Southampion elusiPes in(PisiblesPlus
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# Flavour Symmetries and their Modular Origin 

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## Bethe Forum

Modular Flavor Symmetries


## The Standard Model

## Left-handed

Right-handed

(Including three right-handed neutrinos)

## The Flavour Problem



## Masses



## Mixing

CKM


PMNS
$v_{1}$
$v_{2}$
$v_{3}$

$\nu_{\tau} \quad \square$


## PMNS Lepton mixing matrix



Standard Model states


$$
\left(\begin{array}{l}
\nu_{e} \\
\nu_{\mu} \\
\nu_{\tau}
\end{array}\right)=\left(\begin{array}{lll}
U_{e 1} & U_{e 2} & U_{e 3} \\
U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\
U_{\tau 1} & U_{\tau 2} & U_{\tau 3}
\end{array}\right)\left(\begin{array}{l}
\nu_{1} \\
\nu_{2} \\
\nu_{3}
\end{array}\right)
$$



PMNS Lepton mixing matrix

$$
U_{P M N S}=\left(\begin{array}{ccc}
1 & 0 & 0 \\
0 & c_{23} & s_{23} \\
0 & -s_{23} & c_{23}
\end{array}\right)\left(\begin{array}{ccc}
c_{13} & 0 & s_{13} e^{-i \delta} \\
0 & 1 & 0 \\
-s_{13} e^{i \delta} & 0 & c_{13}
\end{array}\right)\left(\begin{array}{ccc}
c_{12} & s_{12} & 0 \\
-s_{12} & c_{12} & 0 \\
0 & 0 & 1
\end{array}\right)\left(\begin{array}{ccc}
1 & 0 & 0 \\
0 & e^{i \frac{\alpha_{21}}{2}} & 0 \\
0 & 0 & e^{i \frac{\alpha_{31}}{2}}
\end{array}\right)
$$

Atmospheric
Reactor
Solar
Majorana

$$
\begin{array}{r}
\left(\begin{array}{ccc}
c_{12} c_{13} & s_{12} c_{13} & s_{13} e^{-i \delta} \\
-s_{12} c_{23}-c_{12} s_{13} s_{23} e^{i \delta} & c_{12} c_{23}-s_{12} s_{13} s_{23} e^{i \delta} & c_{13} s_{23} \\
s_{12} s_{23}-c_{12} s_{13} c_{23} e^{i \delta} & -c_{12} s_{23}-s_{12} s_{13} c_{23} e^{i \delta} & c_{13} c_{23}
\end{array}\right) \\
\times \operatorname{diag}\left(1, e^{i \alpha_{21} / 2}, e^{i \alpha_{31} / 2}\right)
\end{array}
$$



## Who ordered all of that?



Isidor Issac Rabi

## SM Yukawa couplings

$$
y_{i j} H \bar{\psi}_{L i} \psi_{R j}
$$



Is there a symmetry at work?

## Family/Flavour Symmetry

Basic idea is to distinguish the families by some quantum numbers

## under a new "horizontal" <br> family/flavour symmetry

The symmetry is assumed to be spontaneously broken by
"flavons"


## Family/Flavour Symmetry



## Example: U(1) Family/Flavour Symmetry

Consider a $U(1)$ family symmetry spontaneously broken by a flavon vev $\quad\langle\phi\rangle \neq 0$

Suppose $\mathrm{U}(1)$ charges are $\mathrm{Q}\left(\psi_{3}\right)=0, \mathrm{Q}\left(\psi_{2}\right)=1, \mathrm{Q}\left(\psi_{1}\right)=3, \mathrm{Q}(\mathrm{H})=0, \mathrm{Q}(\phi)=-1$

Then the lowest order allowed Yukawa coupling is $\mathrm{H} \psi_{3} \psi_{3}$

$$
Y=\left(\begin{array}{lll}
0 & 0 & 0 \\
0 & 0 & 0 \\
0 & 0 & 1
\end{array}\right)
$$

The other Yukawa couplings are generated from higher order operators which respect $\mathrm{U}(1)$ family symmetry due to flavon $\phi$ insertions:

$$
\frac{-1+0+1+0=0}{M} H \psi_{2} \psi_{3}+\left(\frac{\phi}{M}\right)^{2} H \psi_{2} \psi_{2}+\left(\frac{\phi}{M}\right)^{3} H \psi_{1} \psi_{3}+\left(\frac{\phi}{M}\right)^{4} H \psi_{1} \psi_{2}+\left(\frac{\phi}{M}\right)^{6} H \psi_{1} \psi_{1}
$$

When the flavon gets its VEV it generates small effective Yukawa couplings in terms
of an expansion parameter

$$
\varepsilon=\frac{\langle\phi\rangle}{M}
$$

$\underset{\text { texture zero }}{\text { Approximate }} \xrightarrow[Y]{ }=\left(\begin{array}{lll}\varepsilon^{6} & \varepsilon^{4} & \varepsilon^{3} \\ \varepsilon^{4} & \varepsilon^{2} & \varepsilon \\ \varepsilon^{3} & \varepsilon & 1\end{array}\right)$

## Froggatt-Nielsen Mechanism (1979)



What is the origin of the higher order operators?
Froggatt and Nielsen took their inspiration from the see-saw mechanism

$$
\frac{H^{2}}{M_{v_{R}}} v_{L} \nu_{L}
$$



$$
\frac{\phi}{M_{\chi}} H \psi_{2} \psi_{3}
$$

Where $\chi$ are heavy fermion messengers c.f. heavy RH neutrinos

## Froggatt-Nielsen Mechanism (1979)

There may be Higgs messengers or fermion messengers

$\psi_{2} \quad \psi_{3}$
Fermion messengers may be $\operatorname{SU}(2)_{\llcorner }$doublets or singlets


## Neutrinos motivate new family/flavour symmetries

## CKM Matrix

PMNS Matrix


Froggatt-Nielsen tends to predict small mixing


What symmetry gives this?

Mu-Tau Symmetry $\nu_{\mu} \leftrightarrow \nu_{\tau}^{*}$


Basic Idea:
Two rows have equal magnitudes Z.z.Xing and S.Zhou, 0804.3512

$$
\rightarrow \quad \theta_{13} \neq 0, \quad \theta_{23}=45^{\circ}, \quad \delta_{\mathrm{CP}}= \pm 90^{\circ}
$$

Mu-Tau Symmetry $\nu_{\mu} \leftrightarrow \nu_{\tau}^{*}$

$$
\begin{gathered}
\mathrm{Bi}^{( }\left\{\begin{array}{c}
- \\
\rightarrow
\end{array}\right) \begin{array}{c}
\text { Basic Idea: } \\
\text { Two rows have }
\end{array} \\
\theta_{13} \neq 0, \quad \theta_{23}=45^{\circ}, \quad \begin{array}{c}
\text { equal magnitudes } \\
\text { e.z.xing and s.zhou, o804.3512 }
\end{array} \\
\delta_{\mathrm{CP}}= \pm 90^{\circ}
\end{gathered}
$$

#  Tri-Bimaximal-Reactor 


$\sin \theta_{12}=\frac{1}{\sqrt{3}}$
Allowed at 3 sigma
$\sin \theta_{23}=\frac{1}{\sqrt{2}} \quad \sin \theta_{13}=\frac{\lambda}{\sqrt{2}}$
Allowed at
3 sigma

Huge literature e.g. Antusch and SFK, hep-ph/0508044; I.Girardi, S.T.Petcov and A.V.Titov,1410.8056, .. Charged lepton corrections

Charged lepton rotation Tri-bimaximal neutrinos

$$
\begin{aligned}
& U_{\text {PMNS }}=\left(\begin{array}{ccc}
c_{12}^{e} & s_{12}^{e} e^{-i \delta_{12}^{e}} & 0 \\
-s_{12}^{e} e^{i \delta_{12}} & c_{12}^{e} & 0 \\
0 & 0 & 1
\end{array}\right)\left(\begin{array}{ccc}
\sqrt{\frac{2}{3}} & \frac{1}{\sqrt{3}} & 0 \\
-\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \\
\frac{1}{\sqrt{6}} & -\frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}}
\end{array}\right) \\
& \begin{array}{r}
=\left(\begin{array}{ccc}
\cdots & \cdots & \left.\frac{s_{12}^{e}}{\sqrt{2}} e^{-i \delta_{1}^{e}}\right) \\
\cdots & \ldots & \frac{c_{12}^{e}}{1 / 2} \\
\frac{1}{\sqrt{6}} & -\frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}}
\end{array}\right)
\end{array}
\end{aligned}
$$

Tri-maximal Mixing

$$
\begin{aligned}
& (T M 2)^{T r i} \\
& \left\{\left(\begin{array}{ll}
0 & \mathrm{TMI} \\
0 & { }^{0}
\end{array}\right) \not U_{\mathrm{TM} 1} \approx\left(\begin{array}{ccc}
\sqrt{\frac{2}{3}} & - & - \\
-\frac{1}{\sqrt{6}} & - & - \\
\frac{1}{\sqrt{6}} & - & -
\end{array}\right)\right. \\
& \text { Second column of TBM } \\
& U_{\mathrm{TM} 2} \approx\left(\begin{array}{ccc}
- & \frac{1}{\sqrt{3}} & - \\
- & \frac{1}{\sqrt{3}} & - \\
- & -\frac{1}{\sqrt{3}} & -
\end{array}\right) \\
& \text { First column of TBM }
\end{aligned}
$$

C.H.Albright and W.Rodejohann, 0812.0436; C.H.Albright, A.Dueck and W.Rodejohann, 1004.2798

## Tri-maximal Mixing


C.H.Albright and W.Rodejohann, 0812.0436; C.H.Albright, A.Dueck and W.Rodejohann, 1004.2798

## Tri-maximal Mixing



## Littlest Seesaw



- Fit includes effects of RG corrections
- Determines the RHN masses!

SFK,1304.6264; 1512.07531 SFK, Molina Sedgwick, Rowley, 1808.01005

4 real input parameters
Describes:
3 neutrino masses ( $m_{1}=0$ ), 3 mixing angles, 1 Dirac CP phase, 2 Majorana phases (1 zero) 1 BAU parameter $\mathrm{Y}_{\mathrm{B}}$ = 10 observables of which 7 are constrained

| Predictions | $1 \sigma$ range |
| :--- | ---: |
| $\theta_{12} /^{\circ}$ | $34.254 \rightarrow 34.350$ |
| $\theta_{13} /{ }^{\circ}$ | $8.370 \rightarrow 8.803$ |
| $\theta_{23} /^{\circ}$ | $45.405 \rightarrow 45.834$ |
| $\Delta m_{12}{ }^{2} / 10^{-5} \mathrm{eV}^{2}$ | $7.030 \rightarrow 7.673$ |
| $\Delta m_{31}{ }^{2} / 10^{-3} \mathrm{eV}^{2}$ | $2.434 \rightarrow 2.561$ |
| $\delta /{ }^{\circ}$ | $-88.284 \rightarrow-86.568$ |
| $Y_{B} / 10^{-10}$ | $0.839 \rightarrow 0.881$ |

##  Non-Abelian Family Symmetry

## Traditionally used for TB mixing

S.F.K. and G.G. Ross,
hep-ph/0108112; hep-ph/0307190
E.Ma and G.Rajasekaran, hep-ph/0106291;
K.S.Babu, E.Ma, J.W.F.Valle, hep-ph/0206292;
G.Altarelli and F.Feruglio, hep-ph/0504165,hep-ph/0512103
I.de Medeiros Varzielas
S.F.K. and G.G. Ross,
hep-ph/0512313;
hep-ph/0607045


These days can explain charged lepton corrections, TMI,TM2, Littlest seesaw,...

## Reviews

F.Feruglio and A.Romanino, 1912.06028 S.F.K. and C.Luhn, 1301.1340
S.F.K., A.Merle, S.Morisi, Y.Shimizu and M.Tanimoto, 1402.4271


| $S^{2}=T^{3}=U^{2}=(S T)^{3}=(S U)^{2}=(T U)^{2}=(S T U)^{4}=1$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| $S_{4}$ | $A_{4}$ | $S$ | $T$ | $U$ |
| $\mathbf{1}, \mathbf{1}^{\prime}$ | $\mathbf{1}$ | 1 | 1 | $\pm 1$ |
| $\mathbf{2}$ | $\binom{\mathbf{1}^{\prime \prime}}{\mathbf{1}^{\prime}}$ | $\left(\begin{array}{ll}1 & 0 \\ 0 & 1\end{array}\right)$ | $\left(\begin{array}{cc}\omega & 0 \\ 0 & \omega^{2}\end{array}\right)$ | $\left(\begin{array}{ll}0 & 1 \\ 1 & 0\end{array}\right)$ |
| $\mathbf{3}, \mathbf{3}^{\prime}$ | $\mathbf{3}$ | $\frac{1}{3}\left(\begin{array}{ccc}-1 & 2 & 2 \\ 2 & -1 & 2 \\ 2 & 2 & -1\end{array}\right)$ | $\left(\begin{array}{ccc}1 & 0 & 0 \\ 0 & \omega^{2} & 0 \\ 0 & 0 & \omega\end{array}\right)$ | $\mp\left(\begin{array}{lll}1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0\end{array}\right)$ |

Diagonalised by TB matrix

#  $S_{4}$ vacuum alignments ${ }^{\substack{\text { ciatachise }}}$ 

$\left\langle\phi_{3^{\prime}}^{\nu}\right\rangle=\left(\begin{array}{l}1 \\ 1 \\ 1\end{array}\right)$ preserves $\mathbb{S}, \cup \quad\left\langle\phi_{3^{\prime}}^{l}\right\rangle=\left(\begin{array}{l}1 \\ 0 \\ 0\end{array}\right)$ preserves $T$

| residual <br> symmetry | $U$ | $S$ | $S U$ |
| :---: | :---: | :---: | :---: |
| $\mathbf{1}$ | - | - | - |
| $\mathbf{1}^{\prime}$ | - | 1 | - |
| $\mathbf{2}$ | - | $(1,-1)^{T}$ | - |
| $\mathbf{3}$ | $(0,1,-1)^{T}$ | $(1,1,1)^{T}$ | $(2,-1,-1)^{T}$ |
| $\mathbf{3}^{\prime}$ | $(1,0,0)^{T}$ | - | $(0,1,-1)^{T}$ |

$S_{4}$ flavour symmetry


$T M^{E} T=M^{E}$
$S M^{\nu} S=M^{\nu}$
$U M^{\nu} U=M^{\nu}$
$\frac{\sqrt{2}}{3}$
$-\frac{1}{\sqrt{6}}$
$\frac{1}{\sqrt{6}}$

$$
\left.\begin{array}{cc}
\frac{1}{\sqrt{3}} & 0 \\
\frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \\
-\frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}}
\end{array}\right)
$$

TB mixing excluded so need to break S,T,U
$\mathrm{S}_{4}$ flavour symmetry

$\mathrm{S}_{4}$ flavour symmetry

$\mathrm{S}_{4}$ flavour symmetry




Y.Koide,0705.2275; T.Banks and N.Seiberg,1011.5120;
Y.L.Wu,1203.2382; A.Merle and R.Zwicky,1110.4891;
B.L.Rachlin and T.W.Kephart,1702.08073; C. Luhn, 1101.2417

## Origin of flavour symmetry

## Break SO(3) using large Higgs reps E.g. 7-plet

| irrep | $\underline{1}$ | $\underline{3}$ | $\underline{5}$ | 7 |
| :--- | :--- | :--- | :--- | :--- |
| subgroups | $S O(3)$ | $S O(2)$ | $Z_{2} \times Z_{2}$ | 1 |
|  |  | $S O(3)$ | $S O(2)$ | $A_{4}$ |
|  |  |  | $S O(3)$ | $Z_{3}$ |
|  |  |  |  | $D_{4}$ |
|  |  |  | $S O(2)$ |  |
|  |  |  | $S O(3)$ |  |

A4 preserving direction of 7-PletVEV
$\left\langle\xi_{123}\right\rangle \equiv \frac{v_{\xi}}{\sqrt{6}}, \quad\left\langle\xi_{111}\right\rangle=\left\langle\xi_{112}\right\rangle=\left\langle\xi_{113}\right\rangle=\left\langle\xi_{133}\right\rangle=\left\langle\xi_{233}\right\rangle=\left\langle\xi_{333}\right\rangle=0$
S.F.K. and Ye-Ling Zhou, 1809.10292


## Modular Symmetry



General modular transformation

$$
\left.\begin{array}{ccc}
\tau \rightarrow \gamma \tau=\frac{a \tau+b}{c \tau+d} & \begin{array}{c}
\text { Integers a,b,c,d } \\
a d-b c=1
\end{array} & \gamma=\left(\begin{array}{cc}
a & b \\
c & d
\end{array}\right)
\end{array} \begin{array}{c}
\text { Infinite group } \\
\Gamma \equiv S L(2, \mathbb{Z})
\end{array}\right] \begin{array}{cc} 
\\
S: \tau \mapsto-\frac{1}{\tau}, \quad T: \tau \mapsto \tau+1 & S=\left(\begin{array}{cc}
0 & 1 \\
-1 & 0
\end{array}\right) \quad T=\left(\begin{array}{ll}
1 & 1 \\
0 & 1
\end{array}\right)
\end{array}
$$

## From Infinite to Finite Modular Symmetry

$$
\Gamma \equiv S L(2, \mathbb{Z}) \quad S^{2}=-\mathbb{1}_{2}, \quad S^{4}=(S T)^{3}=\mathbb{1}_{2}, \quad S^{2} T=T S^{2}
$$

$\longrightarrow \bar{\Gamma} \equiv \operatorname{PSL}(2, \mathbb{Z}) \quad S^{2}=(S T)^{3}=(T S)^{3}=1 \quad$ Infinite
Finite
level N
$\Gamma_{N} \quad S^{2}=(S T)^{3}=(T S)^{3}=1 \quad$ and $\quad T^{N}=1$

| $\Gamma_{2} \approx S_{3}$ |
| :--- |
| $\Gamma_{3} \approx A_{4}$ |
| $\Gamma_{4} \approx S_{4}$ |
| $\Gamma_{5} \approx A_{5}$ |
| $\Gamma_{7} \approx \Sigma(168)$ |

Yukawa coupling transforms as an irrep of $\Gamma_{\mathrm{N}}$ and as a modular form

$$
\begin{aligned}
& Y(\tau) \rightarrow Y(\gamma \tau)=(c \tau+d)^{k_{Y}} \rho_{\mathbf{r}_{Y}}(\gamma) Y(\tau) \\
& Y(\tau) \phi_{1} \phi_{2} \phi_{3} \quad \phi_{1} \rightarrow(c \tau+d)^{k_{1}} \rho_{1}(\gamma) \phi_{1}
\end{aligned}
$$

$k_{Y}=k_{1}+k_{2}+k_{3}$ modular weights balance
$\rho_{\mathbf{r}_{Y}} \times \rho_{1} \times \rho_{2} \times \rho_{3}=1+\ldots \quad$ contains singlet

|  | Leptons | Quarks | $S \circlearrowleft(5)$ | $S O(10)$ |
| :---: | :---: | :---: | :---: | :---: |
| $N=2, S 3$ | T.Kobayashi, K.Tanaka and T.H.Tatsuishi, 1803.10391,... |  | T.Kobayashi, <br> Y. Shimizu, K.Takagi, <br> M.Tanimoto and <br> T.H.Tatsuishi <br> 1906.10341,... |  |
| $N=3, A_{4}$ | F.Feruglio,1706.08749 <br> J.C.Criado and <br> F.Feruglio, 1807. 01125 <br> G.J.Ding,S.F.King and <br> X.G.Liu, 1907.11714,... | H.Okada, M.Tanimoto, 1812.09677; <br> 1905.13421; <br> S.J.D. King,S.F.King, $2002.00969, \ldots$ | F.J.de Anda, S.F.King, E.Perdomo, 1812.05620; P.Chen, G.J.Ding and S.F.King, 2101.12724, ... | G.J.Ding, S.F.King, J.N.Lu, 2108.09655 |
| $N=4, S_{4}$ | J.T.Penedo,S.T.Petcov, 1806.11040;P.P.Novichkov <br> J.T.Penedo,S.T.Petcov, <br> A.V.Titov,1811.04933, <br> J.C.Criado,F.Feruglio, <br> S.J.D.King, 1908.11867.... |  | Y. Zhao and H.H.Zhang, 2101.02266; G.J.Ding, S.F.King and C.Y.Yao, 2103.16311 , ... |  |
| $N=5, A_{5}$ | P.P.Novichkov, <br> J.T.Penedo, S.T.Petcov, <br> A.V.Titov, 1812.02158; <br> G.J.Ding, S.F.King, <br> X.G.Liu, 1903.12588,... |  |  |  |
| $\mathrm{N}=6, \mathrm{~S}_{3} \times \mathrm{A}_{4}$ |  |  |  |  |
| $N=7, \quad \sum(168)$ | G.J.Ding, S.F.King, C.C.Li, Y.L.Zhou, 2004.12662 |  |  |  |

For integer/fractional/CP/eclectic/stabilisation... see other talks

See Ferruccio's talk F.Feruglio, 1706.08749

## Example: Level $\mathrm{N}=3 \sim \mathrm{~A}_{4}$

Yukawa couplings involving twisted states whose modular weights do not add up to zero are modular forms

A 4 triplet 3
Weight ky $=2 \quad Y=\left(\begin{array}{l}Y_{1}(\tau) \\ Y_{2}(\tau) \\ Y_{3}(\tau)\end{array}\right)=\left(\begin{array}{c}1+12 q+36 q^{2}+12 q^{3}+84 q^{4}+72 q^{5}+\ldots \\ -6 q^{1 / 3}\left(1+7 q+8 q^{2}+18 q^{3}+18 q^{4}+\ldots\right) \\ -18 q^{2 / 3}\left(1+2 q+5 q^{2}+4 q^{3}+8 q^{4}+\ldots\right)\end{array}\right)$
Notation $Y=Y_{3}(2)$

$$
q \equiv e^{i 2 \pi \tau} \pi \text { modulus vev }
$$

Weinberg $\frac{1}{\Lambda}$
operator
$\Lambda$$H_{u} H_{u} L L Y Y\left(\begin{array}{ccc}2 Y_{1} & -Y_{3} & -Y_{2} \\ -Y_{3} & 2 Y_{2} & -Y_{1} \\ -Y_{2} & -Y_{1} & 2 Y_{3}\end{array}\right) \frac{v_{u}^{2}}{\Lambda}$ A ${ }_{4}$ rep: 333
Modular weights k: I I 2
no flavons (apart from tau)

\section*{Example with weighton: Level $\mathrm{N}=3 \sim \mathrm{~A}_{4}$ <br> |  | $L$ | $e_{3}^{c}$ | $e_{2}^{c}$ | $e_{1}^{c}$ | $N^{c}$ | $H_{u, d}$ | $\phi$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $A_{4}$ | $\mathbf{3}$ | $\mathbf{1}^{\prime}$ | $\mathbf{1}^{\prime \prime}$ | $\mathbf{1}$ | $\mathbf{3}$ | $\mathbf{1}$ | $\mathbf{1}$ |
| $k_{I}$ | 1 | 0 | -1 | -3 | 1 | 0 | 1 | <br> \[

W_{d r i v}=\chi\left(Y_{1}^{(4)} \frac{\phi^{4}}{M_{f l}^{2}}-M^{2}\right)

\] <br> \[

\tilde{\phi} \equiv \frac{\langle\phi\rangle}{M_{f l}} \sim\left(M / M_{f l}\right)^{1 / 2} \quad small
\]}

$$
W_{e}=\alpha_{e} e_{1}^{c} \tilde{\phi}^{4}\left(L Y_{\mathbf{3}}^{(2)}\right)_{\mathbf{1}} H_{d}+\beta_{e} e_{2}^{c} \tilde{\phi}^{2}\left(L Y_{\mathbf{3}}^{(2)}\right)_{\mathbf{1}^{\prime}} H_{d}+\gamma_{e} e_{3}^{c} \tilde{\phi}\left(L Y_{\mathbf{3}}^{(2)}\right)_{\mathbf{1}^{\prime \prime}} H_{d}
$$

$\left(\begin{array}{lll}\alpha_{e} \tilde{\varphi}^{4} Y_{1} & \alpha_{e} \tilde{\phi}^{4} Y_{3} & \alpha_{e} \tilde{\varphi}^{4} Y_{2} \\ \beta_{e} Y_{2}\end{array}\right) \quad$ Natural explanation of charged lepton hierarchy c.f. FN

Unlike the FN flavon, the weighton phi does not break the flavour symmetry

## Stabilizers and Fixed points

$$
\begin{aligned}
& \gamma_{0} \tau_{0}=\tau_{0} \quad \text { e.g. } \quad S \tau_{S}=\tau_{S} \quad \longrightarrow \quad \tau_{S}=i \\
& \text { Invariant under } \quad S: \tau \mapsto-\frac{1}{\tau}
\end{aligned}
$$

## Alignments from fixed points

Modular transformation
$Y_{I_{Y}}(\gamma \tau)=(c \tau+d)^{2 k_{Y}} \rho_{I_{Y}}(\gamma) Y_{I_{Y}}(\tau)$
Fixed point relations

$$
\gamma \tau_{\gamma}=\tau_{\gamma} \quad Y_{I}\left(\gamma \tau_{\gamma}\right)=Y_{I}\left(\tau_{\gamma}\right)
$$

Eigenvalue equation gives alignments directly

$$
\rho_{I}(\gamma) Y_{I}\left(\tau_{\gamma}\right)=\left(c \tau_{\gamma}+d\right)^{-2 k} Y_{I}\left(\tau_{\gamma}\right)
$$

## Example

Eigenvalue equation

$$
\begin{gathered}
\rho(S) Y\left(\tau_{S}\right)=Y\left(\tau_{S}\right) \\
\frac{1}{3}\left(\begin{array}{ccc}
-1 & 2 & 2 \\
2 & -1 & 2 \\
2 & 2 & -1
\end{array}\right)
\end{gathered}
$$

Eigenvector

## Level 3 fixed points and alignments <br> G.J.Ding, S.F.K., X.G.Liu and J.N.Lu, 1910.03460

| The alignments of triplet modular forms $Y_{\mathbf{3 , 3}}\left(\gamma \tau_{S}\right)$ of level 3 up to weight 6 |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathbf{A}_{\mathbf{4}} \gamma$ | $\gamma \tau_{S}$ | $Y_{\mathbf{3}}^{(2)}\left(\gamma \tau_{S}\right), Y_{3, I}^{(6)}\left(\gamma \tau_{S}\right)$ | $Y_{\mathbf{3}}^{(4)}\left(\gamma \tau_{S}\right)$ | $Y_{\mathbf{3}, I I}^{(6)}\left(\gamma \tau_{S}\right)$ |  |
| $\{1, S\}$ | $i$ | $(1,1-\sqrt{3}, \sqrt{3}-2)$ | $(1,1,1)$ | $(1,-2-\sqrt{3}, 1+\sqrt{3})$ |  |
| $\{T, T S\}$ | $1+i$ | $\left(1,(1-\sqrt{3}) \omega,(\sqrt{3}-2) \omega^{2}\right)$ | $\left(1, \omega, \omega^{2}\right)$ | $\left(1,(-2-\sqrt{3}) \omega,(1+\sqrt{3}) \omega^{2}\right)$ |  |
| $\{S T, S T S\}$ | $\frac{-1+i}{2}$ | $\left(1,(1+\sqrt{3}) \omega,(-2-\sqrt{3}) \omega^{2}\right)$ | $\left(1, \omega, \omega^{2}\right)$ | $\left(1,(\sqrt{3}-2) \omega,(1-\sqrt{3}) \omega^{2}\right)$ |  |
| $\left\{T^{2}, T^{2} S\right\}$ | $2+i$ | $\left(1,(1-\sqrt{3}) \omega^{2},(-2+\sqrt{3}) \omega\right)$ | $\left(1, \omega^{2}, \omega\right)$ | $\left(1,(-2-\sqrt{3}) \omega^{2},(1+\sqrt{3}) \omega\right)$ |  |
| $\left\{S T^{2}, S T^{2} S\right\}$ | $\frac{-2+i}{5}$ | $\left(1,(1+\sqrt{3}) \omega^{2},(-2-\sqrt{3}) \omega\right)$ | $\left(1, \omega^{2}, \omega\right)$ | $\left(1,(\sqrt{3}-2) \omega^{2},(1-\sqrt{3}) \omega\right)$ |  |
| $\left\{T^{2} S T, T S T^{2}\right\}$ | $\frac{3+i}{2}$ | $(1,1+\sqrt{3},-2-\sqrt{3})$ | $(1,1,1)$ | $(1, \sqrt{3}-2,1-\sqrt{3})$ |  |

The alignments of triplet modular forms $Y_{\mathbf{3 , 3 ^ { \prime }}}\left(\gamma \tau_{S T}\right)$ of level 3 up to weight 6

| The alignments of triplet modular forms $Y_{\mathbf{3}, \mathbf{3}^{\prime}}\left(\gamma \tau_{S T}\right)$ of level 3 up to weight 6 |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| $\gamma$ | $\gamma \tau_{S T}$ | $Y_{\mathbf{3}}^{(2)}\left(\gamma \tau_{S T}\right), Y_{\mathbf{3}, I}^{(6)}\left(\gamma \tau_{S T}\right)$ | $Y_{\mathbf{3}}^{(4)}\left(\gamma \tau_{S T}\right)$ | $Y_{\mathbf{3}, I I}^{(6)}\left(\gamma \tau_{S T}\right)$ |
| $\left\{1, S T, T^{2} S\right\}$ | $\frac{-1+i \sqrt{3}}{2}$ | $\left(1, \omega, \frac{-1}{2} \omega^{2}\right)$ | $\left(1, \frac{-1}{2} \omega, \omega^{2}\right)$ | $\left(1,-2 \omega,-2 \omega^{2}\right)$ |
| $\left\{T, S T^{2} S, S\right\}$ | $\frac{1+i \sqrt{3}}{2}$ | $\left(1, \omega^{2},-\frac{1}{2} \omega\right)$ | $\left(1,-\frac{1}{2} \omega^{2}, \omega\right)$ | $\left(1,-2 \omega^{2},-2 \omega\right)$ |
| $\left\{T S, T^{2}, T^{2} S T\right\}$ | $2+\omega$ | $\left(1,1,-\frac{1}{2}\right)$ | $\left(1,-\frac{1}{2}, 1\right)$ | $(1,-2,-2)$ |
| $\left\{S T S, S T^{2}, T S T^{2}\right\}$ | $\frac{-3+i \sqrt{3}}{6}$ | $(0,0,1)$ | $(0,1,0)$ | $(1,0,0)$ |

The alignments of triplet modular forms $Y_{3,3^{\prime}}\left(\gamma \tau_{T S}\right)$ of level 3 up to weight 6

| $\gamma$ | $\gamma \tau_{T S}$ | $Y_{\mathbf{3}}^{(2)}\left(\gamma \tau_{T S}\right), Y_{3, I}^{(6)}\left(\gamma \tau_{T S}\right)$ | $Y_{\mathbf{3}}^{(4)}\left(\gamma \tau_{T S}\right)$ | $Y_{3, I I}\left(\gamma \tau_{T S}\right)$ |
| :---: | :---: | :---: | :---: | :---: |
| $\left\{1, T S, S T^{2}\right\}$ | $\frac{1+i \sqrt{3}}{2}$ | $\left(1, \omega^{2},-\frac{1}{2} \omega\right)$ | $\left(1,-\frac{1}{2} \omega^{2}, \omega\right)$ | $\left(1,-2 \omega^{2},-2 \omega\right)$ |
| $\left\{T, T^{2} S, T S T^{2}\right\}$ | $\frac{3+i \sqrt{3}}{2}$ | $\left(1,1,-\frac{1}{2}\right)$ | $\left(1,-\frac{1}{2}, 1\right)$ | $(1,-2,-2)$ |
| $\left\{S T, S T^{2} S, T^{2} S T\right\}$ | $\frac{(-1)^{5 / 6}}{\sqrt{3}}$ | $(0,0,1)$ | $(0,1,0)$ | $(1,0,0)$ |
| $\left\{S T S, T^{2}, S\right\}$ | $2+\omega$ | $\left(1, \omega, \frac{-1}{2} \omega^{2}\right)$ | $\left(1, \frac{-1}{2} \omega, \omega^{2}\right)$ | $\left(1,-2 \omega,-2 \omega^{2}\right)$ |

The alignments of triplet modular forms $Y_{\mathbf{3}, \mathbf{3}^{\prime}}\left(\gamma \tau_{T}\right)$ of level 3 up to weight 6

| $\gamma$ | $\gamma \tau_{T}$ | $Y_{\mathbf{3}}^{(2)}\left(\gamma \tau_{T}\right), Y_{\mathbf{3}, I}^{(6)}\left(\gamma \tau_{T}\right), Y_{\mathbf{3}}^{(4)}\left(\gamma \tau_{T}\right)$ | $Y_{\mathbf{3}, I I}^{(6)}\left(\gamma \tau_{T}\right)$ |
| :---: | :---: | :---: | :---: |
| $\left\{1, T, T^{2}\right\}$ | $i \infty$ | $(1,0,0)$ | $(0,1,0)$ |
| $\left\{S T, S T^{2}, S\right\}$ | 0 | $(1,-2,-2)$ | $\left(1,-\frac{1}{2}, 1\right)$ |
| $\left\{T S, S T^{2} S, T S T^{2}\right\}$ | 1 | $\left(1,-2 \omega,-2 \omega^{2}\right)$ | $\left(1,-\frac{1}{2} \omega, \omega^{2}\right)$ |
| $\left\{S T S, T^{2} S, T^{2} S T\right\}$ | -1 | $\left(1,-2 \omega^{2},-2 \omega\right)$ | $\left(1,-\frac{1}{2} \omega^{2}, \omega\right)$ |

# Level 4 fixed points and alignments 

G.J.Ding, S.F.K., X.G.Liu and J.N.Lu, 1910.03460


| The alignments of triplet modular forms $Y_{3,3^{\prime}}\left(\gamma \tau_{T}\right)$ of level 4 up to weight 6 |  |  |  |
| :---: | :---: | :---: | :---: |
| $\gamma$ | $\gamma \tau_{T}$ | $Y_{\mathbf{3}}^{(2)}\left(\gamma \tau_{T}\right), Y_{\mathbf{3}}^{(4)}\left(\gamma \tau_{T}\right), Y_{\mathbf{3} \mathbf{I}}^{(6)}\left(\gamma \tau_{T}\right), Y_{\mathbf{3}, \mathbf{I}}^{(6)}\left(\gamma \tau_{T}\right)$ | $Y_{3^{\prime}}^{(4)}\left(\gamma \tau_{T}\right), Y_{3^{\prime}}^{(6)}\left(\gamma \tau_{T}\right)$ |
| $\left\{1, T, T^{2}, T^{3}\right\}$ | $i \infty$ | $\left(1, \omega^{2}, \omega\right)$ | $(0,0,0)$ |
| $\left\{S T^{2} S, S T^{2} S T,\left(S T^{2}\right)^{2}, T S T^{2} S\right\}$ | $-\frac{1}{2}$ |  |  |
| $\frac{\left\{S T,(T S)^{2}, S, S T^{2}\right\}}{\left\{T^{2} S T, T^{2} S T^{3} T^{2} S T^{2} T^{2} S\right\}}$ | 0 | $\left(1, \omega, \omega^{2}\right)$ |  |
| $\left\{T^{2} S T, T^{2} S T^{3}, T^{2} S T^{2}, T^{2} S\right\}$ | 2 |  |  |
| $\left\{T S, T S T^{2}, T S T^{3}, T S T\right\}$ | -1 | $(1,1,1)$ |  |

## Example with $\operatorname{SU}(5)$ GUT: Level $N=4 \sim S_{4}$

weighton

| Field | $T_{3}$ | $T=\left(T_{2}, T_{1}\right)^{T}$ | $F$ | $N_{\mathrm{a}}$ | $N_{\mathrm{s}}$ | $H_{5}$ | $H_{\overline{5}}$ | $H_{\overline{45}}$ | $\phi$ | $\chi^{0}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $S U(5)$ | $\mathbf{1 0}$ | $\mathbf{1 0}$ | $\overline{\mathbf{5}}$ | $\mathbf{1}$ | $\mathbf{1}$ | $\mathbf{5}$ | $\overline{\mathbf{5}}$ | $\overline{\mathbf{4 5}}$ | $\mathbf{1}$ | $\mathbf{1}$ |
| $S_{4}$ | $\mathbf{1}$ | $\mathbf{2}$ | $\mathbf{3}$ | $\mathbf{1}$ | $\mathbf{1}^{\prime}$ | $\mathbf{1}^{\prime}$ | $\mathbf{1}$ | $\mathbf{1}$ | $\mathbf{1}$ | $\mathbf{1}$ |
| $k_{I}$ | 4 | 1 | 3 | 4 | -1 | -2 | 1 | 1 | 1 | 0 |

$\alpha_{u} \tilde{\phi}^{4} Y_{2}^{(4)}(T T)_{2} H_{5}+\beta_{u} \tilde{\phi}^{2} Y_{2}^{(2)}(T T)_{2} H_{5}+\gamma_{u} Y_{1^{\prime}}^{(6)} T_{3} T_{3} H_{5}+\epsilon_{u} \tilde{\phi} T_{3}\left(T Y_{2}^{(4)}\right)_{1^{\prime}} H_{5}$

$$
\mathcal{Y}_{\mathrm{GUT}}^{u} \approx\left(\begin{array}{ccc}
\alpha_{u} \tilde{\phi}^{4} & 0 & 0 \\
0 & \beta_{u} \tilde{\phi}^{2} & \epsilon_{u} \tilde{\phi} \\
0 & \epsilon_{u} \tilde{\phi} & \gamma_{u}
\end{array}\right) \quad \begin{aligned}
& \mathcal{Y}_{\mathrm{GUT}}^{d}=\left(\mathcal{Y}_{\overline{5}}+\mathcal{Y}_{\overline{45}}\right)^{T} \\
& \mathcal{Y}_{\mathrm{GUT}}^{e}=\mathcal{Y}_{\overline{5}}-3 \mathcal{Y}_{\overline{45}}
\end{aligned}
$$



Littlest Modular Seesaw from fixed point alignments

$$
\begin{gathered}
Y_{\mathbf{3}^{\prime}}^{(6)} \propto\left(\begin{array}{c}
0 \\
1 \\
-1
\end{array}\right), \quad Y_{\mathbf{3}}^{(2)} \propto\left(\begin{array}{c}
1 \\
1+\sqrt{6} \\
1-\sqrt{6}
\end{array}\right) \\
m_{\nu}=m_{a}\left(\begin{array}{ccc}
0 & 0 & m_{s} \\
0 & 1 & -1 \\
0 & -1 & 1
\end{array}\right)+m_{a_{s}} e^{i \eta}\left(\begin{array}{ccc}
1 & 1-\sqrt{6} & 1+\sqrt{6} \\
1-\sqrt{6} & 7-2 \sqrt{6} & -5 \\
1+\sqrt{6} & -5 & 7+2 \sqrt{6}
\end{array}\right)^{0.05}= \\
Y_{\mathbf{3}^{\prime}}^{(6)} Y_{\mathbf{3}^{\prime}}^{(6)} \\
Y_{\mathbf{3}}^{(2)} Y_{\mathbf{3}}^{(2)^{T}}
\end{gathered}
$$

(a)
S.F.K. and Y.L. Zhou,1908.02770

## Example with two groups: Level $\mathrm{N}=4 \sim \mathrm{~S}_{4}$

$S_{4}^{l}, \tau_{l}$
$\frac{\langle\Phi\rangle}{\square}$ C.f.

Family


Use S4 basis: $\quad T=S_{\tau} T_{\tau}, \quad S=T_{\tau}^{2}, \quad U=T_{\tau} S_{\tau} T_{\tau}^{2} S_{\tau}$

Fixed points:
$\left\langle\tau_{\nu}\right\rangle=\tau_{S U}=-\frac{1}{2}+\frac{i}{2}$
$\left\langle\tau_{l}\right\rangle=\tau_{T}=-\frac{1}{2}+i \frac{\sqrt{3}}{2}$
$Y\left(\tau_{S U}\right) \propto\left(\begin{array}{c}2 \\ -1 \\ -1\end{array}\right)$
$U_{\mathrm{TM}_{1}}=\left(\begin{array}{ccc}\frac{2}{\sqrt{6}} & - & - \\ -\frac{1}{\sqrt{6}} & - & - \\ -\frac{1}{\sqrt{6}} & - & -\end{array}\right)$

## Summary

- Flavour problem motivates family/flavour symmetry
- U(1) with FN for hierarchies and small mixing
- Neutrino mass and mixing motivates non-Abelian
- TBM, TM1/TM2, Littlest Seesaw...enforced by $S_{4}$ and flavon alignments...gauged or modular origin


## - Large literature on bottom-up modular models

- Weightons for charged fermion hierarchies
- Stabilizers/fixed points for Yukawa alignments
- SU(5) GUT with $S_{4}$ and Littlest Modular Seesaw
- Twin modular $S_{4}$ symmetries for TM1

