### Deflated Multigrid Multilevel Monte Carlo

#### Gustavo Ramirez-Hidalgo (Andreas Frommer, Mostafa Nasr Khalil)

Bergische Universität Wuppertal Department of Applied Mathematics

August 11, 2022



BERGISCHE UNIVERSITÄT WUPPERTAL





European Joint Doctorates

Gustavo Ramirez-Hidalgo

Lattice 2022

August 11, 2022 1 / 14

## Short recap

- 2 Some preliminary results in LQCD
- 3 Skipping levels in multigrid multilevel Monte Carlo
- Exact deflation in multigrid multilevel Monte Carlo
- 5 A (short) note on eigensolving



э

イロト イポト イヨト イヨト

Short recap

<□▶ < @▶ < E▶ < E▶ E - のへで 2/14

The problem at hand is:  $tr(A^{-1})$ . The method presented here works for any f(A), in particular any  $f(A) = (\Gamma A)^{-1}$ : but we focus on  $f(A) = A^{-1}$ :

We start with the decomposition:

$$\operatorname{tr}(A^{-1}) = \sum_{\ell=1}^{L-1} \operatorname{tr}\left(\hat{P}_{\ell} A_{\ell}^{-1} \hat{R}_{\ell} - \hat{P}_{\ell+1} A_{\ell+1}^{-1} \hat{R}_{\ell+1}\right) + \operatorname{tr}(\hat{P}_{L} A_{L}^{-1} \hat{R}_{L}).$$

Why should this reduce the variance?

$$\mathbb{V}[x^{H}A^{-1}x] = \|\text{offdiag}(A^{-1})\|_{F}^{2} = \sum_{i=1}^{n} \sigma_{i}^{-2} - \sum_{i=1}^{n} |(A^{-1})_{ii}|^{2}$$



3

ヘロト 人間 ト 人 ヨト 人 ヨトー

A bit more specifically:

• two levels: 
$$\operatorname{tr}(A_1^{-1}) = \operatorname{tr}(A_1^{-1}) - \operatorname{tr}(P_1A_2^{-1}R_1) + \operatorname{tr}(P_1A_2^{-1}R_1)$$

• multilevel 
$$(R_{\ell}P_{\ell} = I)$$
:  
 $\operatorname{tr}(A^{-1}) = \sum_{\ell=1}^{L-1} \operatorname{tr}(A_{\ell}^{-1} - P_{\ell}A_{\ell+1}^{-1}R_{\ell}) + \operatorname{tr}(A_{L}^{-1})$ 

We can connect this to inexact/approximate deflation:

$$A^{-1} = A^{-1}(I - \Pi) + A^{-1}\Pi$$
$$\Pi = \hat{U}_k (\hat{V}_k^* A \hat{U}_k)^{-1} \hat{V}_k^* A, \hat{U}_k, \hat{V}_k \in \mathbb{C}^{n \times k}$$

which is discussed in *Multigrid deflation for Lattice QCD* (Romero, Stathopoulos and Orginos), and already in Lüscher's local coherence paper.



BERGISCHE UNIVERSITÄT WUPPERTAL

4/14

#### Short recap Multigrid multilevel Monte Carlo for $tr(A^{-1})$



Some preliminary results in LQCD

◆□▶ ◆舂▶ ◆臣▶ ◆臣▶ 臣 - 釣�� 5/14

# Some preliminary results in LQCD Wilson and twisted mass

Wilson:  $64 \times 32^3$  (work in collaboration with Jose Jimenez-Merchan). Twisted mass:  $96 \times 48^3$ .



	measurement	Mult. MC	Hutchinson
$\ell = 1$	# estimates	63	1164
	time (seconds)	17,010.0	221,160.0
$\ell = 2$	# estimates	172	-
	time (seconds)	15,480.0	-
$\ell = 3$	# estimates	903	-
	time (seconds)	9,030.0	-
$\ell = 4$	# estimates	7715	-
	time (seconds)	1,543.0	-
total	# estimates	-	1164
total	time (seconds)	43,063.0	221,160.0



#### Skipping levels in multigrid multilevel Monte Carlo

### Skipping levels in multigrid multilevel Monte Carlo Skip the second level

Let's fully write a four-level expansion:

$$\operatorname{tr}(A^{-1}) = \operatorname{tr}\left(A_{1}^{-1} - P_{1}A_{2}^{-1}R_{1}\right) + \operatorname{tr}\left(P_{1}A_{2}^{-1}R_{1} - P_{1}P_{2}A_{3}^{-1}R_{2}R_{1}\right) + \operatorname{tr}\left(P_{1}P_{2}A_{3}^{-1}R_{2}R_{1} - P_{1}P_{2}P_{3}A_{4}^{-1}R_{3}R_{2}R_{1}\right) + \operatorname{tr}\left(P_{1}P_{2}P_{3}A_{4}^{-1}R_{3}R_{2}R_{1}\right) + \operatorname{tr}\left(P_{1}$$

Skip second level-difference:

$$\begin{aligned} \operatorname{tr}(\mathcal{A}^{-1}) &= \operatorname{tr}\left(\mathcal{A}_{1}^{-1} - \mathcal{P}_{1}\mathcal{P}_{2}\mathcal{A}_{3}^{-1}\mathcal{R}_{2}\mathcal{R}_{1}\right) + \\ \operatorname{tr}\left(\mathcal{P}_{1}\mathcal{P}_{2}\mathcal{A}_{3}^{-1}\mathcal{R}_{2}\mathcal{R}_{1} - \mathcal{P}_{1}\mathcal{P}_{2}\mathcal{P}_{3}\mathcal{A}_{4}^{-1}\mathcal{R}_{3}\mathcal{R}_{2}\mathcal{R}_{1}\right) + \operatorname{tr}\left(\mathcal{P}_{1}\mathcal{P}_{2}\mathcal{P}_{3}\mathcal{A}_{4}^{-1}\mathcal{R}_{3}\mathcal{R}_{2}\mathcal{R}_{1}\right) \end{aligned}$$

Then, reduce further:

$$\operatorname{tr}(A^{-1}) = \operatorname{tr}\left(A_{1}^{-1} - P_{1}P_{2}A_{3}^{-1}R_{2}R_{1}\right) + \operatorname{tr}\left((A_{3}^{-1} - P_{3}A_{4}^{-1}R_{3})(R_{2}R_{1}P_{1}P_{2})\right) \\ + \operatorname{tr}\left(A_{4}^{-1}(R_{3}R_{2}R_{1}P_{1}P_{2}P_{3})\right)$$

Exact deflation in multigrid multilevel Monte Carlo

#### Exact deflation in MGMLMC Deflating non-deflated modes

Mostafa (talk on Th. @ 10:00) already talked about inexact deflation on top of multigrid multilevel Monte Carlo. Let's talk about exact deflation.



# Exact deflation in MGMLMC

Deflating non-deflated modes

Exact deflation in "plain" Hutchinson:

$$\begin{aligned} & \operatorname{tr}(A^{-1}) = A^{-1}(I - U_1 U_1^H) + \\ & \operatorname{tr}(U_1^H \Sigma_1^{-1} V_1) \end{aligned}$$

In LQCD, to compute the singular vectors we can use:

 $A = U\Sigma V \Rightarrow Q = \Gamma_5 A$ 

 $Q = X\Lambda X^{H}$  $\Rightarrow A = (\Gamma_{5}X \operatorname{sign}(\Lambda))|\Lambda|V$  Exact deflation in multigrid multilevel Monte Carlo:

instead of *A*, we have a more complicated operator to deflate from in multigrid multilevel Monte Carlo:

$$M_{\ell} = A_{\ell}^{-1} - P_{\ell} A_{\ell+1}^{-1} P_{\ell}^{H}$$



3

프 에 에 프 어

Can we construct a symmetrized difference-level operator?

To do this easily, the multigrid hierarchy needs to be built via a  $\Gamma_5\text{-}\text{compatible}$  aggregation:

$$\Gamma_5^\ell P_\ell = P_\ell \Gamma_5^{\ell+1}, \ \ P_\ell^H \Gamma_5^\ell = \Gamma_5^{\ell+1} P_\ell^H$$

and then:

$$\begin{aligned} J_{\ell} &= M_{\ell} \Gamma_{5}^{\ell} \\ &= A_{\ell}^{-1} \Gamma_{5}^{\ell} - P_{\ell} A_{\ell+1}^{-1} P_{\ell}^{H} \Gamma_{5}^{\ell} = Q_{\ell}^{-1} - P_{\ell} A_{\ell+1}^{-1} \Gamma_{5}^{\ell+1} P_{\ell}^{H} \\ &= Q_{\ell}^{-1} - P_{\ell} Q_{\ell+1}^{-1} P_{\ell}^{H} \end{aligned}$$



## Exact deflation in MGMLMC

Deflating non-deflated modes

 $m_0 = -0.1320$ , Schwinger. Single core, Python, laptop.

Mathad	exec.	eigsolv.
Method	time	time
Hutch. + exact defl. (384)	1663.74	119.92
MGMLMC	2701.05	0.0
MGMLMC + exact defl. (32,32,32)	1067.77	1544.12
MGMLMC + exact defl. (32,32) + skip	530.26	948.35



A (short) note on eigensolving

We need to compute large modes of the operator:

$$J_{\ell} = Q_{\ell}^{-1} - P_{\ell}Q_{\ell+1}^{-1}P_{\ell}^{H}$$

Done naively (as we have in our first tests with exact deflation), this is extremely expensive. It's probably better to transform the eigenproblem into a generalized Hermitian eigenproblem:

$$(Q_{\ell}^{-1} - P_{\ell}Q_{\ell+1}^{-1}P_{\ell}^{H})v = \lambda v$$
  
$$\Rightarrow (I - Q_{\ell}P_{\ell}Q_{\ell+1}^{-1}P_{\ell}^{H})v = \lambda Q_{\ell}v$$



# A (short) note on eigensolving Using Jacobi-Davidson

$$(Q_{\ell}^{-1} - P_{\ell}Q_{\ell+1}^{-1}P_{\ell}^{H})v = \lambda v$$
  
 $\Rightarrow (I - Q_{\ell}P_{\ell}Q_{\ell+1}^{-1}P_{\ell}^{H})v = \lambda Q_{\ell}v$ 

With the GHEP at hand, when using Jacobi-Davidson (which can be realized efficiently and scalably for Hermitian systems), we will need to apply the following operators at every iteration:

$$I-Q_\ell P_\ell Q_{\ell+1}^{-1} P_\ell^H \ Q_\ell$$

and (very approximately e.g.  $10^{-1}$ ) solve a linear system with the operator:

$$(I - Q_\ell P_\ell Q_{\ell+1}^{-1} P_\ell^H) - \theta Q_\ell$$



3

#### Thank you!



ъ

BERGISCHE UNIVERSITÄT WUPPERTAL

のへで 14/14

・ロト ・ 御 ト ・ ヨト・