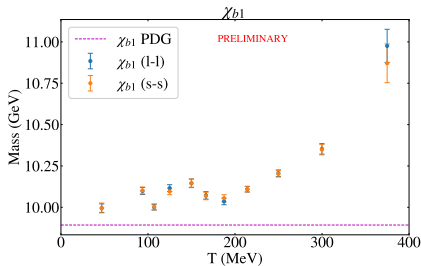
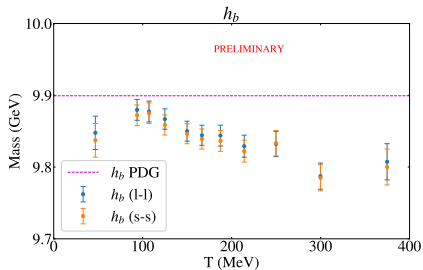
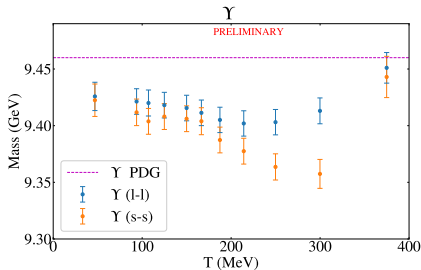
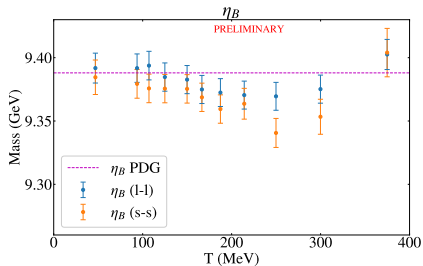


Figure: Laplace-shift in action using correlation functions for the χ_{b1} .

Results from the lattice - $b\bar{b}$ sector



Results from the lattice – Υ

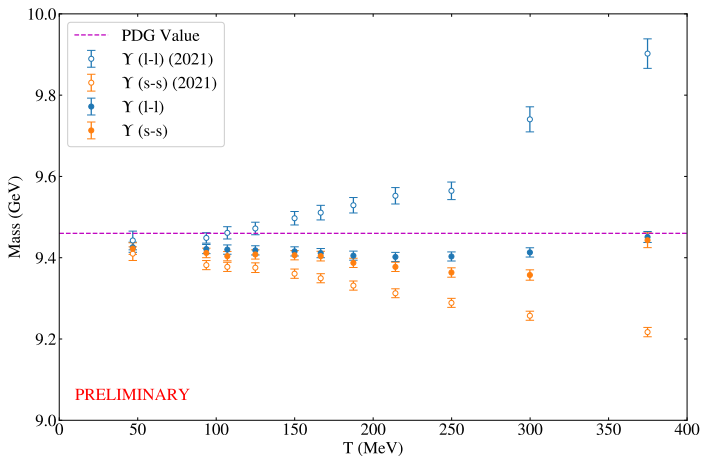


Figure: Υ mass comparison for local and smeared sources, 2021 v. 2022

Results from the lattice – Υ

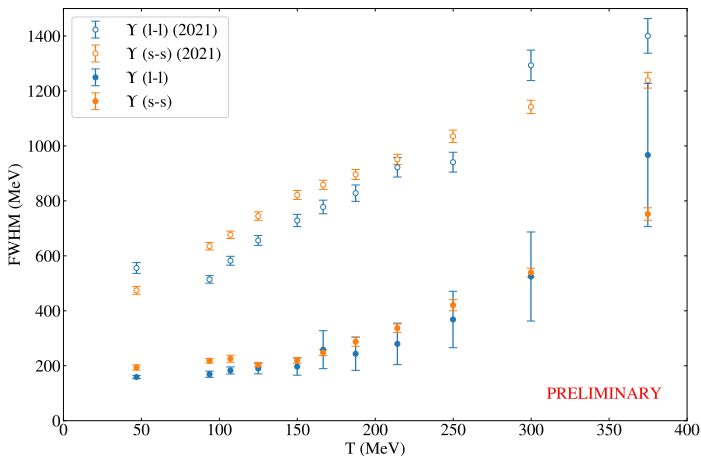


Figure: Υ FWHM comparison for local (l-l) and smeared (s-s) sources

Results from the lattice - Υ

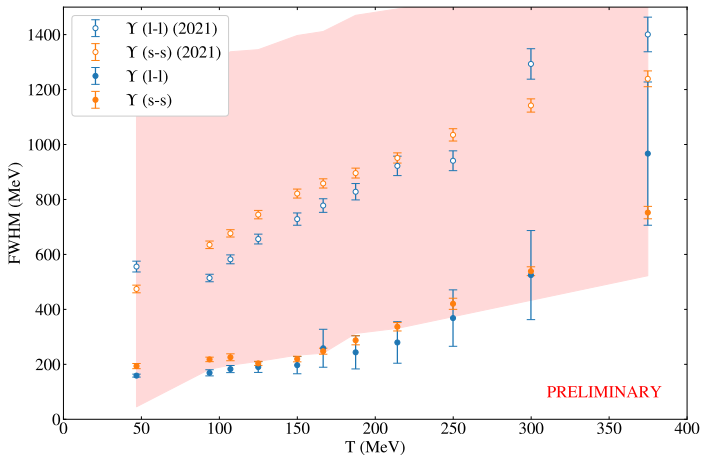


Figure: Υ FWHM comparison with resolution band

- $c_\tau(\omega_0)$ calculation requires inversion of an ill-conditioned matrix \mathcal{K}
- \mathcal{K} may also be conditioned via Tikhonov-like approach with the covariance matrix Σ :

$$\mathcal{K}(\alpha) = \alpha\mathcal{K} + (1 - \alpha)\text{Cov}[G] \quad (9)$$

Question:

\implies How does the condition of $\text{Cov}[G]$ change with Δ ?

Studies into the covariance matrix

- $\text{Cov}[G]$ transforms as $e^{\Delta \cdot \tau} \text{Cov}[G]_{\tau\tau'} e^{\Delta \cdot \tau'}$ under Laplace shift
- For τ, τ' large, expect $\text{Cov}[G]$ scales as

$$\text{Cov}[G]_{\tau\tau'} \sim G(\tau) \cdot G(\tau') \sim e^{-M_{\text{ground}}(\tau+\tau')} \quad (10)$$

where M_{ground} is the mass of the ground state.

Principle:

\implies When $\Delta = M_{\text{ground}}$, $\text{Cov}[G]_{\tau\tau'} \sim \mathcal{O}(1)$ and $\text{cond}(\text{Cov}[G]) \rightarrow \infty$

Studies into the covariance matrix

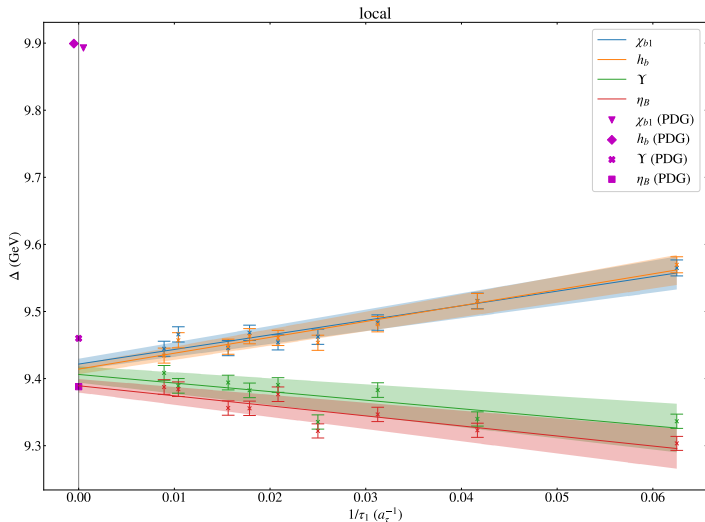


Figure: Covariance mass prediction from local sources

Studies into the covariance matrix

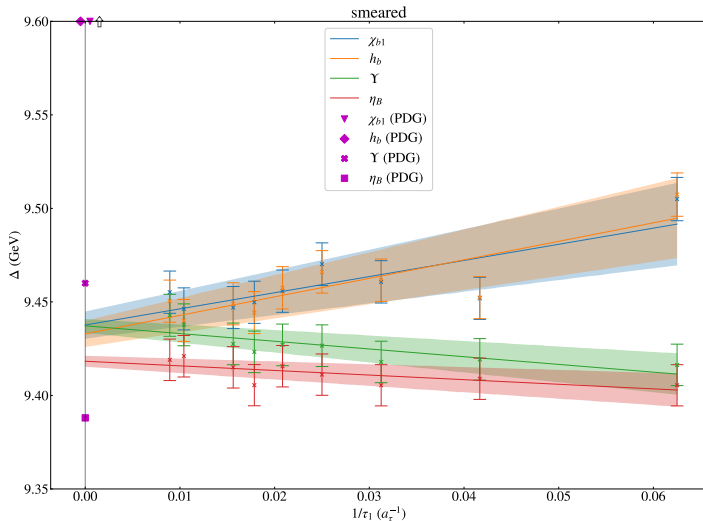


Figure: Covariance mass prediction from smeared sources

Studies into the covariance matrix

In general, $\Delta \neq M_{\text{ground}}$ but instead $\Delta \leq M_{\text{ground}}$

\implies Parisi-Lepage scaling⁵

e.g. for Υ , $\text{Cov}[G]$ includes $\Upsilon\Upsilon$ correlation terms – conservation of angular momentum allows construction of $\eta_b\eta_b$ terms which scale with mass $M_{\eta_B} < M_{\Upsilon}$.

\implies We can probe M_{η_b} from S/P-wave vector/pseudoscalar quantum numbers!

\implies Laplace shift limited by smallest sector mass

⁵Parisi 1983, Lepage 1989

Conclusions and closing comments

- New results from Tikhonov-conditioned Backus-Gilbert for multiple heavy sector meson channels
- Laplace shifting gives improved results when paired with Backus-Gilbert
- Change in resolving power shows systematic uncertainties are still not controlled
- Studies into Laplace shifting $\text{Cov}[G]$ support Parisi-Lepage scaling for heavy sector
- **Question:** Does Laplace shifting offer improvements for other recon. techniques?

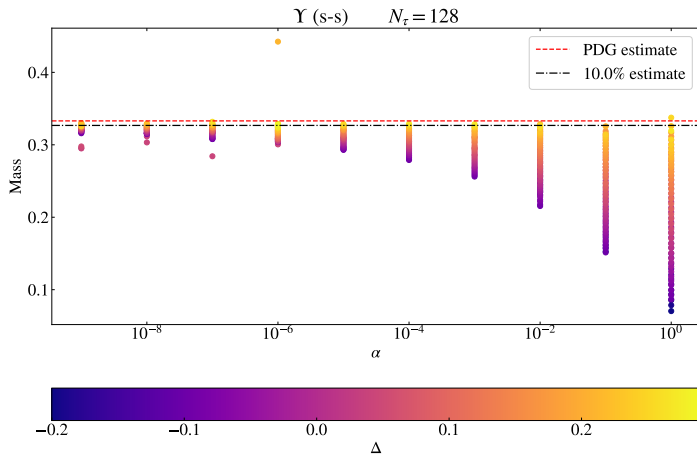
Thank you for your time.

Questions are welcome!

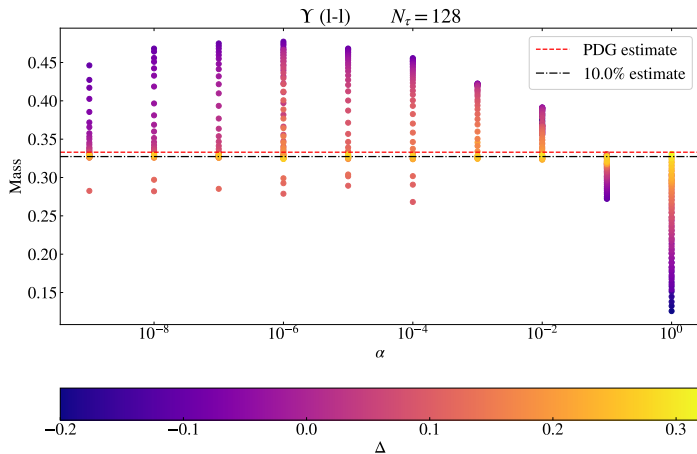
or contact me via:
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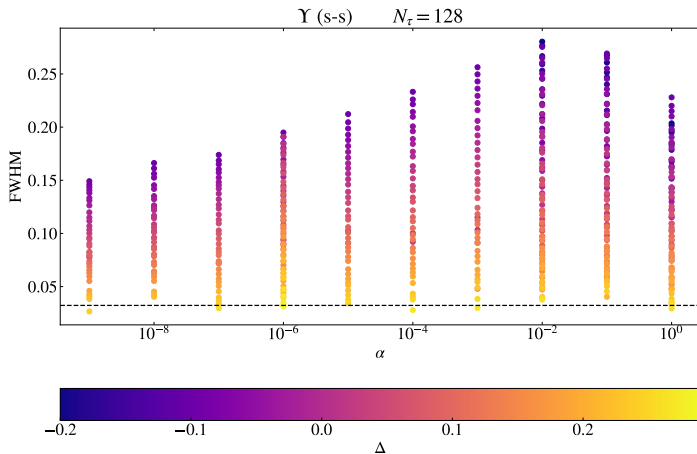
Fitting the mass



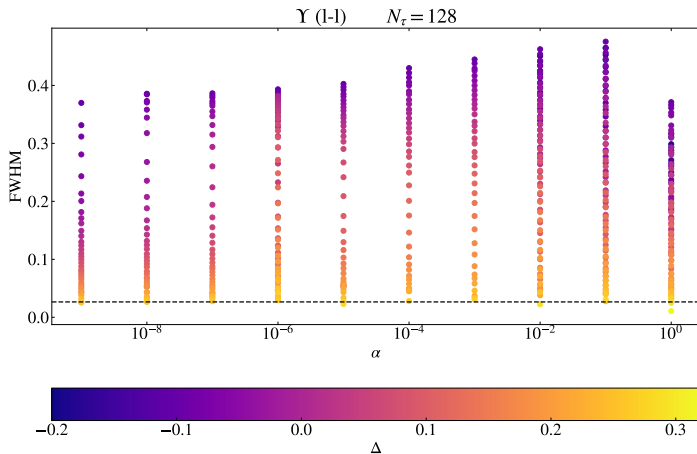
Fitting the mass



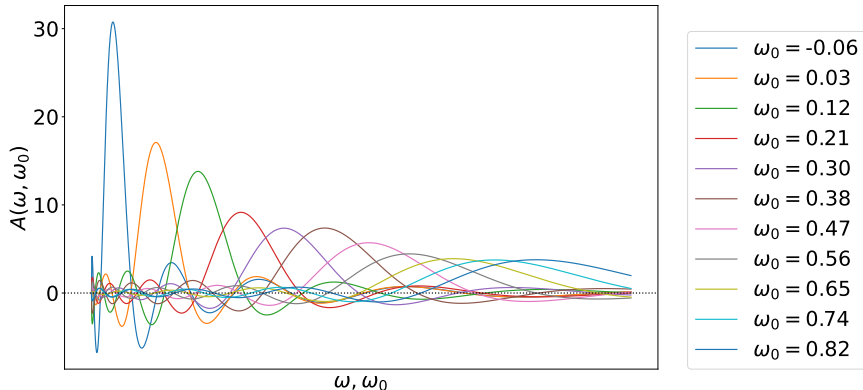
Fitting the width



Fitting the width



Averaging functions



χ_{b1} widths

