

## Halo EFT: News and Views

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#### Outline



- Halo nuclei and Halo EFT
- Efimov physics in halo nuclei
   Zhang, Fu, Guo, HWH, Phys. Rev. C 108, 044304 (2023)
- Nuclear reactions with neutrons HWH, Son, Proc. Nat. Acad. Sci. 118, e2108716118 (2021)
- Summary and Outlook

### Halo Nuclei



#### Low separation energy of valence nucleons: B<sub>valence</sub> « B<sub>core</sub>, E<sub>ex</sub>

 $\longrightarrow$  close to "nucleon drip line"  $\longrightarrow$  scale separation  $\longrightarrow$  EFT



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#### EFT for halo nuclei

(Bertulani, HWH, van Kolck, 2002; Bedaque, HWH, van Kolck, 2003; ...)

## **Halo Effective Field Theory**



Separation of scales:

 $1/k = \lambda \gg R_{core}$ 

- Limited resolution at low energy:
  - $\longrightarrow$  expand in powers of  $kR_{core}$
  - $\longrightarrow$  contact interactions
- Short-distance physics not resolved
  - $\longrightarrow$  capture in low-energy constants using renormalization
  - $\longrightarrow$  include long-range physics explicitly if present
- Nucleon degrees of freedom: pionless EFT
- Exploit cluster substructures Halo EFT



#### **EFT Framework**



- Exploit scale separation in EFT framework
- Here: S-wave case, higher L states can also be treated
- Effective Lagrangian



## Limit Cycle



- **RG** invariance  $\implies$  running coupling  $H(\Lambda) = q_3 \Lambda^2 / (9q_2^2)$ 
  - $H(\Lambda)$  periodic: limit cycle

 $\Lambda \rightarrow \Lambda e^{n\pi/s_0} \approx \Lambda (22.7)^n$ 

(cf. Wilson, 1971)

Anomaly: scale invariance broken to discrete subgroup



$$H(\Lambda) \approx \frac{\cos(s_0 \ln(\Lambda/\Lambda_*) + \arctan(s_0))}{\cos(s_0 \ln(\Lambda/\Lambda_*) - \arctan(s_0))}, \quad s_0 \approx 1.00624$$

(Bedaque, HWH, van Kolck, 1999)

## Limit Cycle: Efimov Effect



Universal spectrum of three-body states (Efimov, 1970)



- Discrete scale invariance for fixed angle
- Geometrical spectrum for  $1/a \rightarrow 0$

$$B_3^{(n)}/B_3^{(n+1)} \xrightarrow{1/a \to 0} (e^{\pi/s_0})^2 = 515.035...$$

■ Ultracold atoms ⇒ variable scattering length ⇒ loss resonances

## **Efimov Physics in Halo Nuclei**



- Efimov effect in halo nuclei? (Fedorov, Jensen, Riisager, 1994) ⇒ excited states obeying scaling relations
- Correlation plot:  $E_{nn} \leftrightarrow S_{1n}$  (Amorin, Frederico, Tomio, 1997)



HWH, Ji, Phillips, J. Phys. G 44, 103002 (2017)

Alternative ways to observe Efimov physics in halo nuclei?

### **Neutron Scattering**



- Neutron scattering off  $J^P = 1/2^+$  one-neutron halos (<sup>11</sup>Be, <sup>15</sup>C, <sup>19</sup>C)
- *J* = 0 channel (three-body force suppressed)



• J = 1 channel (no three-body force)



Zhang, Fu, Guo, HWH, Phys. Rev. C 108, 044304 (2023)

## J = 1 channel



#### S-wave scattering amplitude (Zhang, Fu, Guo, HWH, Phys. Rev. C 108, 044304 (2023))



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## J = 0 channel



#### S-wave scattering amplitude (Zhang, Fu, Guo, HWH, Phys. Rev. C 108, 044304 (2023))



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## Efimov physics in halo nuclei?



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#### Alternative ways to see Efimov physics in halo nuclei?

- Pole in  $p \cot \delta_0$  due to virtual Efimov state close to threshold
- **B** Real excited state appears for  $1/a_{\sigma}$  smaller than some critical value



(Zhang, Fu, Guo, HWH, Phys. Rev. C 108, 044304 (2023))

#### Pole position directly correlated with virtual state energy

- $\implies$  pole position determined by Efimov physics
- Also present in deuteron-halo scattering?

#### **Reactions with neutrons**



 High-energy nuclear reaction with multi-neutron final state (HWH, Son, Proc. Nat. Acad. Sci. 118, e2108716118 (2021))



- Assumption: energy scale of primary reaction  $\gg E_{\mathcal{U}} \frac{\mathbf{p}^2}{2M_{\mathcal{U}}} = E_n^{cms}$
- Factorization:  $\frac{d\sigma}{dF} \sim |\mathcal{M}_{primary}|^2 \operatorname{Im} G_{\mathcal{U}}(E_{\mathcal{U}}, \boldsymbol{p})$
- Reproduces Watson-Migdal treatment of FSI for 2n (Watson, Phys. Rev. 88, 1163 (1952); Migdal, Sov. Phys. JETP 1, 2 (1955))

## **Unitary limit**



Spin-1/2 Fermions with zero-range interactions  $(|a| \gg r_e)$ 



Renormalization group equation:

$$\Lambda rac{d}{d\Lambda} ilde{g}_2 = ilde{g}_2 (1 + ilde{g}_2)$$

Two fixed points:

 $- \, ilde{g}_2 = 0 \,\, \Leftrightarrow \,\, a = 0 \,\,\,\,\, \Rightarrow \,\,$  no interaction, free particles

 $- \, { ilde g}_2 = -1 \, \Leftrightarrow \, 1/a = 0 \quad \Rightarrow \, {
m unitary \, limit}$ 

⇒ conformal/Schrödinger symmetry

(Mehen, Stewart, Wise, PLB 474, 145 (2000); Nishida, Son, PRD 76, 086004 (2007); ...)

Neutrons:  $a \approx -18.6$  fm,  $r_e \approx 2.8$  fm

 $\Rightarrow$  neutrons are close to the unitary limit

## **Conformal field theory**



Two-point function of primary field operator U ("unnucleus")

$$G_{\mathcal{U}}(t, \mathbf{x}) = -i \langle T \mathcal{U}(t, \mathbf{x}) \mathcal{U}^{\dagger}(0, \mathbf{0}) 
angle = C \frac{\theta(t)}{(it)^{\Delta}} \exp\left(\frac{iM\mathbf{x}^2}{2t}\right)$$

- Determined by symmetry up to overall constant C
- Two-point function in momentum space

$$G_{\mathcal{U}}(\omega, \boldsymbol{p}) = -\frac{C}{M} \left(\frac{2\pi}{M}\right)^{3/2} \Gamma\left(\frac{5}{2} - \Delta\right) \left(\frac{\boldsymbol{p}^2}{2M} - \omega - i\epsilon\right)^{\Delta - \frac{5}{2}}$$

- pole only for Δ = 3/2 (free field)
   branch cut for Δ > 3/2
- General unnucleus/unparticle does not behave like a particle ⇒ continuous energy spectrum

#### **Reactions with neutrons**



#### Two ways to do experiments

(a) detect recoil particle B

$$rac{d\sigma}{dE} \sim (E_0 - E_B)^{\Delta - 5/2}, \qquad E_0 = (1 + M_B/M_U)^{-1} E_{\rm kin}$$

(b) detect all final state particles including neutrons

$$\frac{d\sigma}{dE} \sim (E_n^{cms})^{\Delta-5/2}$$

(HWH, D.T. Son, Proc. Nat. Acad. Sci. 118, e2108716118 (2021))

- Consistent with previous experiments for <sup>3</sup>H(π<sup>-</sup>, γ)3n (Miller et al., Nucl. Phys. A 343, 347 (1980))
- Two few events in recent tetraneutron experiment: <sup>4</sup>He(<sup>8</sup>He,<sup>8</sup>Be)4n (Kisamori et al., Phys. Rev. Lett. **116**, 052501 (2016))

## **Reaction calculations**



Two-neutron spectrum for  ${}^{6}\text{He}(p, p\alpha)2n$  (Göbel et al., Phys. Rev. C **104**, 024001 (2021))



Can be understood from dimer propagator (Δ = 2)

$$G_d(E_{nn}, \mathbf{0}) \sim \frac{1}{1/a + i\sqrt{mE_{nn}}} \quad \Rightarrow \quad \operatorname{Im} G_d(E_{nn}, \mathbf{0}) \sim \frac{\sqrt{E_{nn}}}{(ma^2)^{-1} + E_{nn}}$$

#### **Reaction calculations**



#### Radiative muon/pion capture on the triton (AV18 + UIX)



#### New experiments



- New experiments in complete kinematics at RIBF/RIKEN
- Measurement of a<sub>nn</sub> in <sup>6</sup>He(p, pα)2n (T. Aumann et al., NP2012-SAMURAI55R1 (2020))
- Search for tetraneutron resonances in <sup>8</sup>He(p, pα)4n (S. Paschalis et al., NP1406-SAMURAI19R1 (2014))



#### **New experiments**





M. Duer et al., Nature 606 (2022) 678

 $E_R = 2.37 \pm 0.38(stat) \pm 0.44(sys)$  MeV,  $\Gamma_R = 1.75 \pm 0.22(stat) \pm 0.30(sys)$  MeV

- Genuine resonance or dineutron correlations/initial state effect?
  - R. Lazauskas et al., Phys. Rev. Lett. 130, 102501 (2023)  $\implies$  dineutron correlations

## **Summary and Outlook**





- Efimov physics in halo nuclei
- High-energy nuclear reactions with final state neutrons
  - ⇒ (approximate) conformal symmetry
  - $\Rightarrow$  power law behavior of observables determined by  $\Delta$
- Model-independent constraints on nuclear reactions
- Connection between reactions & properties of trapped particles

## **Summary and Outlook**





- Efimov physics in halo nuclei
- High-energy nuclear reactions with final state neutrons
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- Model-independent constraints on nuclear reactions
- Connection between reactions & properties of trapped particles
- Other applications & extensions
  - Two-component Fermions in ultracold atom physics
  - Neutral charm mesons (Braaten, HWH, Phys. Rev. Lett. 128, 032002 (2022), Phys. Rev. D 107, 034017 (2023))
  - Systems with the Efimov effect?
    - $\Rightarrow$  bosonic atoms, nucleons,  $\alpha$  particles
    - $\Rightarrow$  complex scaling dimensions
    - $\Rightarrow$  scale symmetry broken

## **Additional Slides**



## **Conformal field theory**



Imaginary part of propagator

$$\operatorname{Im} \mathsf{G}_{\mathcal{U}}(\omega, \mathbf{p}) \sim \begin{cases} \delta\left(\omega - \frac{\mathbf{p}^2}{2M}\right), & \Delta = \frac{3}{2}, \\ \left(\omega - \frac{\mathbf{p}^2}{2M}\right)^{\Delta - \frac{5}{2}} \theta\left(\omega - \frac{\mathbf{p}^2}{2M}\right), & \Delta > \frac{3}{2} \end{cases}$$

#### Examples of unnuclei

- free field:  $\mathcal{U}=\psi$ ,  $M=m_\psi$ ,  $\Delta=3/2$
- **D** N free fields:  $U = \psi_1 \dots \psi_N$ ,  $M = Nm_{\psi}$ ,  $\Delta = 3N/2$
- N interacting fields:  $U = \psi_1 \dots \psi_N$ ,  $M = Nm_{\psi}$ ,  $\Delta > 3/2$

# • In our case: unnucleus is strongly interacting multi-neutron state with $1/(ma^2) \sim 0.1 \text{ MeV} \ll E_n^{cms} \ll 1/(mr_e^2) \sim 5 \text{ MeV}$

Corrections from finite *a* and *r*<sub>0</sub> (S. Dutta, R. Mishra, D.T. Son, arXiv:2309.15177)

$$\operatorname{Im} G_{\mathcal{U}}(\omega, 0) \sim \omega^{\Delta - \frac{5}{2}} \theta(\omega) \left( 1 + \frac{c_1}{a \sqrt{m\omega}} + c_2 r_0 \sqrt{m\omega} \right), \quad c_2 = 0$$

#### Scaling dimension





#### ■ How to calculate scaling dimension △?

- (1)  $\Delta$  can be obtained from field theory calculation
- (2)  $\Delta$  can be obtained from operator state correspondence

 $\Delta$  of primary operator = (Energy of state in HO)/ $\hbar\omega$ 

(Nishida, Son, Phys. Rev. D 76, 086004 (2007))

Ν	S	L	$\mathcal{O}$	Δ
2	0	0	$\psi_1\psi_2$	2
3	1/2	1	$\psi_1\psi_2 abla_j\psi_2$	4.27272
3	1/2	0	$\psi_1 \nabla_j \psi_2 \nabla_j \psi_2$	4.66622
4	0	0	$\psi_1\psi_2\nabla_i\psi_1\nabla_j\psi_2$	5.07(1)
5	1/2	1		7.6(1)

#### $\Rightarrow$ connection between $\Delta$ and energy of particles in a trap