

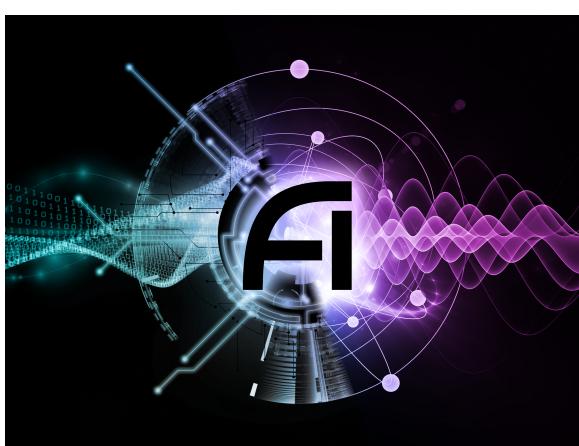
Gauge-equivariant flow models for sampling in LFT with pseudofermions

Fernando Romero-López

Lattice 2022 @ Bonn
August 8th



Massachusetts
Institute of
Technology

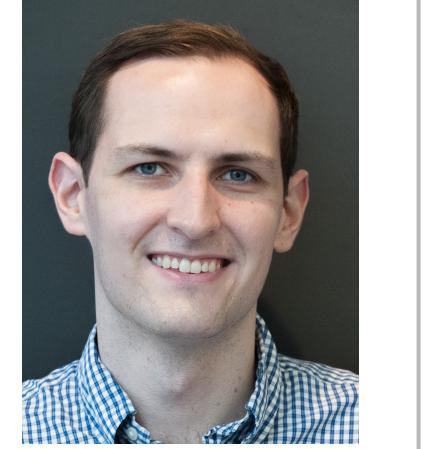
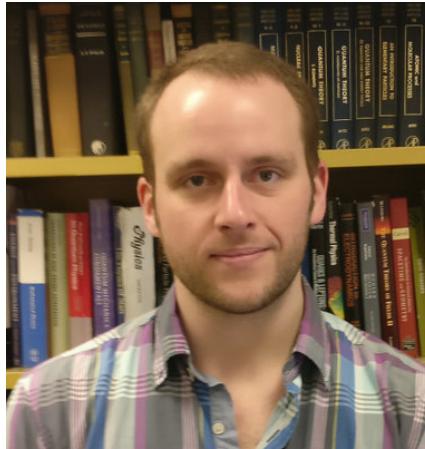


fernando@mit.edu

Collaboration (non-exhaustive)



Massachusetts Institute of Technology



Argonne NATIONAL LABORATORY





UNIVERSITÄT HEIDELBERG ZUKUNFT SEIT 1386



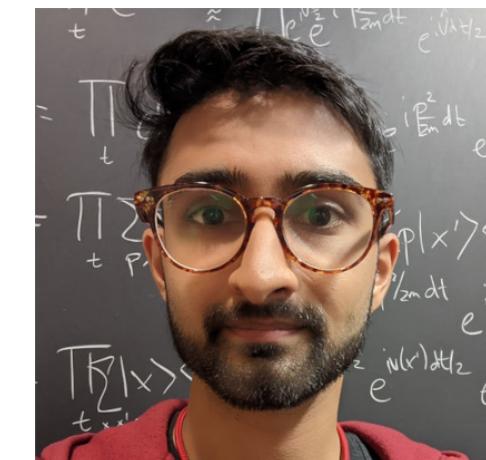
• Julian Urban



NEW YORK UNIVERSITY



• Kyle Cranmer
• Michael Albergo



• Gurtej Kanwar



DeepMind



• Sébastien Racanière
• Danilo Rezende

This talk is based on:

Flow-based sampling in the lattice Schwinger model at criticality

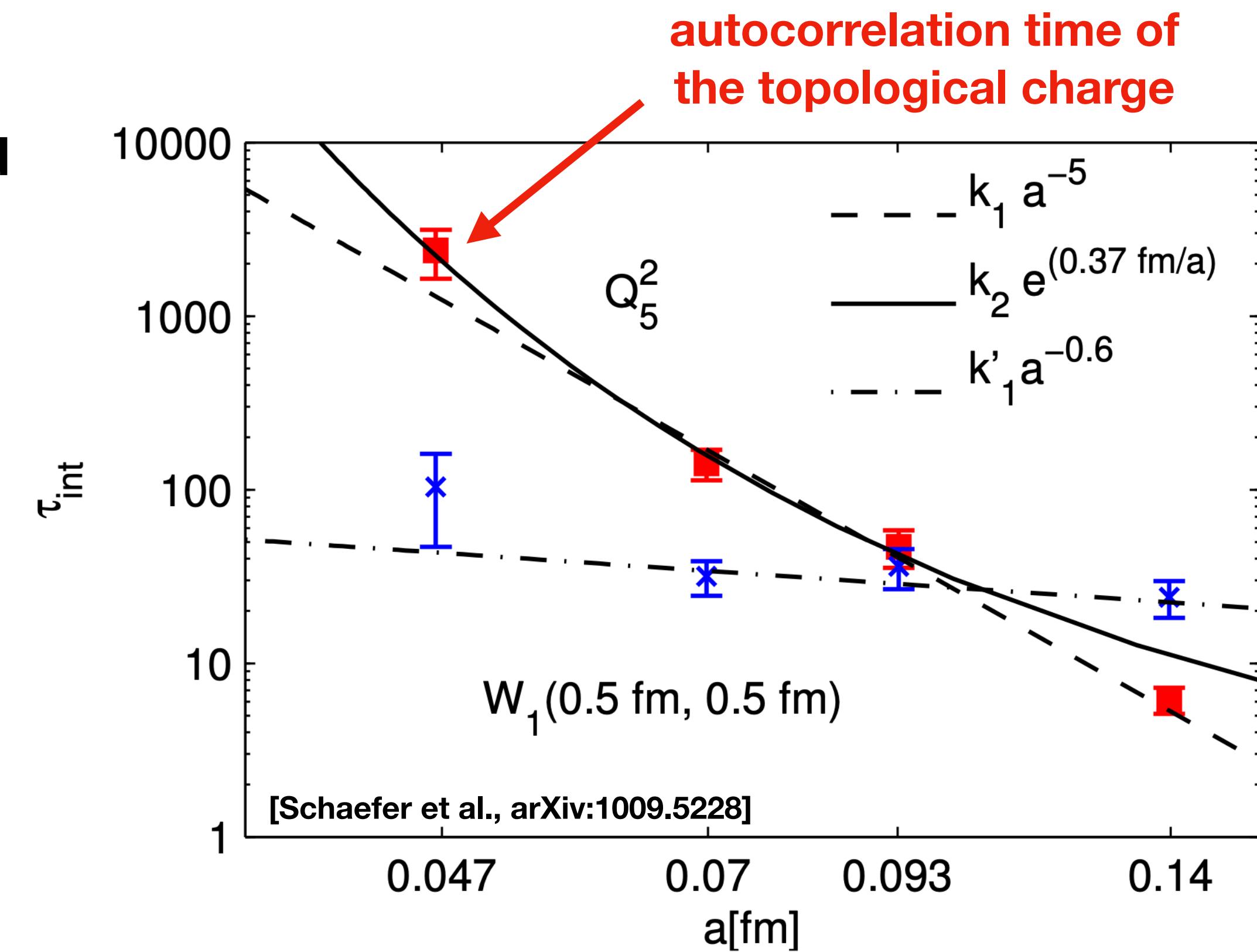
[Albergo, Boyda, Cranmer, Hackett, Kanwar, Racanière, Rezende, FRL, Shanahan, Urban 2202.11712]

Gauge-equivariant flow models for sampling in lattice field theories with pseudofermions

[Abbott, Albergo, Boyda, Cranmer, Hackett, Kanwar, Racanière, Rezende, FRL, Shanahan, Tian, Urban 2207.08945]

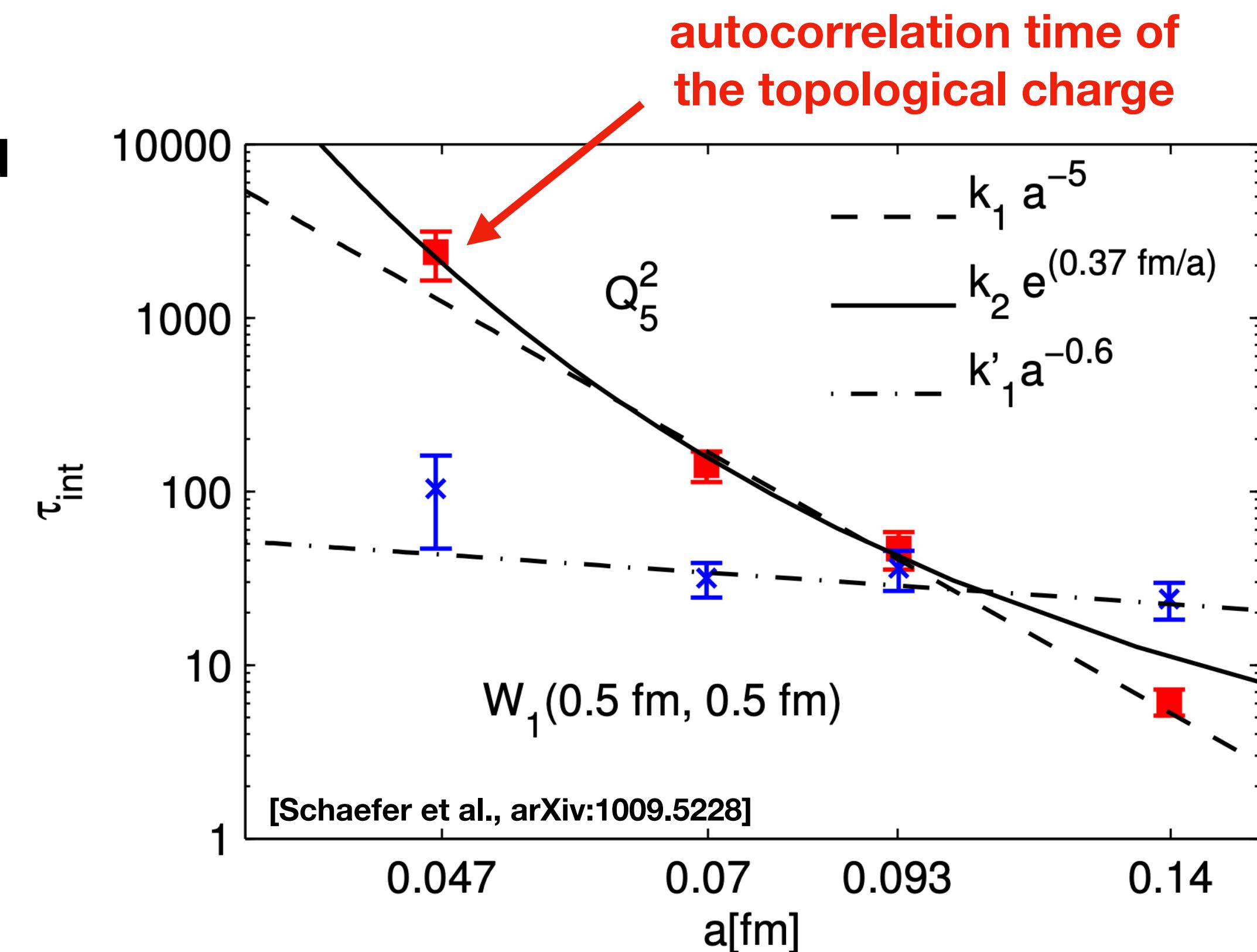
Introduction

- Critical slowing down is a well-known obstacle to extend the reach of state-of-the-art lattice QCD calculations
- Flow models are a promising tool, and have been able to overcome critical slowing down in some 2D theories
[Kanwar, Albergo, Boyda, Cranmer, Hackett, Racanière, Rezende, Shanahan 2003.06413]
[Finkenrath 2201.02216]
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- Several technical challenges remain for at-scale QCD
 - ▶ Compare: HMC has been optimized for the last ~30 years!



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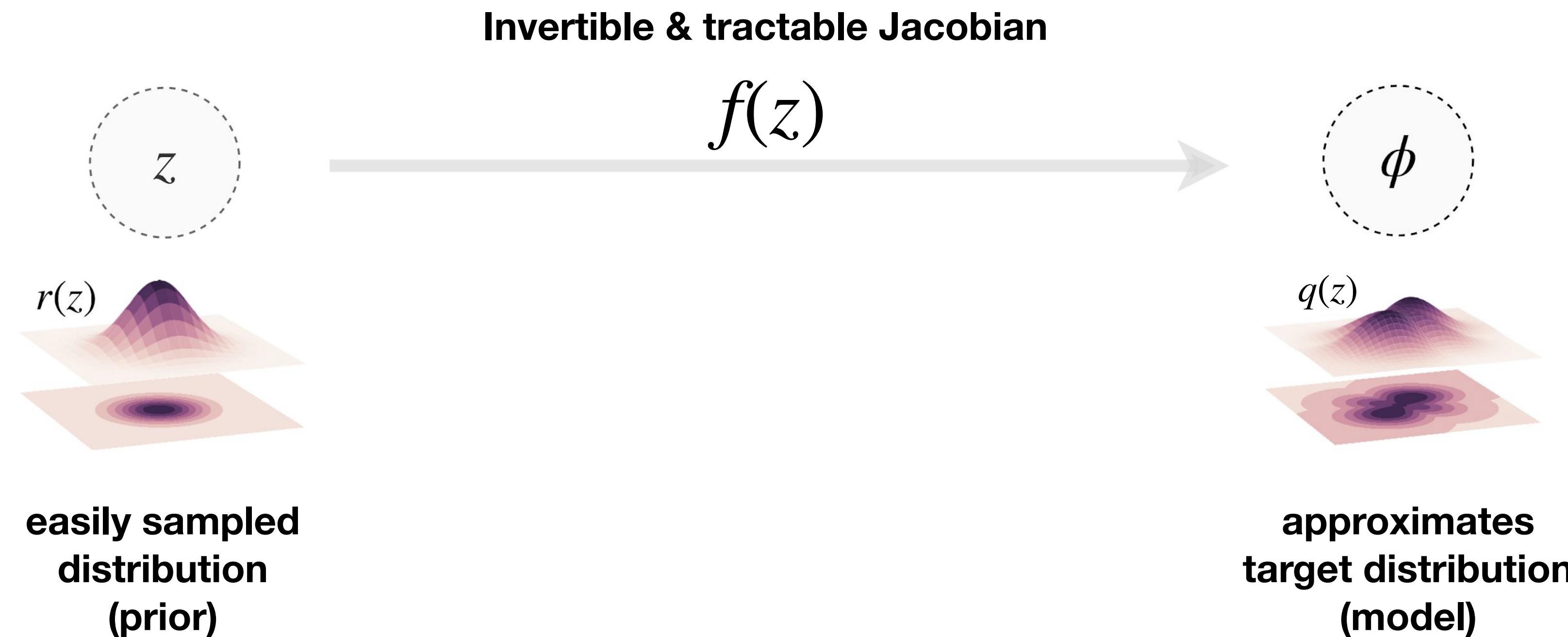


This talk:

Sampling in gauge theories
using pseudofermions

Generative flow models

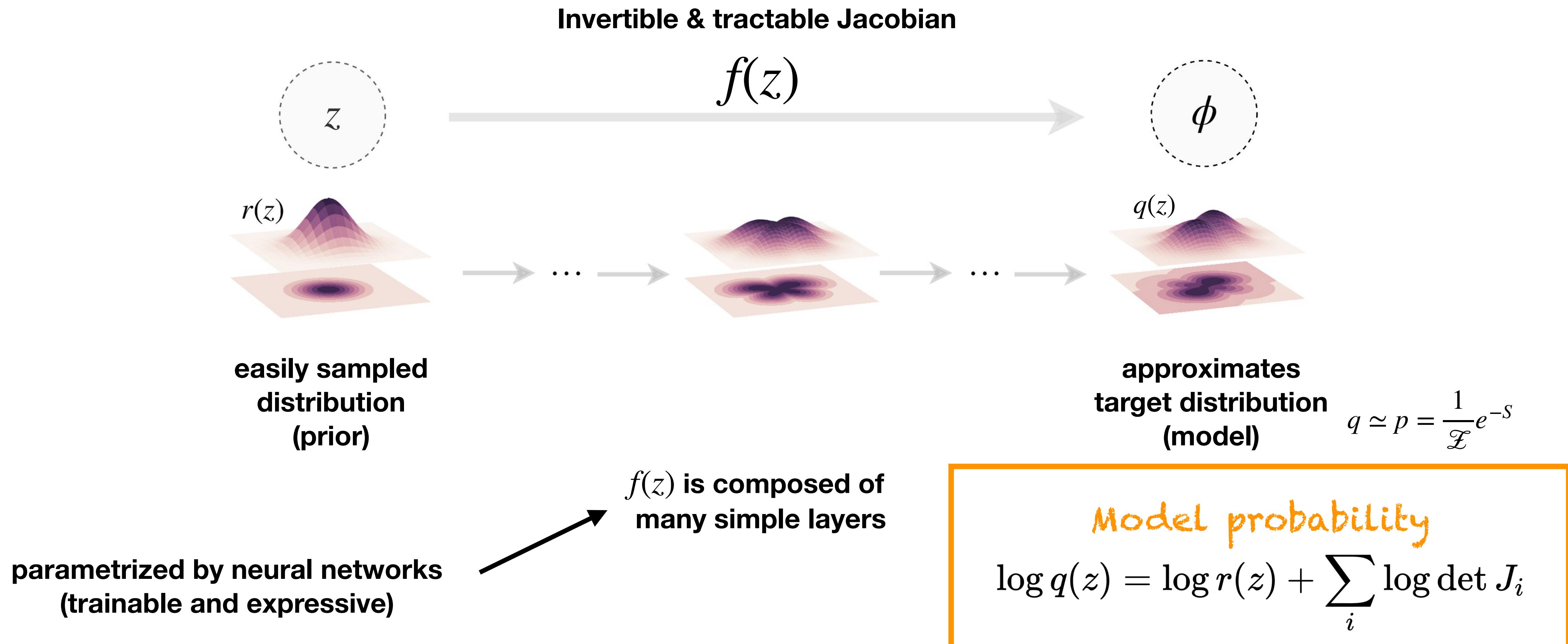
[Rezende, Mohamed, arXiv:1505.05770]



$$q \simeq p = \frac{1}{\mathcal{Z}} e^{-S}$$

Generative flow models

[Rezende, Mohamed, arXiv:1505.05770]



Training \neq using flows

- Train to minimize e.g. Kullback-Leibler divergence:

$$D_{\text{KL}}(q||p) = \int d\phi q(\phi)[\log q(U) + S(\phi)] + (\text{const})$$

- Self-training:
 1. Draw samples from the model to measure sample mean of $[\log q(U) + S(\phi)]$
 2. Gradient-based methods to optimize model parameters (e.g. Adam optimizer)

[Kingma, Ba, arXiv:1412.6980]

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! Trained models are not perfect but exactness is essential.

Recover exactness by e.g. forming a Markov Chain with accept/reject steps

Flows for fermionic gauge theories

- Several examples of flow models for pure gauge theories

[\[Kanwar et al, 2003.06413\]](#) [\[Boyda et al, 2008.05456\]](#)

- For fermions: straightforward approach consists on integrating them out

$$S_E(U) = -\beta \sum_{\text{(Plaquette)}} \text{Re } P(x) - \log \det D[U]^\dagger D[U]$$

(fermion determinant)

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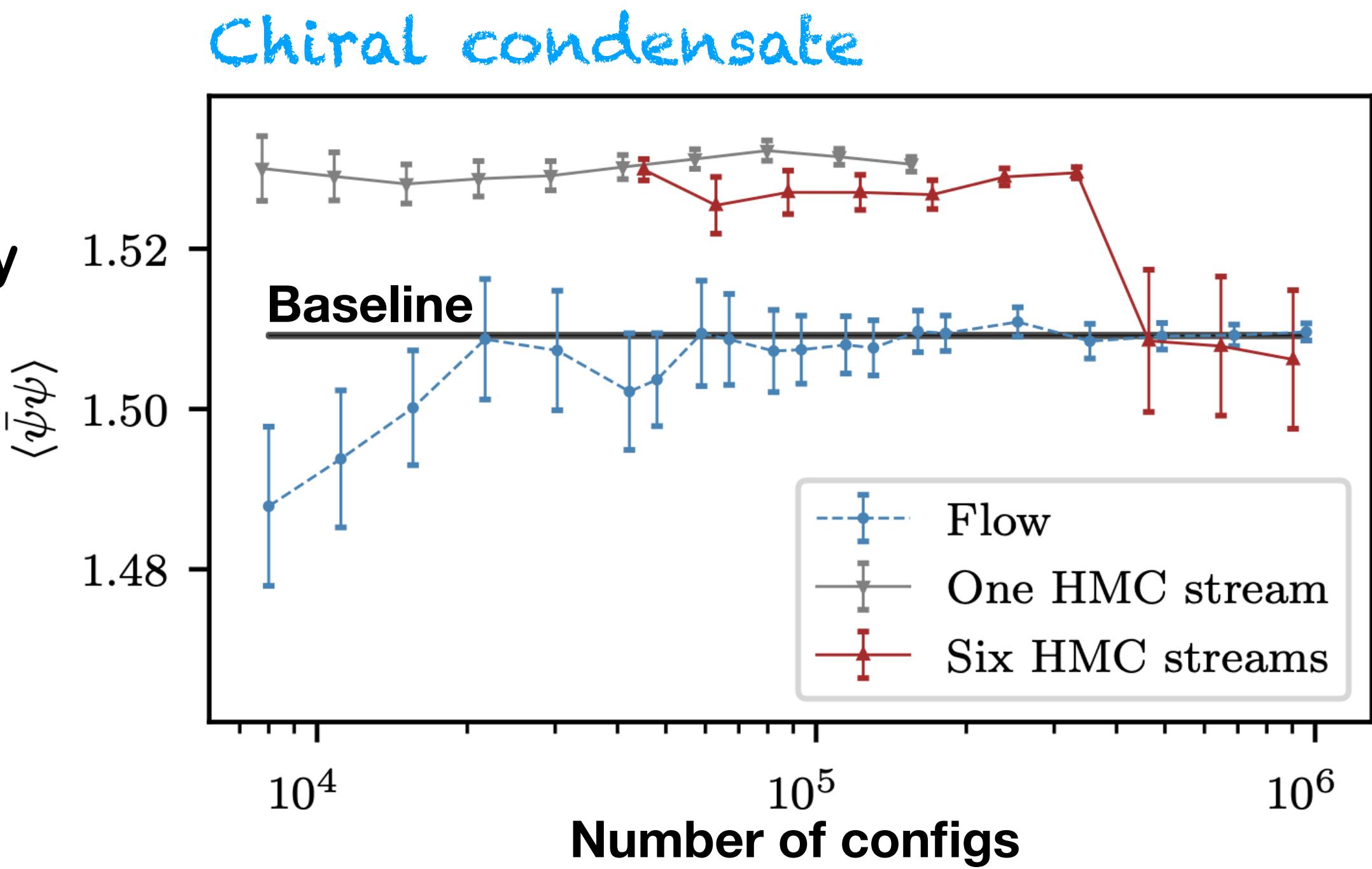
(fermion determinant)

- Successfully applied to the Schwinger model at criticality

[Albergo, Boyda, Cranmer, Hackett, Kanwar, Racanière, Rezende, FRL, Shanahan, Urban 2202.11712]

- ✗ HMC shows biased results with underestimated errors

- ✓ Flow-based sampling provides correct results



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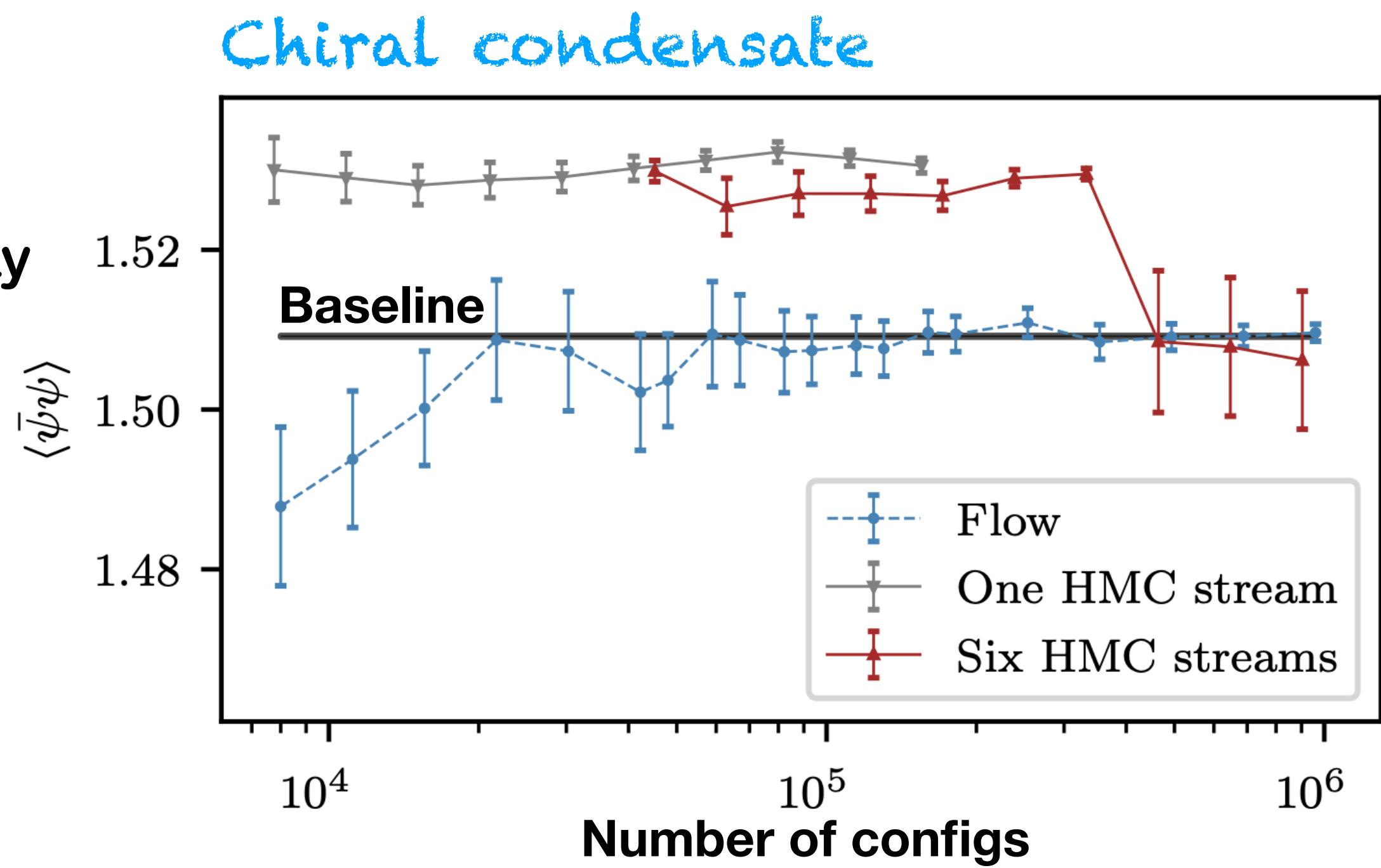
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- ! Evaluation of the fermion determinant is expensive.

Not feasible for QCD-scale calculations!

- Scalable approach: use stochastic determinant estimators!

→ Pseudofermions!



Recap of Pseudofermions

- Action:

$$S(\psi, \bar{\psi}, U) = S_g(U) + S_F(\psi, \bar{\psi}, U)$$

Gauge fields  **Fermions** 

$$S_F(\psi, \bar{\psi}, U) = \sum_{f=1}^{N_f} \bar{\psi}_f D_f(U) \psi_f \quad \longrightarrow \quad \int \mathcal{D}[\psi, \bar{\psi}] e^{-S_F(\psi, \bar{\psi}, U)} = \prod_{f=1}^{N_f} \det D_f(U)$$

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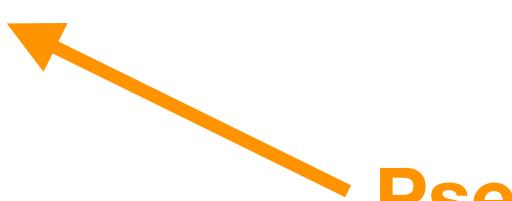
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- Evaluate determinant using auxiliary degrees of freedom

Assuming $N_f=2$ 

$$\det DD^\dagger = \frac{1}{Z_N} \int \mathcal{D}\phi e^{-\phi^\dagger (DD^\dagger)^{-1} \phi}$$

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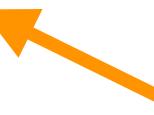
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 **Pseudofermions**

- Joint target distribution

$$p(U, \phi) = \frac{1}{Z} e^{-S_g(U) - S_{\text{pf}}(U, \phi)} \quad \text{with} \quad S_{\text{pf}}(\phi, U) = \phi^\dagger [D(U) D^\dagger(U)]^{-1} \phi^\dagger$$

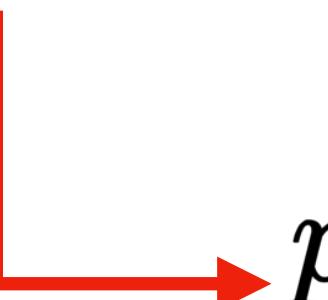
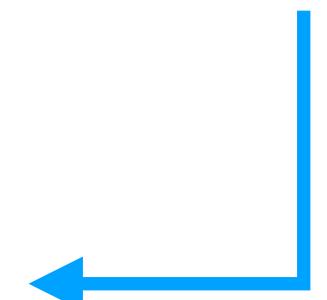
- Only need the Dirac operator applied to the PF field: **scales linearly with the lattice volume**

Joint flow models

$$p(U, \phi) = p(U)p(\phi | U)$$

$$p(U) \propto \det DD^\dagger(U) e^{-S_g(U)}$$

“marginal”



$$p(\phi|U) \propto \frac{1}{\det DD^\dagger(U)} e^{-S_{PF}(\phi|U)}$$

“conditional”

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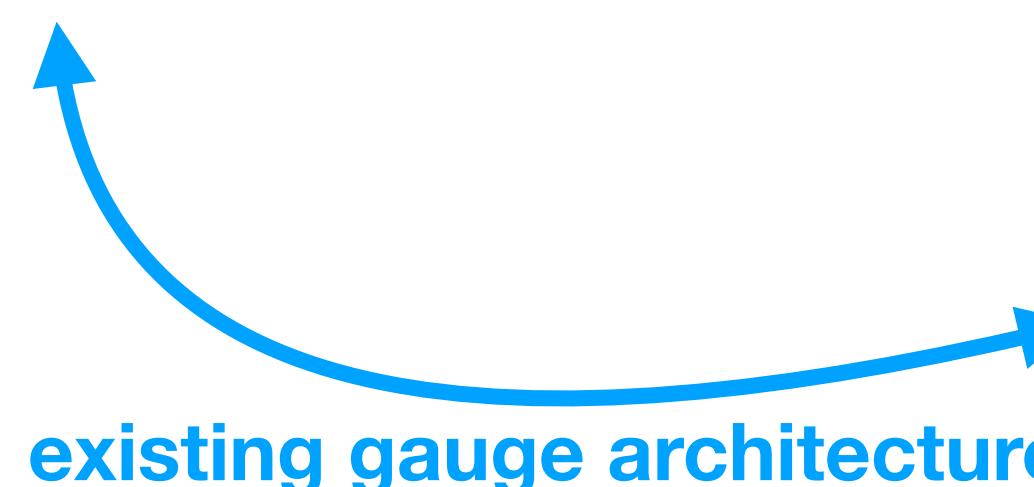
Different flow models to approximate
marginal and **conditional** distributions

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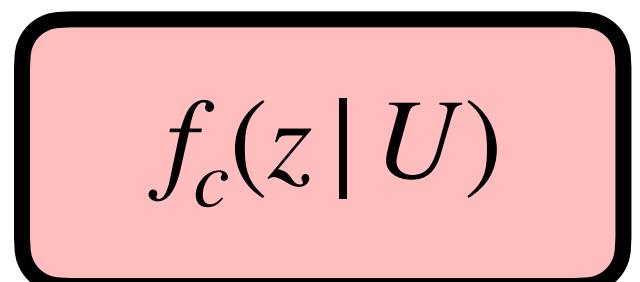


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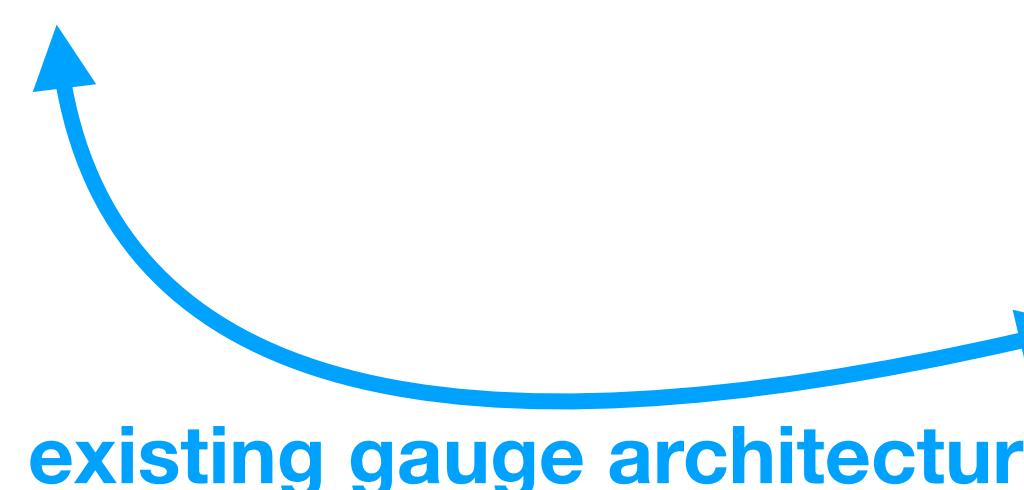
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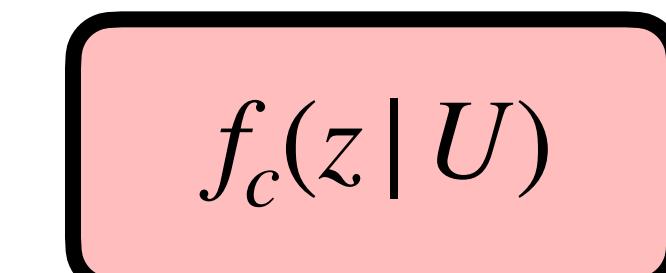


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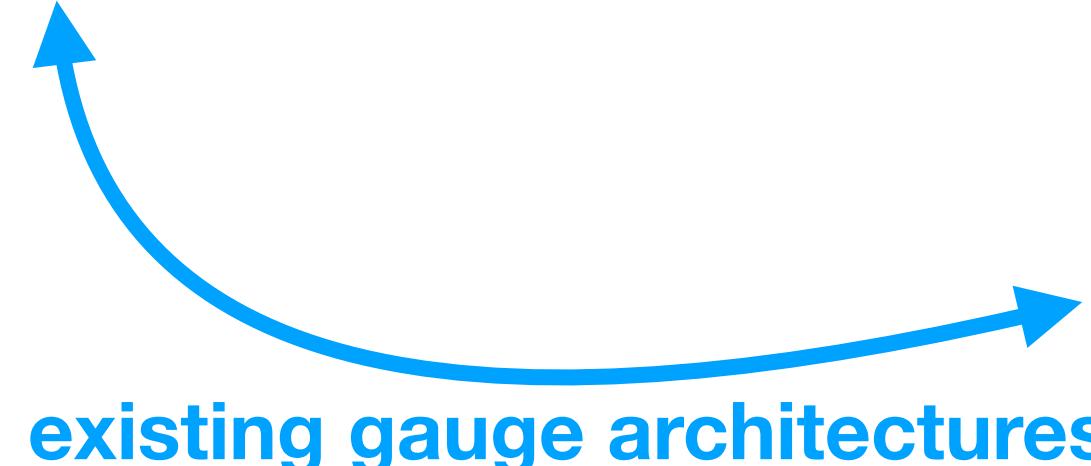
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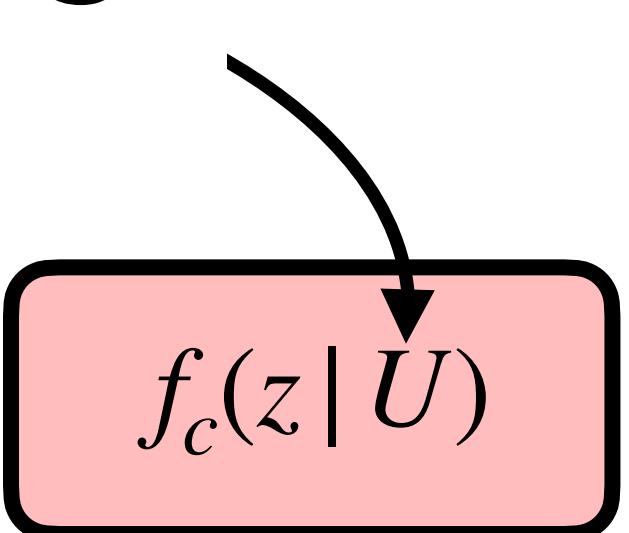
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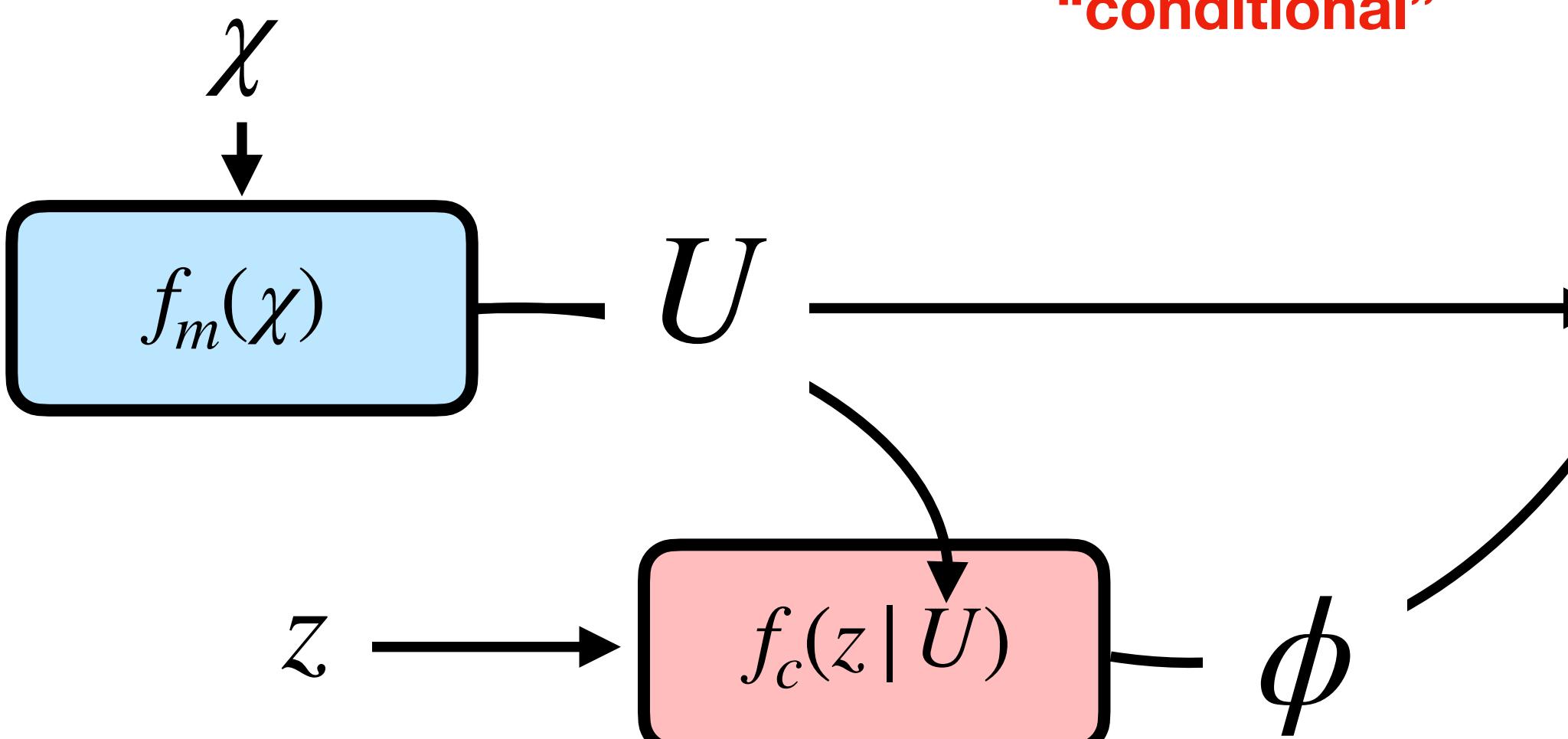
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“marginal”

existing gauge architectures
[Kanwar et al, 2003.06413]
[Boyda et al, 2008.05456]



proposed configuration
 $q(U)q(\phi | U)$

Different flow models to approximate
marginal and **conditional** distributions

new conditional architectures

Gauge-equivariant conditional models

- Train a flow to map an uncorrelated gaussian into a correlated one:

$$r(z) \propto e^{-z^\dagger z} \xrightarrow{f_c(z \mid U)} q(\phi \mid U) \propto e^{-\phi^\dagger A(U)\phi} \simeq e^{-\phi^\dagger (D(U)D^\dagger(U))^{-1}\phi}$$

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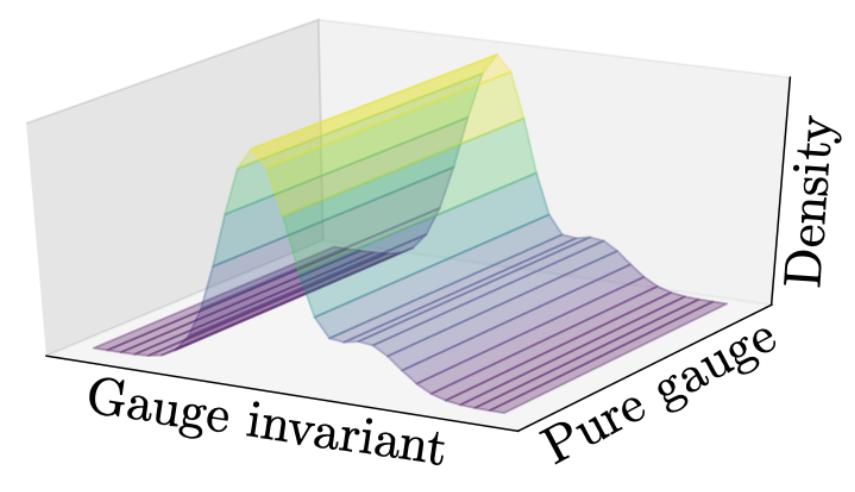
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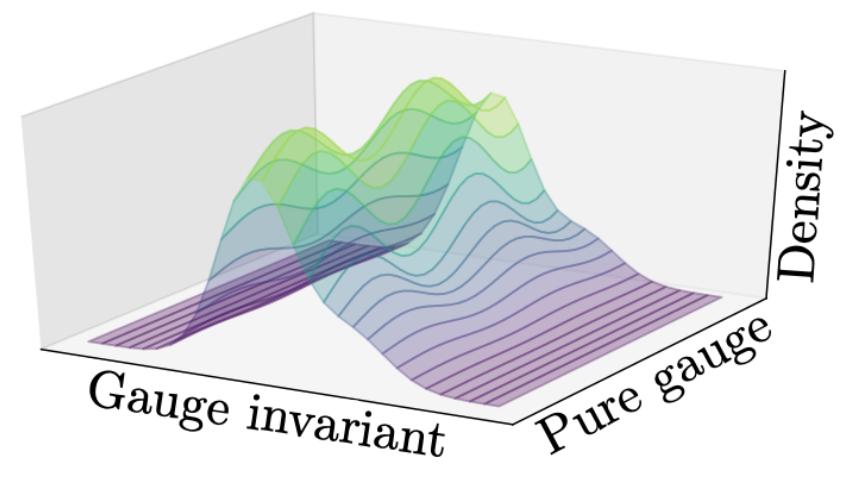
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- Gauge equivariance simplifies the problem:



True distribution



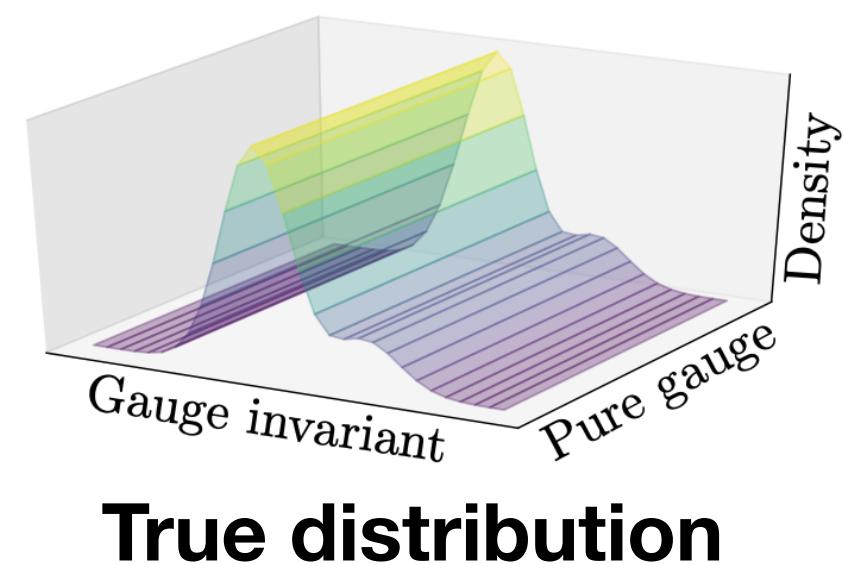
Learned by naive ML

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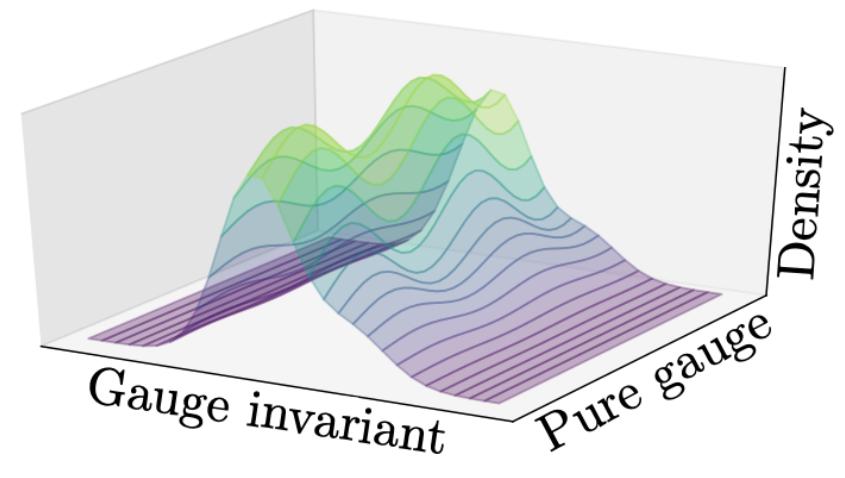
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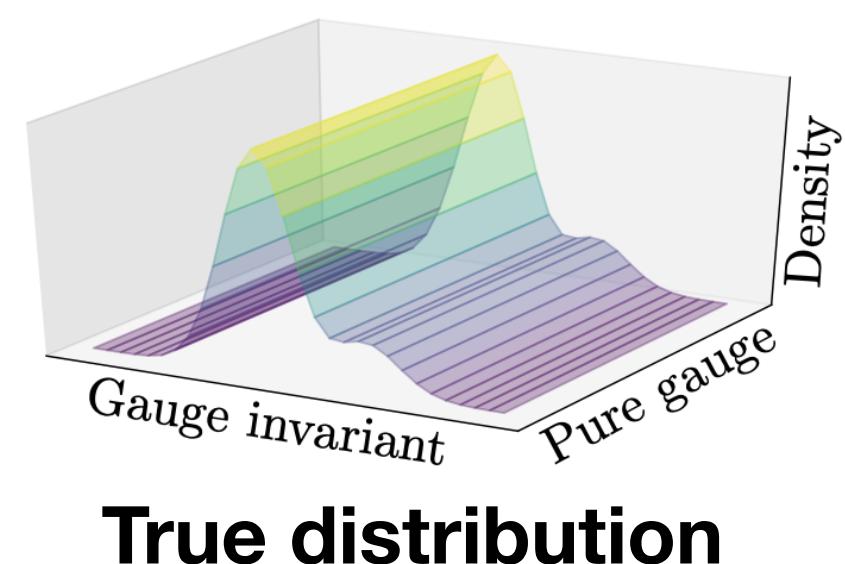
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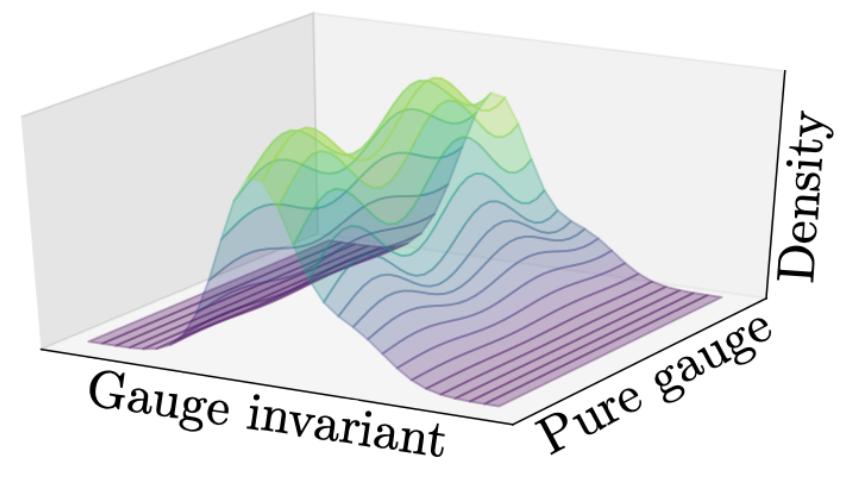
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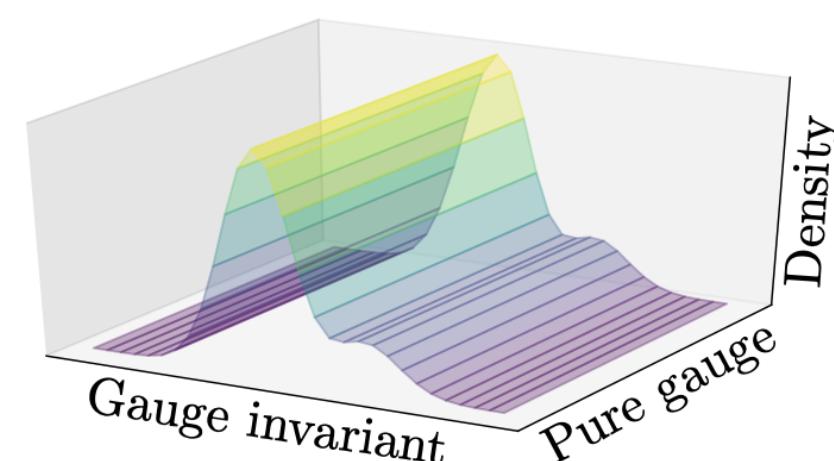
parallel-transported
neighbor

Gauge-equivariant conditional models

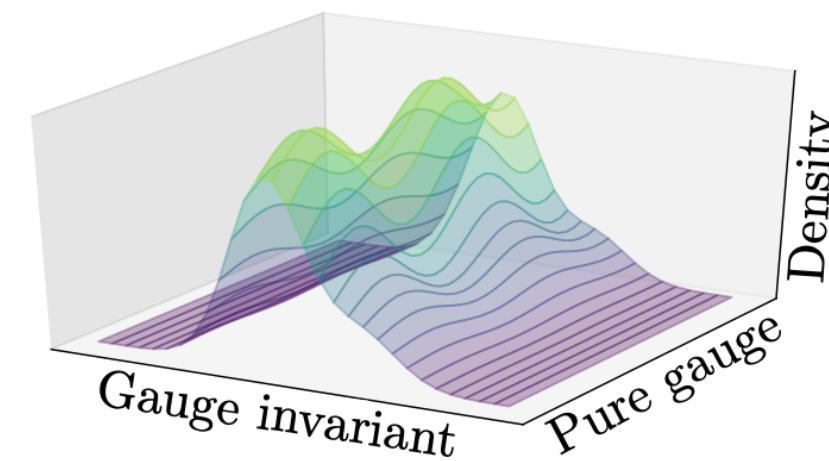
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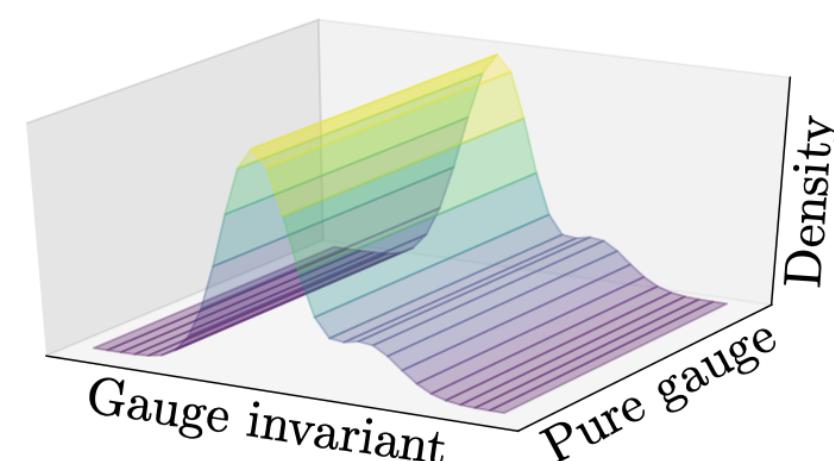
NN outputs
(gauge-invariant inputs)

Gauge-equivariant conditional models

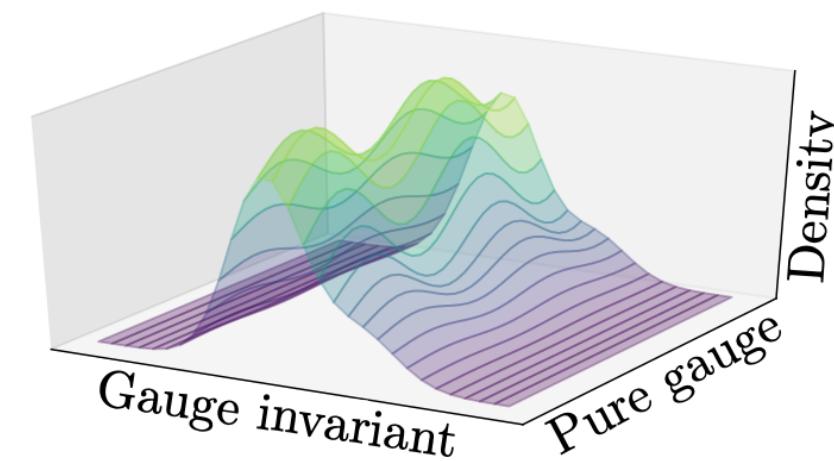
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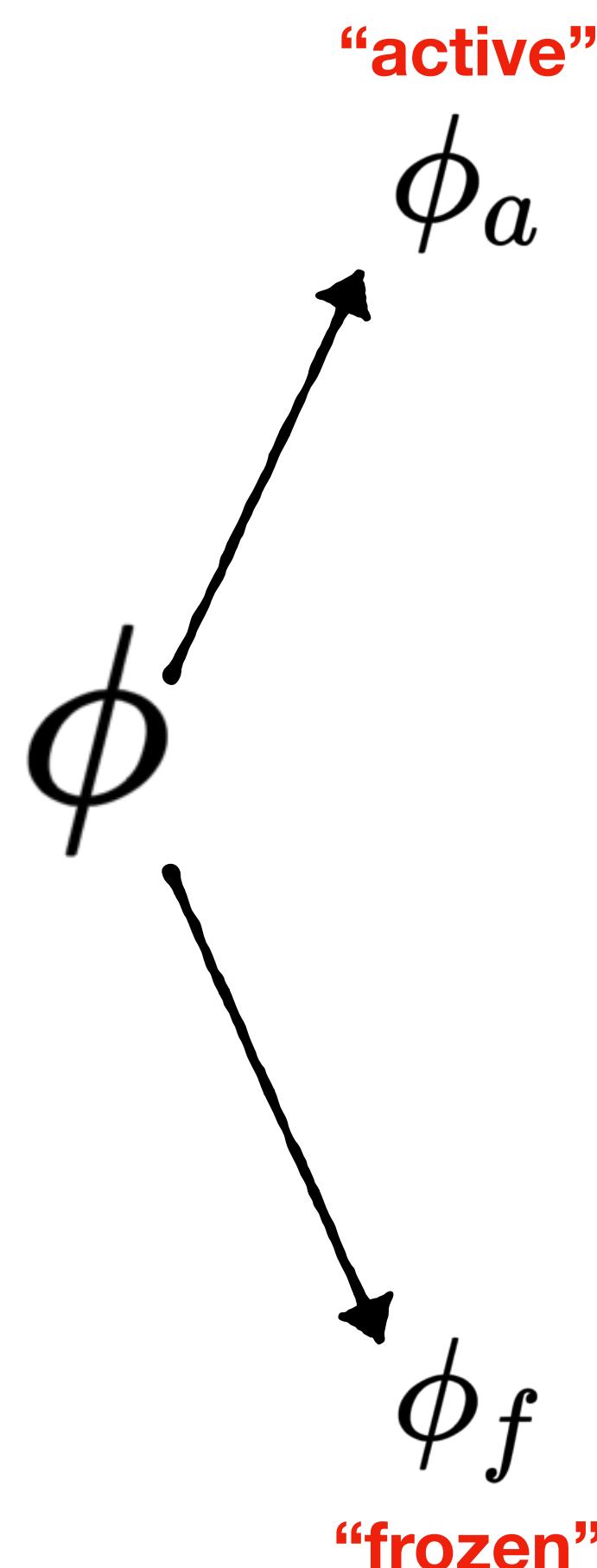
(gauge equivariance)

$A(U)$ and $B(U)$ are circled in red. Red arrows point from these circles to the text "NN outputs (gauge-invariant inputs)" and "parallel-transported neighbor" respectively.

NN outputs (gauge-invariant inputs)

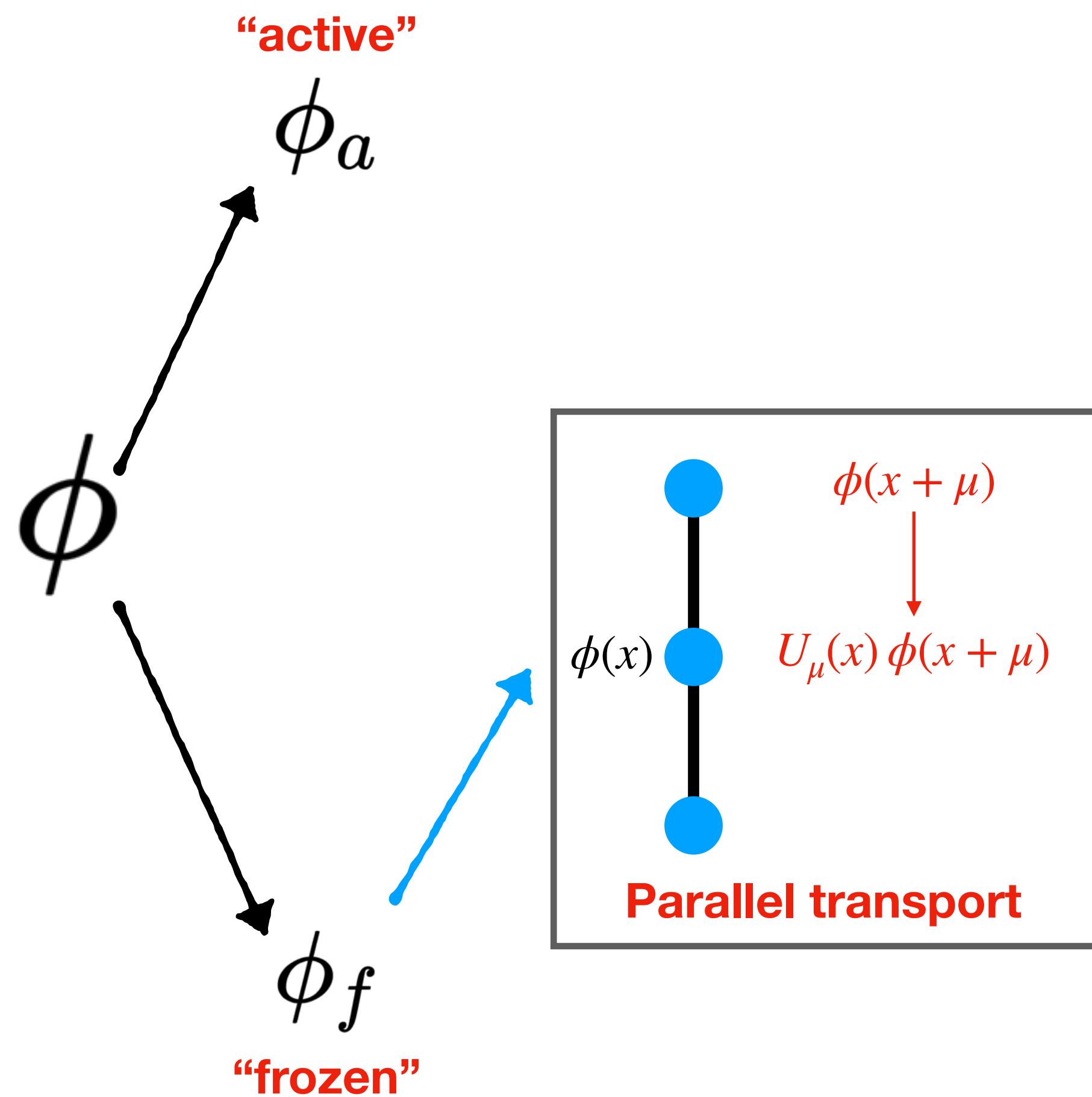
parallel-transported neighbor

Sketch of pseudofermion Layers



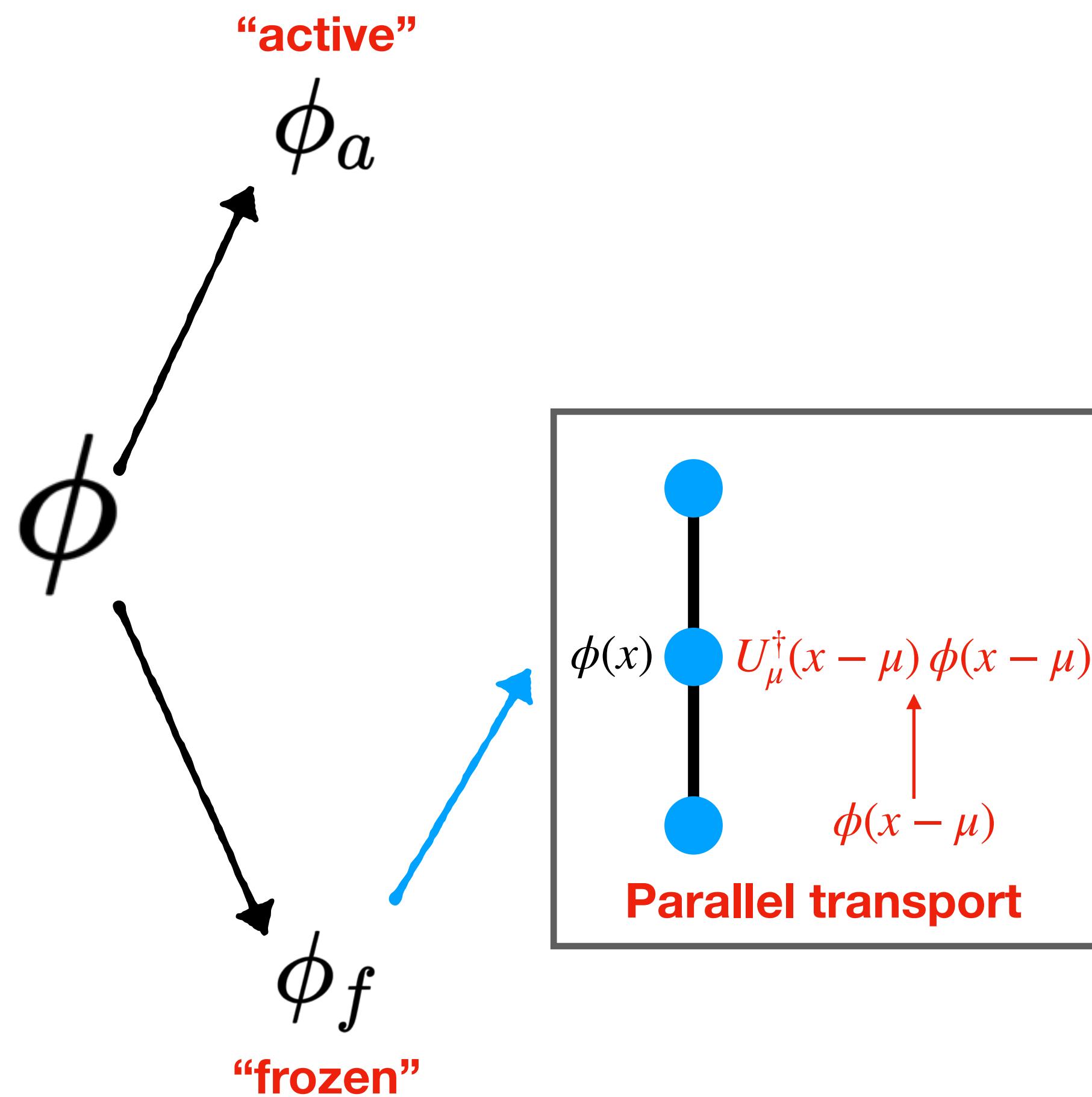
variable partitioning
ensures tractible Jacobian

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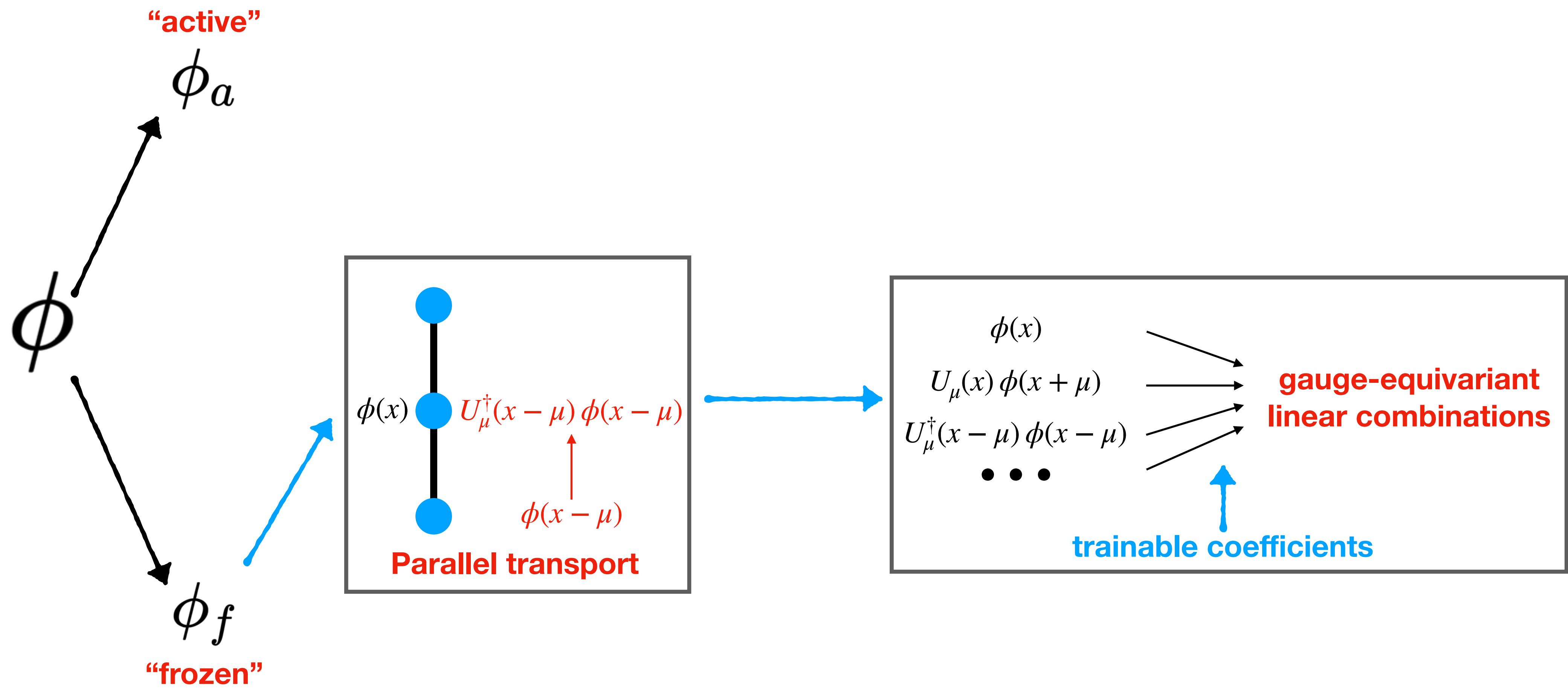
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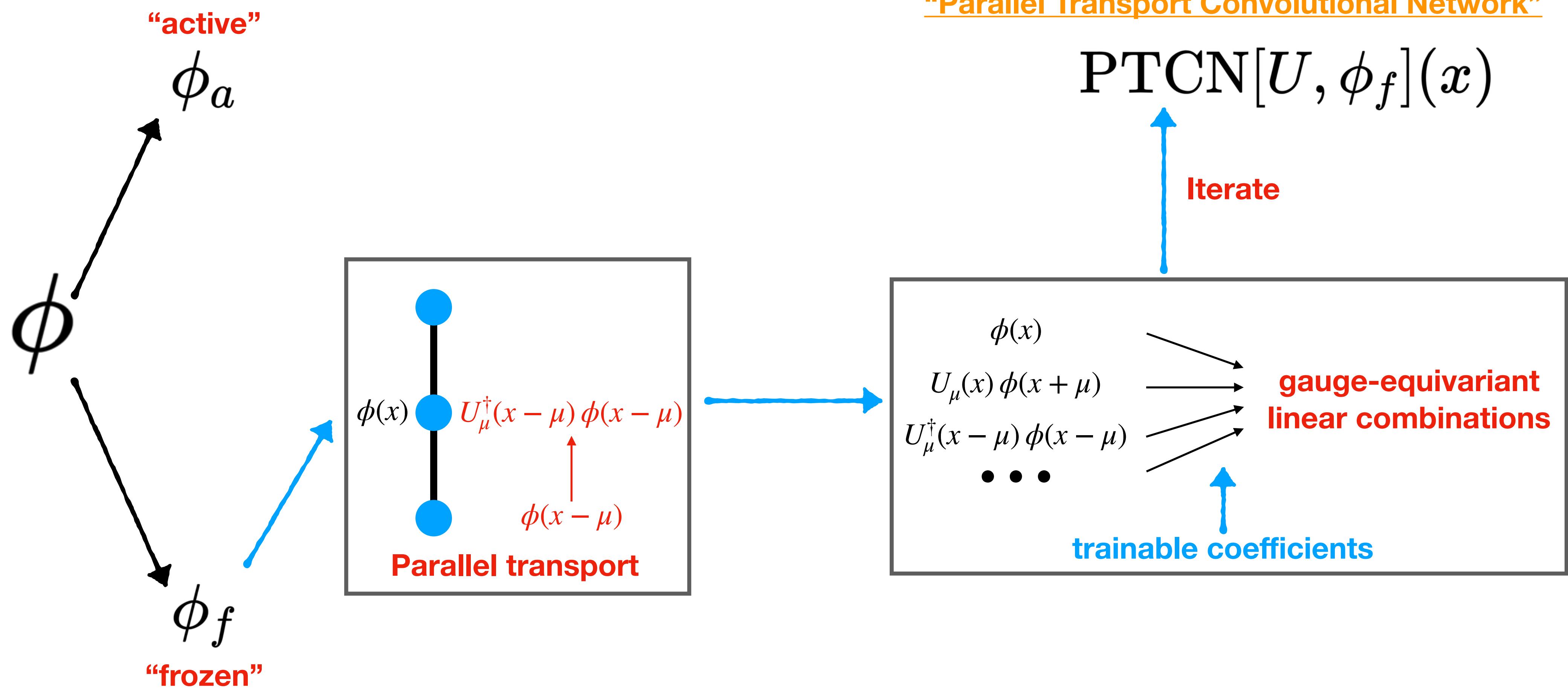
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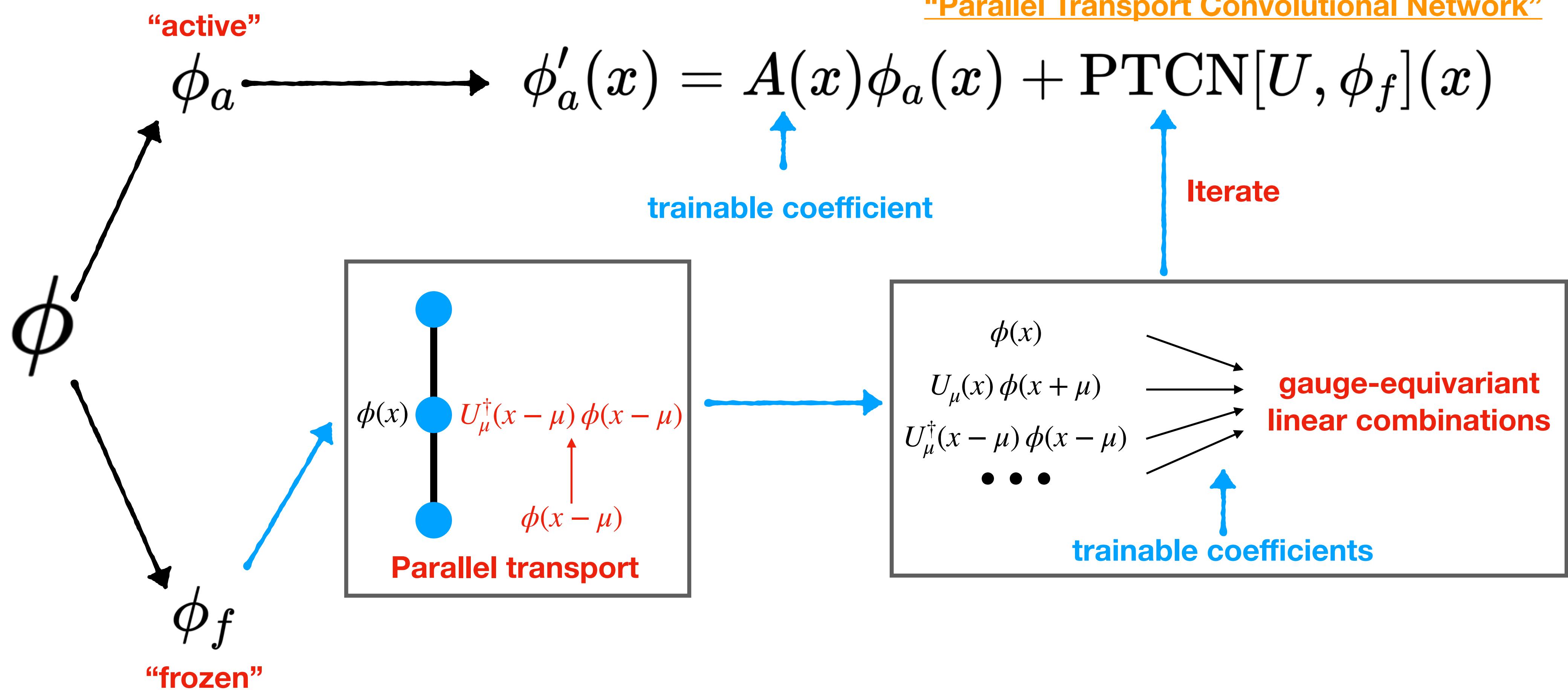
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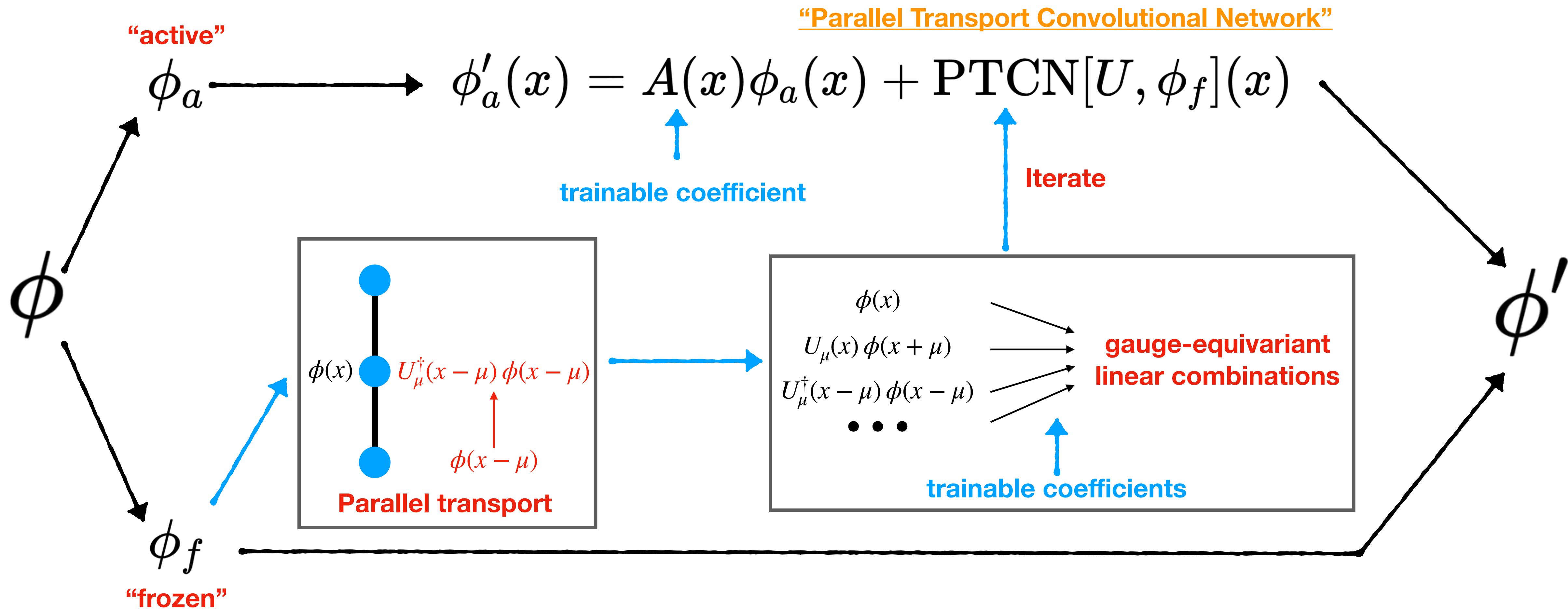


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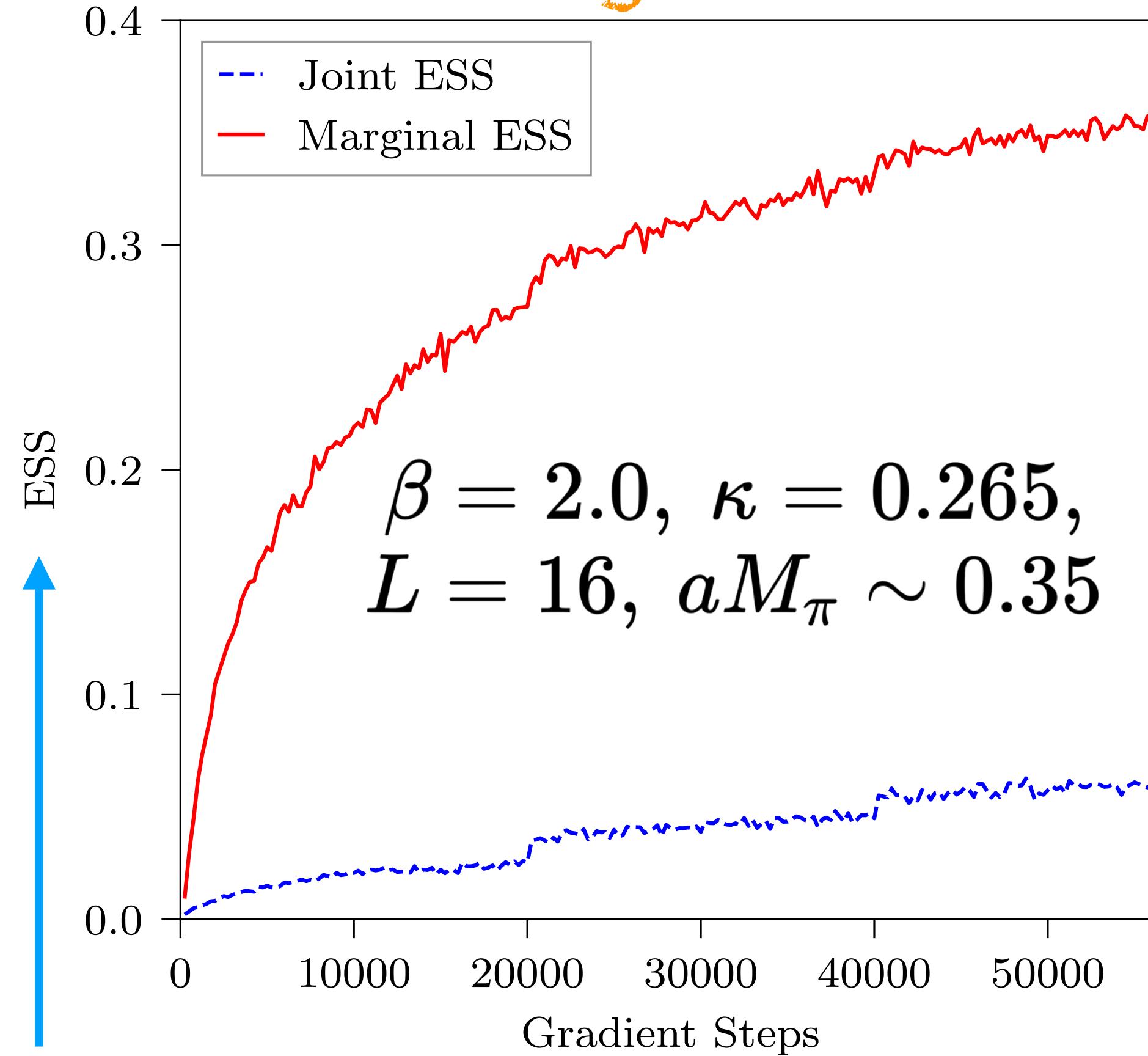
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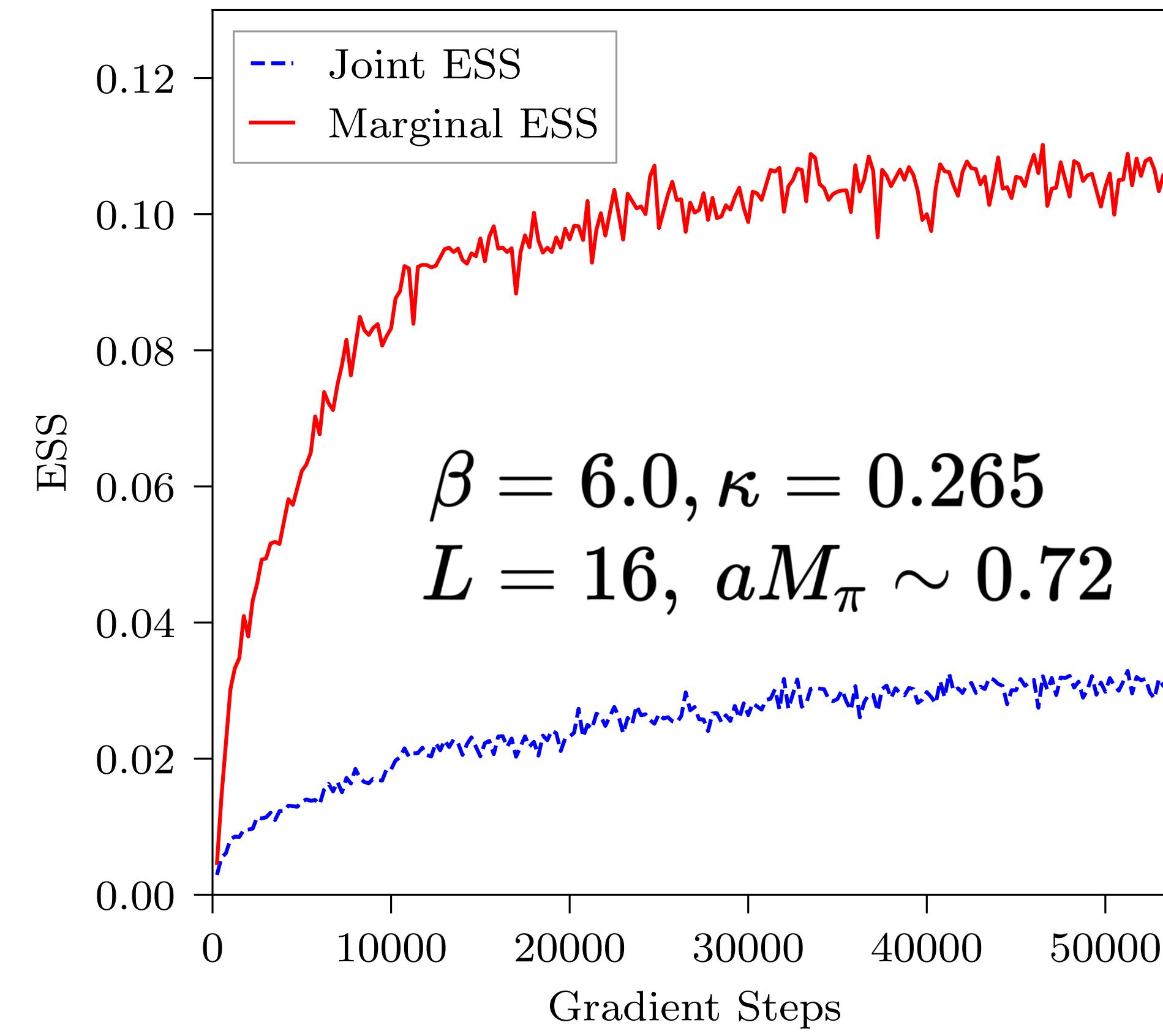
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Examples of flow models

Schwinger Model



2D SU(3) with $N_f=2$



"Effective sample size" \equiv model quality
(ESS=1 for a perfect model)

[Abbott, Albergo, Boyda, Cranmer, Hackett, Kanwar, Racanière, Rezende, FRL, Shanahan, Tian, Urban 2207.08945]

Use of preconditioners

○ Use joint models in combination with preconditioners

[Abbott, Albergo, Boyda, Cranmer, Hackett, Kanwar, Racanière, Rezende, FRL, Shanahan, Tian, Urban 2207.08945]

► Even/odd (EO) preconditioning

$$D = \begin{pmatrix} \mathbb{1} & D_{eo} \\ D_{oe} & \mathbb{1} \end{pmatrix} \quad \det D = \det (\mathbb{1} - D_{eo}D_{oe}) \equiv \det D_{sc}$$

◆ Less PF variables in target

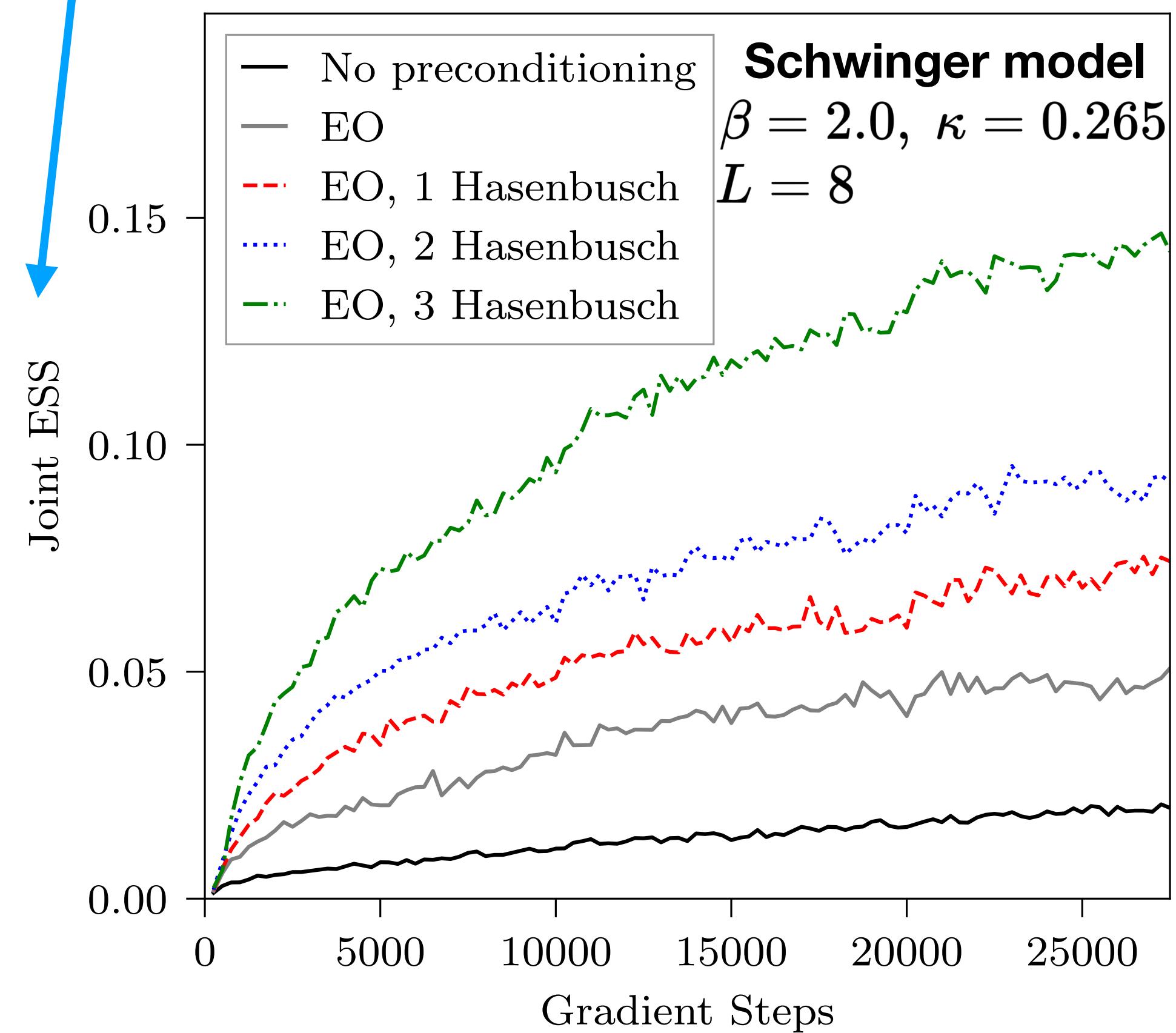
► Hasenbusch factorization

$$\det M = \left[\frac{\det M}{\det (M + \mu)} \right] \det (M + \mu)$$

◆ Reduces the condition number of target

“Effective sample size” \equiv model quality

Significant improvement by using preconditioners!



Improving determinant estimators

- Use joint models more efficiently by drawing multiple pseudofermion samples at fixed gauge field

$$w_{N_{\text{pf}}}(U) = \frac{1}{N_{\text{pf}}} \sum_{i=1}^{N_{\text{pf}}} \frac{p(\phi^{(i)}, U)}{q(\phi^{(i)}, U)}$$

 weight of each config for reweighting or Markov Chain

- This does not require reevaluating the observables

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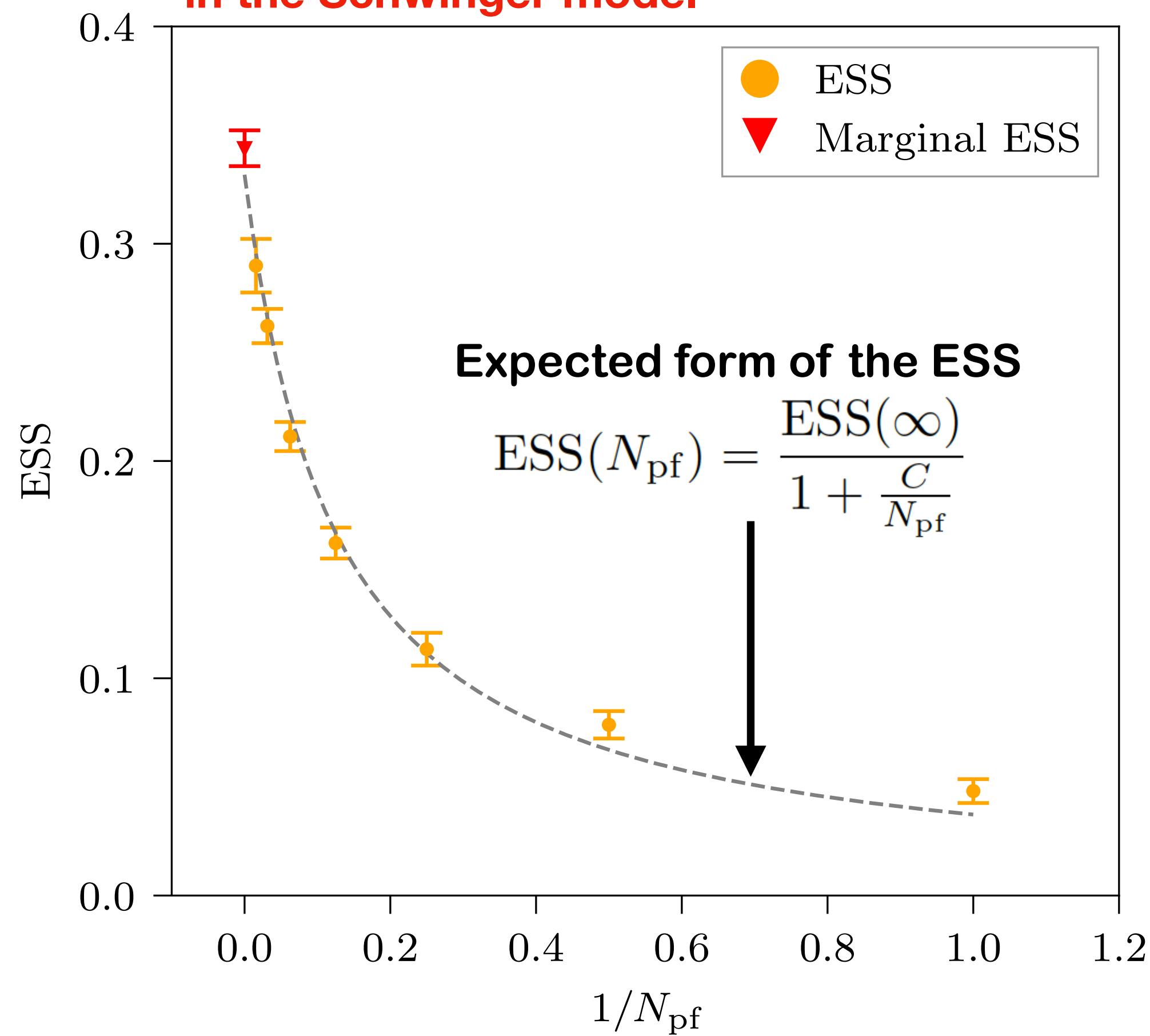
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↑
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Numerical demonstration
in the Schwinger model



Summary & Outlook

- Pseudofermions can be used in flows as a **scalable approach** to the fermion determinant
- We have developed **architectures** for flow models using pseudofermions in gauge theories
 - ▶ **Parallel Transport Convolutional Network** is the central element
[Abbott, Albergo, Boyda, Cranmer, Hackett, Kanwar, Racanière, Rezende, FRL, Shanahan, Tian, Urban 2207.08945]
- Successful proof of concept **applications to 2D theories**: **Schwinger model** and **SU(3) in 2D**
 - ▶ Improve model quality with **preconditioners**
 - ▶ Can use **multiple pseudofermion draws** to better estimate $\det DD^\dagger[U]$
- No fundamental limitation to apply this technology to QCD in 4D
 - ▶ Further engineering still needed for high quality models
[See talk by Phiala Shanahan later in this session!]

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Thanks!

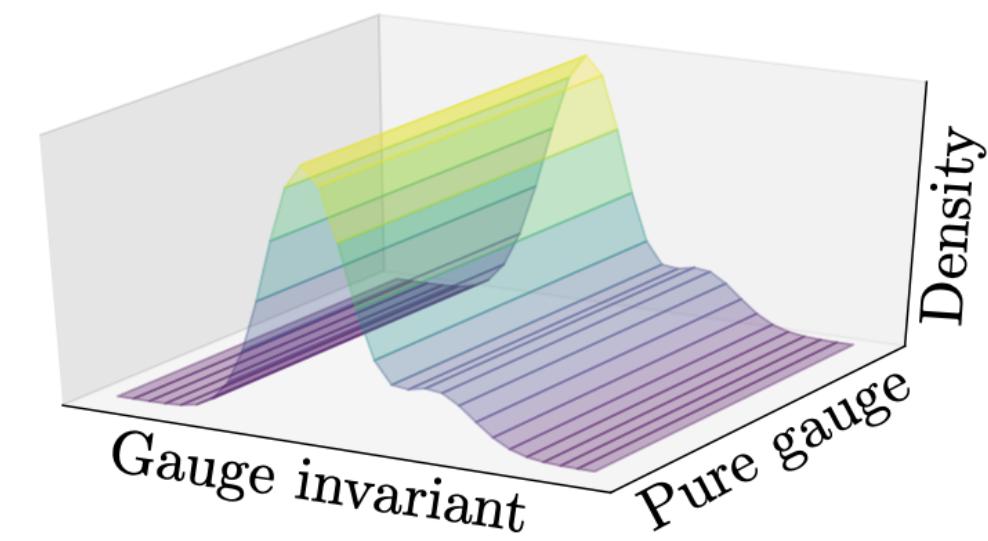
Back-up slides

Incorporating symmetries

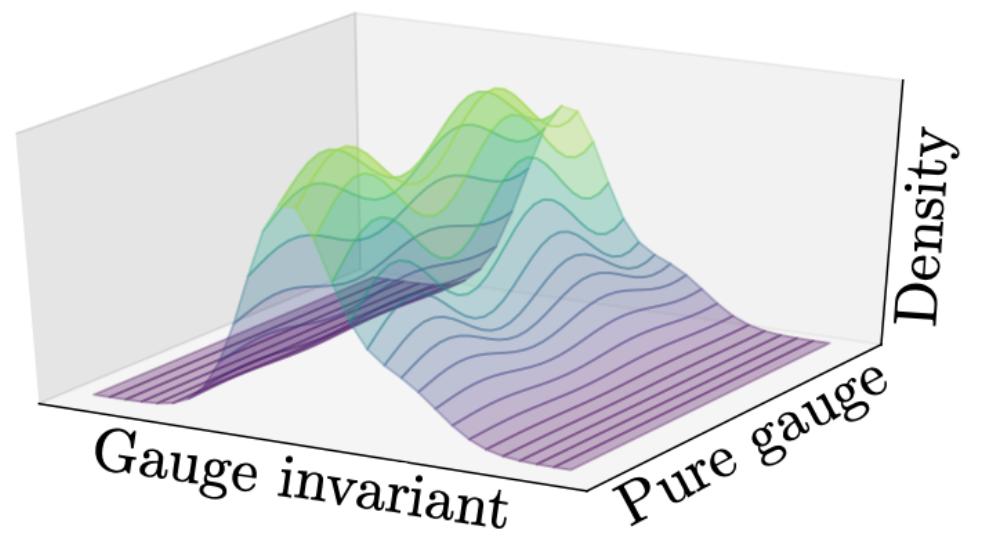
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- ✓ Reduces complexity of training
- ✓ Reduces parameter count

Gauge symmetries:



True distribution



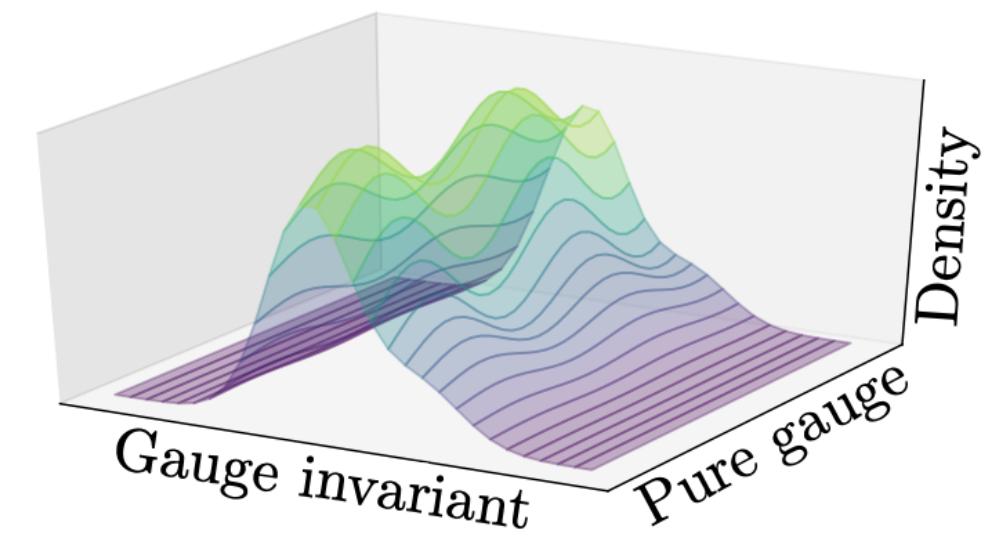
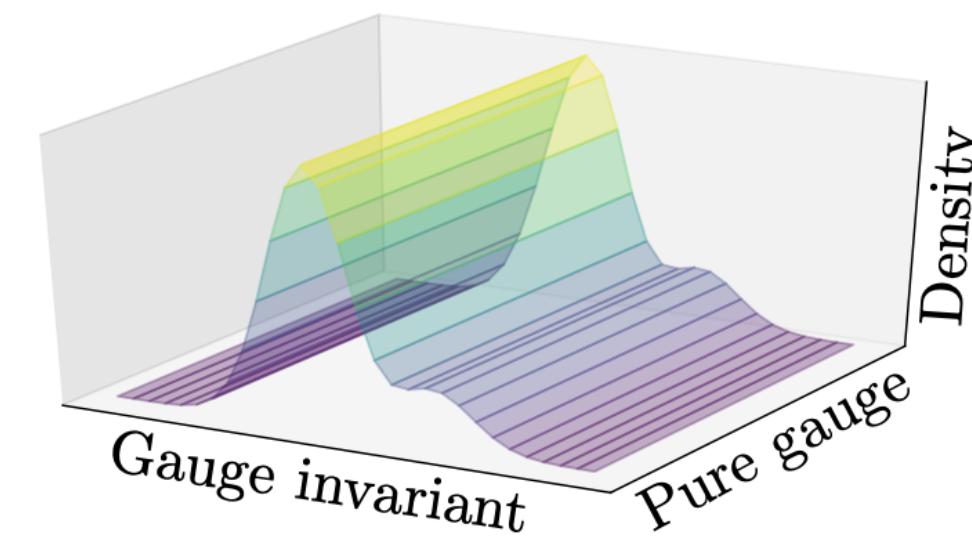
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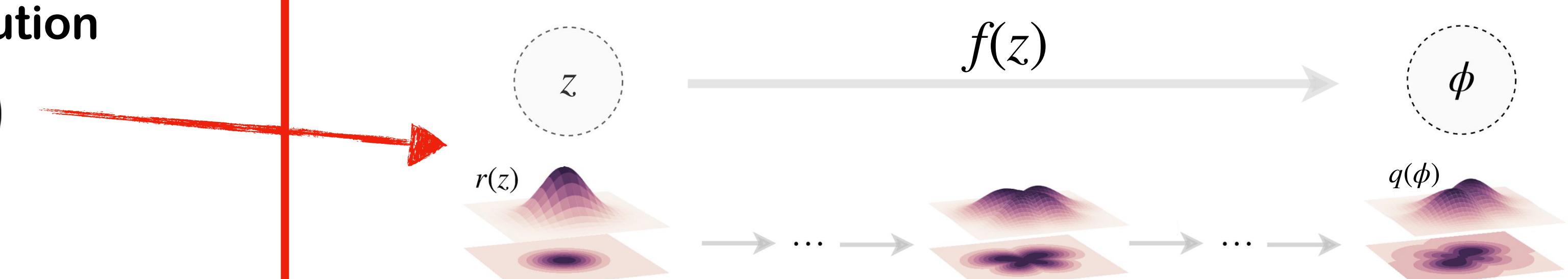
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General approach:

1. Invariant base distribution

$$r(z) = r(t(z))$$

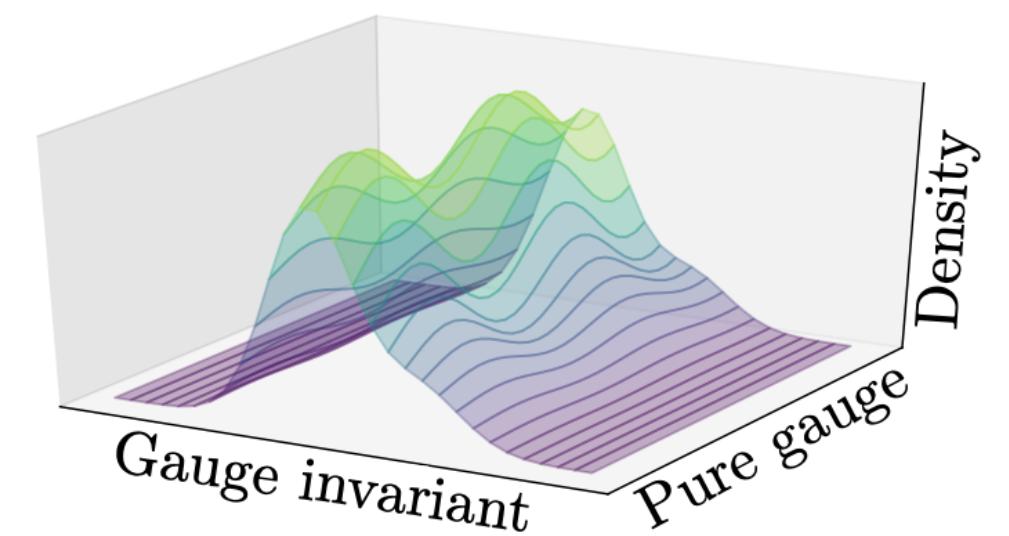
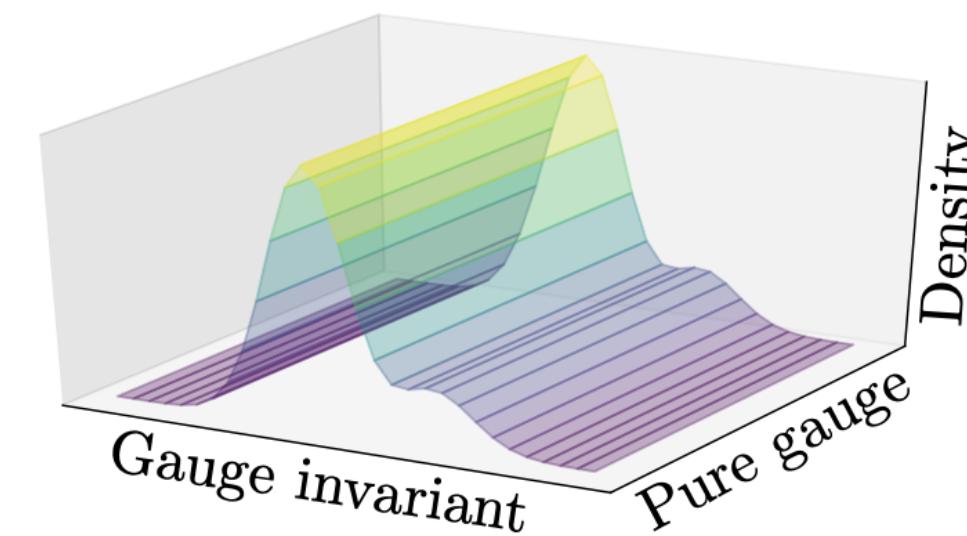


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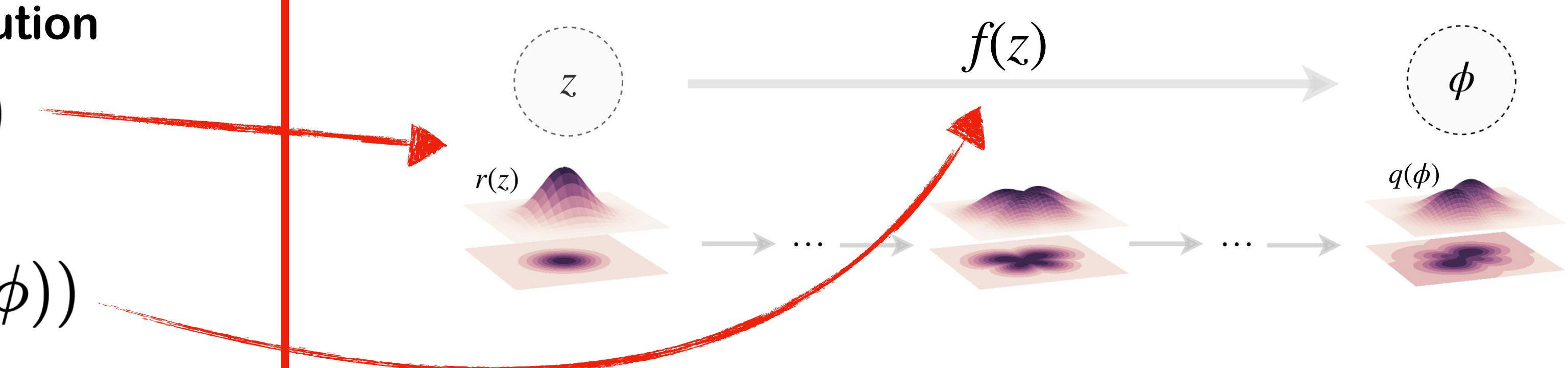
General approach:

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2. Equivariant flow

$$f(t(\phi)) = t(f(\phi))$$



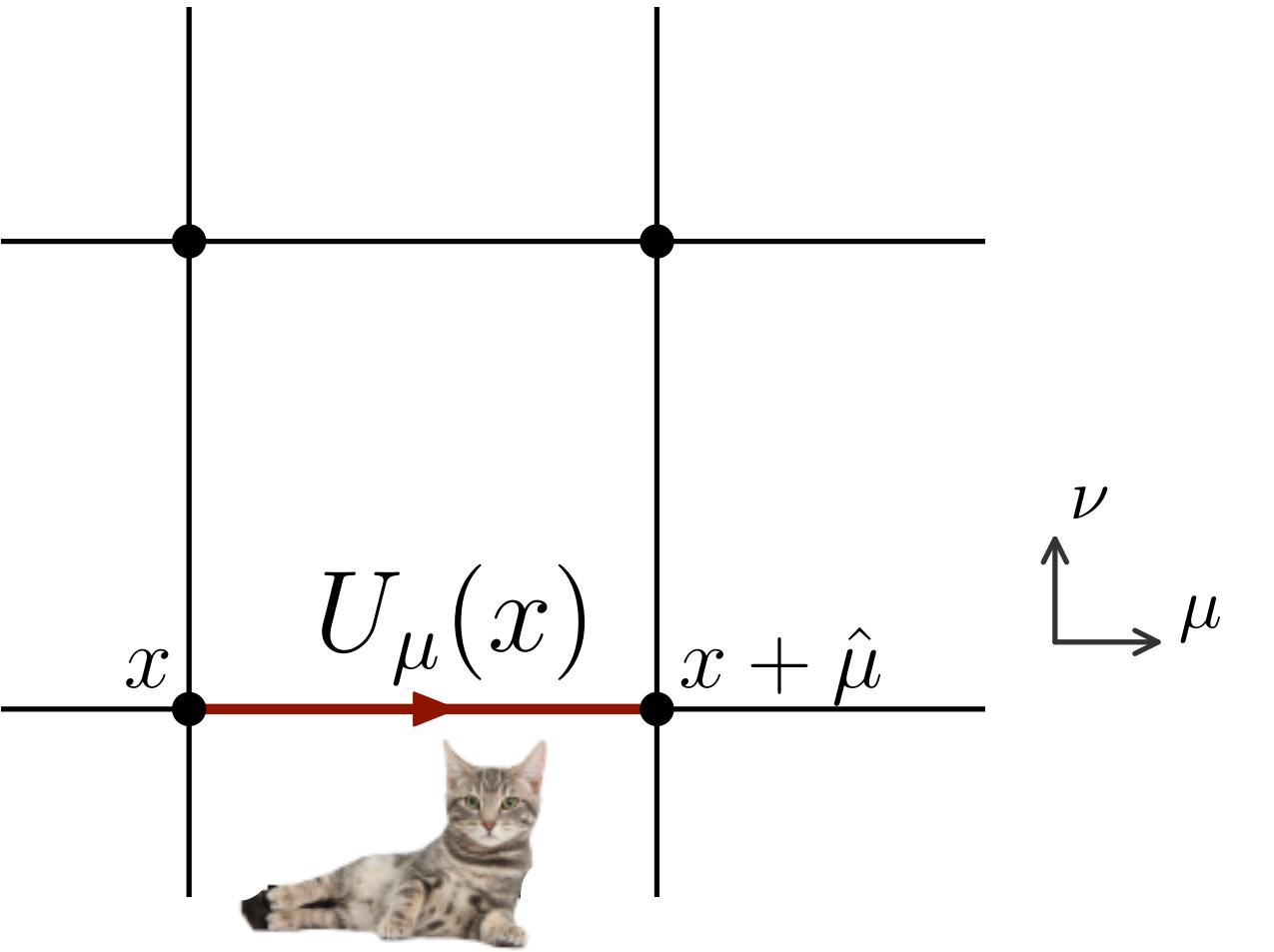
Lattice U(1) gauge symmetry

- Gauge variables are the gauge links

$$U_\mu(x) \in \mathrm{U}(1)$$

$$U_\mu(x) = e^{i a g A_\mu(x)}$$

$$a g A_\mu(x) \in [0, 2\pi)$$



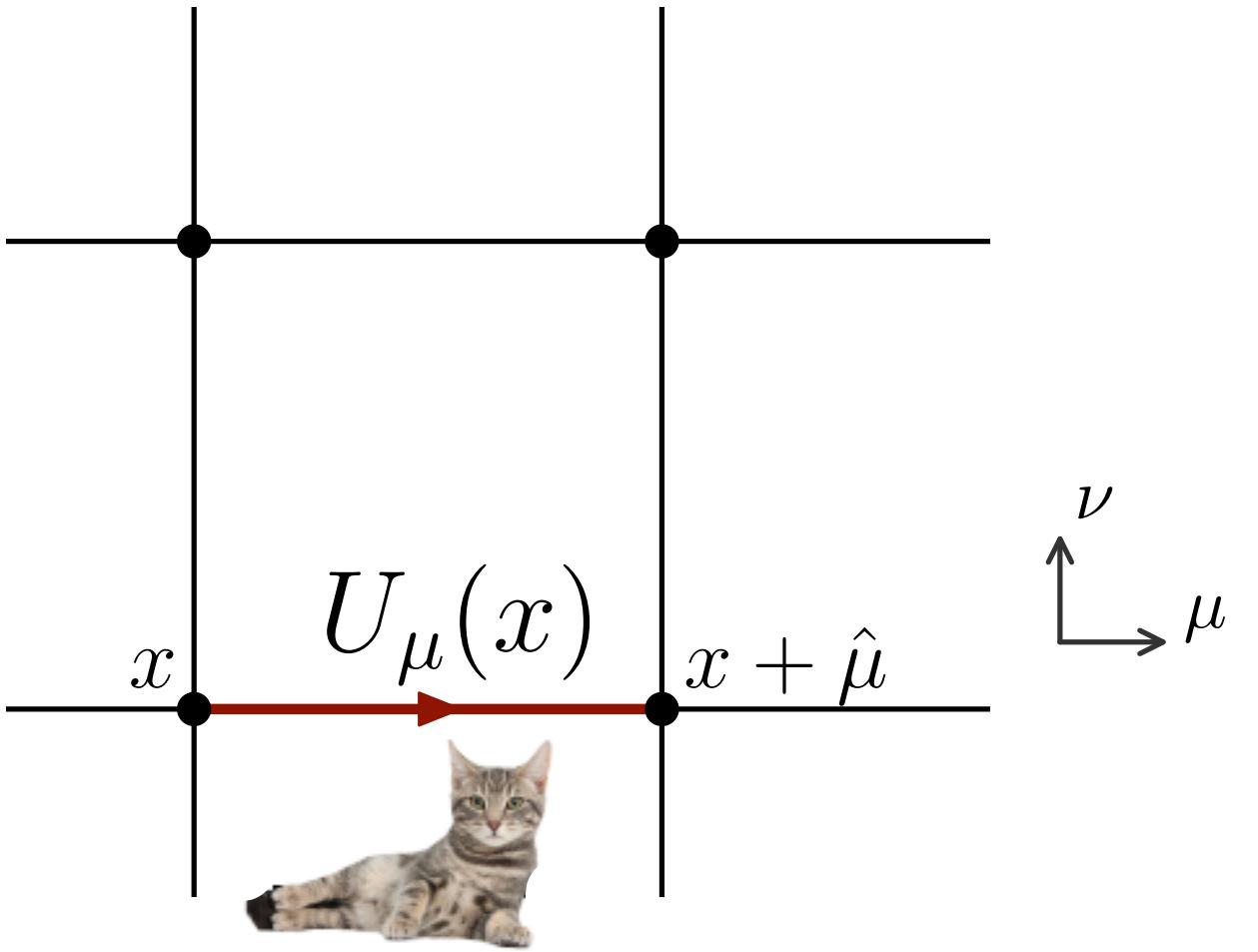
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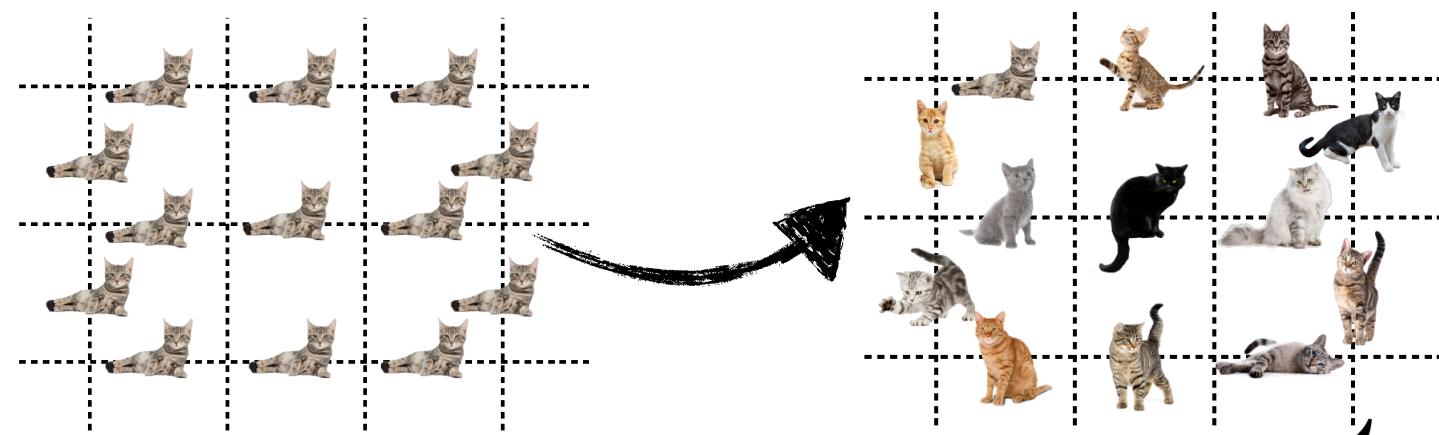
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- A gauge transformation is:

$$U_\mu(x) \rightarrow \Omega(x) U_\mu(x) \Omega^\dagger(x + \hat{\mu})$$

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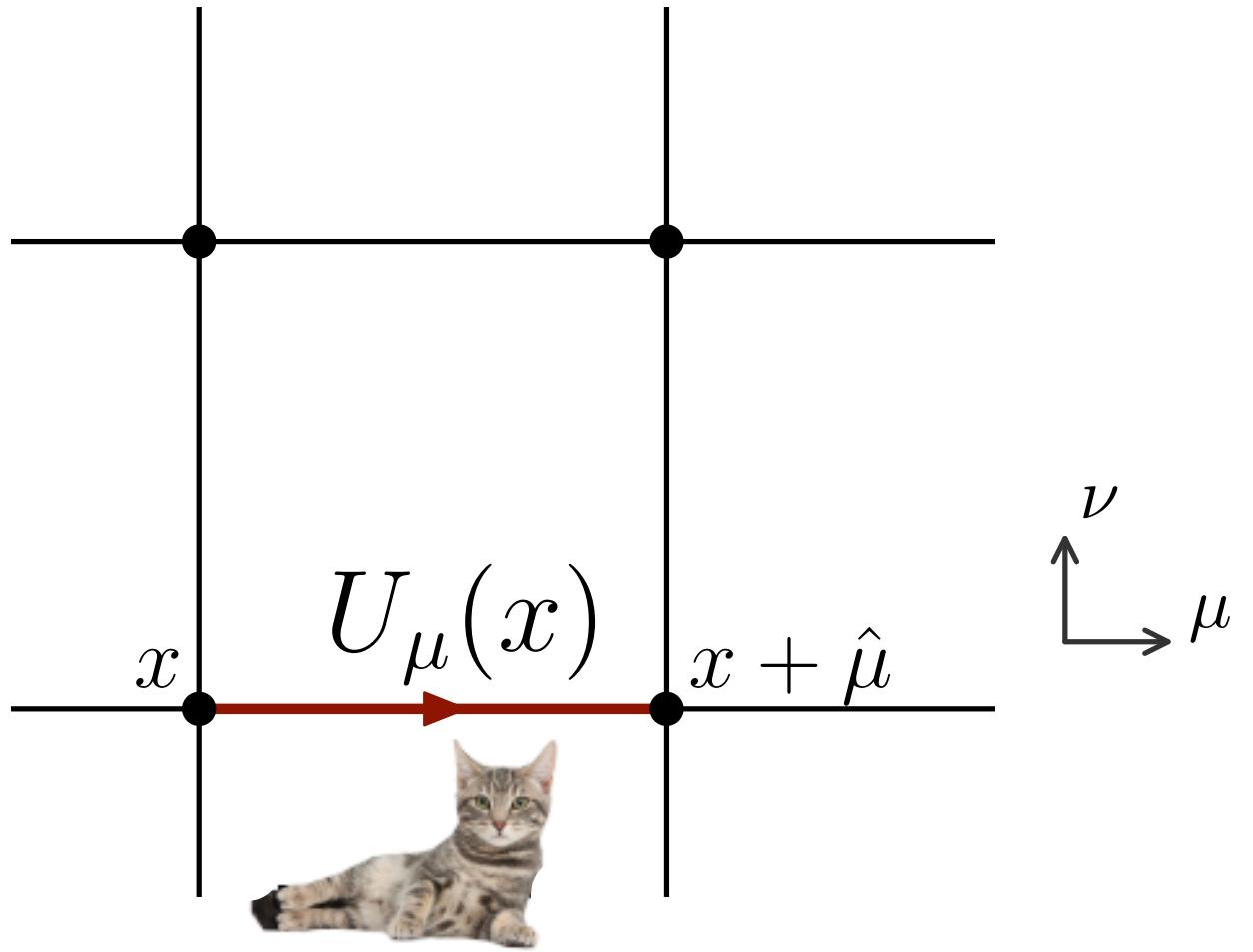
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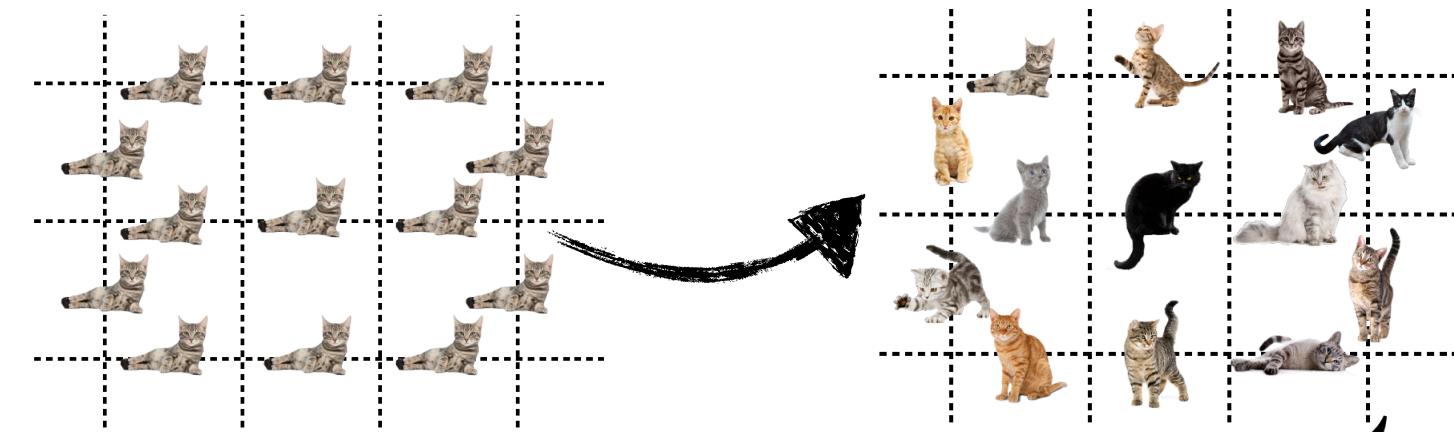
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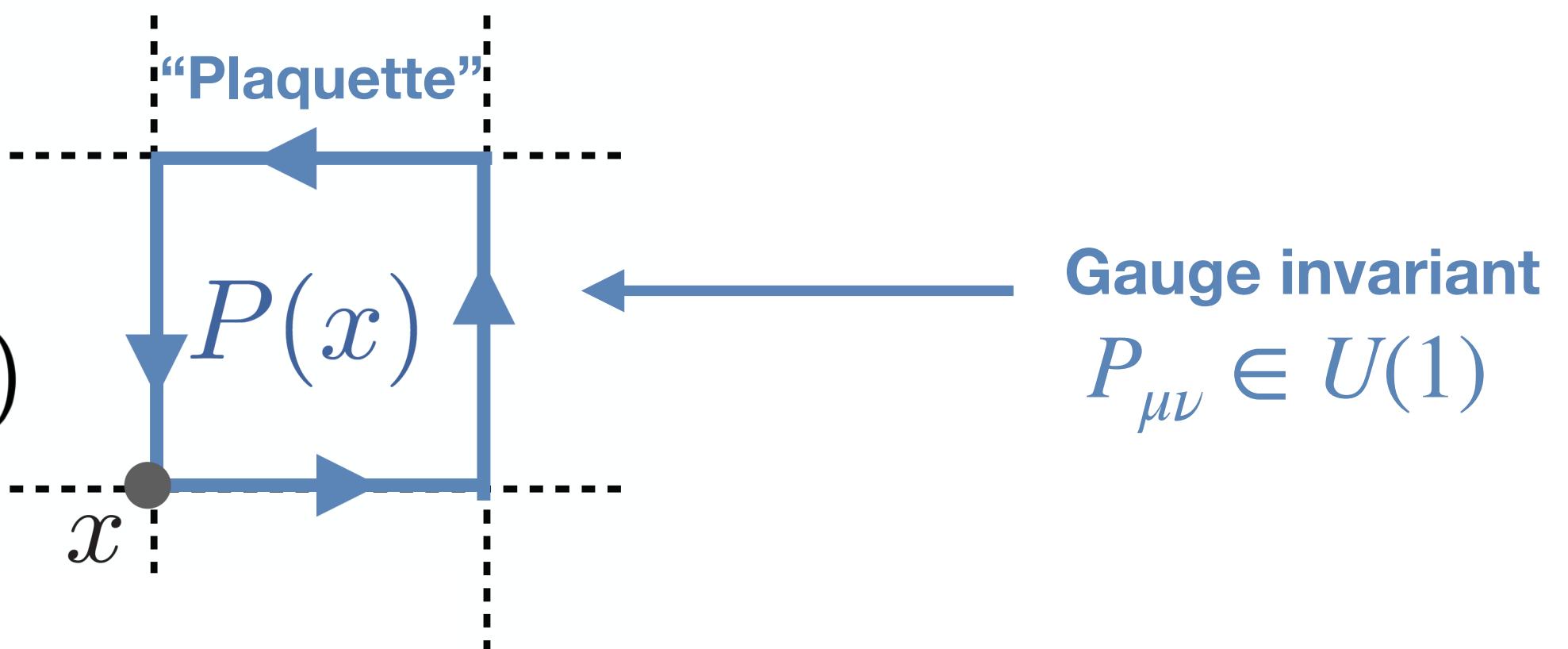
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- Pure gauge action:

$$S_E(U) = -\beta \sum_x \operatorname{Re} P_{\mu\nu}(x)$$

inverse
gauge coupling

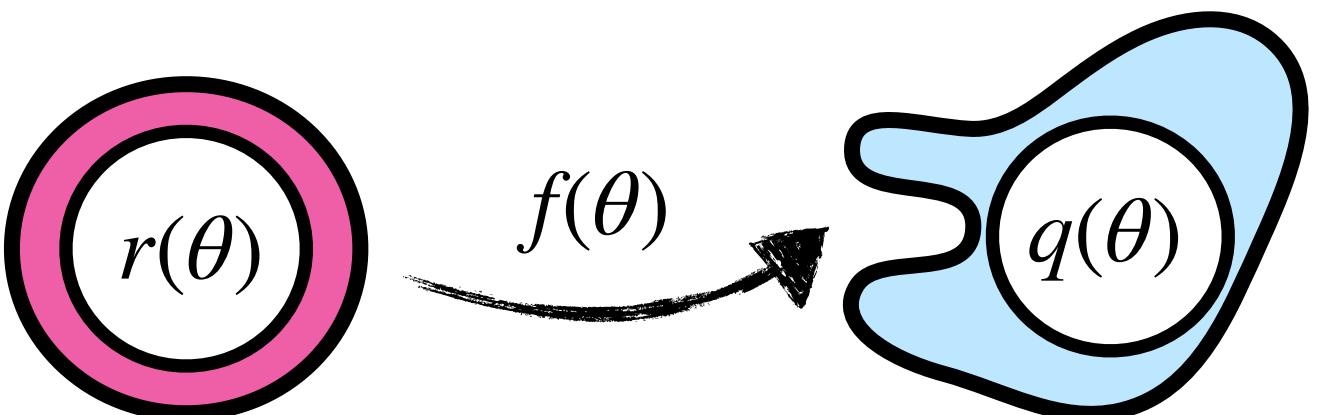


Flows on compact variables

- Flows on compact connected manifolds:

$$U_\mu(x) = \exp(i\theta) \in \mathrm{U}(1)$$

[\[Rezende, Papamakarios, Racanière , Albergo, Kanwar, Shanahan, Cranmer 2002.02428\]](#)

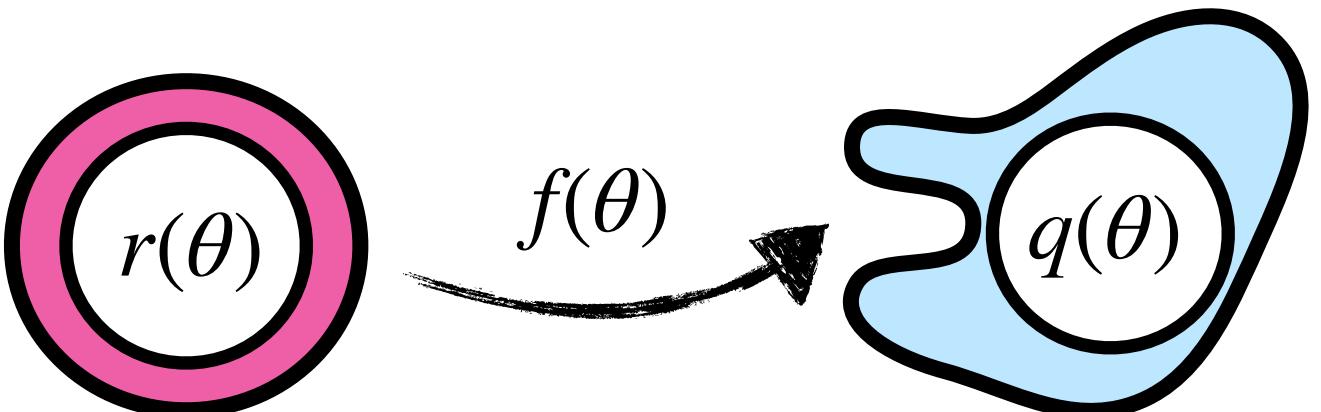


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Diffeomorphism if:

$$f(0) = 0,$$

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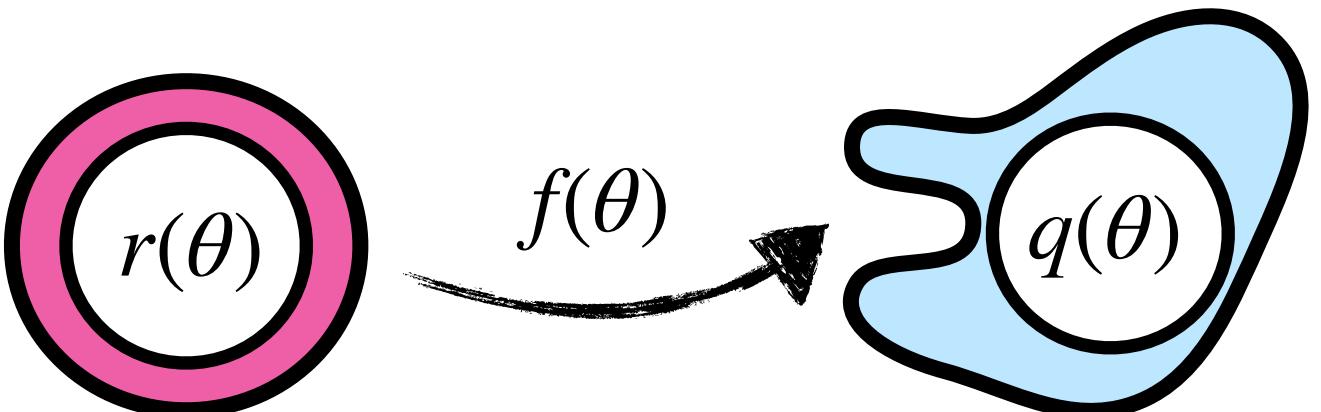
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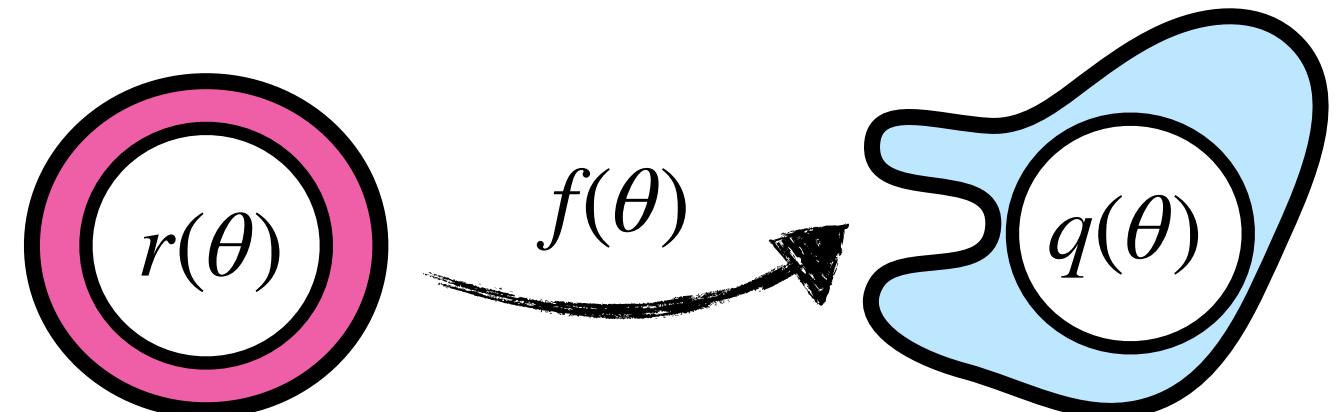
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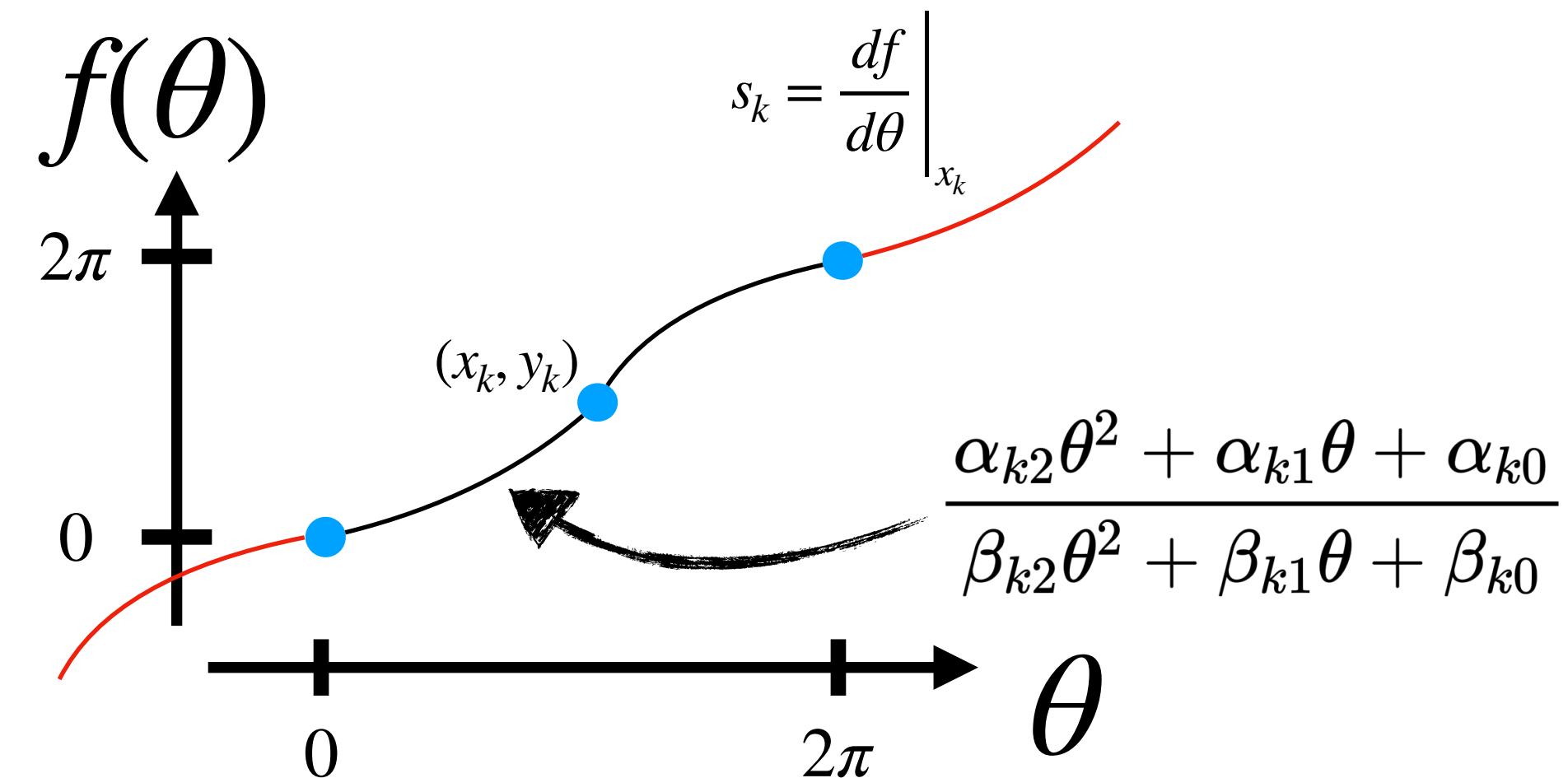
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Circular splines:

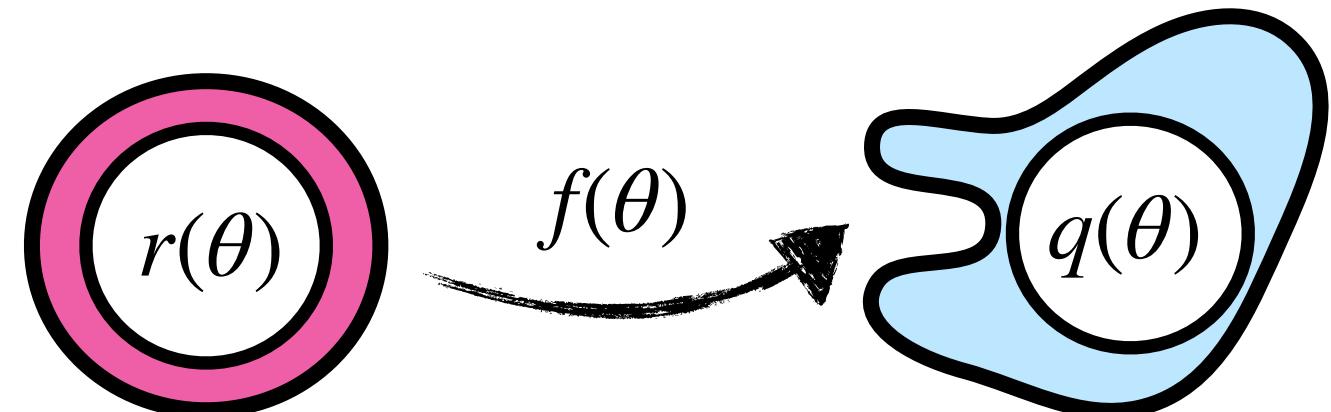


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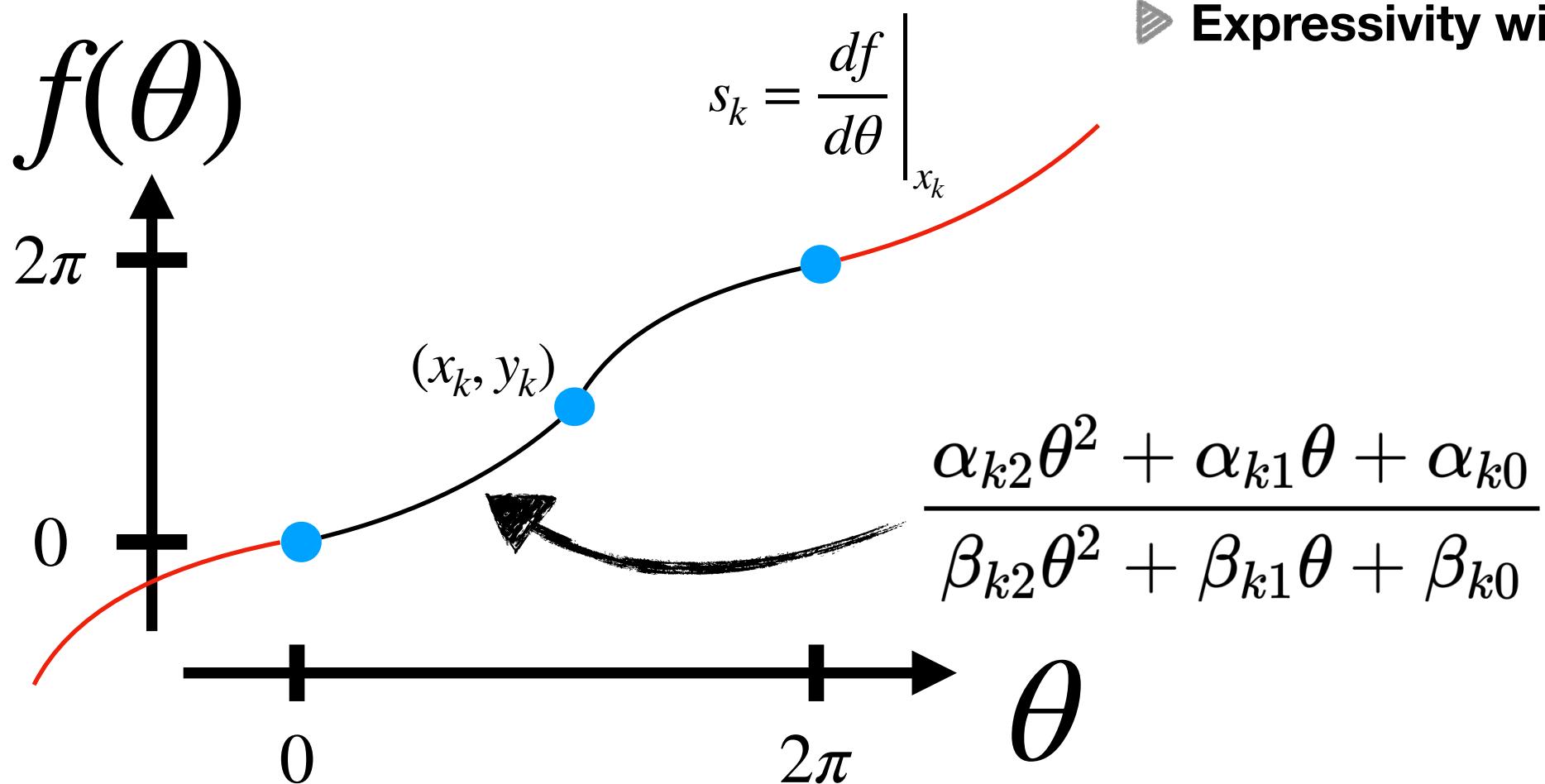


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monotonic,
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Circular splines:



- Trainable positions and slopes
- Expressivity with more knots

Flows with gauge variables

- Gauge invariant prior + gauge-equivariant transformations:

$$\theta \sim \text{Uniform}(0, 2\pi)$$

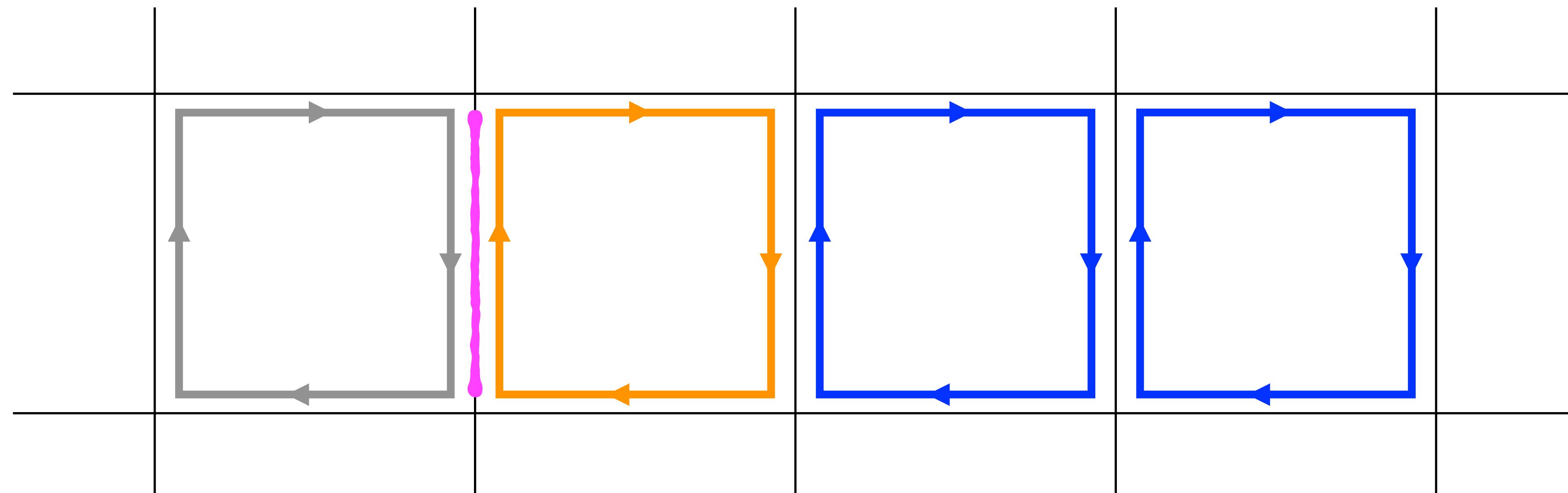
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Link to be updated

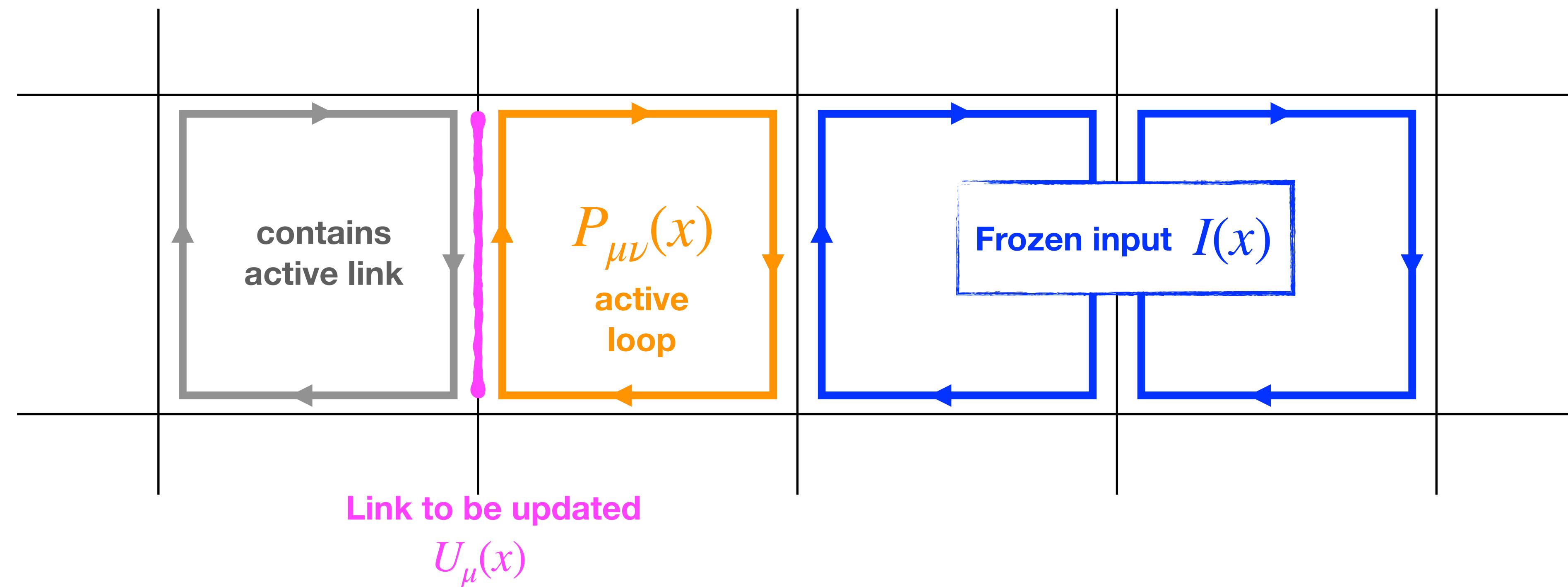
$U_\mu(x)$

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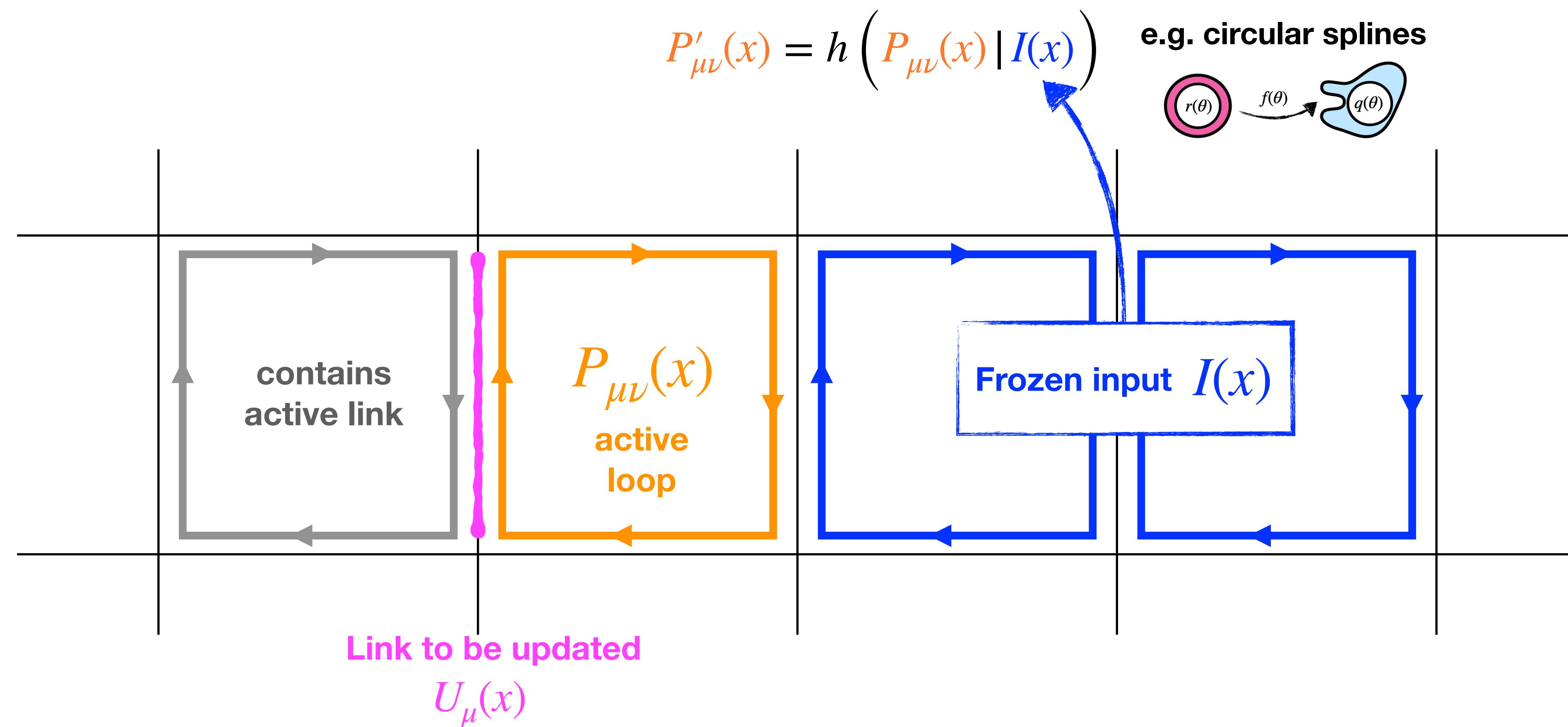


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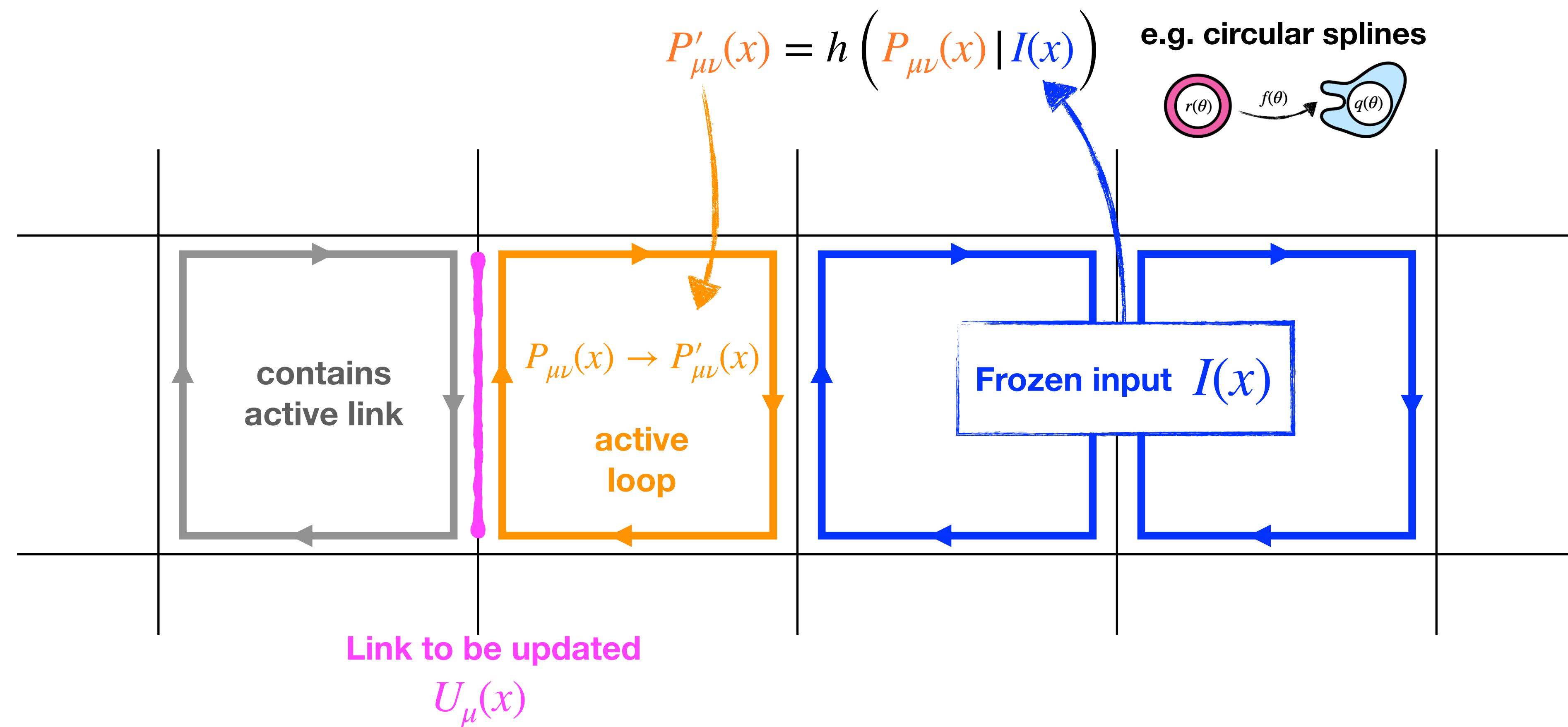


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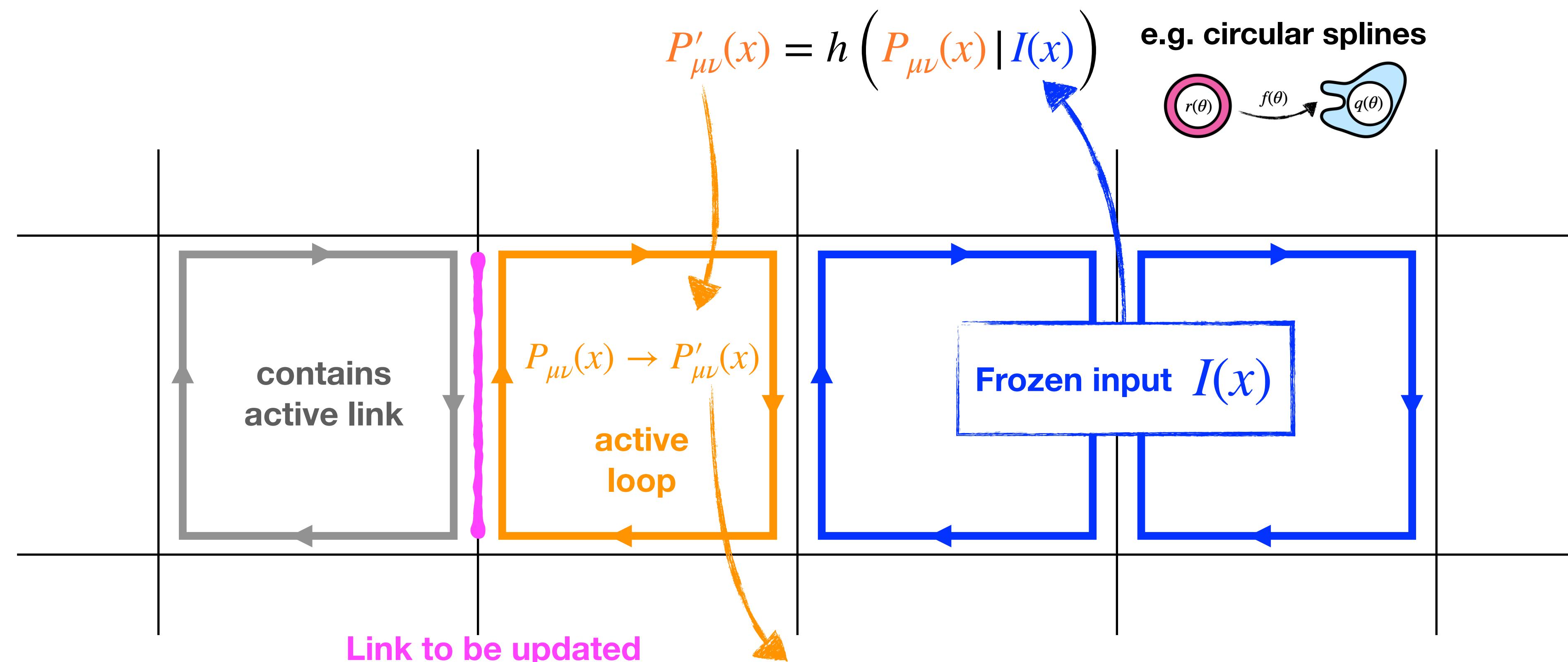


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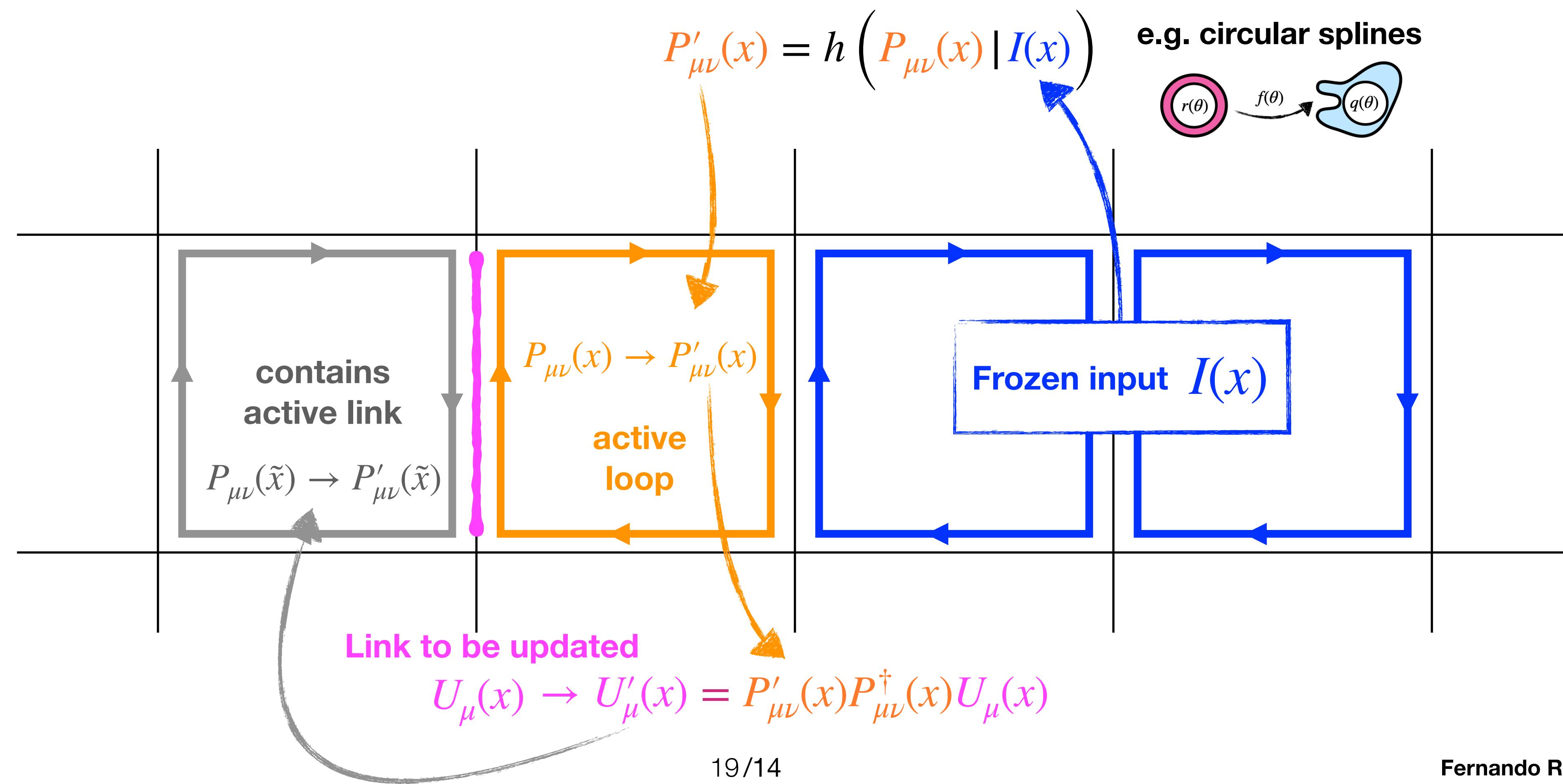
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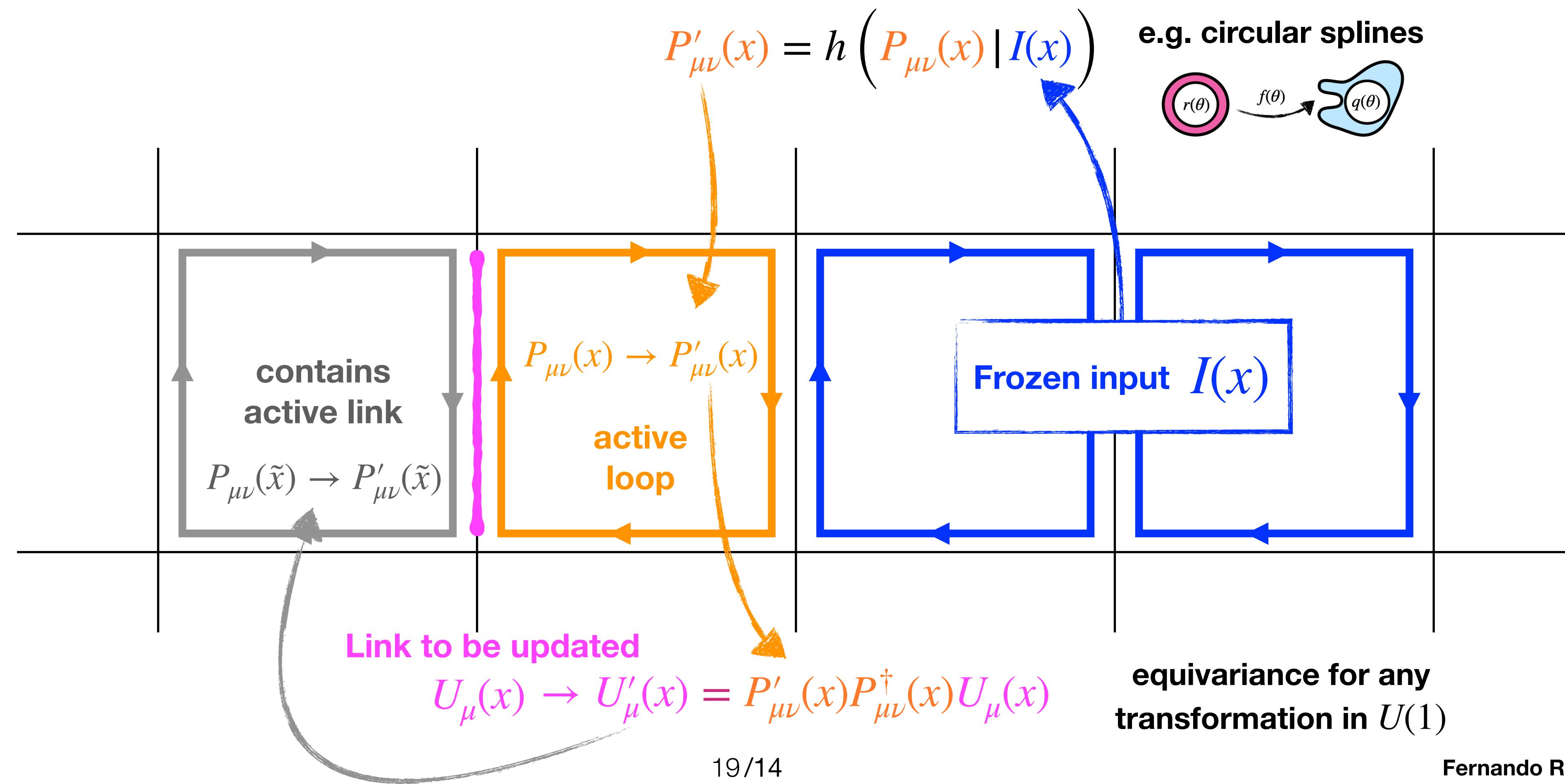


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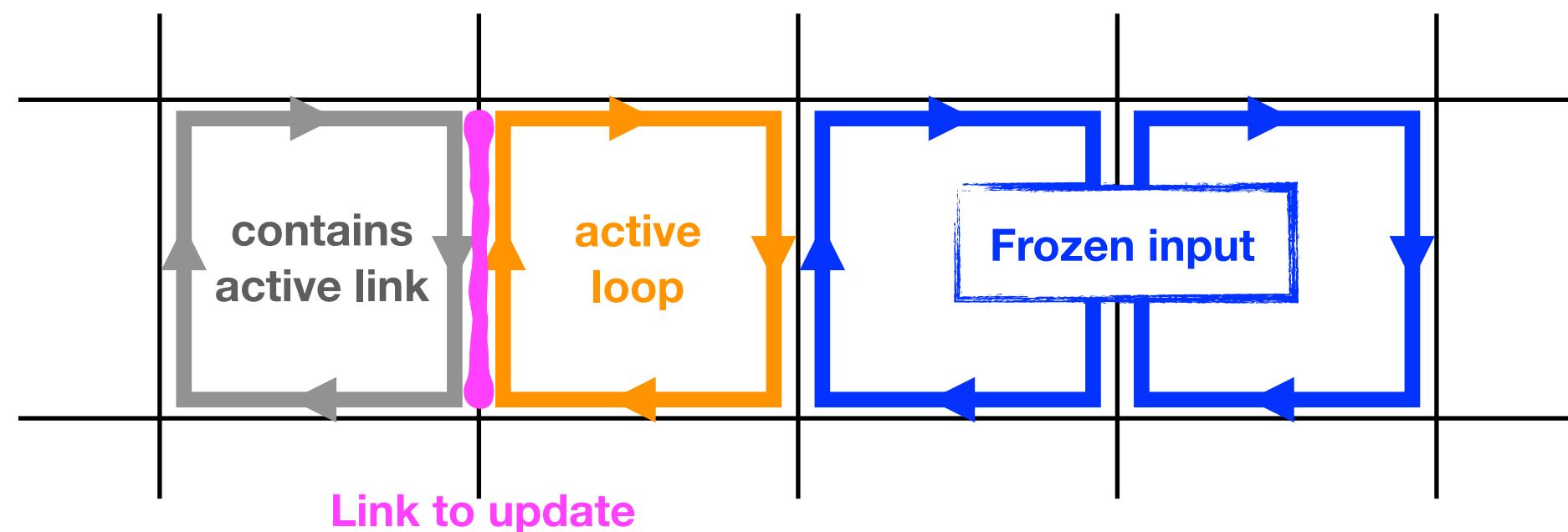
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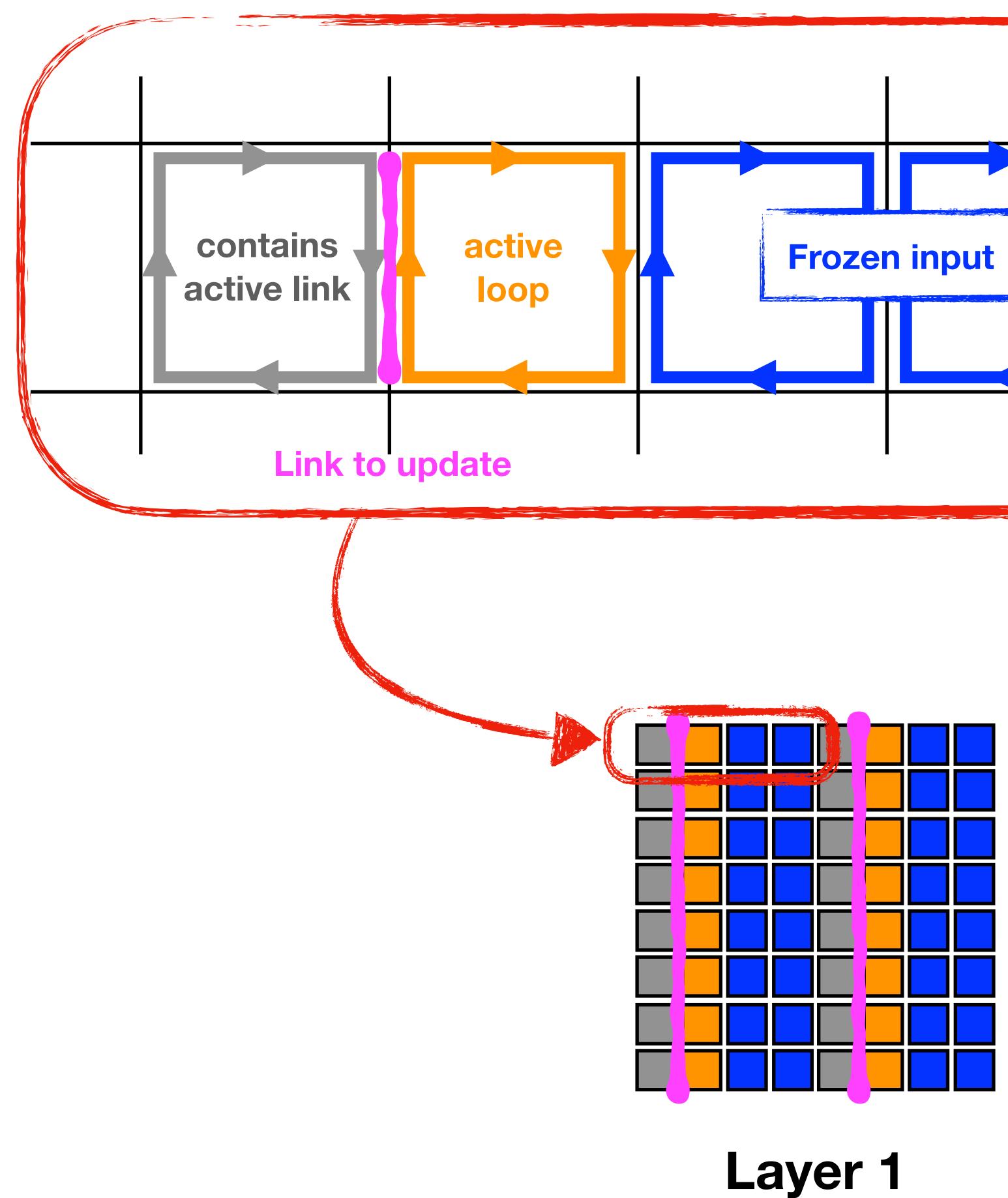
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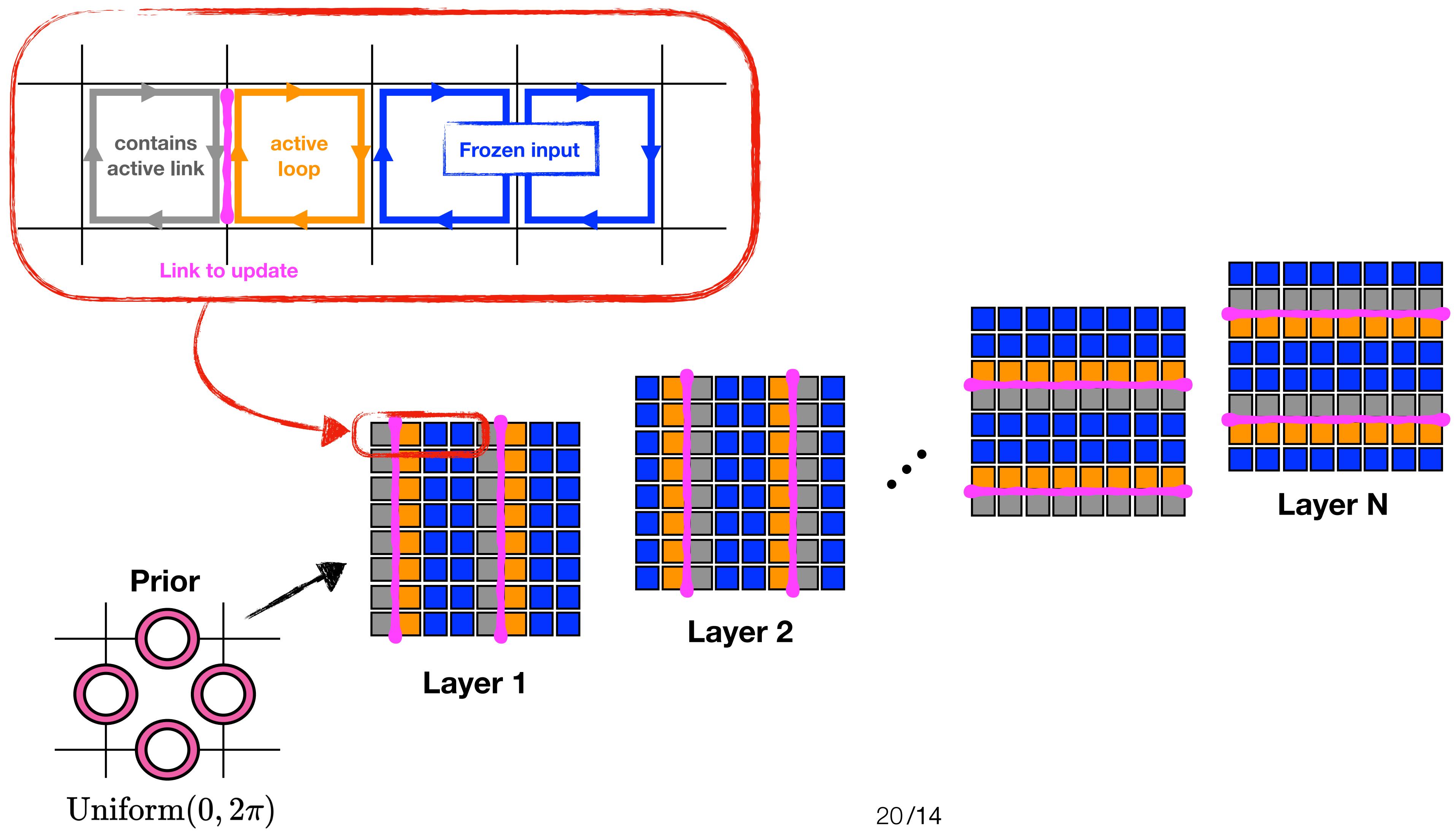
Gauge-equivariant models



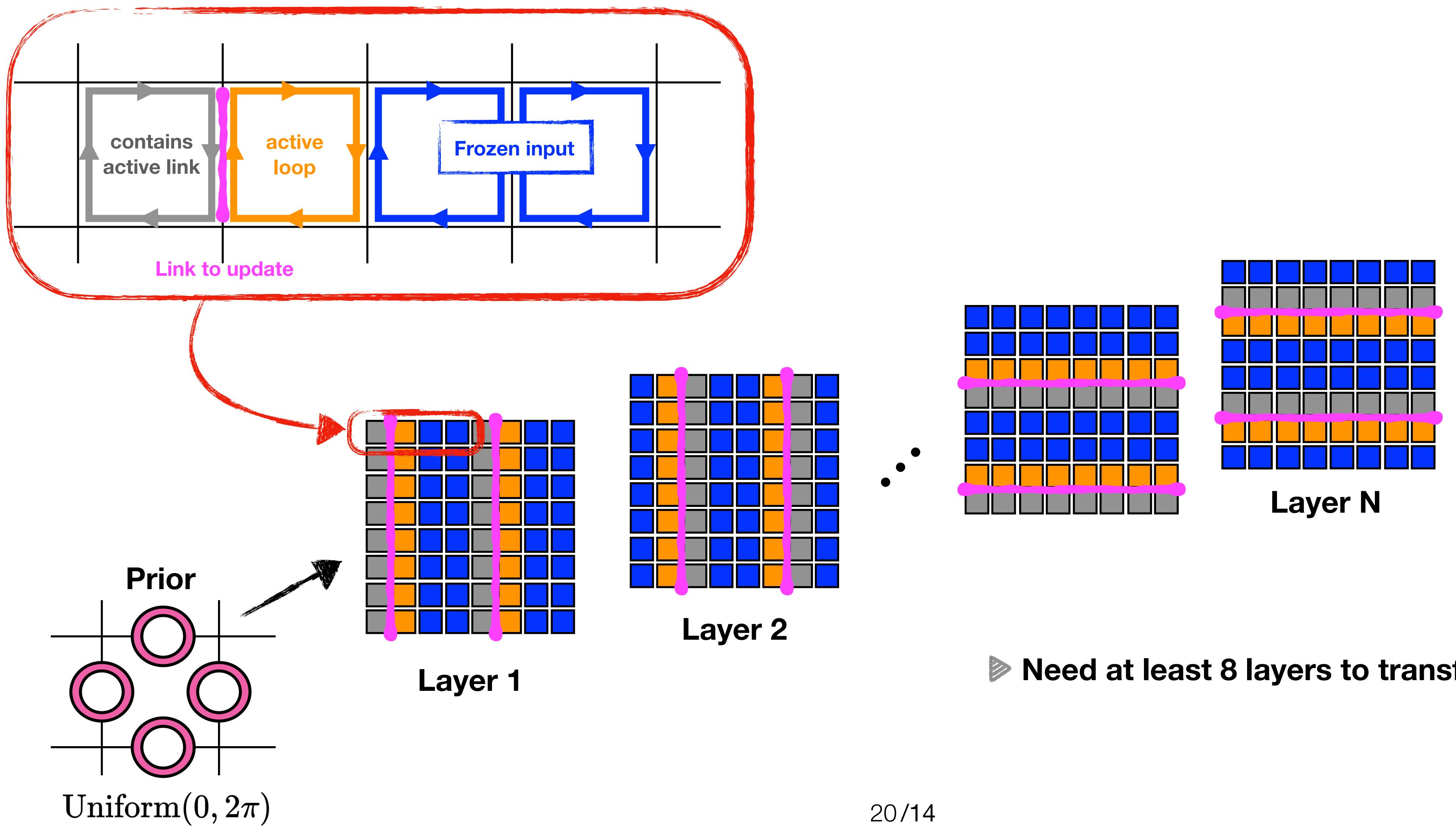
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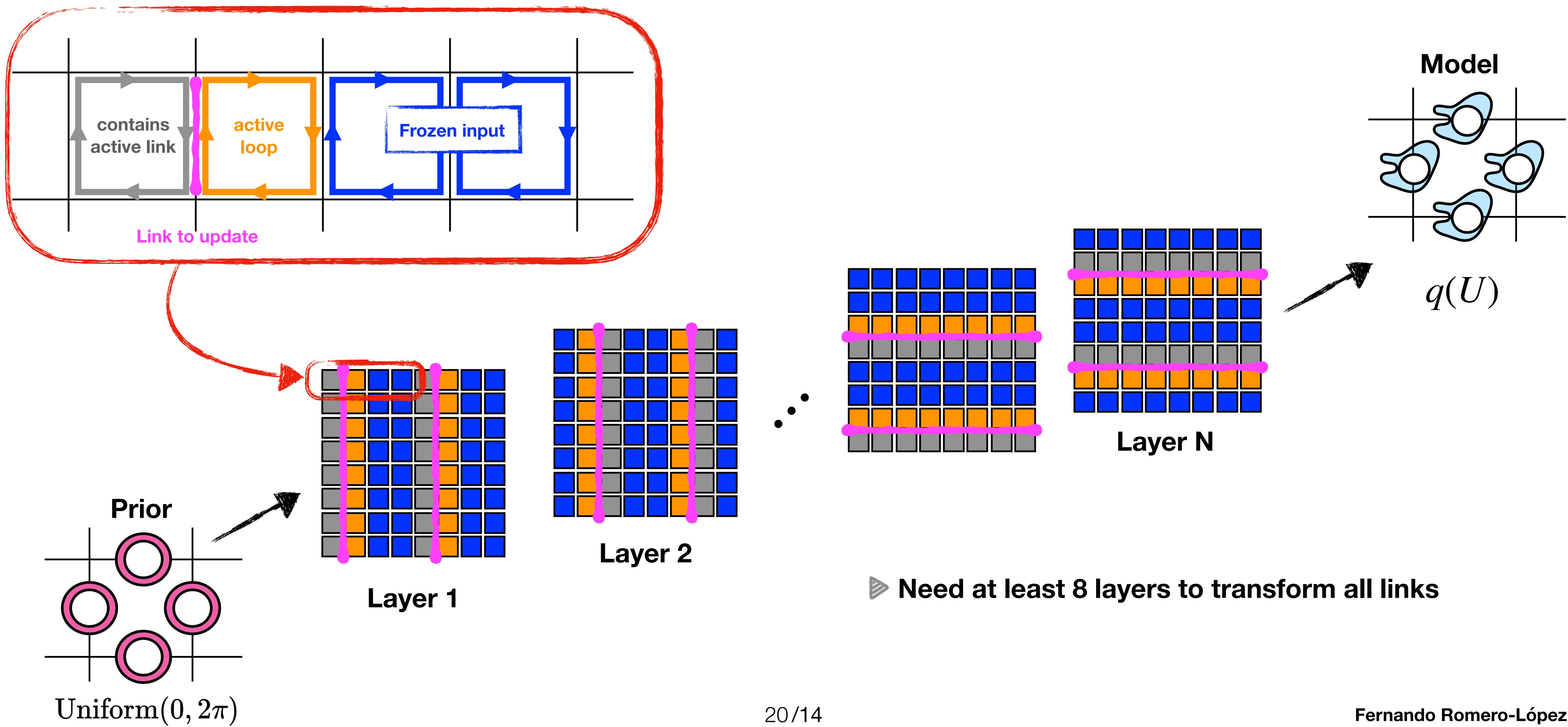
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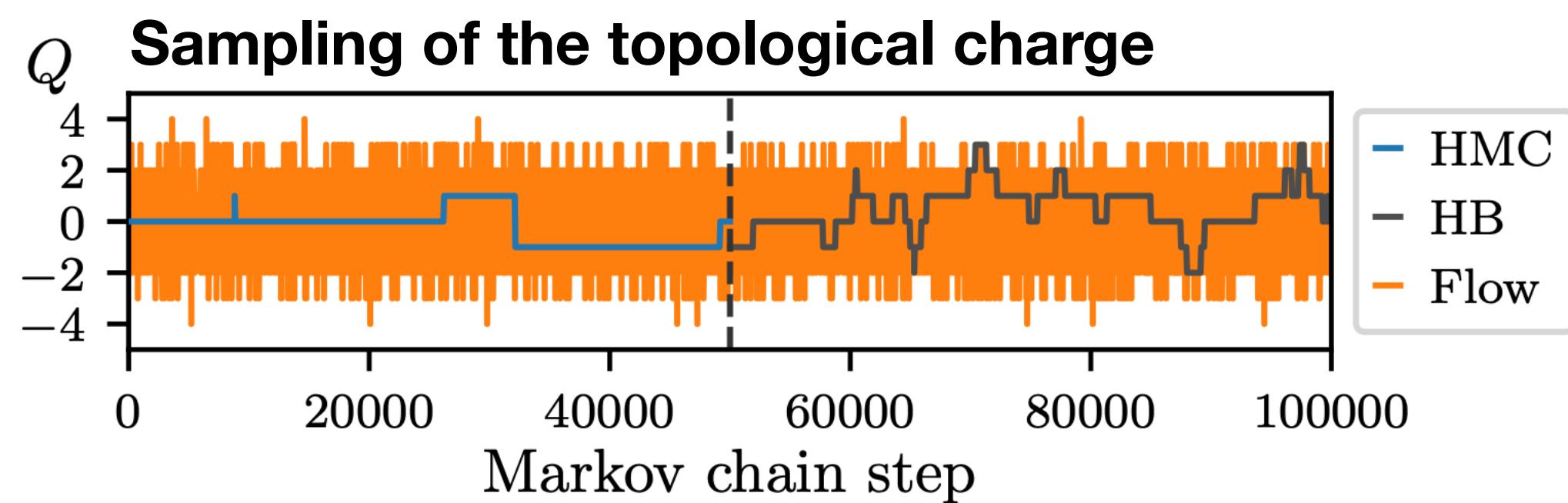
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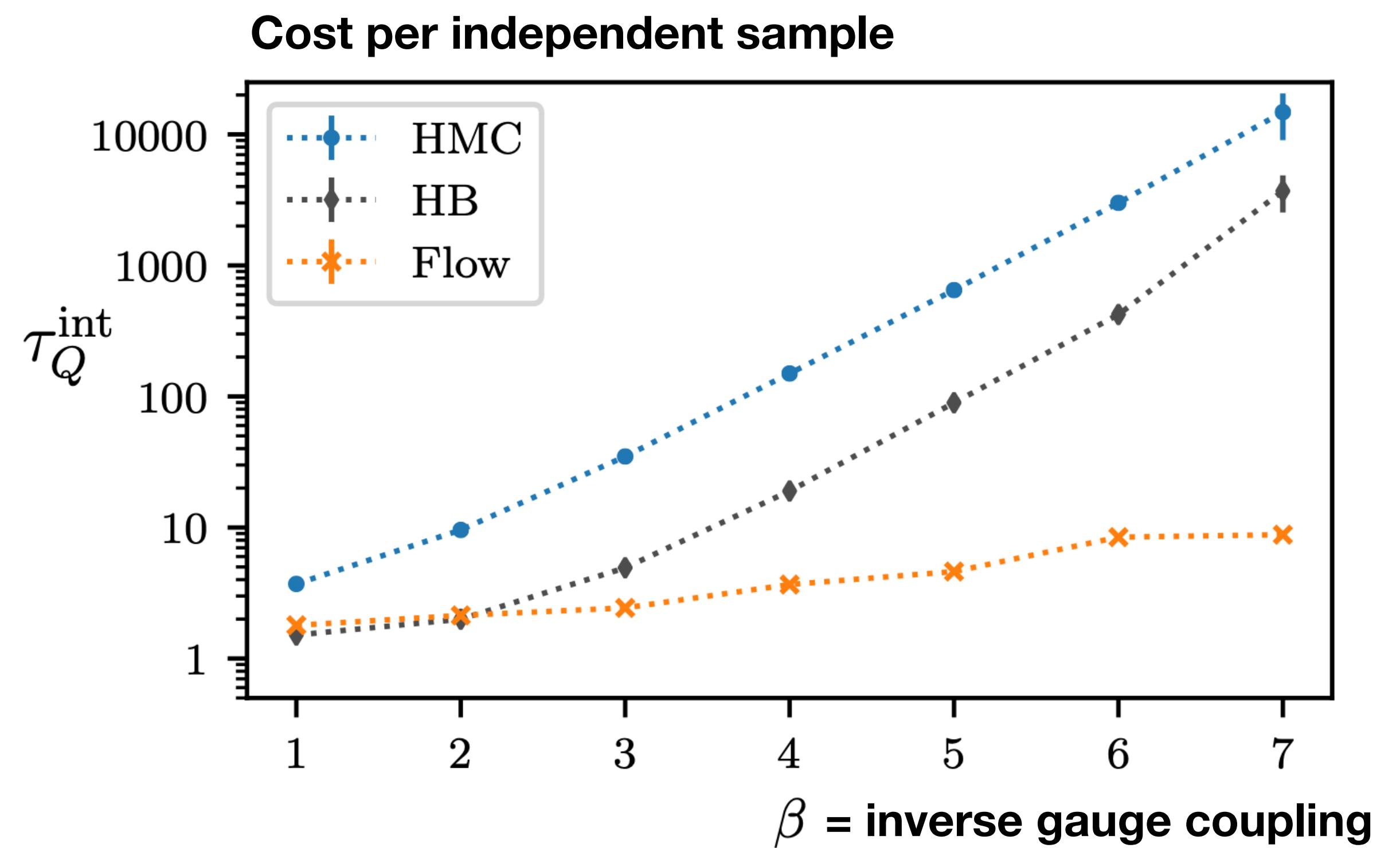
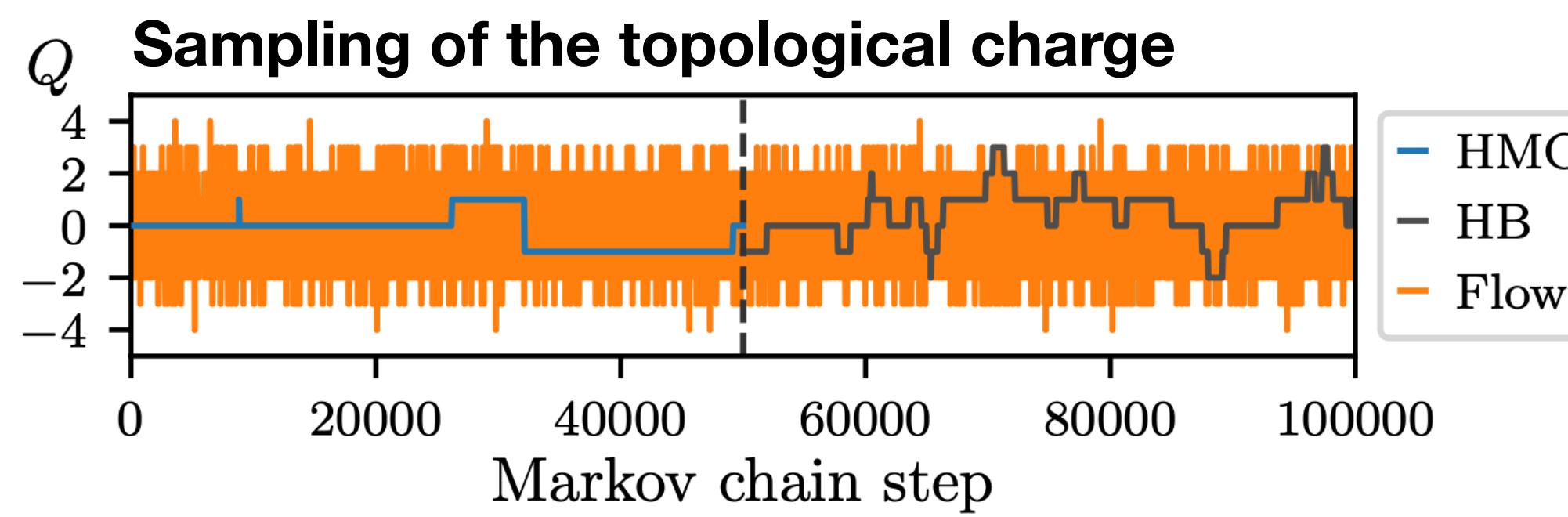
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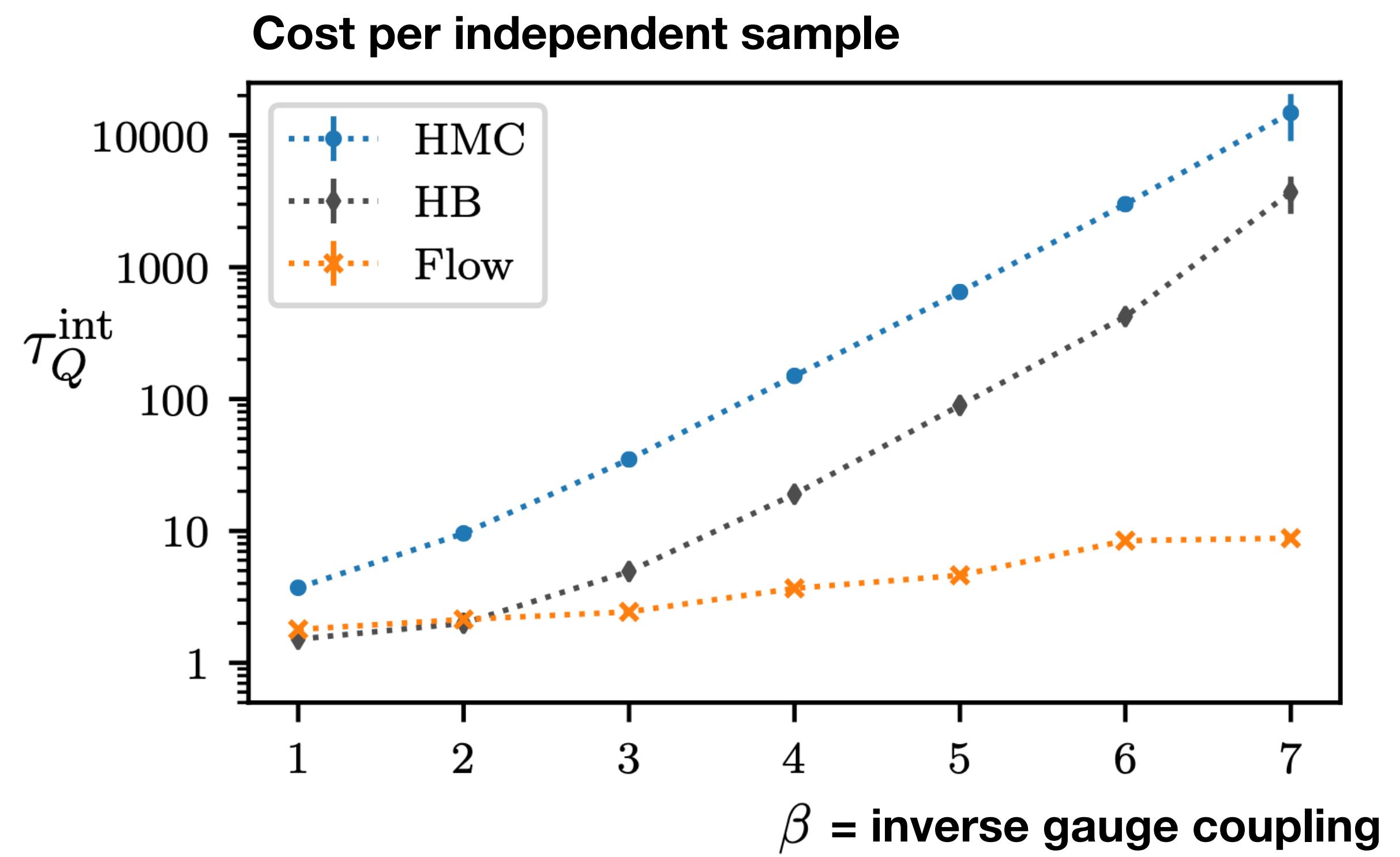
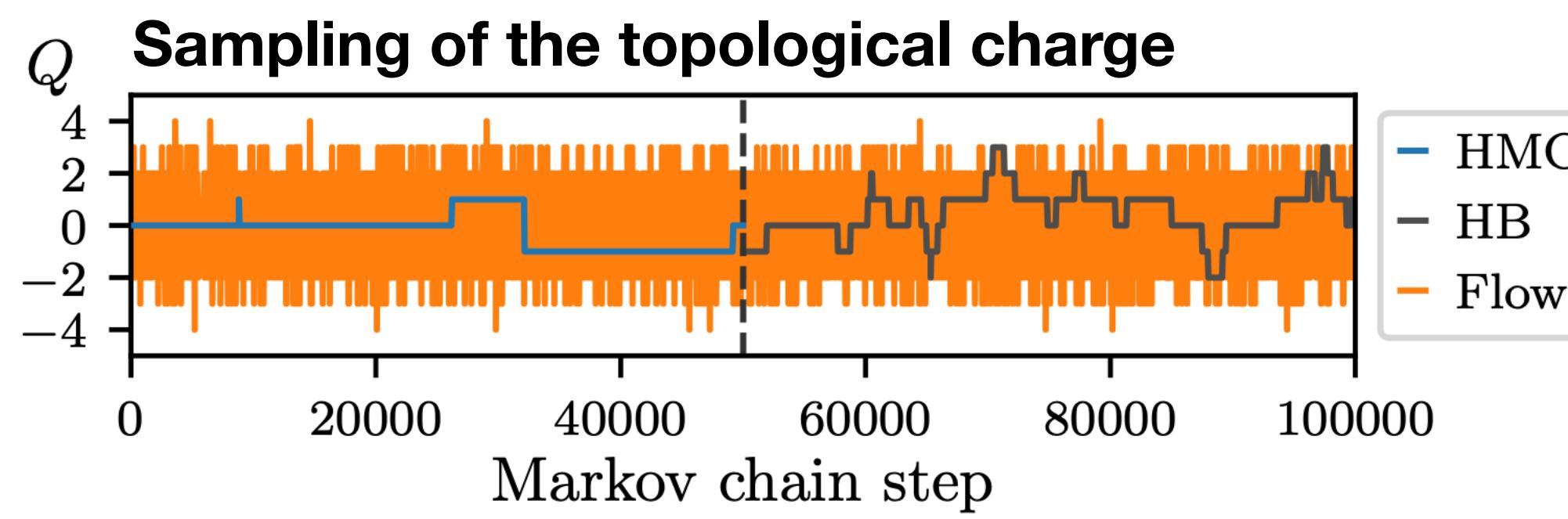
Results for pure U(1) theory



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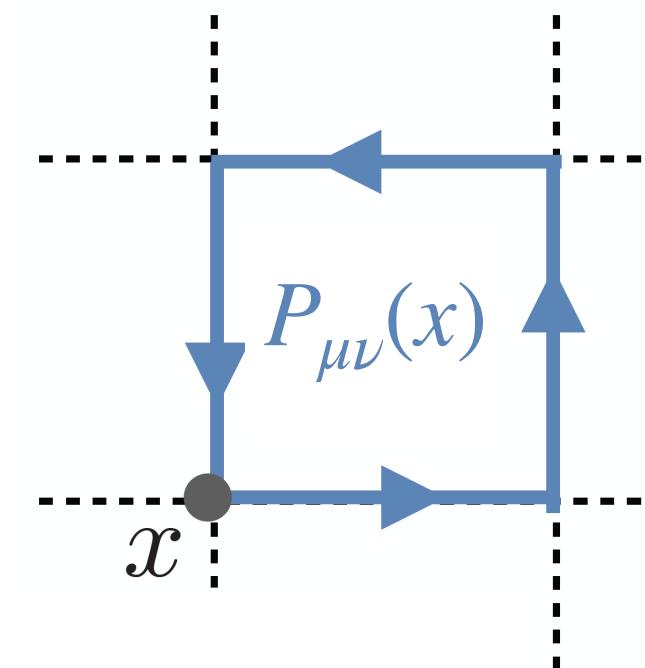


✓ Flow-based sampling has no topological freezing

[Kanwar, Albergo, Boyda, Cranmer, Hackett, Racanière, Rezende, Shanahan 2003.06413]

Flows on SU(3)

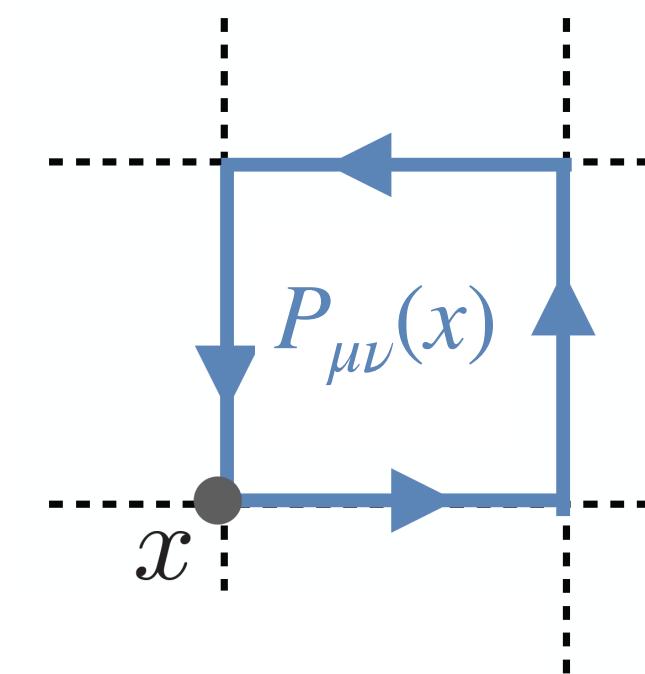
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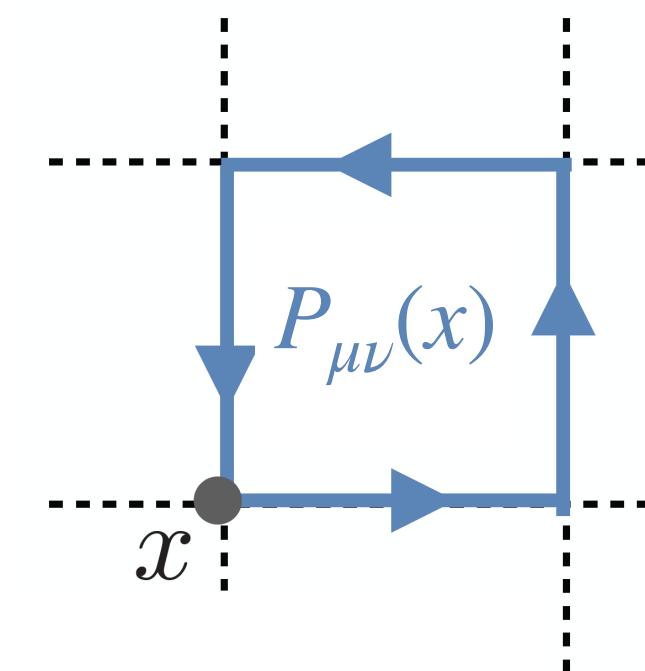


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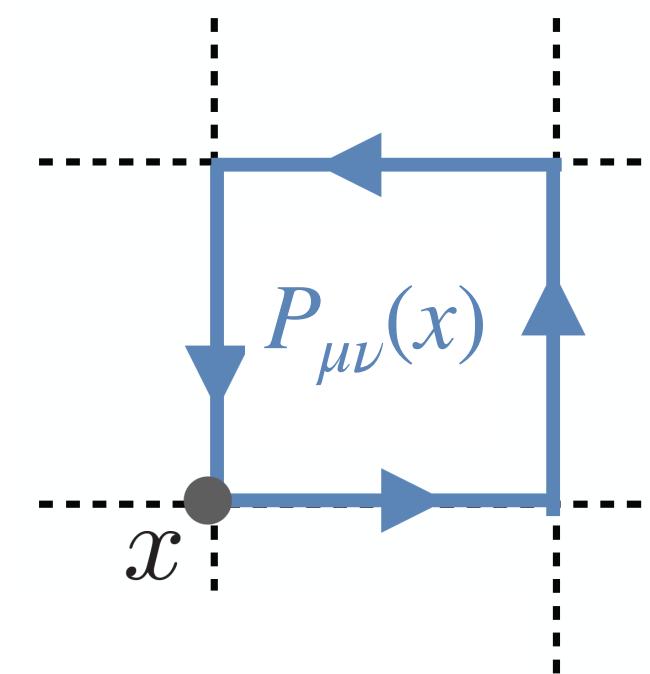
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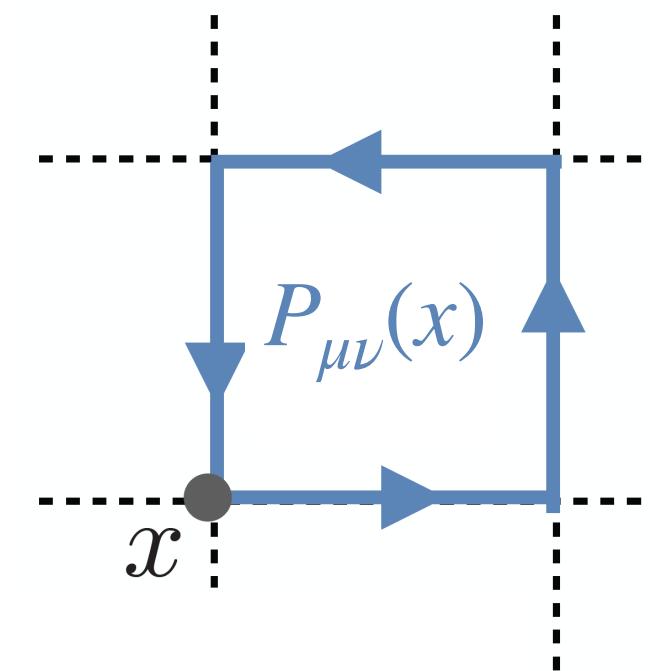
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- Place eigenvalues in canonical ordering
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- Undo canonical ordering

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- ✓ Successfully applied to SU(3) in 2D

[Kanwar, Albergo, Boyda, Cranmer, Hackett, Racanière, Rezende, Shanahan 2003.06413]

The road to at-scale QCD

(1+1)d real scalar field theory

[\[Albergo, Kanwar, Shanahan 1904.12072\]](#)

[\[Hackett, Hsieh, Albergo, Boyda, JW Chen, KF Chen, Cranmer, Kanwar, Shanahan 2107.00734\]](#)

(1+1)d Abelian gauge theory

[\[Kanwar, Albergo, Boyda, Cranmer, Hackett, Racanière, Rezende, Shanahan 2003.06413\]](#)

(1+1)d non-Abelian gauge theory

[\[Boyda, Kanwar, Racanière, Rezende, Albergo, Cranmer, Hackett, Shanahan 2008.05456\]](#)

(1+1)d Yukawa model

i.e. real scalar field theory + fermions

[\[Albergo, Kanwar, Racanière, Rezende, Urban, Boyda, Cranmer, Hackett, Shanahan 2106.05934\]](#)

Schwinger model

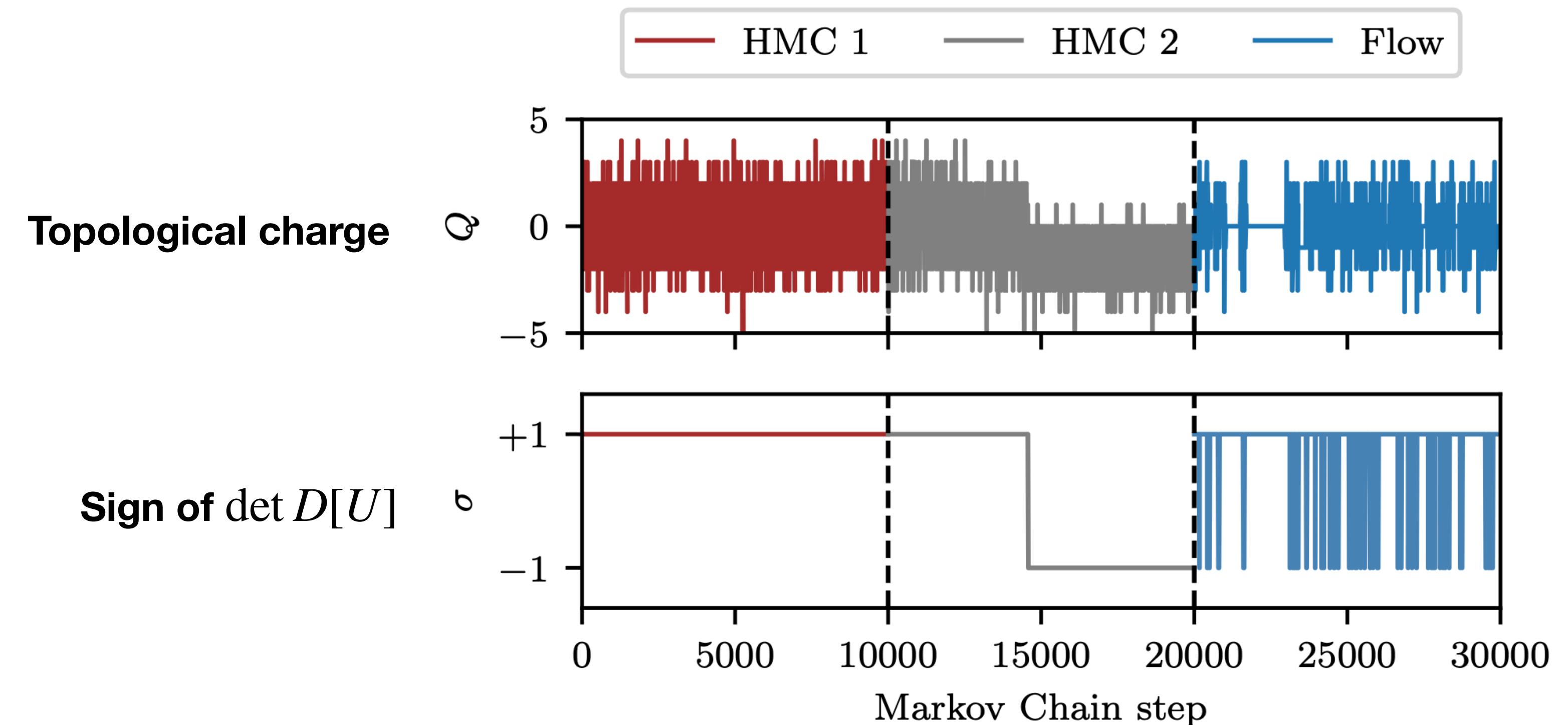
i.e. (1+1)d QED

[\[Albergo, Boyda, Cranmer, Hackett, Kanwar, Racanière, Rezende, FRL, Shanahan, Urban 2202.11712\]](#)

Schwinger and 2D QCD with Pseudofermions

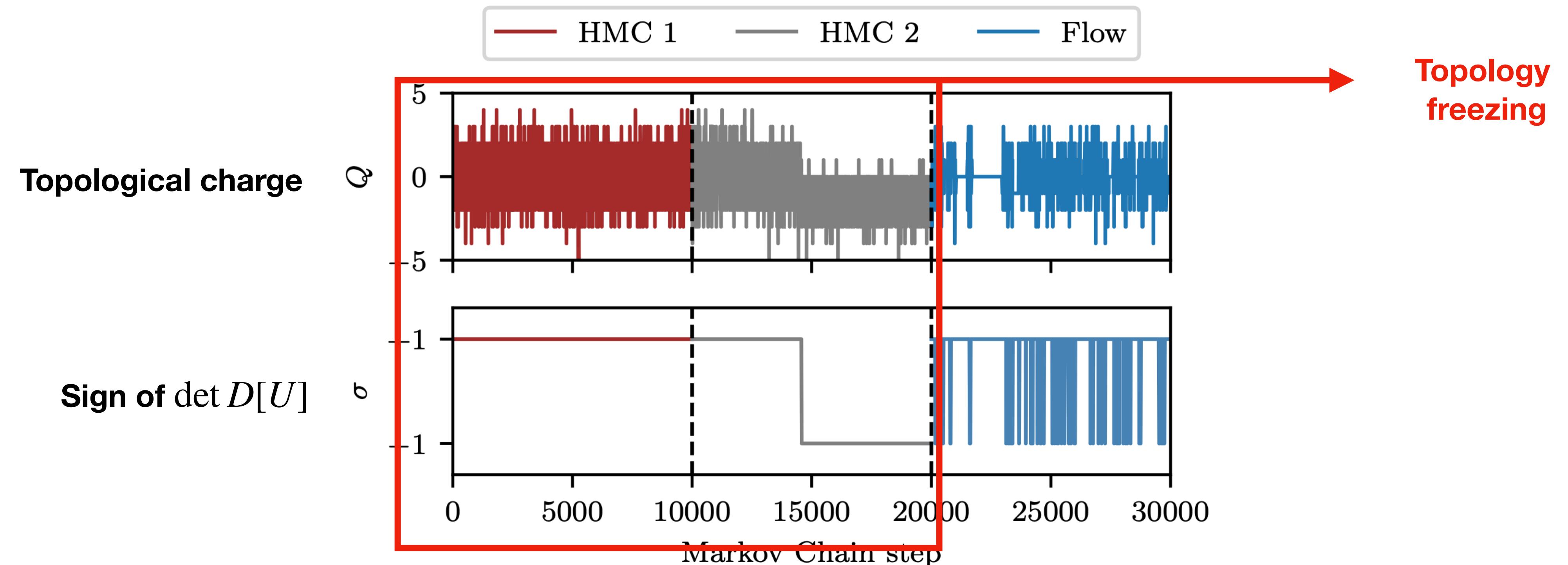
[\[Abbott ,Albergo, Boyda, Cranmer, Hackett, Kanwar, Racanière, Rezende, FRL, Shanahan, Urban 2207.08945\]](#)

The Schwinger model at criticality



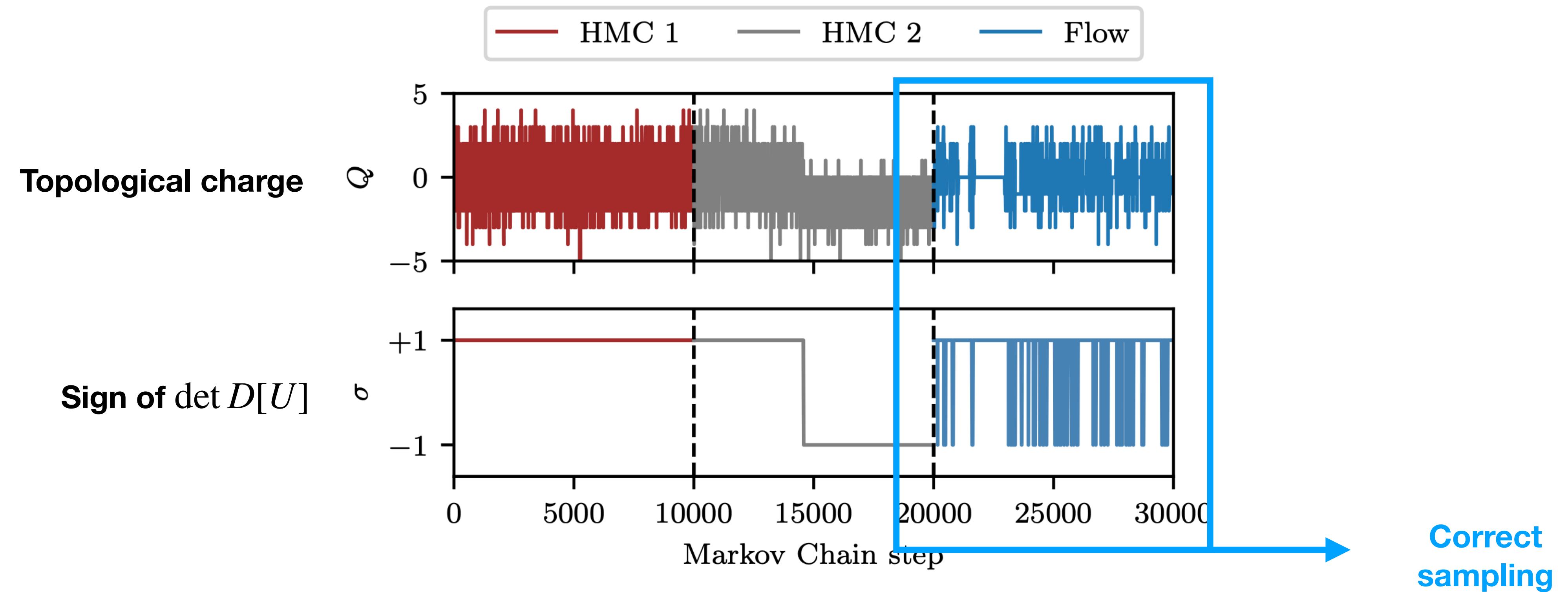
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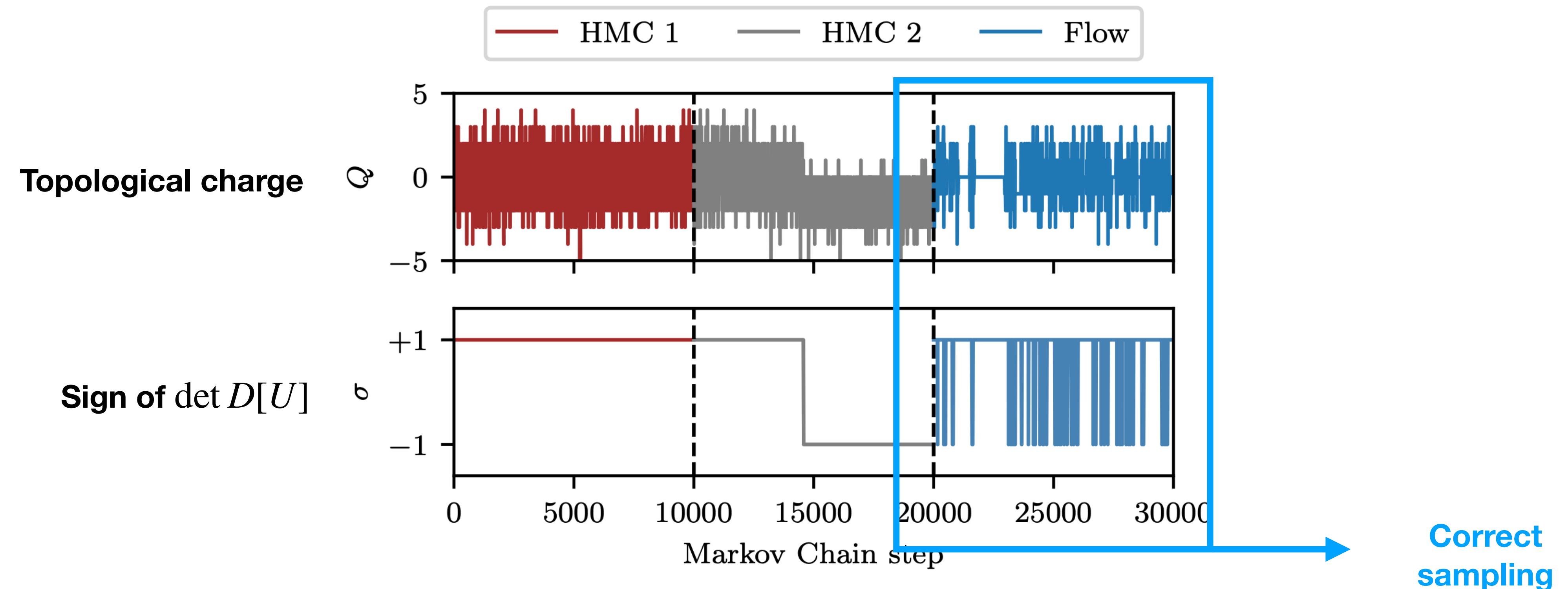
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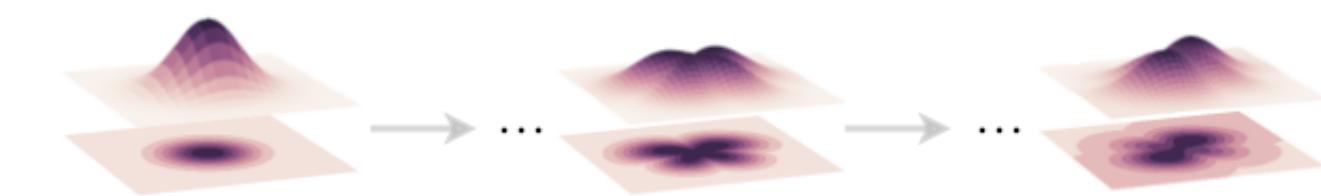
The Schwinger model at criticality



✓ Flow-based sampling mitigates topology freezing even at criticality

[Albergo, Boyda, Cranmer, Hackett, Kanwar, Racanière, Rezende, FRL, Shanahan, Urban 2202.11712]

Lattice QFT via flow models

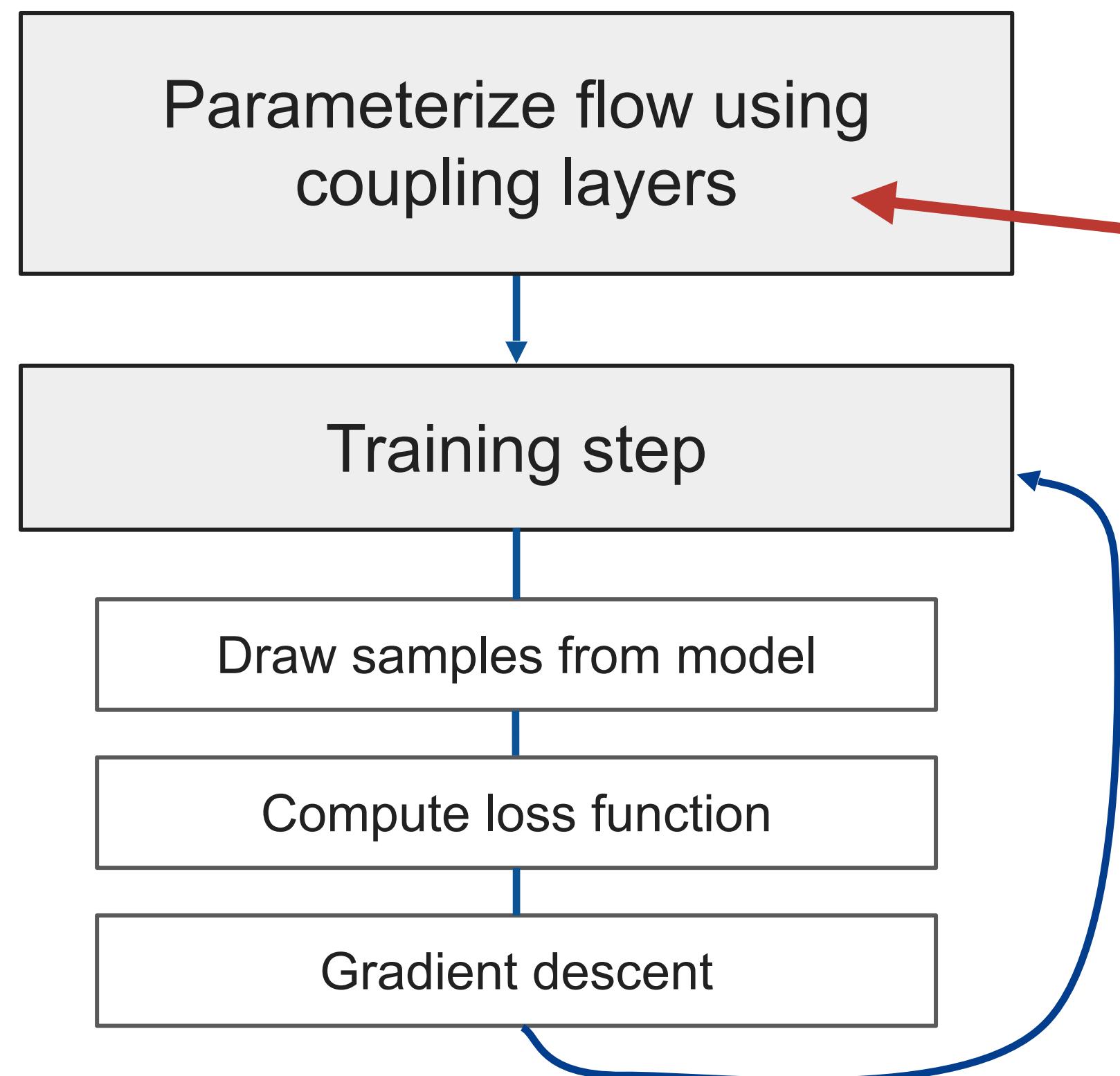
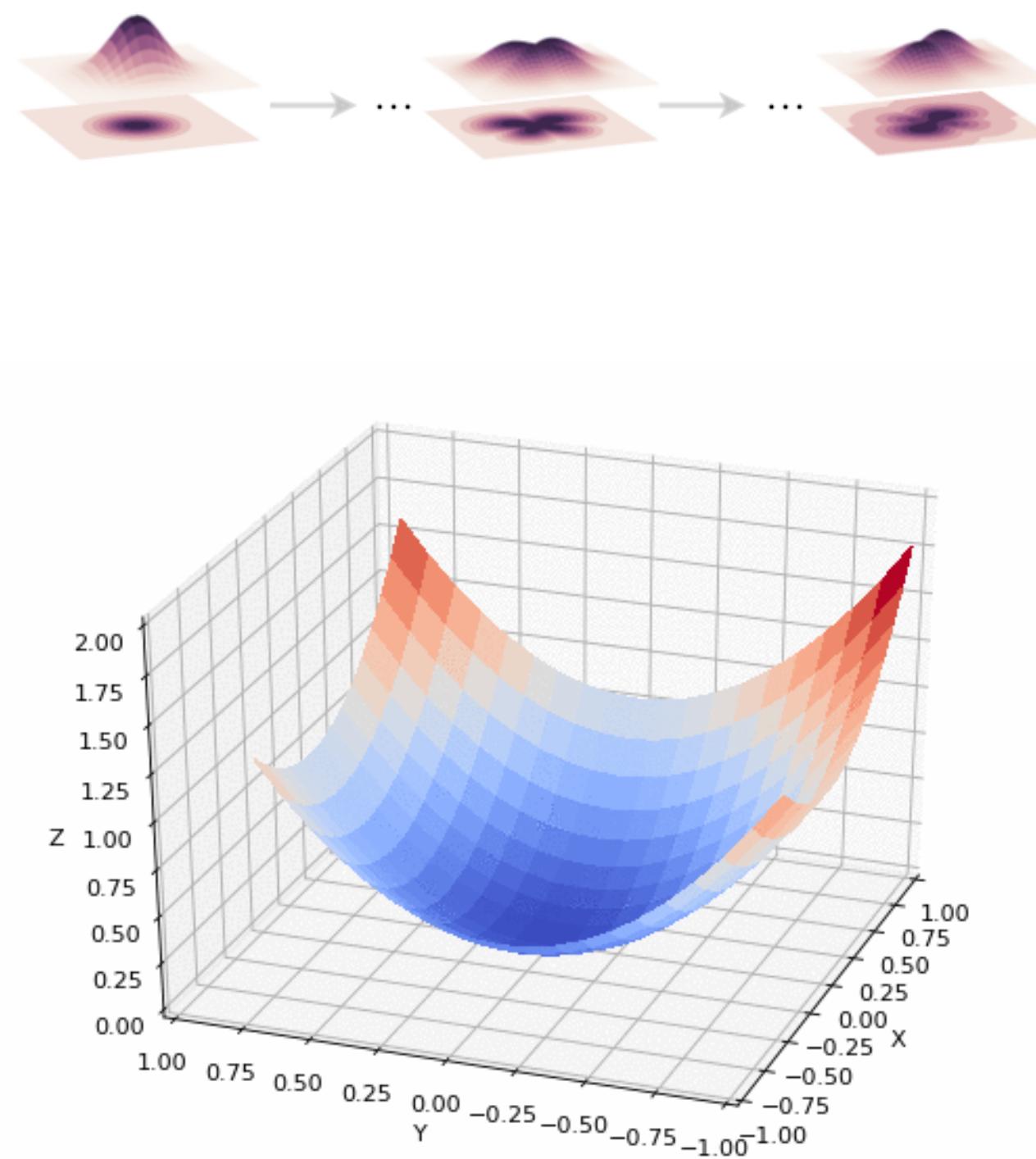


Parameterize flow using
coupling layers

Each layer contains
arbitrary neural nets

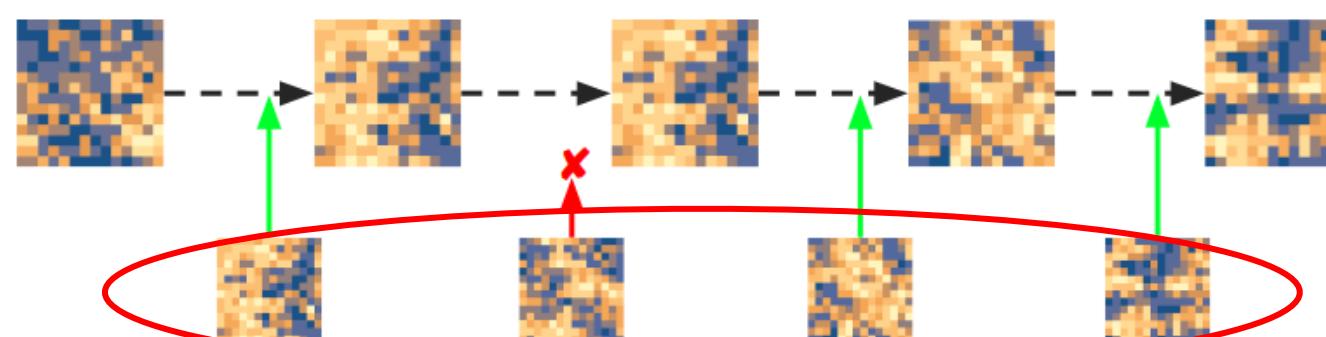
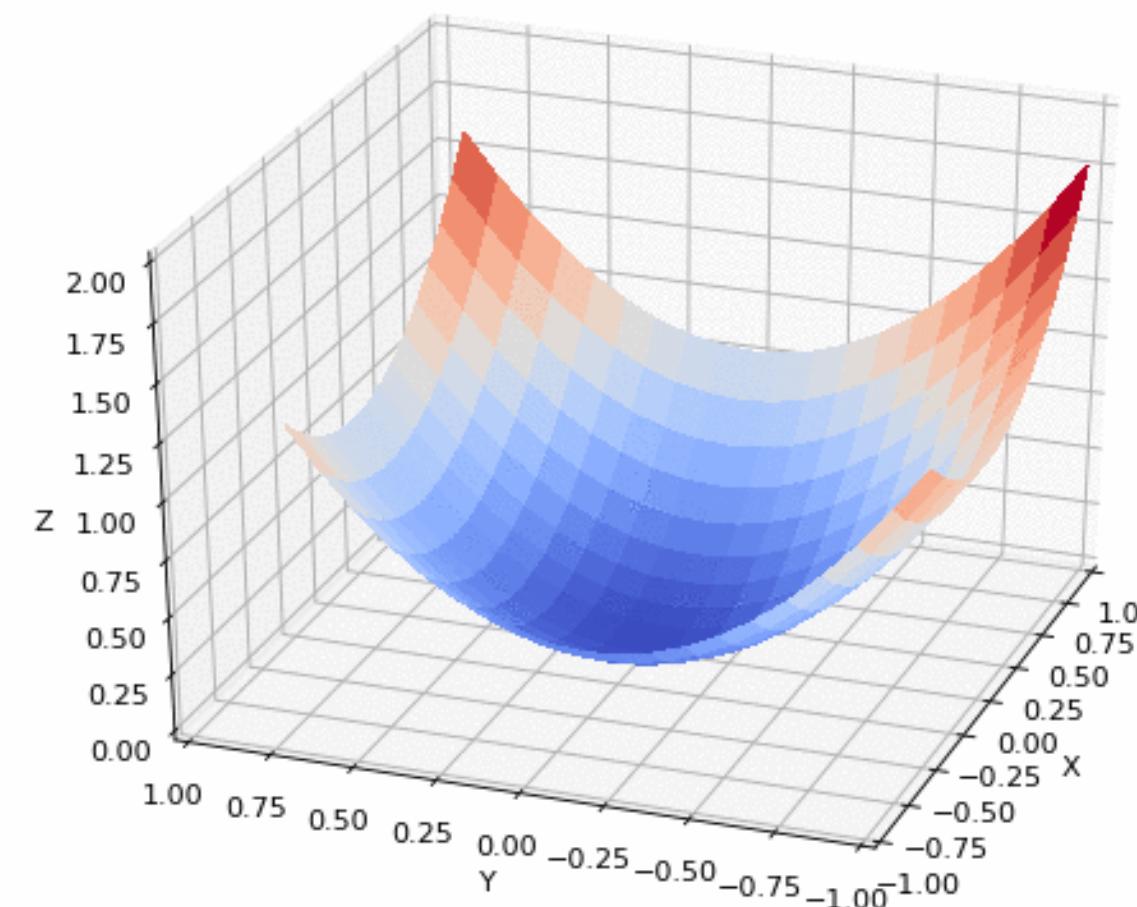
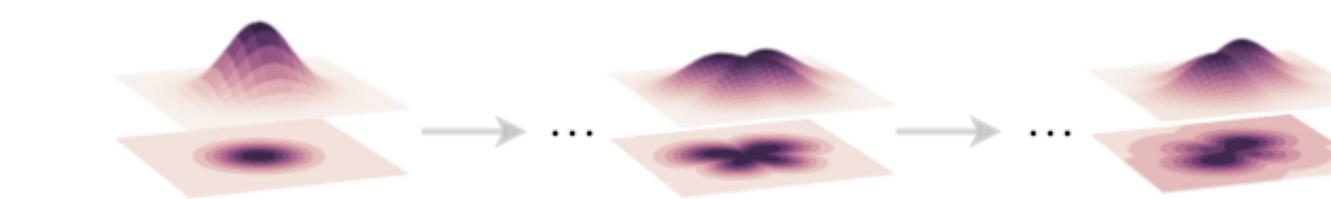
[credits: G. Kanwar]

Lattice QFT via flow models

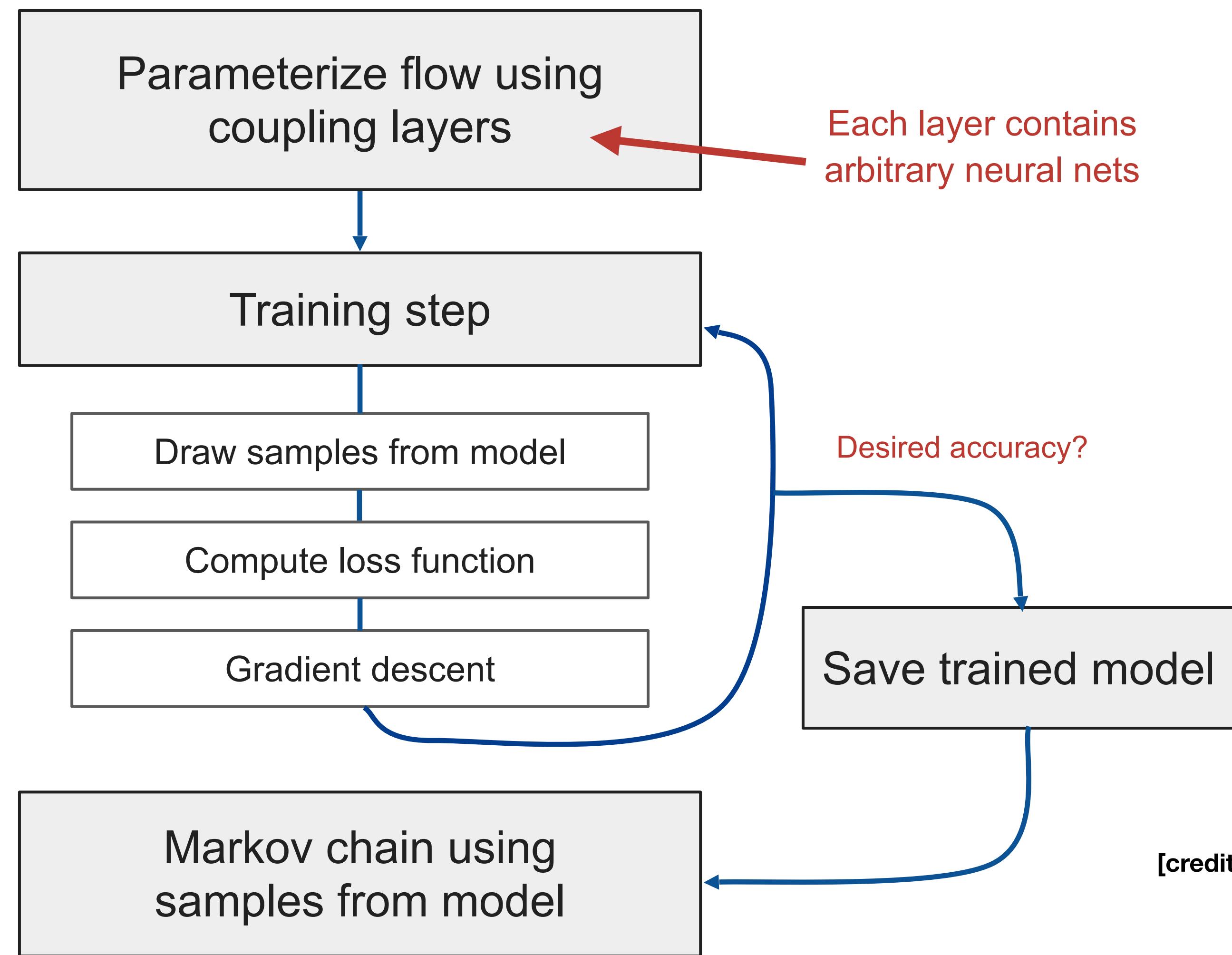


[credits: G. Kanwar]

Lattice QFT via flow models



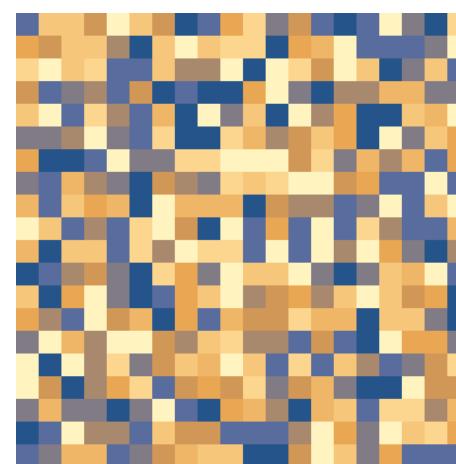
generating samples is
"embarrassingly parallel"



[credits: G. Kanwar]

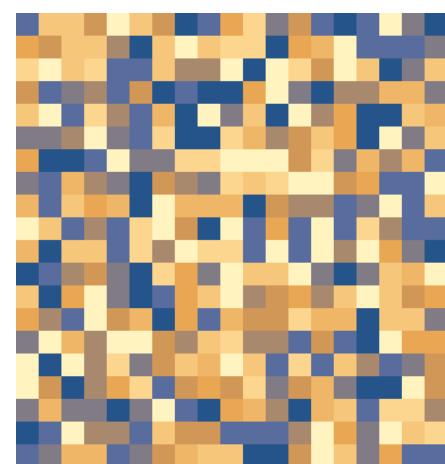
Implementing flows for scalar theory

$$\phi(x) \in (-\infty, +\infty)$$



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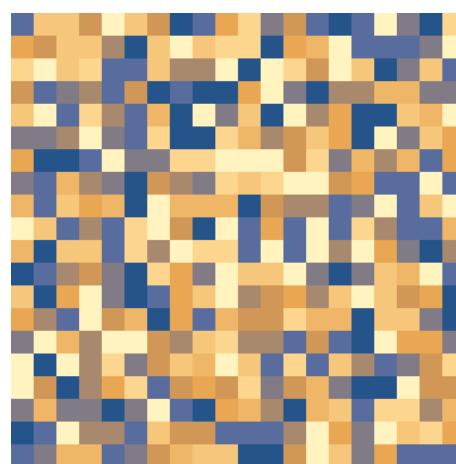
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Idea: Each layer acts on a **subset** of components, conditioned only on the complimentary subset.

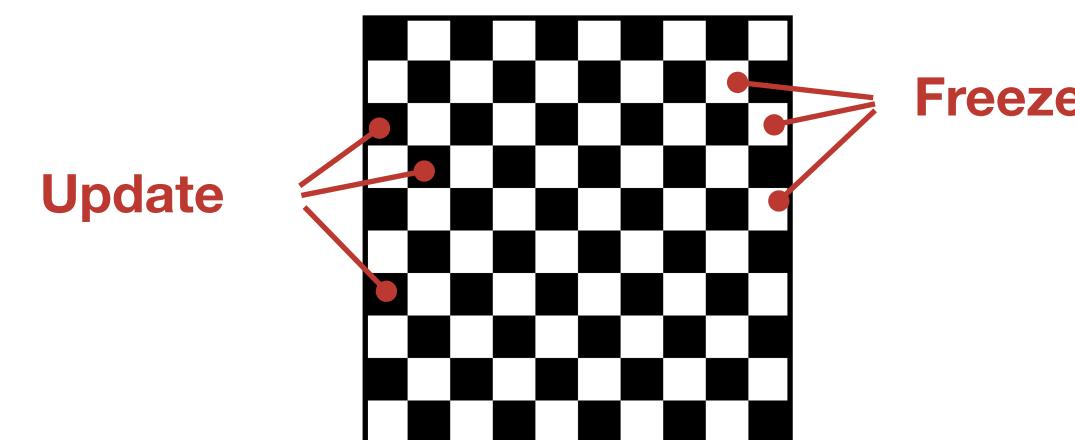
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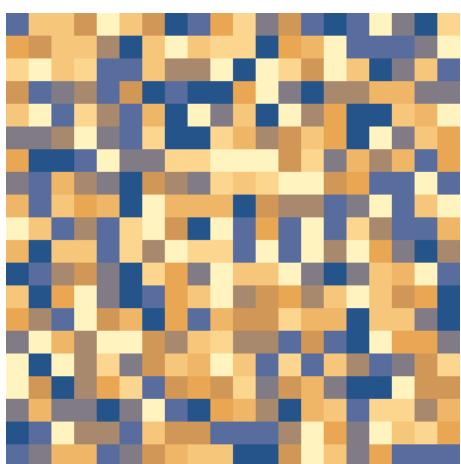
“Masking pattern” m defines subsets.



$$\phi' = \begin{cases} \phi'_{\text{frozen}} = \phi_{\text{frozen}} \\ \phi'_{\text{active}} = h(\phi_{\text{frozen}}) \times \phi_{\text{active}} + t \end{cases}$$

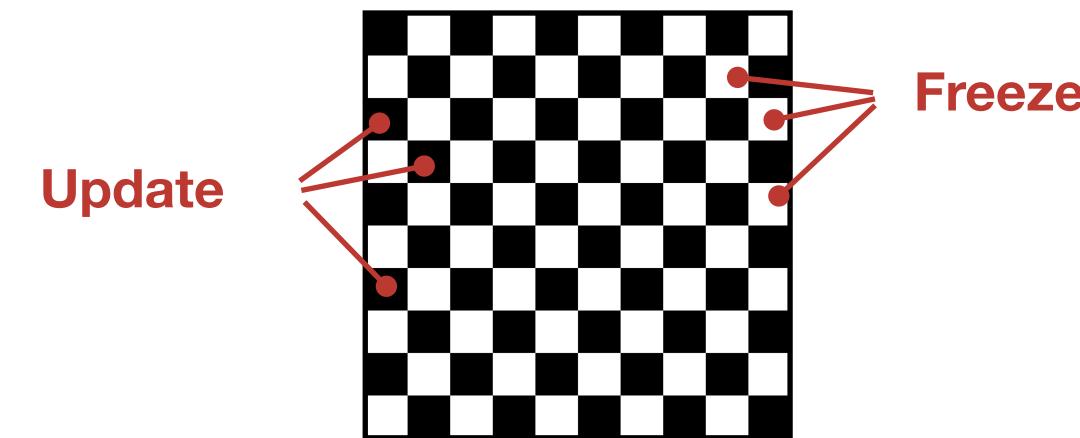
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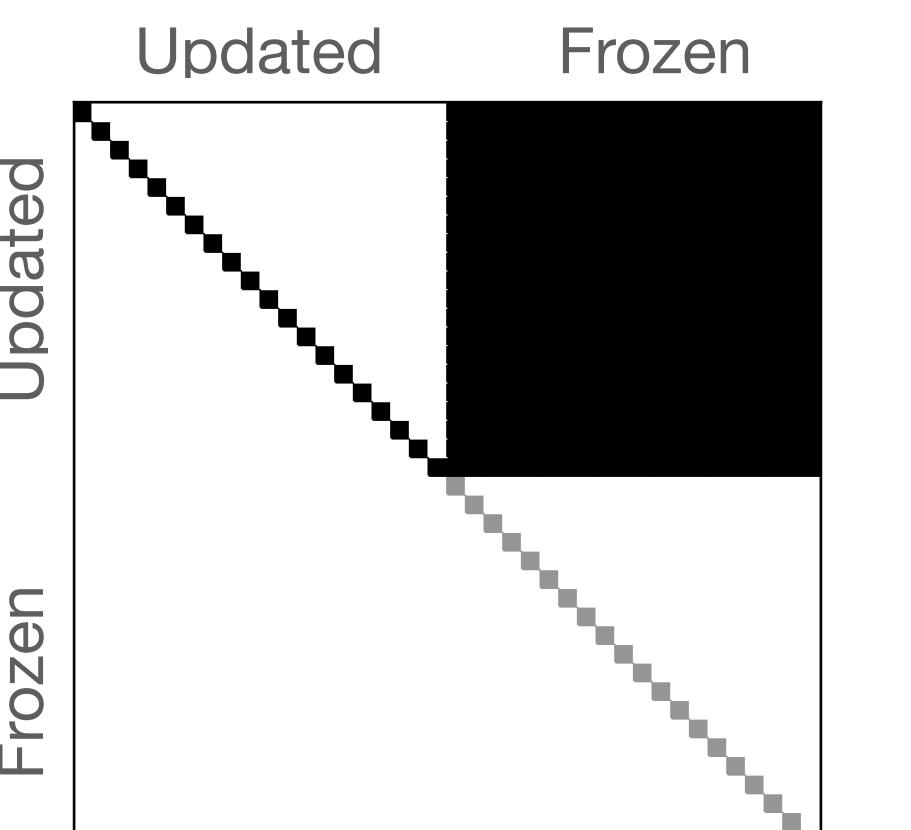
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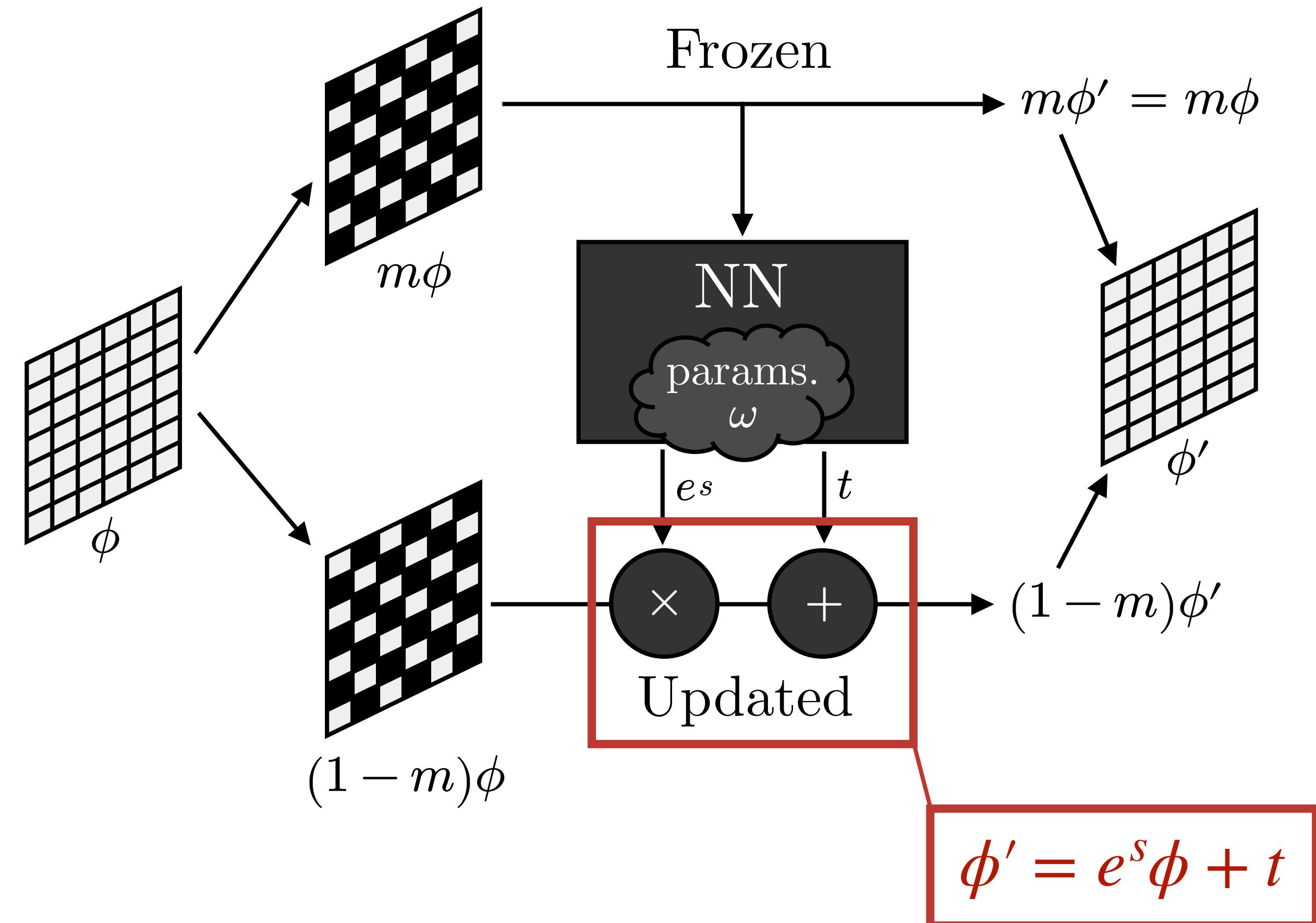
$$\text{Jacobian} = \frac{\partial \phi'}{\partial \phi} =$$



Implementing flows for scalar theory

Non-volume presering coupling layer:

[Dinh, Sohl-Dickstein, Bengio 1605.08803]



Flows for fermionic gauge theories

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- Straightforward approach consists on integrating out fermions:

$$\int \mathcal{D}[U] \mathcal{D}[\psi, \bar{\psi}] e^{-S_{\text{gauge}}(U)} e^{-S_{\text{ferm}}(\psi, \bar{\psi}, U)}$$

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- Full action can be expressed only in terms of gauge variables:

$$S_E(U) = -\beta \sum \operatorname{Re} P(x) - \log \det D[U]^\dagger D[U]$$

"exact determinant"

- One can use gauge-equivariant architectures from pure gauge models:

[\[Kanwar, Albergo, Boyda, Cranmer, Hackett, Racanière, Rezende, Shanahan 2003.06413\]](#)

[\[Boyda, Kanwar, Racanière, Rezende, Albergo, Cranmer, Hackett, Shanahan 2008.05456\]](#)

Training and using the models

- Train to minimize e.g. Kullback-Leibler divergence:

$$D_{\text{KL}}(q||p) = \int d\phi q(\phi)[\log q(U) + S(\phi)] + (\text{const})$$

- Self-training:
 1. Draw samples from the model to measure sample mean of $[\log q(U) + S(\phi)]$
 2. Gradient-based methods to optimize model parameters (e.g. Adam optimizer)

[Kingma, Ba, arXiv:1412.6980]

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! Trained models are not perfect but exactness is essential.

Exactness by forming a Markov chain with accept/reject Metropolis-Hastings steps

Acceptance probability

$$\longrightarrow A(\phi^{(i-1)}, \phi') = \min \left(1, \frac{\frac{q(\phi^{(i-1)})}{p(\phi')}}{\frac{p(\phi^{(i-1)})}{q(\phi')}} \right)$$

True distribution

Model distribution

Improving determinant estimators

- For a fixed U , the fermion determinant can be estimated as

$$\det M(U) \propto \int d\phi e^{-S_{\text{pf}}(\phi, U)} = \int d\phi q(\phi | U) \frac{e^{-S_{\text{pf}}(\phi, U)}}{q(\phi | U)} = \left\langle \frac{e^{-S_{\text{pf}}(\phi, U)}}{q(\phi | U)} \right\rangle_q$$

- By drawing more PF samples at fixed gauge field,
one can improve the determinant estimator

$$w_{N_{\text{pf}}} (U) = \frac{1}{N_{\text{pf}}} \sum_{i=1}^{N_{\text{pf}}} \frac{p(\phi^{(i)}, U)}{q(\phi^{(i)}, U)}$$

 weight of each config for reweighting or Markov Chain

- This does not require reevaluating the observables

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Numerical demonstration
in the Schwinger model

