

Estimation of the photon production rate using imaginary momentum correlators

Csaba Török

in collaboration with:

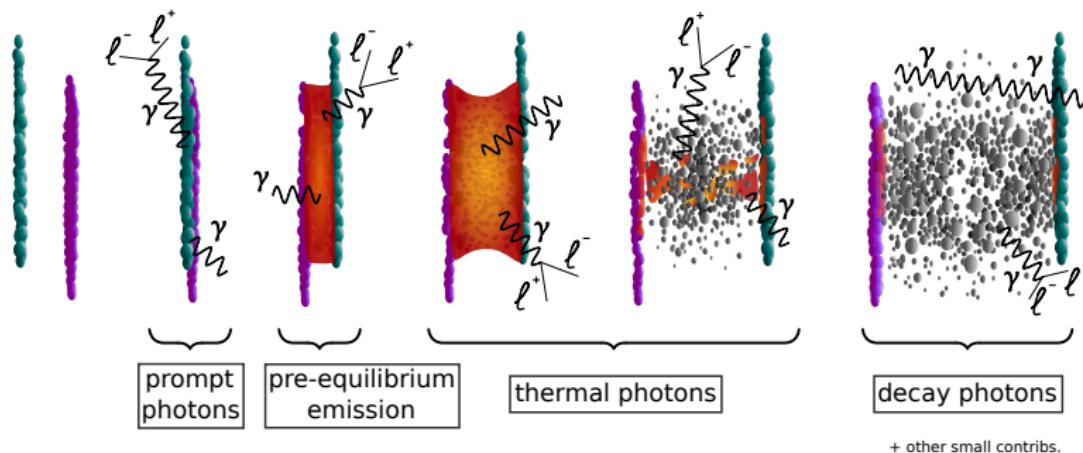
Marco Cé, Tim Harris, Ardit Krasniqi, Harvey Meyer, Samuel Ruhl

39th International Symposium on Lattice Field Theory, Bonn, 2022



Motivation

- electromagnetic probes (photons, dileptons) interact weakly with the QGP medium → on-going experimental research (RHIC, LHC, GSI)



$$\text{direct photons} = \text{total} - [\text{decay photons}]$$

- discrepancies: direct photon excess at low p_T see e.g. Gale, Paquet, Schenke, Shen '21
- dominant contribution at low p_T : **thermal photons**

Correlators at imaginary spatial momentum

- theory input: thermal photon emission rate \longrightarrow thermal photon yield
- model input: hydro

Thermal photon emission rate per unit volume of the QGP:

$$\frac{d\Gamma_\gamma(\omega)}{d\omega} = \frac{\alpha_{\text{em}}}{\pi} \frac{1}{e^{\omega/T} - 1} \omega \sigma(\omega) + \mathcal{O}(\alpha_{\text{em}}^2)$$

where $\sigma(\omega) \equiv \rho_T(\omega, k = \omega)$, $\rho_T = \frac{1}{2}(\delta^{ij} - k^i k^j / k^2)\rho_{ij}(\omega, \vec{k})$

Correlators at imaginary spatial momentum

- theory input: thermal photon emission rate \longrightarrow thermal photon yield
- model input: hydro

Thermal photon emission rate per unit volume of the QGP:

$$\frac{d\Gamma_\gamma(\omega)}{d\omega} = \frac{\alpha_{em}}{\pi} \frac{1}{e^{\omega/T} - 1} \omega \sigma(\omega) + \mathcal{O}(\alpha_{em}^2)$$

where $\sigma(\omega) \equiv \rho_T(\omega, k = \omega)$, $\rho_T = \frac{1}{2}(\delta^{ij} - k^i k^j / k^2)\rho_{ij}(\omega, \vec{k})$

Meyer '18

$$\begin{aligned} H_E(\omega_n) &\equiv G_E^T(\omega_n, k = i\omega_n), \quad \omega_n = 2n\pi T \\ &= -\frac{1}{2} \sum_{i=1}^2 \int_0^\beta dx_0 \int d^3x e^{i\omega_n x_0} e^{-i(i\omega_n)x_3} \langle J_i(x) J_i(0) \rangle \end{aligned}$$

Dispersion relation:

$$H_E(\omega_n) = -\frac{\omega_n^2}{\pi} \int_0^\infty \frac{d\omega}{\omega} \frac{\sigma(\omega)}{\omega^2 + \omega_n^2}$$

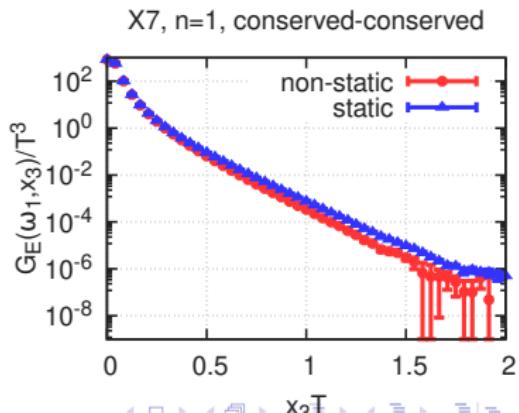
Lattice estimator of H_E

$$\begin{aligned} H_E(\omega_n) &= - \sum_{i=1}^2 \int_0^\beta dx_0 \int d^3x \left(e^{i\omega_n x_0} - e^{i\omega_n x_2} \right) e^{\omega_n x_3} \langle J_i(x) J_i(0) \rangle \\ &= - \int_{-\infty}^{\infty} dx_3 e^{\omega_n x_3} \left[\underbrace{G_s(\omega_n, x_3)}_{\text{static}} - \underbrace{G_{ns}(\omega_n, x_3)}_{\text{non-static}} \right] \end{aligned}$$

- subtraction term: makes the integral finite by power counting
: vanishes in the continuum limit
- $H_E(\omega_n) < 0$

Strategy:

- measure the screening correlators G_{ns} and G_s
- fit their tails with single state fits
- model the tail of their difference



Lattice setup

ensembles:

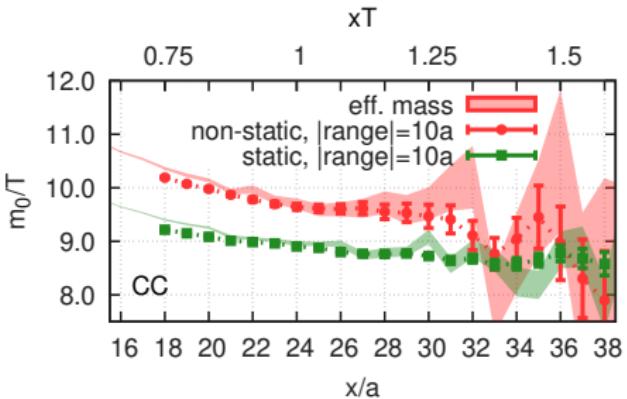
- $N_f = 2$, clover-improved Wilson fermions
- $m_\pi \approx 280$ MeV
- $T \approx 250$ MeV
- three lat. spacings: 0.05, 0.04, 0.033 fm (volume $\sim (3.1 \text{ fm})^3$)
 $64^3 \times 16, 80^3 \times 20, 96^3 \times 24$
- 1500-2000 configurations with 64 point sources

observable:

- $I = 1$ contribution
- different discretizations using the local/conserved current
- focus on the $n = 1$ & $n = 2$ Matsubara sectors

Modelling the tail of the integrand

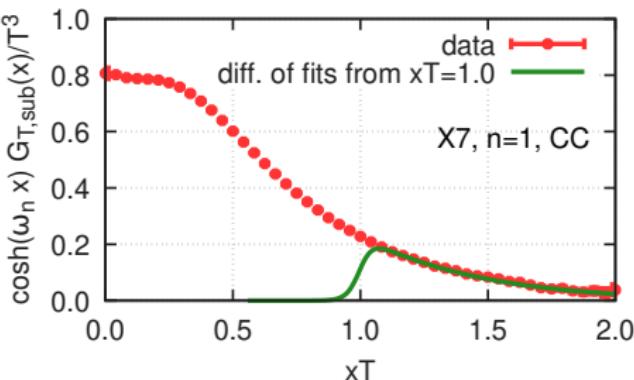
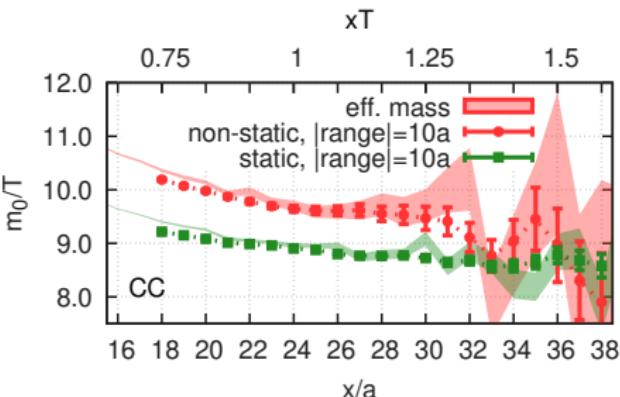
- single state fits on the screening correlators:
 $A \cosh[m_0(x - L/2)]$
- plateau regions consist of only a few points
- three representatives from AIC-based histogram



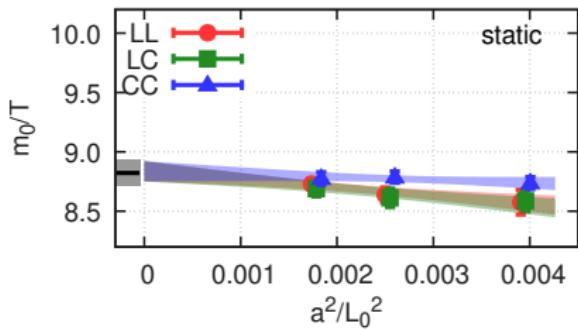
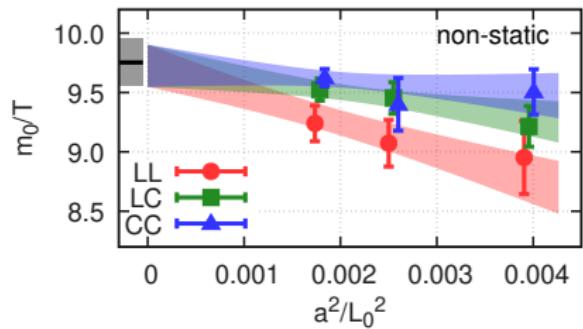
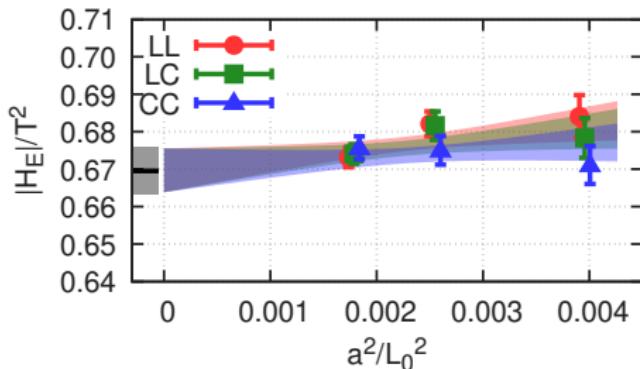
Modelling the tail of the integrand

- single state fits on the screening correlators:
 $A \cosh[m_0(x - L/2)]$
- plateau regions consist of only a few points
- three representatives from AIC-based histogram

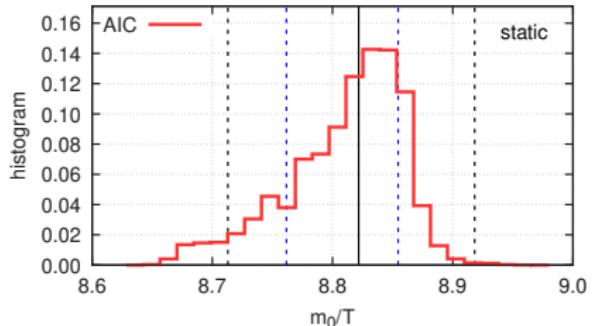
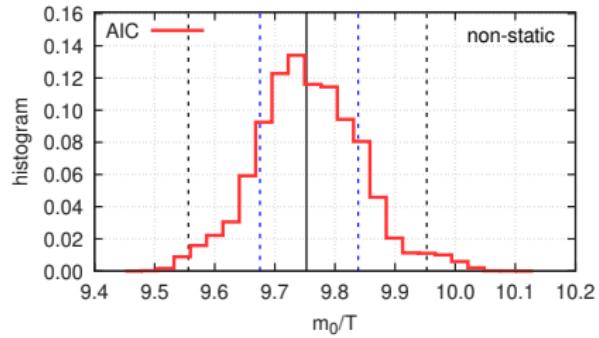
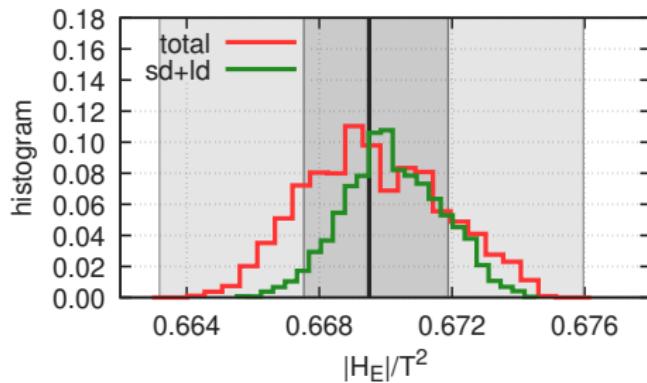
- modelling the tail reduces errors by a factor of ~ 2.5 on the coarsest ensemble
- different switching points give consistent total integral results



Continuum extrapolations



Histograms of continuum extrapolated results



Results for first Matsubara sector ($n = 1$)

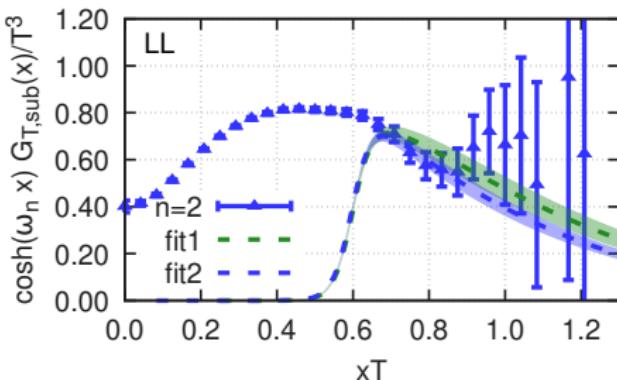
observable	value at the continuum limit
$ H_E /T^2$	0.670(6) _{stat} (2) _{sys}
$\text{sd}(x_w T = 1)$	0.579(3) _{stat} (1) _{sys}
$\text{ld}(x_w T = 1)$	0.091(5) _{stat} (2) _{sys}
$m_0^{\text{non-static}}/T$	9.75(18) _{stat} (9) _{sys}
m_0^{static}/T	8.82(9) _{stat} (6) _{sys}

comparisons:

- free theory: $|H_E|/T^2 = 0.5$ Meyer '18
- $\mathcal{N} = 4$ SYM: $|H_E|/\chi_s^{\text{SYM}} = 0.67$ Meyer '18
while $|H_E|^{(\text{lat})}/\chi_s = 0.75$
- static screening mass for $n = 1$ on our coarsest ensemble
using the dispersion relation $\sqrt{5.76^2 + (2\pi n)^2} = 8.52$

Brandt, Francis, Laine, Meyer '14

Similar steps for $n=2$: more noisy



- data is more noisy
- screening masses are large: $m_0 a \sim 1$ on the coarsest ensemble
- $n = 2$ result is compatible within errors with the $n = 1$ result
- dispersion rel., $\sigma(\omega) > 0$: $\implies |H_E(r)| > |H_E(n)|$ if $r > n$

Summary and outlook

- first analysis of imaginary spatial momentum correlators at finite T in the $n = 1$ & $n = 2$ Matsubara sectors
- simultaneous continuum extrapolation of H_E/T^2 of different discretized correlators using three lattice spacings
- determining H_E/T^2 for $n \geq 2$ (or modelling using e.g. NLO pert. theory) might allow a more educated guess for the photon rate at this temperature

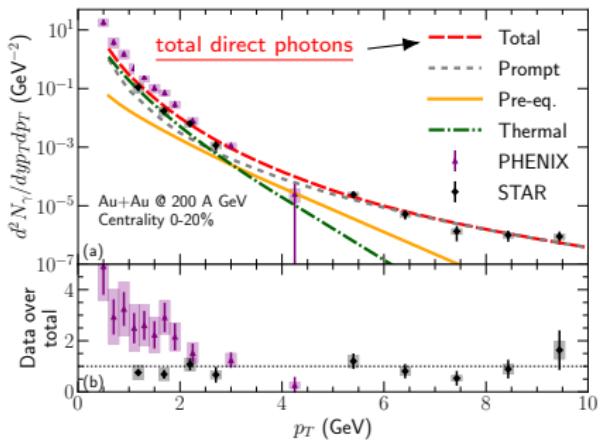
Thank you for your attention!

BACKUP

Discrepancies between experiment and theory

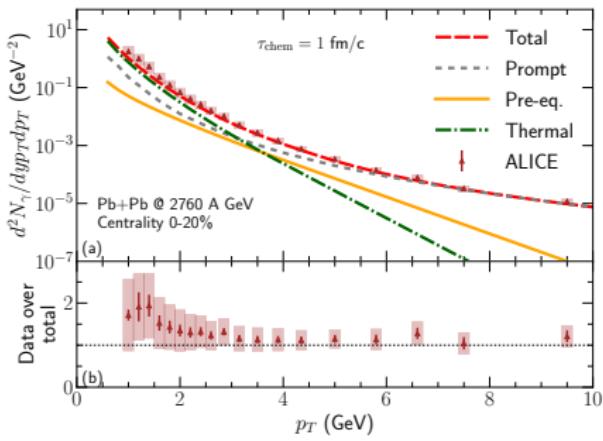
PHENIX, STAR (RHIC)

[1405.3940, 1607.01447]



ALICE (LHC)

[1509.07324]



source of figs: [2106.11216]

- discrepancy between PHENIX & STAR
- discrepancy between PHENIX & theory (/ALICE & theory)
- thermal photon yield: assuming weakly coupled plasma + hydro