



Gradient flow anomalous dimensions for ten-flavor SU(3) gauge theory

Curtis Taylor Peterson, Anna Hasenfratz and Oliver Witzel
On behalf of the Lattice Strong Dynamics (LSD) collaboration

LATICE

Ten-flavor SU(3) gauge theory and the renormalization group (RG)

- Ten-flavor SU(3) gauge theory
 - Interesting for composite Higgs modelling (e.g., LSD 4+6 system)
 - There is ongoing debate over the existence of an infrared fixed point (IRFP)
- Of particular interest are mesonic and baryonic running operator anomalous dimensions

[Hasenfratz, Rebbi, Witzel PRD 101, 114508 (2020)] [Kuti, Fodor, Holland, Wong PoS LATTICE2021 (2021) 396] [LSD PRD 103, 014504 (2020)] [Cacciapaalia, Pica, Sannino Phys. Rep. 877 (2020)]







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per per scale
$$\beta(\mu;g^2) \equiv \mu^2 \frac{\mathrm{d}}{\mathrm{d}\mu^2} g^2(\mu)$$

$$\gamma_{\mathcal{O}}(\mu; g^2) \equiv \mu \frac{\mathrm{d}}{\mathrm{d}\mu} \log Z_{\mathcal{O}}(\mu)$$

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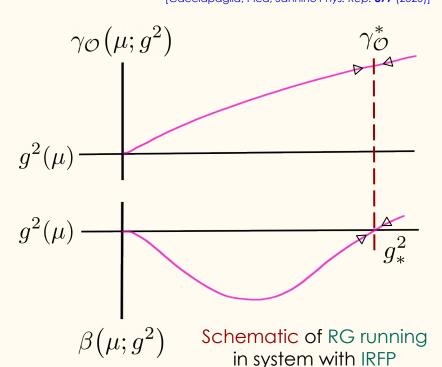
RG beta-functior

Energy scale
$$\beta \left(\mu;g^2\right) \equiv \mu^2 \frac{\mathrm{d}}{\mathrm{d}\mu^2} g^2(\mu)$$

Operator anomalous dimension

$$\gamma_{\mathcal{O}}(\mu; g^2) \equiv \mu \frac{\mathrm{d}}{\mathrm{d}\mu} \log Z_{\mathcal{O}}(\mu)$$

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[Carosso, Hasenfratz, Neil *PRL* **121**, 201601 (2018)]

- Gradient flow describes an RG transformation when combined with a rescaling step in the calculation of expectation values
 - Must take infinite volume limit
 - We can avoid rescaling step for certain quantities using MCRG principles

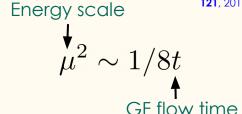






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Partially flowed two-point function
$$G_{\mathcal{O}}(t,\hat{x}_4;g_0^2) \equiv \int_{\hat{\mathbf{x}}} \left\langle \mathcal{O}_t(\hat{\mathbf{x}},\hat{x}_4)\mathcal{O}(\hat{\mathbf{x}},0) \right\rangle_{g_0^2}$$





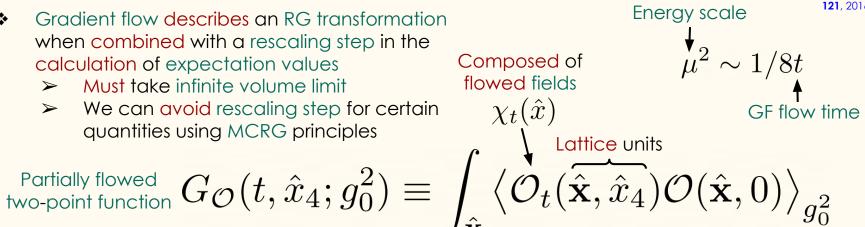




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Partially flowed
$$G_{\mathcal{O}}(t,\hat{x}_4;g_0^2) \equiv$$





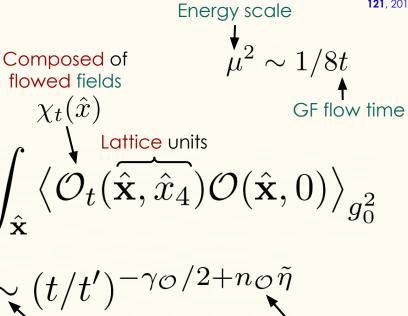


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Partially flowed two-point function $G_{\mathcal{O}}(t,\hat{x}_4;g_0^2) \equiv \int_{\hat{\mathbf{x}}} \langle \mathcal{O}_t(\hat{\mathbf{x}},\hat{x}_4)\mathcal{O}(\hat{\mathbf{x}},0) \rangle_{g_0^2}$

$$\frac{G_{\mathcal{O}}(t, \hat{x}_4; g_0^2)}{G_{\mathcal{O}}(t', \hat{x}_4; g_0^2)}$$



$$\frac{G_{\mathcal{O}}(t,\hat{x}_4;g_0^2)}{G_{\mathcal{O}}(t',\hat{x}_4;g_0^2)} \sim \underbrace{(t/t')^{-\gamma_{\mathcal{O}}/2+n_{\mathcal{O}}\tilde{\eta}}}_{8t/a^2,8t'/a^2\gg 1} \underbrace{\overset{\text{Comes}}{\underset{\text{from }Z_\chi}{\text{cand}}}}_{\text{cand}}$$









GF can be used to define a scheme for the renormalization of local operators

[Carosso, Hasenfratz., Neil PRL 121, 201601 (2018)] [Hasenfratz, Monahan, Rizik, Schindler, Witzel PoS LATTICE2021 (2021) 155]









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Matches onto tree-level
$$\overline{\rm MS}$$
 coupling $g_{
m GF}^2(t;g_0^2)\equiv \mathcal{N}\langle t^2E(t)
angle_{g_0^2}$









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$$\mathcal{R}_{\mathcal{O}}\big(t,\hat{x}_4;g_0^2\big) \equiv \frac{G_{\mathcal{O}}(t,\hat{x}_4;g_0^2)}{G_{\mathcal{V}}(t,\hat{x}_4;g_0^2)^{n_{\mathcal{O}}/n_{\mathcal{V}}}}$$
 Independent of
$$\hat{x}_4 \text{ for} \qquad \qquad \text{Cancels off } n_{\mathcal{O}}\tilde{\eta}$$

$$\hat{x}_4 \gg \sqrt{8t/a}$$









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$$g_{\rm GF}^2 (t;g_0^2) \equiv \mathcal{N} \langle t^2 E(t) \rangle_{g_0^2}$$

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$$\beta_{GF}(t; g_0^2) = -t \frac{\mathrm{d}}{\mathrm{d}t} g_{GF}^2(t; g_0^2)$$
$$\gamma_{\mathcal{O}}(t, \hat{x}_4; g_0^2) = -2t \frac{\mathrm{d}}{\mathrm{d}t} \log \mathcal{R}_{\mathcal{O}}(t, \hat{x}_4; g_0^2)$$

$$\mathcal{R}_{\mathcal{O}} \big(t, \hat{x}_4; g_0^2 \big) \equiv \frac{G_{\mathcal{O}} (t, \hat{x}_4; g_0^2)}{G_{\mathcal{V}} (t, \hat{x}_4; g_0^2)^{n_{\mathcal{O}}/n_{\mathcal{V}}}}$$
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Independent of
$$\hat{x}_4$$
 for $\hat{x}_4 \gg \sqrt{8t}/a$

Cancels off
$$n_{\mathcal{O}}\tilde{\eta}$$

$$\beta_{\rm GF}(t;g_0^2) = -t\frac{\mathrm{d}}{\mathrm{d}t}g_{\rm GF}^2(t;g_0^2)$$

$$\gamma_{\mathcal{O}}(t, \hat{x}_4; g_0^2) = -2t \frac{\mathrm{d}}{\mathrm{d}t} \log \mathcal{R}_{\mathcal{O}}(t, \hat{x}_4; g_0^2)$$

We can trade t dependence for $g_{\mathrm{GF}}^2 \left(t; g_0^2 \right)$ dependence

See also:

[O. Witzel: Wed. 2:00 PM]

[A. Shindler: Wed. 4:30 PM] [C. Monahan: Wed. 4:50]

[C.H. Wong: Thu. 12:10 PM] [R. Harlander: Thu. 12:30 PM]

[A. Harlander: Inv. 12:30 PM]
[A. Hasenfratz: Fri. 8:50 AM]









Simulation details



Ten-flavor simulations are performed using Symanzik gauge action with stout smeared Möbius domain wall fermions (DWF) using GRID

- > Bare gauge couplings $\beta \equiv 6/g_0^2 = 5.00, 6.00, 4.20, 4.60$ 4.10, 4.05
- ightharpoonup Volumes $24^3 \times 64$ and $32^3 \times 64$
- Gradient flow performed with Wilson flow using QLUA
 - We define the gradient flow coupling in finite volume using tree-level normalization
 - Corrects for gauge zero modes and tree-level cutoff effects

[Boyle, Yamaguchi, Cossu, Portelli LATTICE2015 (2015) 023]
[Pochinsky PoS LATTICE2008 (2008) 040]
[Fodor, Holland, Kuti, Mondal, Nogradi, JHEP 09 (2014) 018]

$$\begin{split} g_{\mathrm{GF}}^2 \left(t; L, g_0^2\right) \\ &\equiv \frac{\mathcal{N}}{C(t, L)} \langle t^2 E(t) \rangle_{L, g_0^2} \\ &\stackrel{\mathrm{Tree-level}}{\underset{\mathrm{normalization}}{\underset{\mathrm{factor}}{\sim}}} \end{split}$$

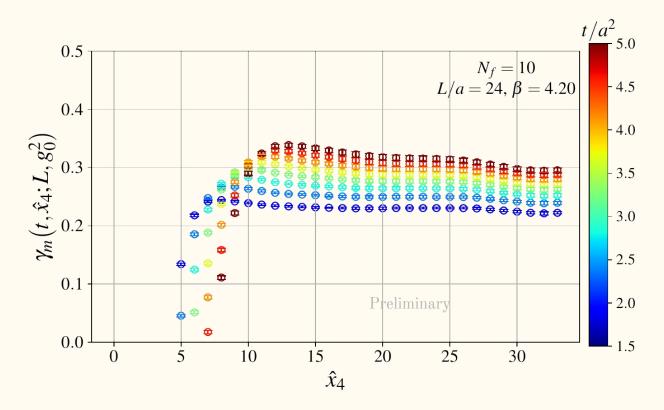








Anomalous dimension \hat{x}_4 -independence



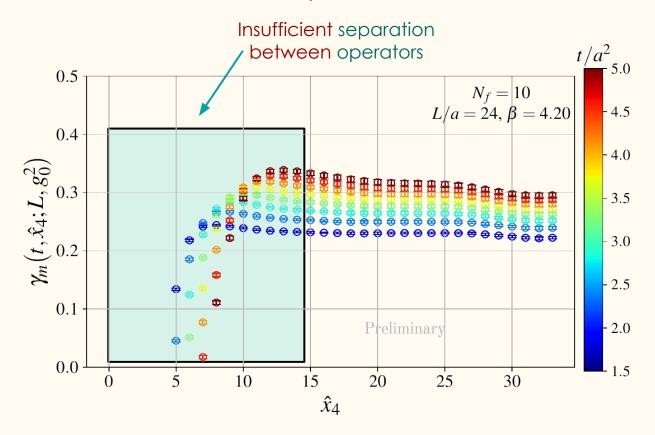








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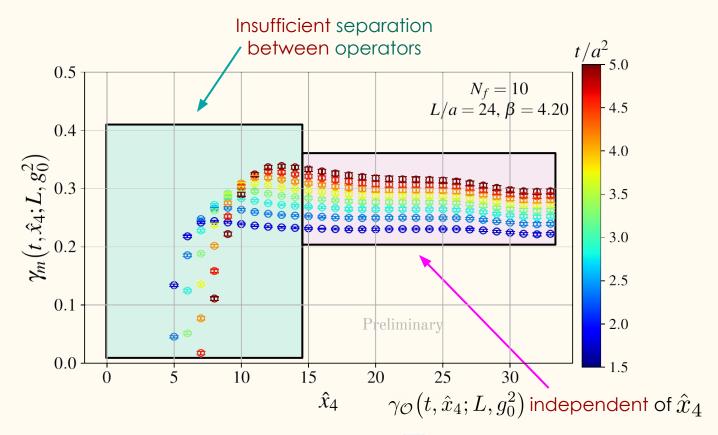








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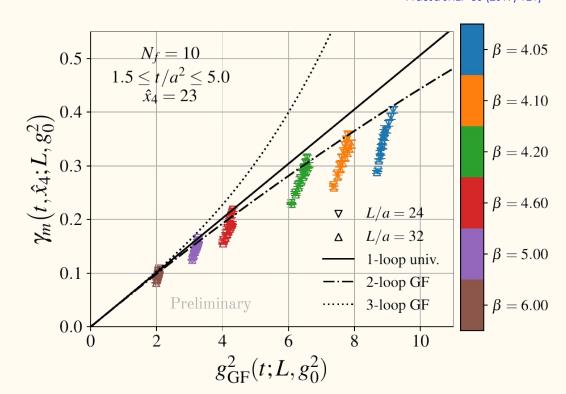




Anomalous dimension (no extrapolations)

[Artz, Harlander, Lange, Neumann, Prausa JHEP **06** (2019) 121]

- Finite volume effects appear to be small
- Each bare gauge coupling appears to approach the 1- and 2-loop GF curves from perturbation theory





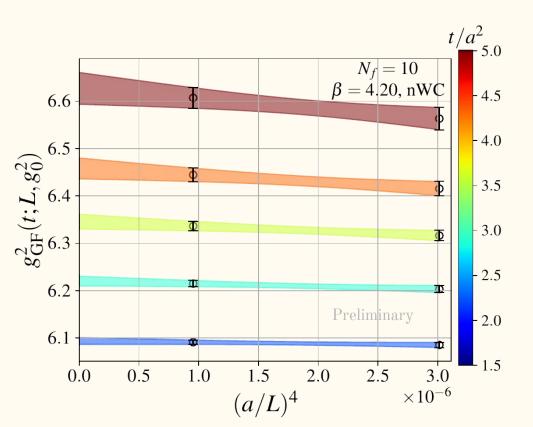








Infinite volume limit (GF coupling)



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Leading-order finite volume effects predicted from scaling of Yang-Mills energy density

$$g_{\text{GF}}^2(t; L, g_0^2) \approx g^2(t; g_0^2) + k(t^2/L^4)$$









Infinite volume limit (anomalous dimension)

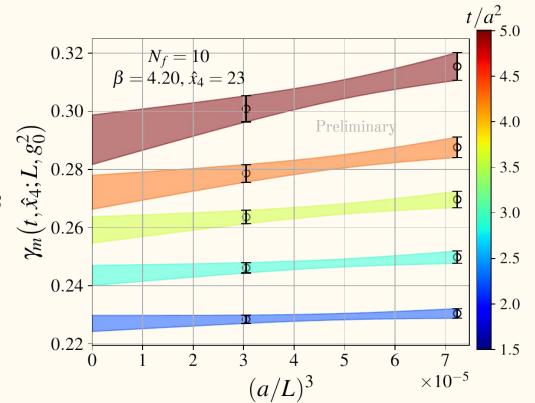
Leading finite volume effects predicted by scaling of two-point function

$$\gamma_{\mathcal{O}}(t,\hat{x}_4;L,g_0^2)$$

$$\approx \gamma_{\mathcal{O}}(t, \hat{x}_4; L, g_0^2) + \kappa (t/L^2)^{\delta_{\mathcal{O}}/2}$$

$$\delta_{\mathcal{O}} pprox 3$$
 (meson)

$$pprox 9/2$$
 (baryon)







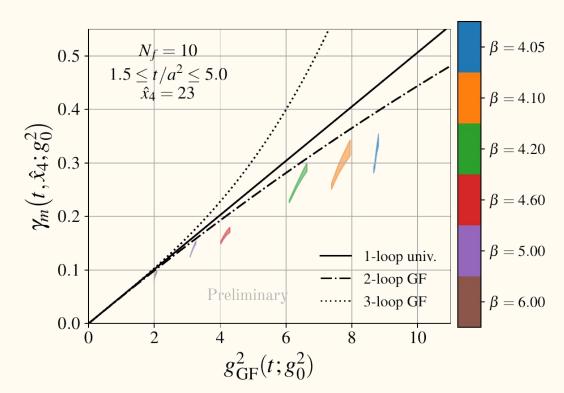




Infinite volume limit (final result)

Bare gauge couplings do not overlap

Must interpolate between bare gauge coupling





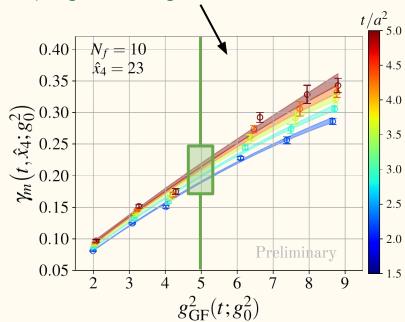






Continuum limit (extrapolation)

Interpolation between bare gauge couplings at fixed gradient flow time





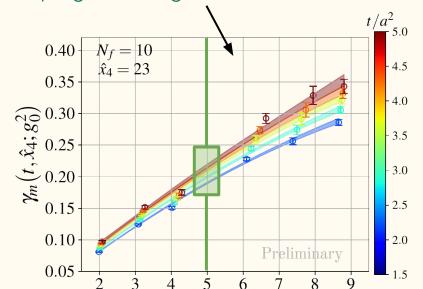






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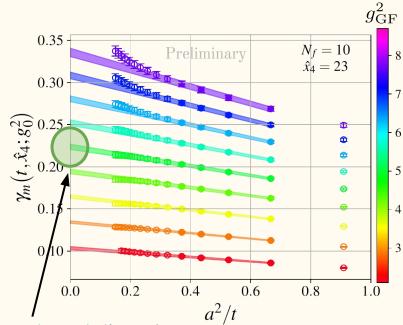
Interpolation between bare gauge couplings at fixed gradient flow time



 $g_{GF}^2(t;g_0^2)$

$$\gamma_{\mathcal{O}}(t, \hat{x}_4; g_0^2)$$

 $\approx \gamma_{\mathcal{O}}(t, \hat{x}_4) + c(a^2/t)$



Continuum extrapolation at fixed gradient flow coupling



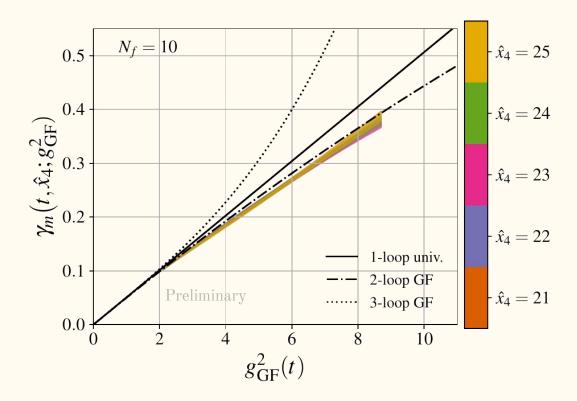






Continuum limit (final result)





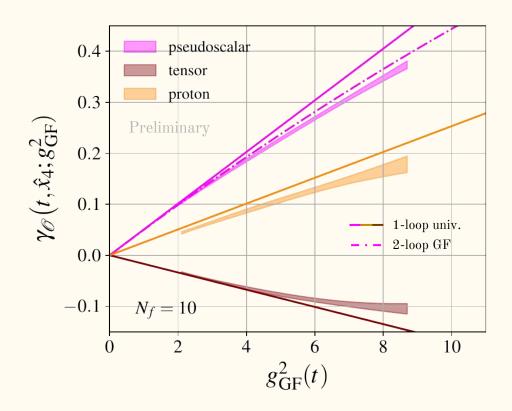
Continuum mass anomalous dimension appears to closely follow 2-loop GF







We have also calculated proton and tensor anomalous dimensions





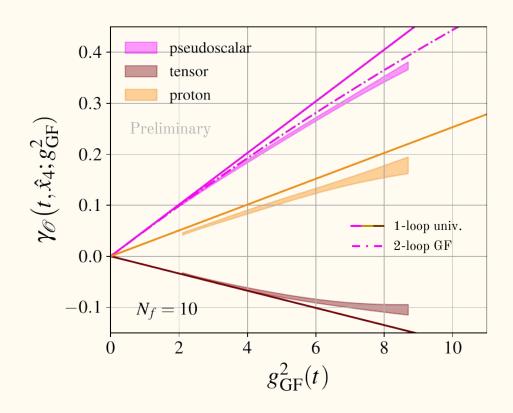








- We have also calculated proton and tensor anomalous dimensions
- Gradient flow anomalous dimension closely follow perturbative curves over our current range of couplings
 - No evidence for non-perturbative enhancement over our current range of couplings





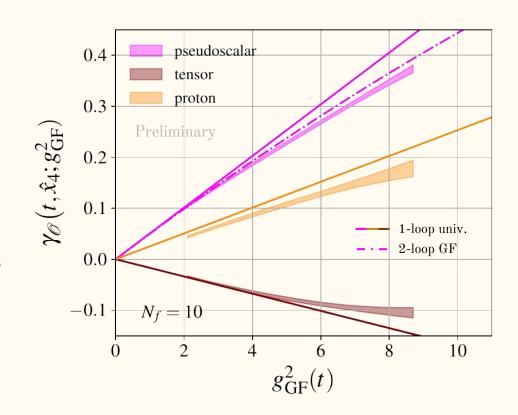








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- Plan to gather more statistics on all bare gauge couplings on L/a = 32





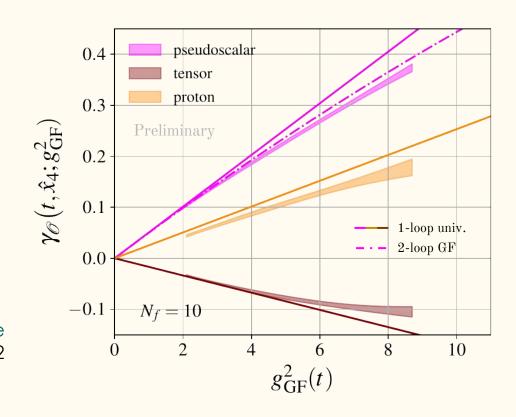








- We have also calculated proton and tensor anomalous dimensions
- Gradient flow anomalous dimension closely follow perturbative curves over our current range of couplings
 - No evidence for non-perturbative enhancement over our current range of couplings
- Plan to gather more statistics on all bare gauge couplings on L/a = 32
- We are currently running one more bare gauge coupling at $\beta=4.03$ for L/a = 32











Acknowledgements

U. Colorado: RMACC Summit

LLNL: Lassen, Boraxo, Quartz

U. Siegen: OMNI

NSF GRFP

LSD Collaboration





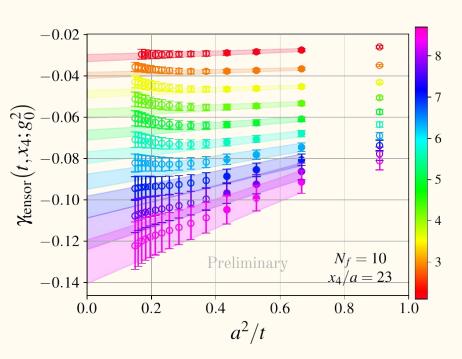


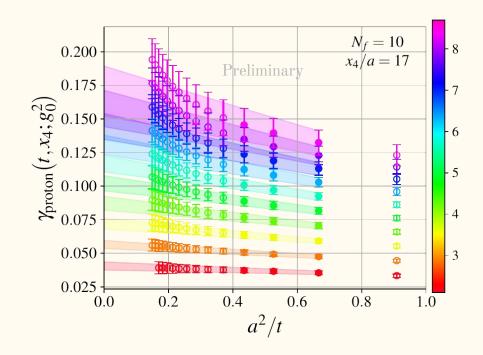






Continuum extrapolation of tensor and proton















Tensor and proton anomalous dimension

