

Non-perturbative study of Yang-Mills theory with four supercharges in two dimensions



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With Raghav G. Jha, Anosh Joseph, and David Schaich

Quick RECAP

Presented preliminary analysis in **Lattice 2021**.

) Some remarks from the talk:

Scalars behaviour

Existence of bound state at finite temperature for $U(N)$ with $N = 2; 4; 8; 12$.

arXiv: [2108.08111](#) [hep-lat] NSD, Jha, Joseph, Schaich

Comparison with 16 supercharge theory

Theory looks to be in different universality class to maximal theory.

Not discussed

Possible 'Spatial Deconfinement' transition.

This talk

of four supercharge theory on .

with maximal theory in regimes.

Signature of transition and its possible .



Two-dimensional

$N = (2; 2)$ **SYM**

Two-dimensional $N = (2;2)$ SYM

Constructed from dimensional reduction of four dimensional theory.

$$N = \dots, d = \dots \quad / \quad N = (\dots, \dots), d = \dots$$

Not a "**maximal**" theory.

No holographic dual "**exists**".

Regularised on lattice using "**twisting**".

Phys. Rep. **484** () -
Catterall, Kaplan, Ünsal

Maximal Supersymmetric theories on Lattice talks:

Goksu Toga: Now TD-I

Angel Sherletov: Monday- : pm

David Schaich: Monday- : pm

Arpith Kumar: Wednesday- : pm

Two-dimensional $N = (2;2)$ SYM

Continuum Action

$$S = \frac{N}{4} \int d^2x \text{Tr} \left[F_{ab} F_{ab} + \bar{D}_a; D_a \right] + \frac{1}{2} \int d^2x \text{Tr} \left[D_a; D_a \right]^2$$

After integrating out auxiliary field

$$S = \frac{N}{4} \int d^2x \text{Tr} \left[\bar{F}_{ab} F_{ab} + \frac{1}{2} \bar{D}_a; D_a \right]^2 + \int d^2x \text{Tr} \left[D_a; D_a \right]^2$$

Two-dimensional $N = (2;2)$ SYM

$$S = \frac{N}{4} \int d^2x \text{Tr} \left[\bar{F}_{ab} F_{ab} + \frac{1}{2} \bar{D}_a; D_a^2 - \frac{1}{2} \bar{D}_a D_a \right]$$

Using geometrical discretization / theory lives on d lattice.

JHEP 11 () 6

Catterall

To control at directions, scalar potential term added to discretized action.

JHEP 11 ()

Catterall, Damgaard, DeGrand, Galvez, Mehta

Discretization used and all the observables studied can be accessed via publicly available software github.com/daschaich/susy.

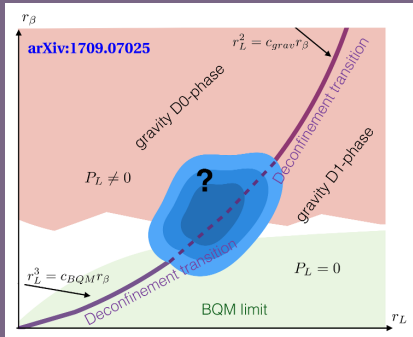


Comparison with Two-dimensional

$N = (8; 8)$ **SYM**

Two-dimensional $N = (8;8)$ SYM

At low temperature and large N) dual to type IIB supergravity.



PRD **97**, 86 () 8

Catterall, Jha, Schaich, Wiseman

Maximal theory prediction from gravity dual

Scalars behave as: $\text{Tr}(X^2) / t$.

JHEP **07** () Wiseman

Energy density $\propto t^2$ (for $t > 1$),
 $\propto t^3$ (for $t < 1$).

JHEP **07** () Wiseman

First order GL Phase transition.

PRL **70**, 8 () Gregory, La amme

Back to Target Theory

) Lattice simulations for four supercharge theory (

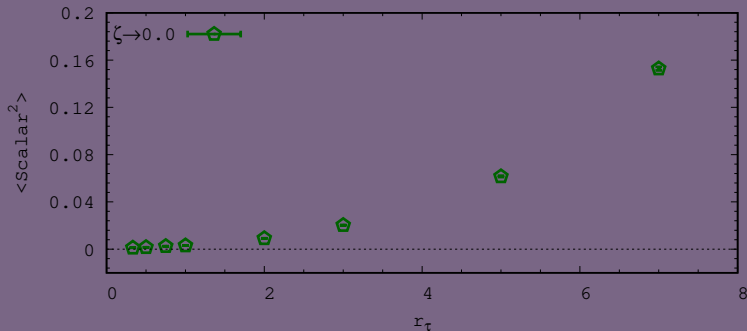
Worked with finite mass deformation parameter ρ ,
 $= \frac{r}{N_t} = \rho a, r = 1=t, 0.33 \quad r = 7:0, 2 (0.2;0.3;0.4;0.5)$.

Different Lattice aspect ratios ρ ,
 $= \frac{N_x}{N} = \frac{r_x}{r} \quad 2 (0.5;1.0;1.5;2.0)$.

Different gauge groups, $2 \quad N = 20$

Anti-periodic boundary conditions for fermions along temporal direction.

Scalar behaviour



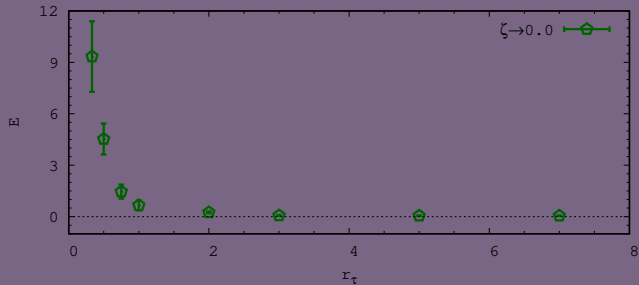
Scalar² \propto Tr(X)²

24 24 lattice,
 $N = 12$.

$r > 1$! r^3 behaviour, Maximal case ! $1=r$.

$r < 1$! r behaviour, Maximal case ??

Energy density



$$E = \frac{3}{24} \frac{1}{24 \text{ lattice, } N = 12} \frac{2}{3N^2} S_B$$

$r > 1$! r^0 behaviour, Maximal case ! $1=r^3$.

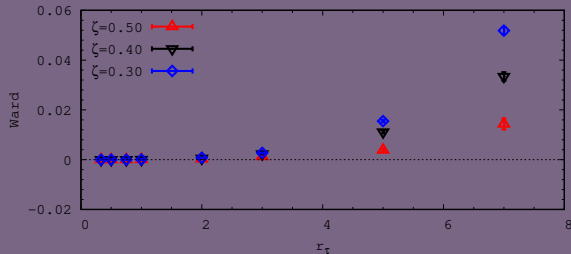
$r < 1$! $1=r^2$ behaviour, Maximal case ! $1=r^2$.

Vanishing energy density at zero temperature ! Preserved SUSY.

PRD **80**, 6 () Hanada, Kanamori

PRD **97**, (8) Catterall, Jha, Joseph

Ward Identity

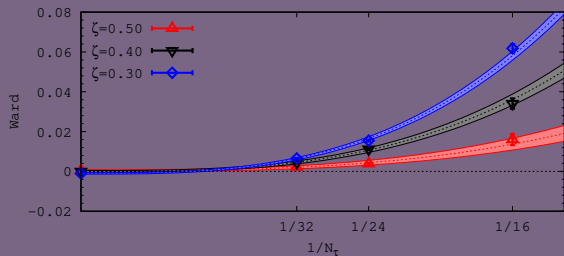


Ward Identity:

$$Q_a^P U_a U_a$$

At larger temperatures ($r < 1$), ward identity satisfied.

At smaller temperatures ($r > 1$), satisfied at larger volume.



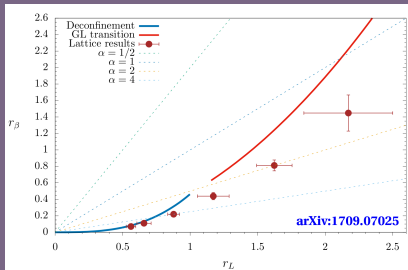
24 \times 24 lattice, $N = 12$
Bottom left plot with $r = 5.0$.

‘Spatial deconfinement’
in two-dimensional
 $N = (2;2)$ **SYM**

Spatial deconfinement

Four supercharge theory:
so far

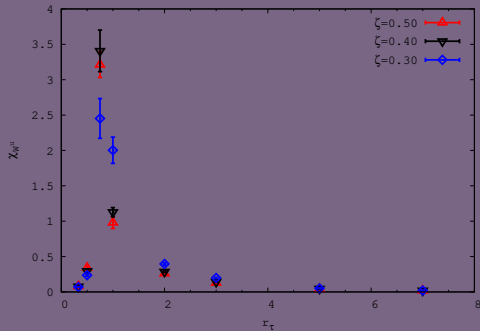
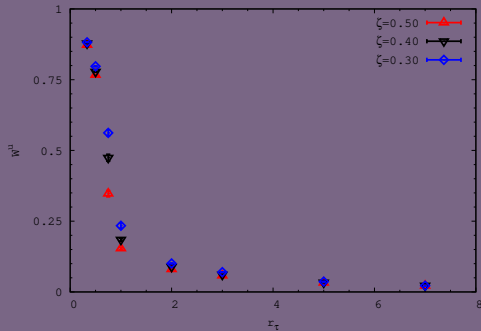
- Preserved SUSY.
- Different behaviour compared with maximal case.
- What about deconfinement transition? which exists in sixteen supercharge theory.



PRD **97**, 86 (8)

Catterall, Jha, Schaich, Wiseman

Spatial deconfinement - Signal



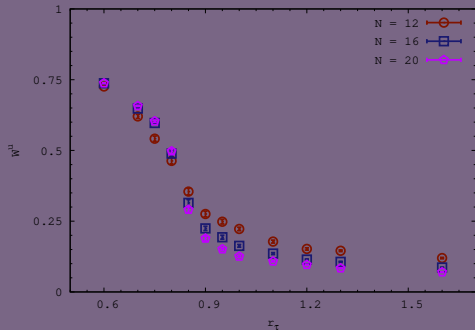
Spatial wilson lines and its susceptibility as order parameter for deconfinement transition.

24 \times 24 lattice, $N = 12$.

Transition around $r_t = 1.0$.

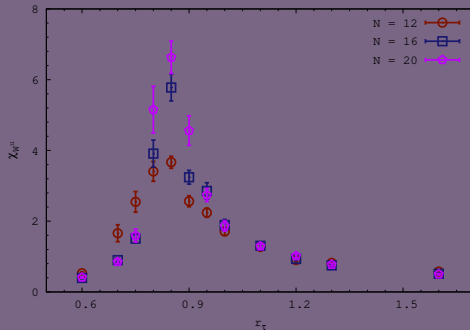
Slight dependence.

Spatial deconfinement - Order



Spatial wilson lines and its susceptibility for different N values.

12 12 lattice, $\beta = 0.30$.

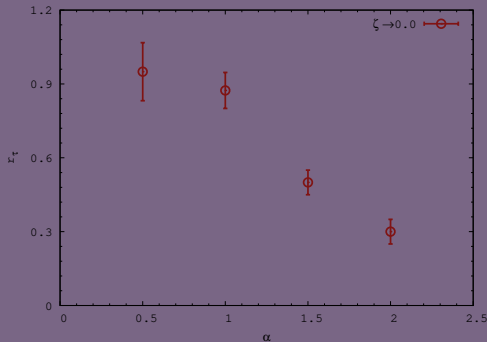


Critical r independent of N .
Hints of second-order transition.

PRL 113, 6 ()

Azuma, Morita, Takeuchi

Phase transition vs Lattice aspect ratio



$N = 12$, Lattices used
12; 24; 12; 12; 24; 16; 24; 12.
 r (critical) has dependence for
1.

Spatial deconfinement transition
similar to maximal theory but
restricted only in weaker coupling
regime.

Conclusions

Spatial deconfinement

Spatial deconfinement phase transition observed in this theory with different lattice volumes.

Weak coupling behaviour

Similar to maximal theory,
with different normalizations

Phase transition observed.

Energy density behaviour same.

Strong coupling behaviour

Different from maximal theory,

No Phase transition.

Scalars behaviour different.

Energy density behaviour
different.

Open question

Holographic dual to two-dimensional $N = (2;2)$ SYM - - - - ???

