

Improved constraints for axion-like particles from 3-photon events at e^+e^- colliders

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October 30, 2023

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EINN2023



QCD axion

- The **absence of CP violating effects in QCD**, leading to the unnaturally small value of the vacuum angle $\theta \lesssim 10^{-11}$, is known as the **strong CP problem** [Phys. Rev. Lett. 124, 081803 \(2020\)](#)
- Spontaneous breaking of Peccei-Quinn symmetry gives rise to a **new pseudoscalar particle** with a linear relation between mass m_a and coupling $g_{a\gamma\gamma}$ which **solves this problem** [Phys. Rev. Lett. 40, 223 \(1978\)](#); [Phys. Rev. Lett. 40, 279 \(1978\)](#)
- **Axion-photon coupling** is widely discussed. Experimental bounds in **MeV-GeV range** mostly from e^+e^- colliders [Phys. Rev. Lett. 125, 161806 \(2020\)](#); [arXiv:2211.12699v3](#); and others

The results are upper limits on $|d_n|$.

VALUE ($10^{-25} e \text{ cm}$)	CL%	DOCUMENT ID
< 0.18	90	¹ ABEL 2020
• • We do not use the following data for averages		
< 0.22	95	² SAHOO 2017
< 0.16	95	GRANER 2016
< 0.30	90	³ PENDLEBURY 2015
< 0.55	90	SEREBROV 2015
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< 0.29	90	⁵ BAKER 2006
< 0.63	90	⁶ HARRIS 1999
< 0.97	90	ALTAREV 1996
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PDG data on neutron electric dipole moment, one of the most precisely measured θ -induced effects ($\vec{d} \sim \theta \cdot 10^{-16} e \text{ cm}$)

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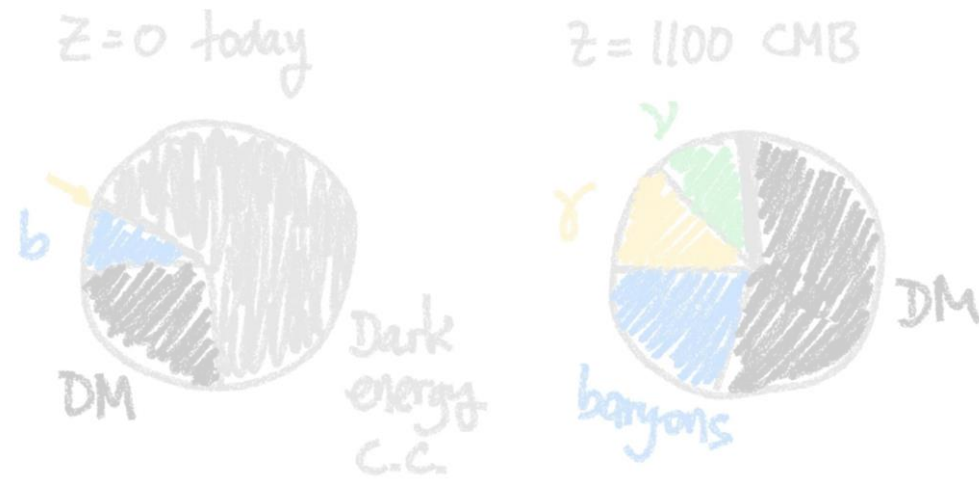
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Axion-like particles



„Cosmological pie“ (by Martin Schmaltz)

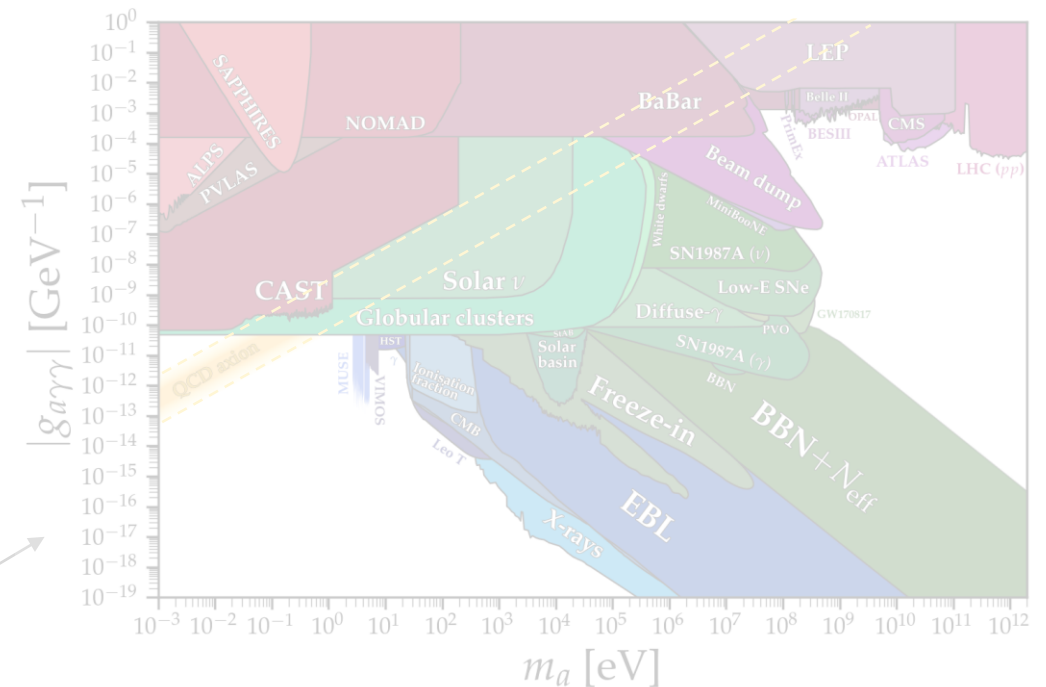
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- Improve upon previous analyses by also including the ALP coupling to electrons
- Significantly extend the ALP search range and impose more stringent restrictions on their couplings

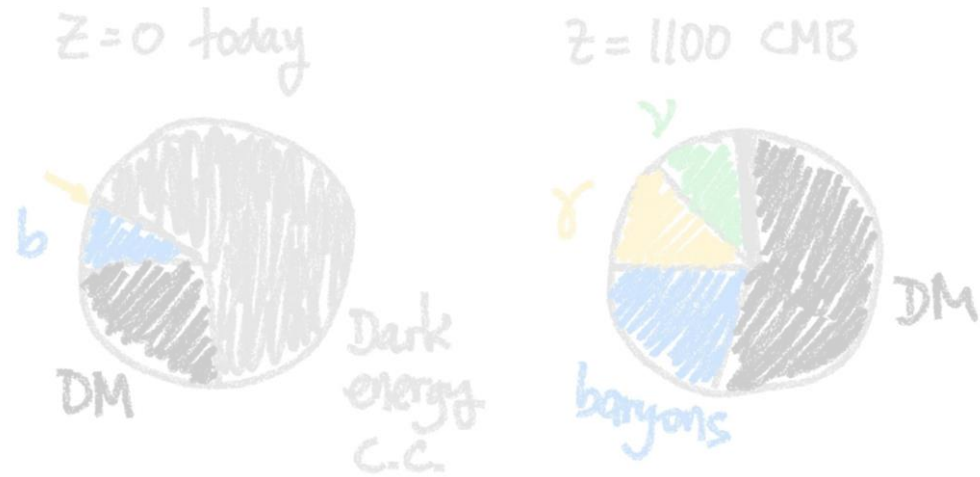
Existing bounds from <https://cajohare.github.io/AxionLimits>

Axion-like particle (ALP):

- Is an **axion generalization**, which appears in many **extensions of the SM**
- Has mass-coupling relation unconstrained with linear relation
- A promising dark matter candidate



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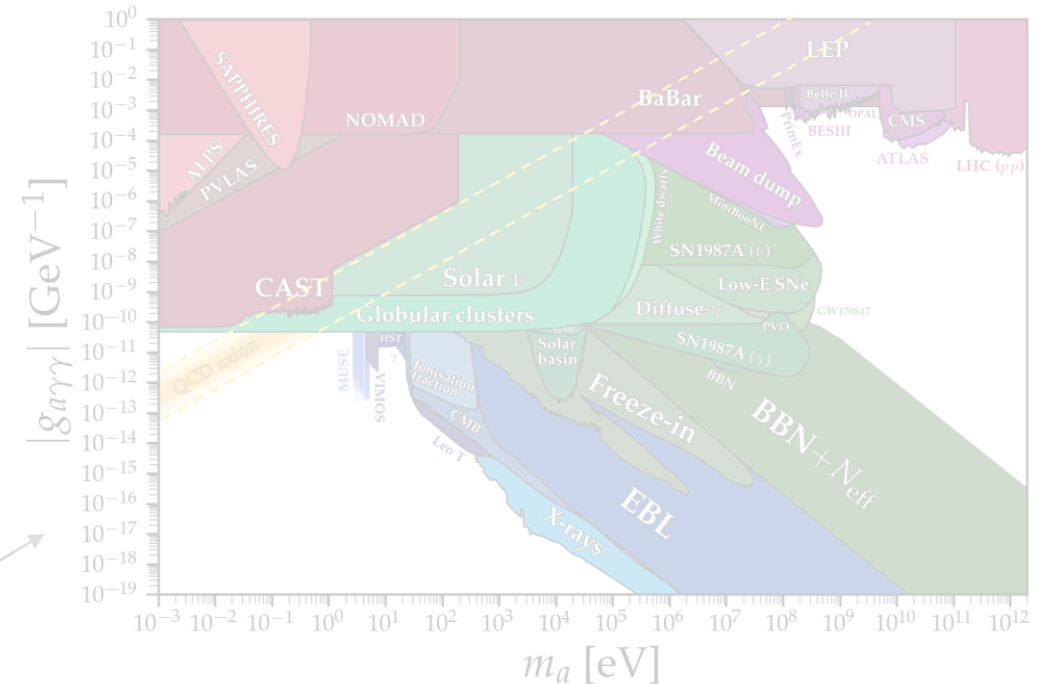
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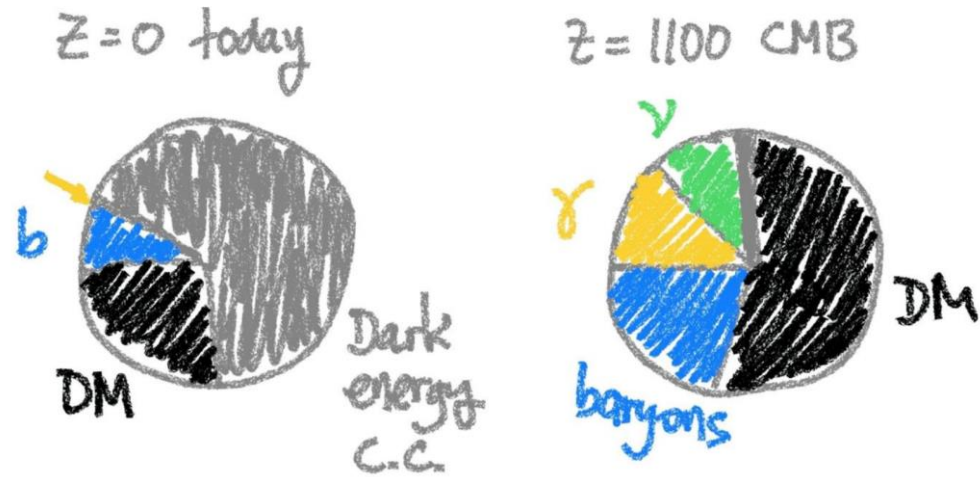
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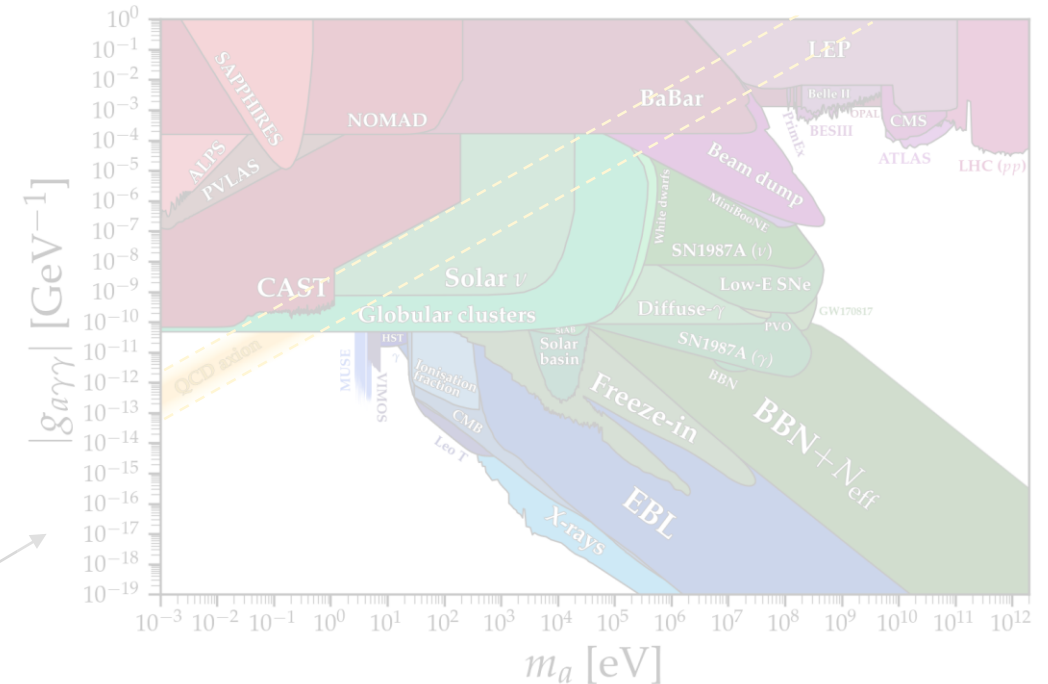
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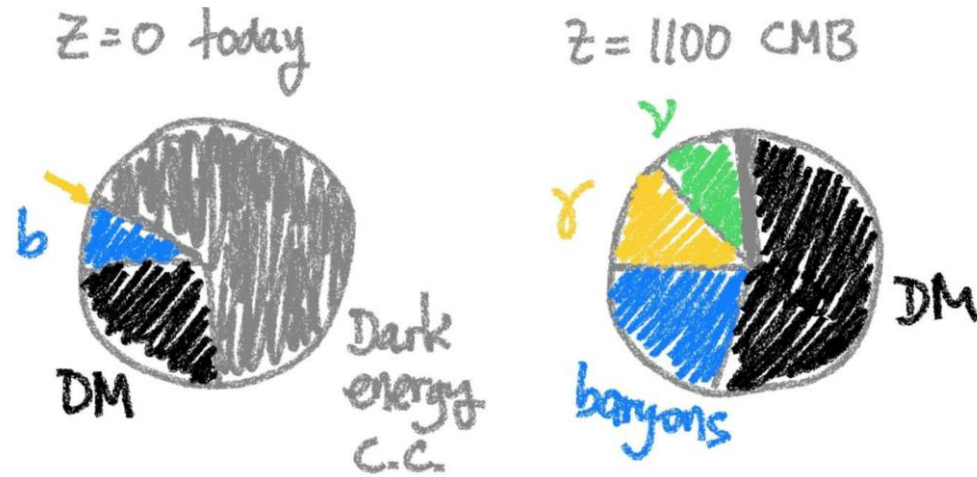
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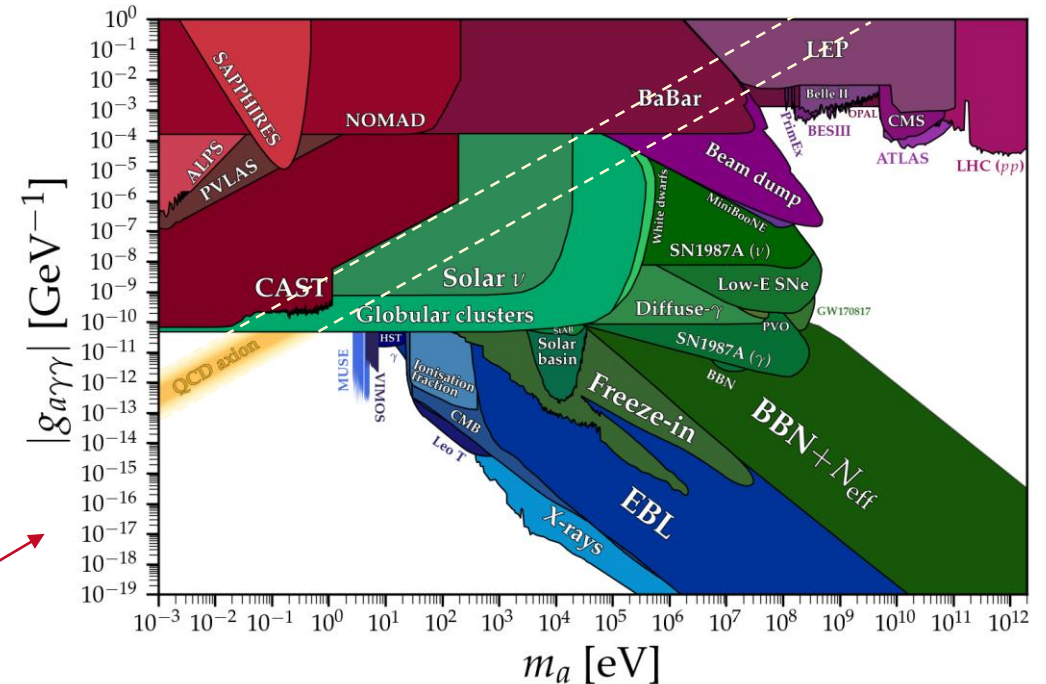
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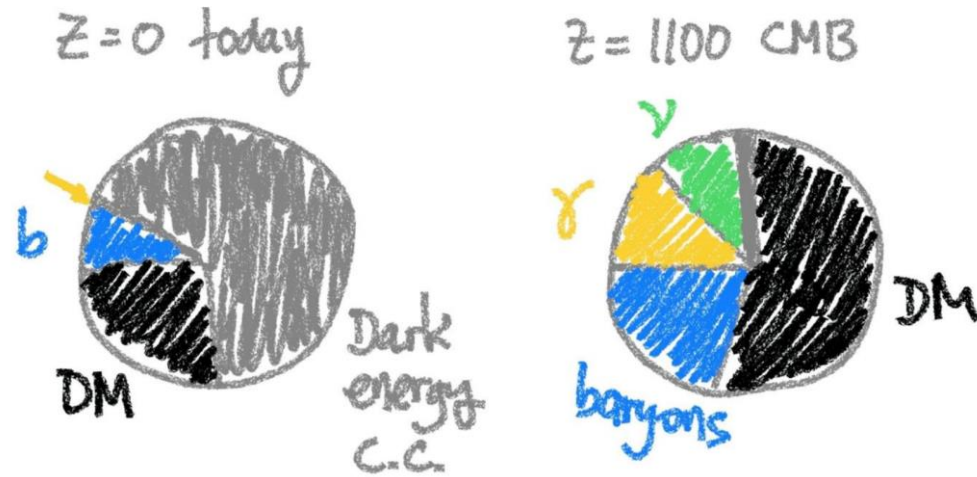
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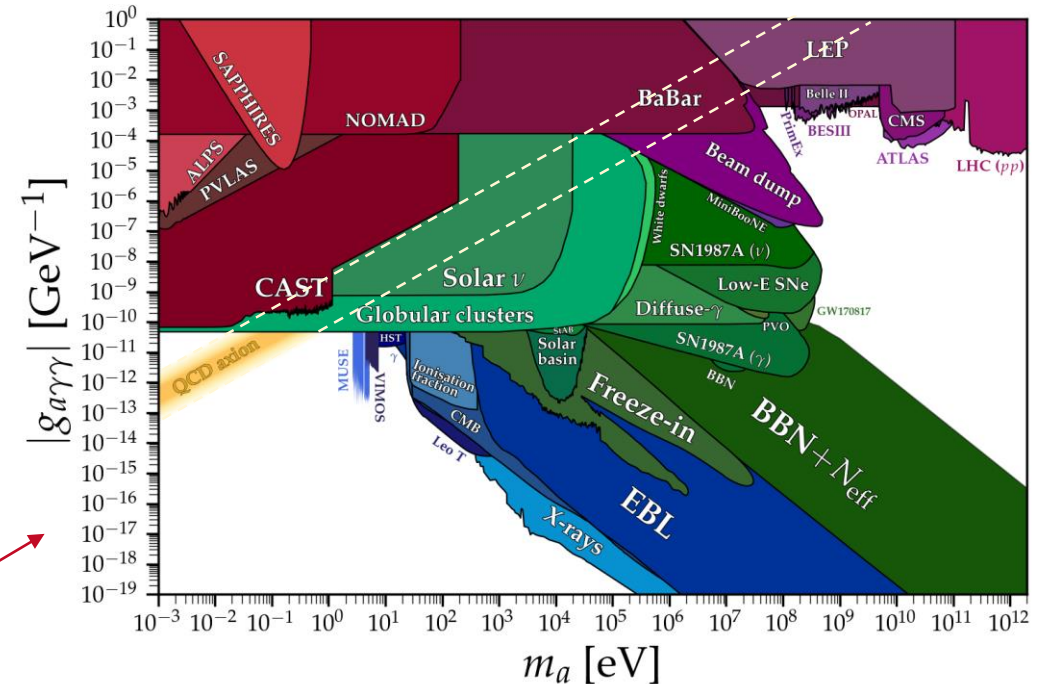
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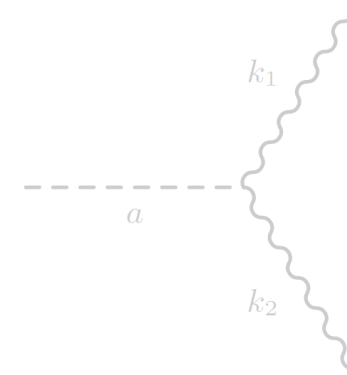
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- e/a - electron/ALP field of the mass $m_{e/a}$
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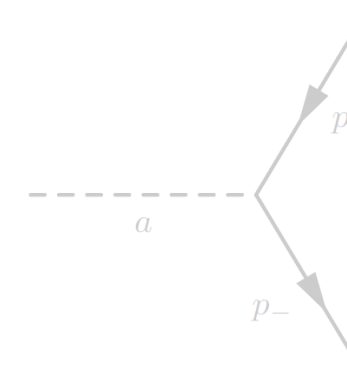
In this work we assume that ALPs in the MeV to GeV mass range **couple predominantly to photons and electrons**, i.e., no enhanced impact to $(g - 2)_\mu$

JHEP, 07:092, 2018



$$\Rightarrow \Gamma_{a\gamma\gamma} = \frac{g_{a\gamma\gamma}^2 m_a^3}{64\pi}$$

ALP-photon interaction was investigated in details by M. J. Dolan and others in: JHEP, 12:094, 2017



$$\Rightarrow \Gamma_{aee} = \frac{g_{aee}^2 m_a}{8\pi}$$

Dominant channel for lower-mass ALP!

In the absence of interaction with other fields, the total ALP decay width is assumed to consist of these two contributions

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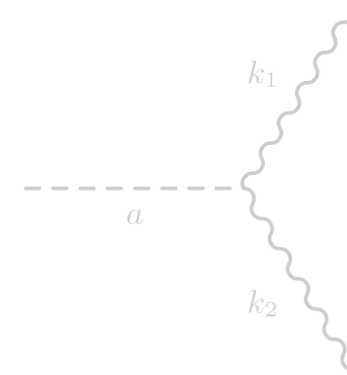
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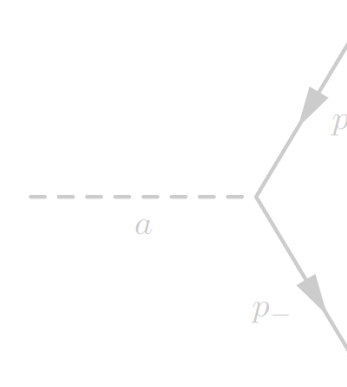
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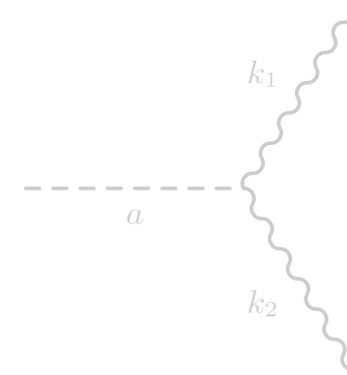
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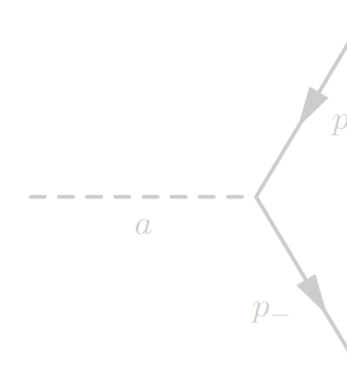
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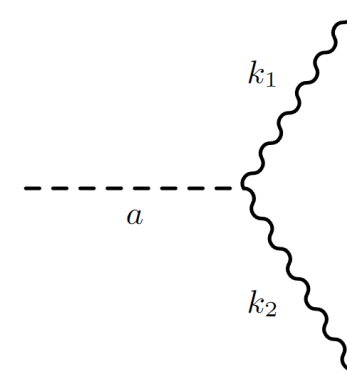
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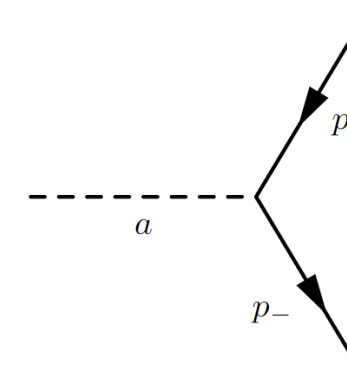
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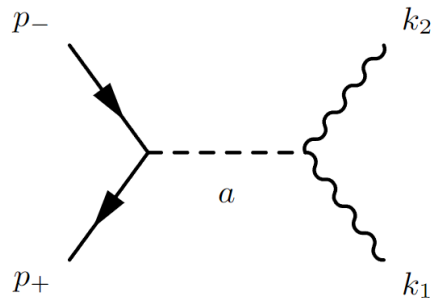
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2- and 3-photon events

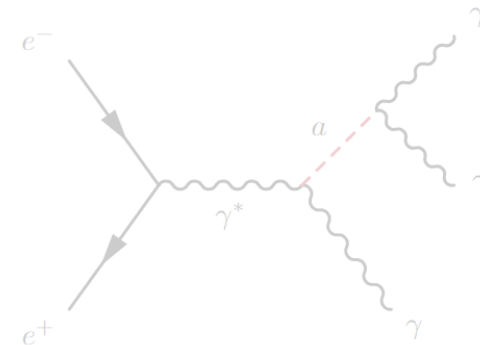
ALP contributes to 2-photon events through the single Feynman diagram:



However, it is resonant and thus has sensitivity only around $m_a^2 \approx s$

For a fixed collider energy, the process does not provide constraints in a broad parameter space

Thus, we proceed to 3-photon events



$$= i \frac{H_{e^+e^- \rightarrow \gamma^* \rightarrow a\gamma}(k_1)}{K_{23}^2 - m_a^2 + im_a\Gamma_a} \times M_{a \rightarrow \gamma\gamma}(k_2, k_3) + \text{crossed}$$

Where the amplitude of $e^+e^- \rightarrow a\gamma_i$ is given by:

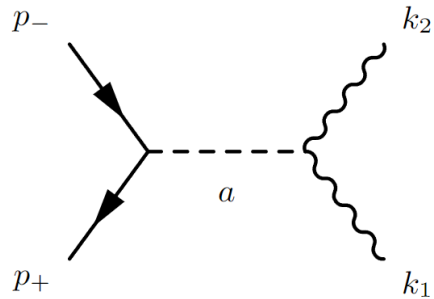
$$H_{e^+e^- \rightarrow \gamma^* \rightarrow a\gamma}(k_i) = -ieg_{a\gamma\gamma} \varepsilon_{\alpha\beta\mu\gamma} q^\alpha k_i^\beta \epsilon^\gamma(k_i, \lambda_i) \times \frac{\bar{v}(p_+, s_+) \gamma^\mu u(p_-, s_-)}{s}$$

And the amplitude of ALP decay to a photon pair:

$$M_{a \rightarrow \gamma\gamma}(k_1, k_2) = -ig_{a\gamma\gamma} k_{1,\kappa} k_{2,\beta} \varepsilon^{\kappa\beta\mu\nu} \times \epsilon_\mu^*(k_1, \lambda_1) \epsilon_\nu^*(k_2, \lambda_2)$$

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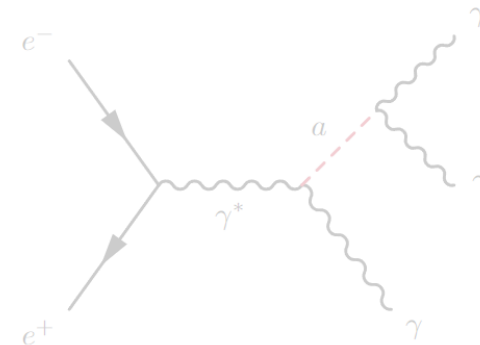
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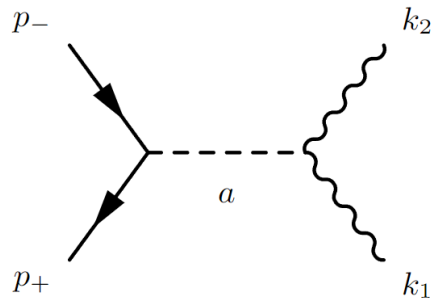
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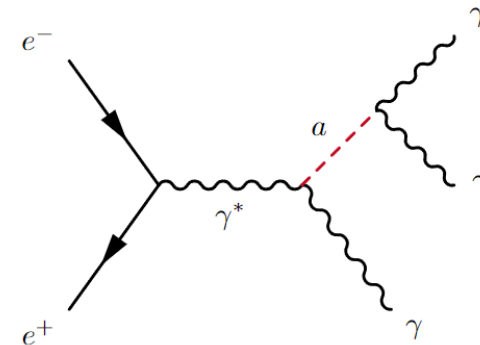
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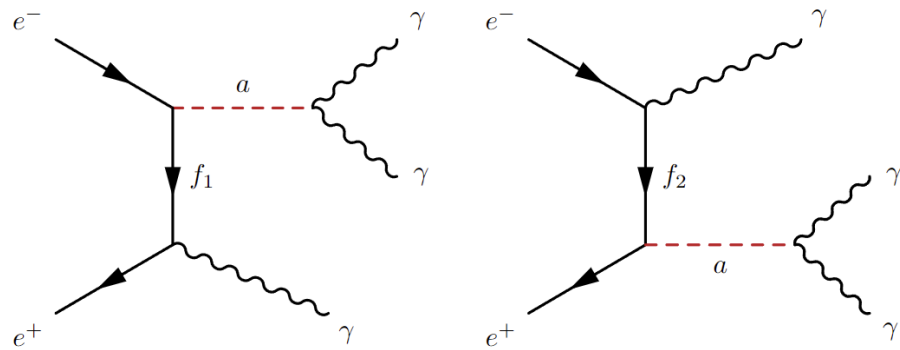
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The next contribution arises from **ALP-electron coupling**:



$$= i \frac{H_{e^+e^- \rightarrow a\gamma,1}(k_1) + H_{e^+e^- \rightarrow a\gamma,2}(k_1)}{K_{23}^2 - m_a^2 + im_a\Gamma_a} \times M_{a \rightarrow \gamma\gamma}(k_2, k_3) + \text{crossed},$$

where $H_{e^+e^- \rightarrow a\gamma,j}(k_i)$ denote amplitudes for the corresponding $2 \rightarrow 2$ process

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With the internal electron momenta:

$$l_i = k_i - p_+, \quad f_i = p_- - k_i$$

Note: no interference between this set of diagrams and diagrams from the previous slide

Background analysis

General assumption: ALPs are very long-lived particles \Rightarrow decay width $\Gamma_a \ll$ **exp. resolution** (interference is suppressed):

$$\frac{1}{(p_a^2 - m_a^2)^2 + (m_a \Gamma_a)^2} \rightarrow \frac{\pi}{m_a \Gamma_a} \delta(p_a^2 - m_a^2)$$

Due to the resonant behavior of the amplitude, one photon is always emitted with a fixed energy:

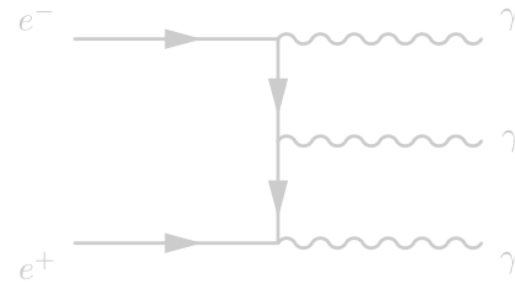
$$\omega = \frac{4E^2 - m_a^2}{4E}$$

Signal detection strategy: a narrow peak in the photon energies distribution, or a narrow peak in the squared mass distribution $m_{\gamma\gamma}$

Total $e^+e^- \rightarrow \gamma\gamma\gamma$ cross section has three terms:

$$\sigma_{ALP+B} = \sigma_{ALP} + \sigma_B + \sigma_{int}$$

➤ The dominant background: QED 3-photon annihilation



Aimed sensitivity is then:

$$\frac{\sigma_{ALP}}{\sigma_{QED}} = \frac{N}{\sqrt{L \cdot \sigma_{QED}}}$$

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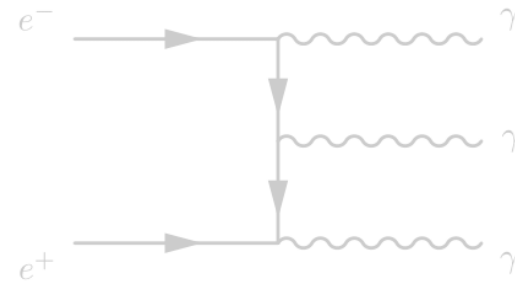
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Background analysis

General assumption: ALPs are very long-lived particles \Rightarrow decay width $\Gamma_a \ll$ **exp. resolution** (interference is suppressed):

$$\frac{1}{(p_a^2 - m_a^2)^2 + (m_a \Gamma_a)^2} \rightarrow \frac{\pi}{m_a \Gamma_a} \delta(p_a^2 - m_a^2)$$

Due to the resonant behavior of the amplitude, one photon is always emitted with a fixed energy:

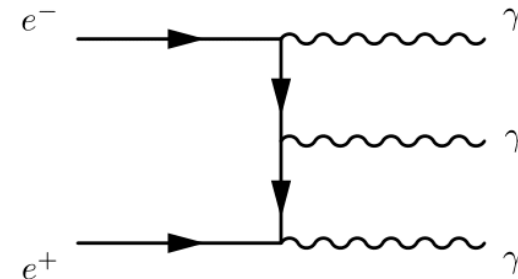
$$\omega = \frac{4E^2 - m_a^2}{4E}$$

Signal detection strategy: a narrow peak in the photon energies distribution, or a narrow peak in the squared mass distribution $m_{\gamma\gamma}$

Total $e^+e^- \rightarrow \gamma\gamma\gamma$ **cross section** has three terms:

$$\sigma_{ALP+B} = \sigma_{ALP} + \sigma_B + \sigma_{int}$$

➤ The dominant background: **QED 3-photon annihilation**



Aimed sensitivity is then:

$$\frac{\sigma_{ALP}}{\sigma_{QED}} = \frac{N}{\sqrt{L \cdot \sigma_{QED}}}$$

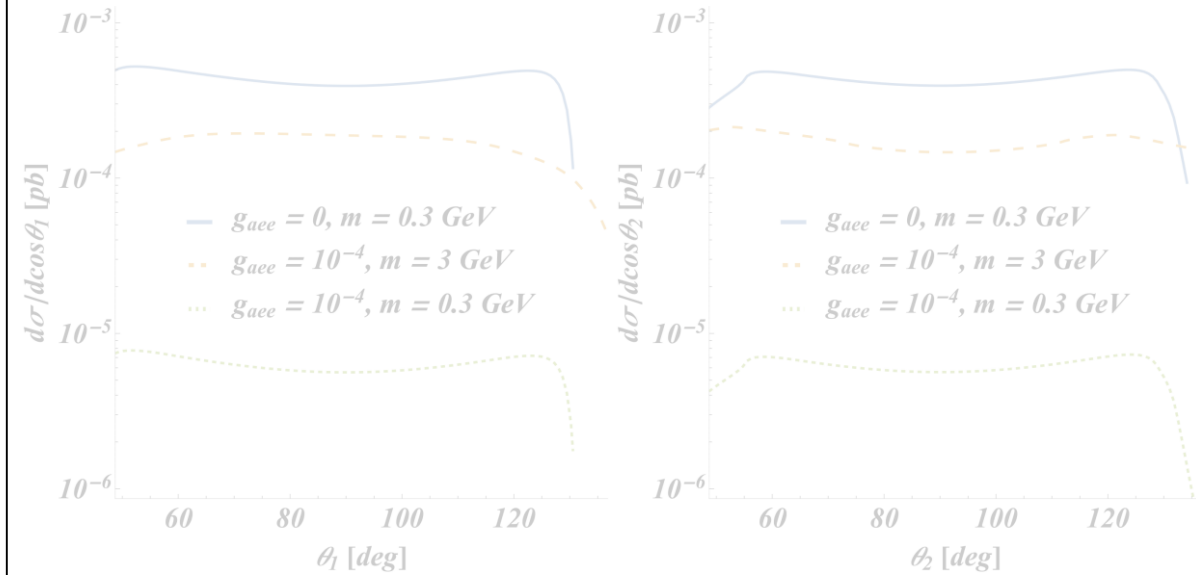
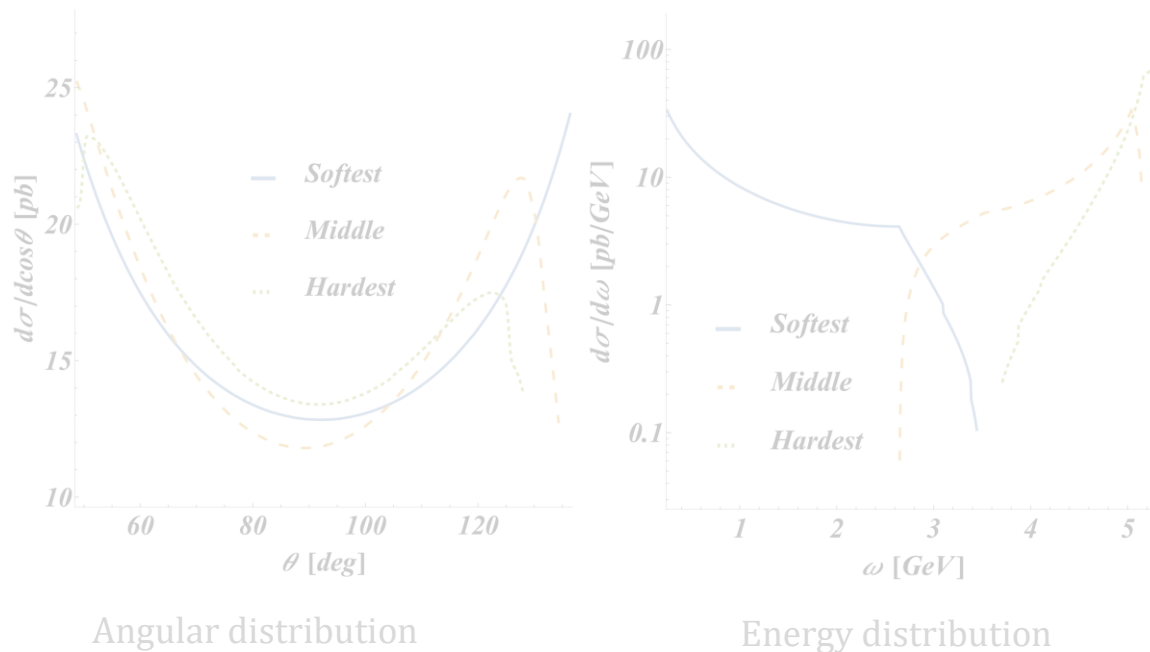
- N – the number of standard deviations that determines whether or not a fluctuation is considered as a signal ($N = 2 \Leftrightarrow 95\%$ c.l.)
- L – integrated luminosity

Distributions

- In the 2018 data run Belle II achieved an integrated luminosity of 445 pb^{-1} , which was used for the ALP searches

Phys. Rev. Lett. 125, 161806 (2020);

- LO QED 3-photon distributions for the softest, middle and hardest photons in this kinematics are shown below:



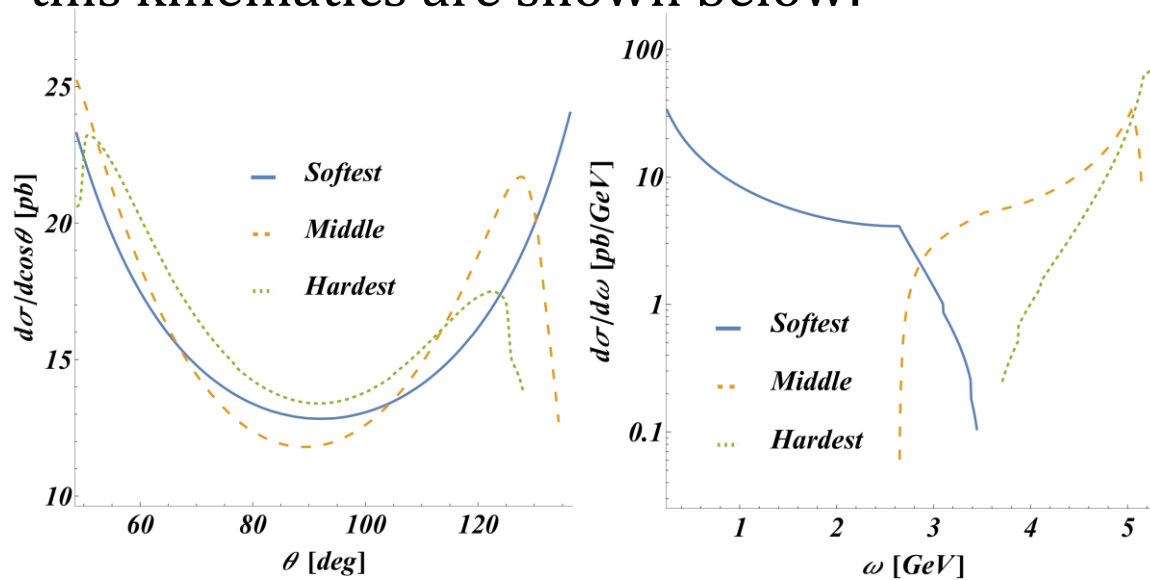
- In contrast to the ALP-related process, which exhibits a rather uniform angular distribution, the QED three-photon annihilation is characterized by an **enhanced angular distribution** in both the forward and backward directions.
- The presence of a distinct peak in the photon energy distribution would serve as an **indication of ALP creation**

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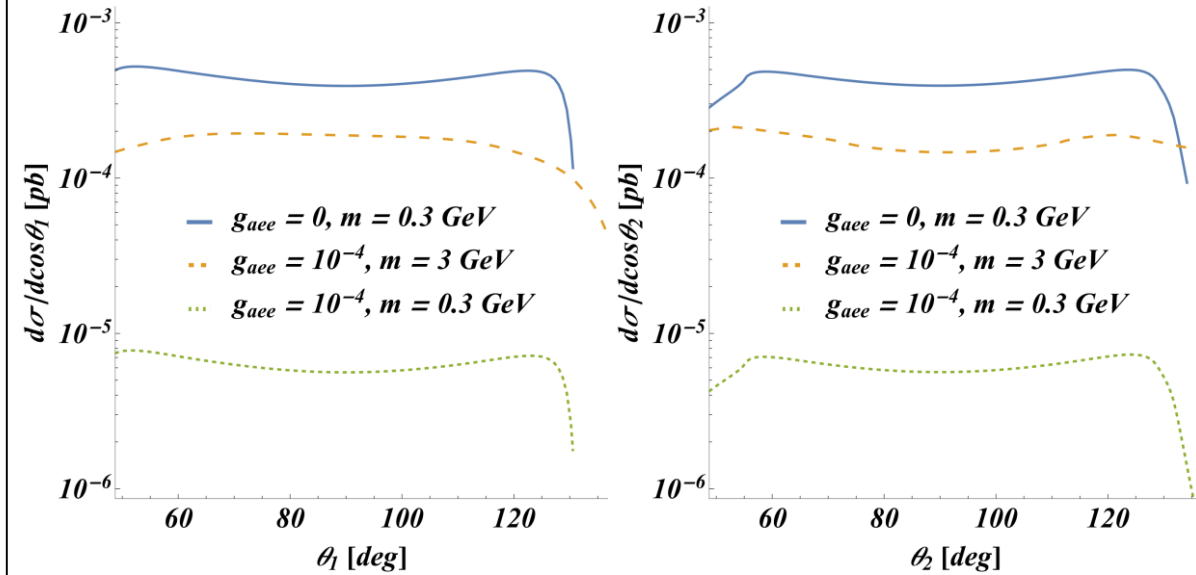
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Angular distribution

Energy distribution



Angular distributions of a photon which **accompanies** the ALP

Angular distributions of a photon **produced from ALP** decay

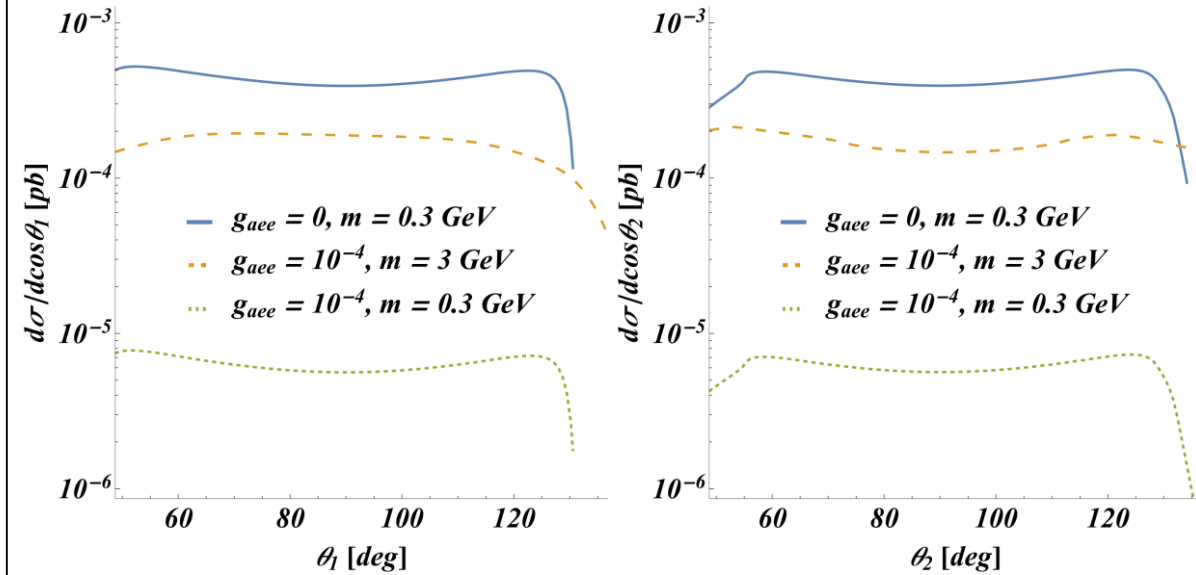
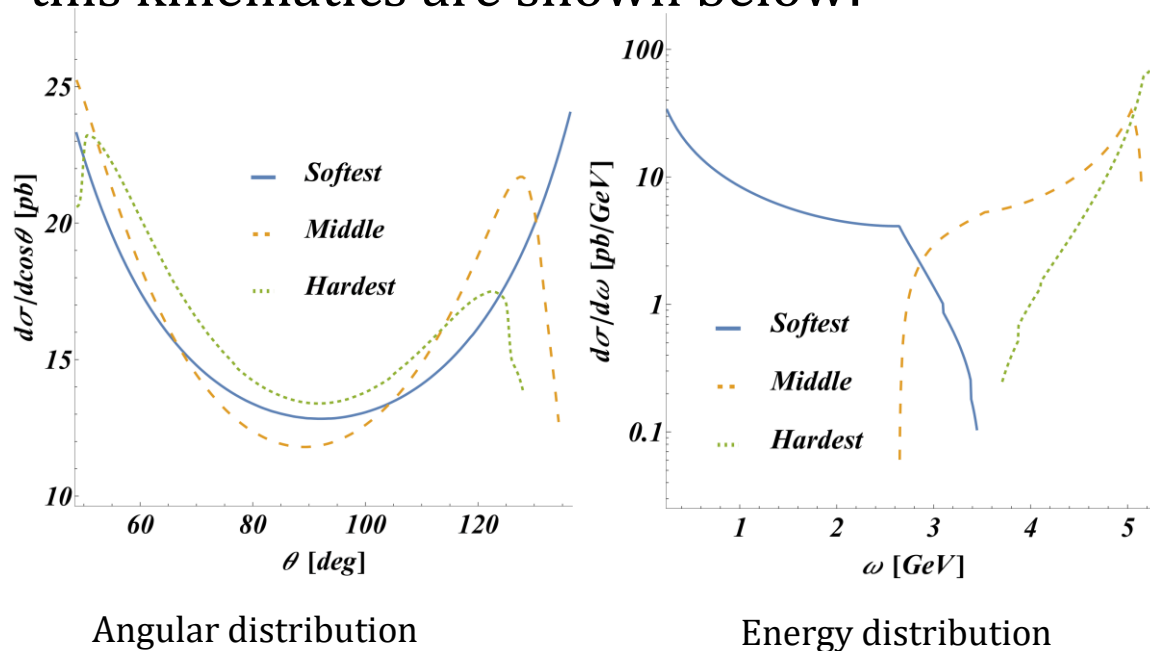
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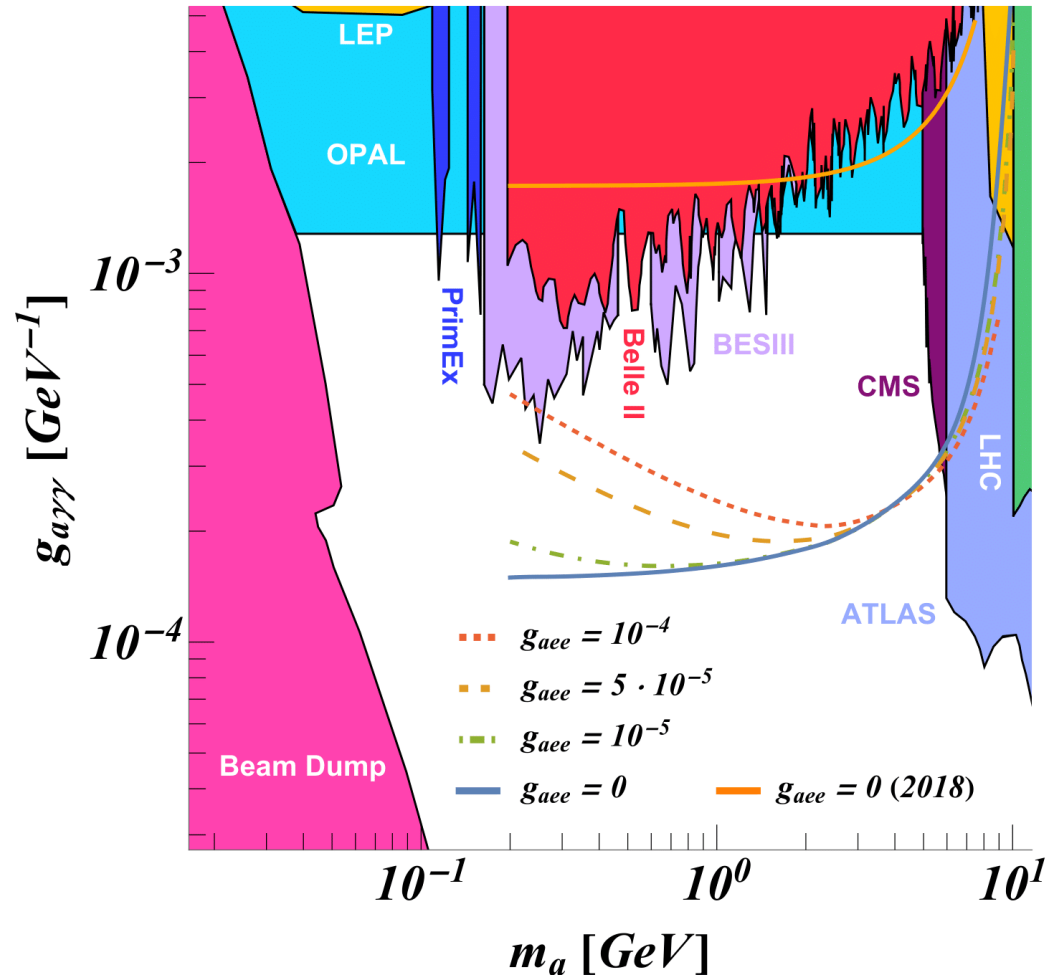
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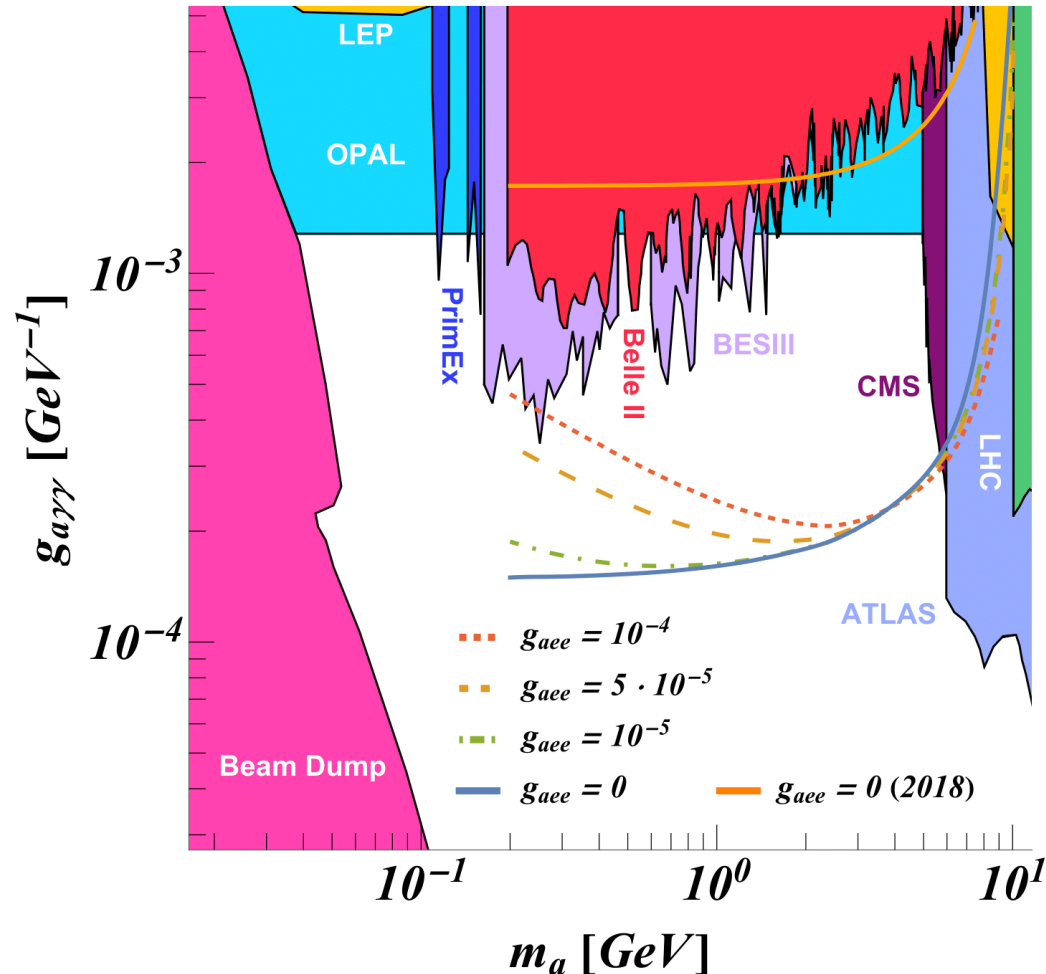
Results



Belle II constraints for $g_{a\gamma\gamma}$ based on the 2018 data set of 445 pb^{-1} and projected results for 50 ab^{-1} of integrated luminosity – improvement by an order of magnitude

- A good agreement in the higher ALP mass region ($m_a \gtrsim 2 \text{ GeV}$)
- In the lower mass region the invariant mass of a photon pair also becomes low (two photons travel in a very narrow cone with each other, oppositely to the third photon)
- This produces very asymmetric kinematics and the QED background becomes suppressed
- The assumption $g_{aee} = 0$ generally leads to an overestimated $g_{a\gamma\gamma}$ limit, which may be incorrect if the ALP has other decay channels besides the 2-photon mode

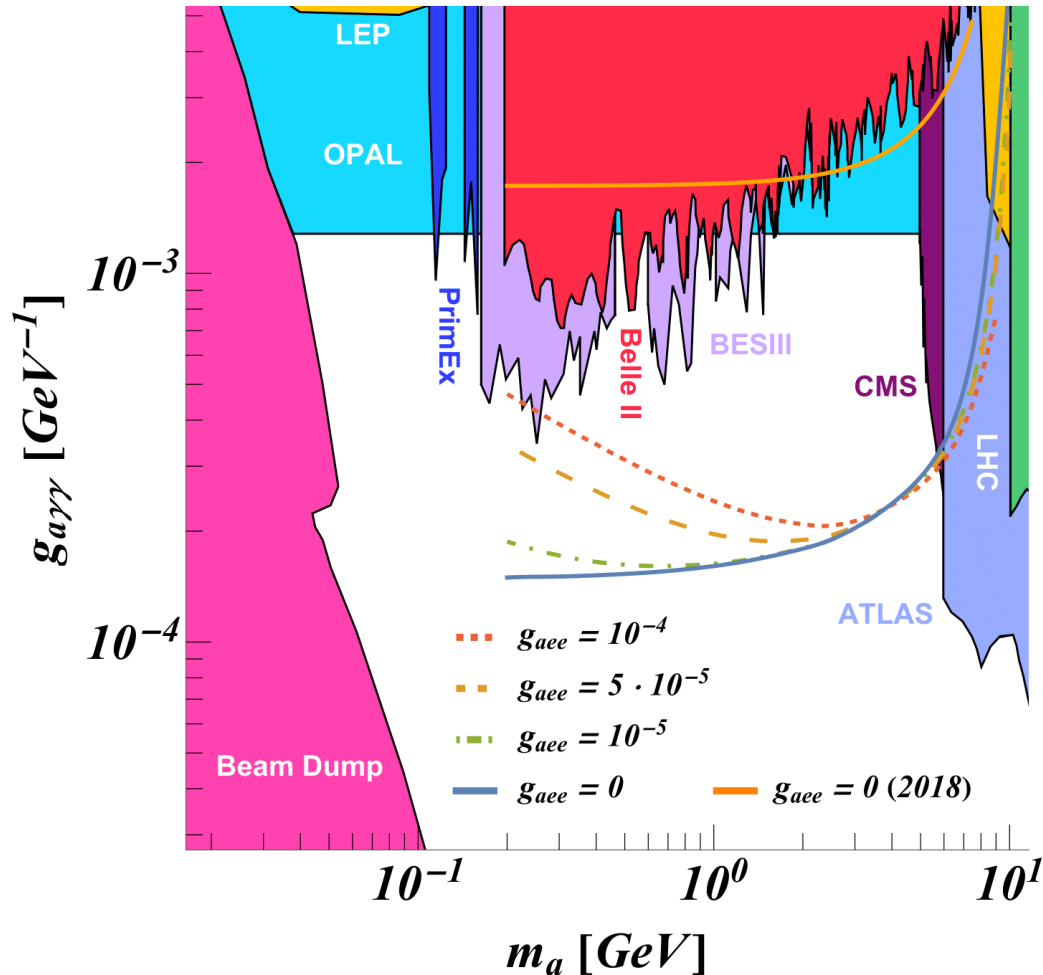
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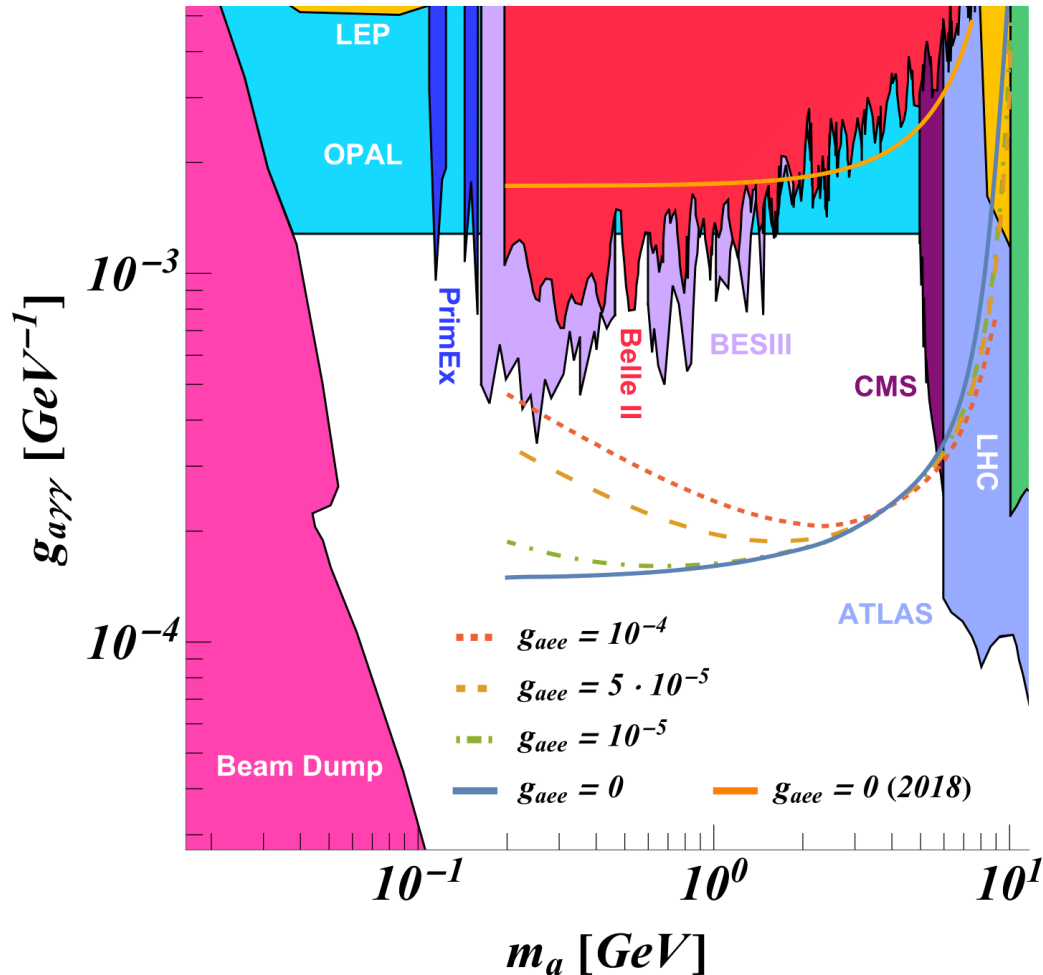
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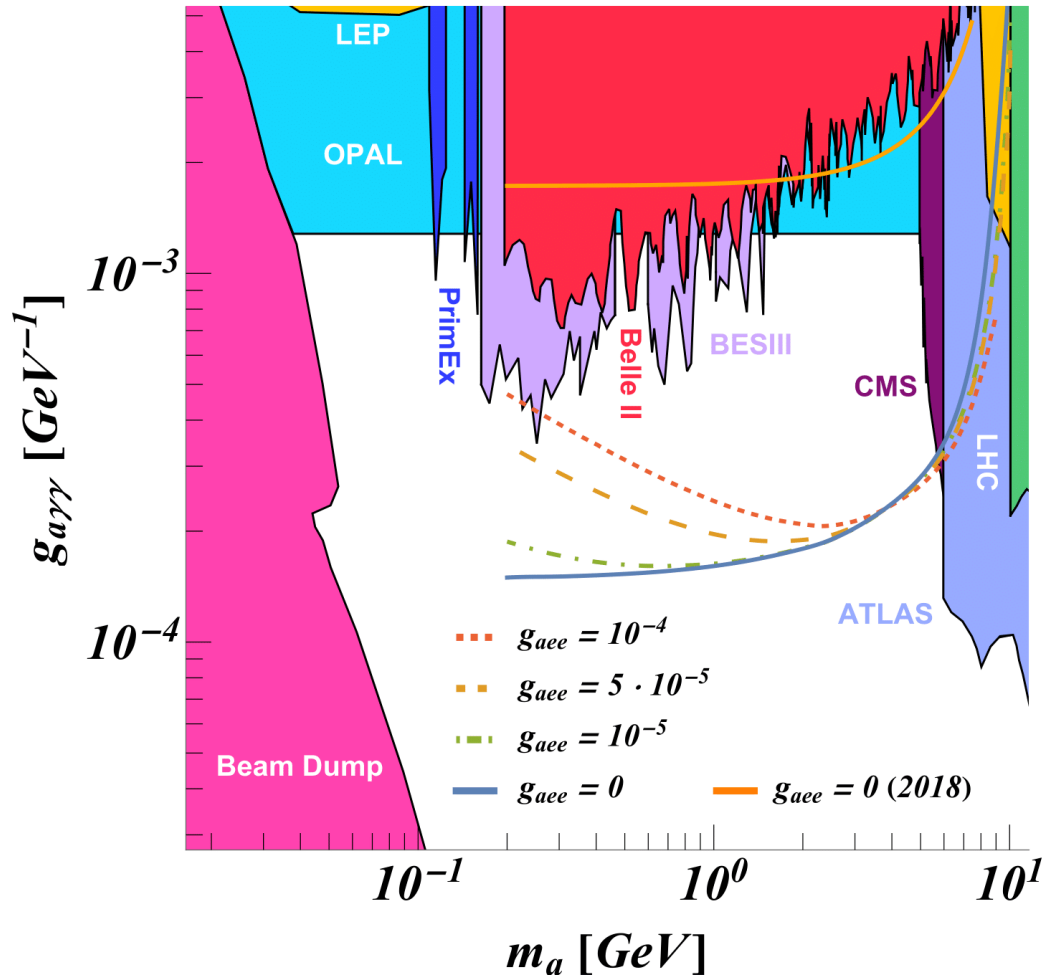
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Conclusion

The contributions of ALP states to $e^+e^- \rightarrow \gamma\gamma\gamma$ annihilation were calculated:

- **New constraints acquired** for ALPs in the MeV to GeV mass range, which can be effectively **applied** in e^+e^- colliders
- Tested on existing data for Belle II kinematics and projected for the forthcoming data collection with the integrated luminosity of 50 ab^{-1}
- Narrowed down the search area for potential ALPs and constrained the possible solution of the strong CP problem in the MeV to GeV mass range

There is **still some room for improvement**, e.g. more sophisticated background analysis? Investigation of the $(g - 2)_\mu$ impact and limitations?

For more details see arXiv: 2309.15106 [hep-ph]

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Thank you very much for your attention!

I wish you a good time during EINN conference and hope to see you during the poster session!

JOHANNES GUTENBERG
UNIVERSITÄT MAINZ



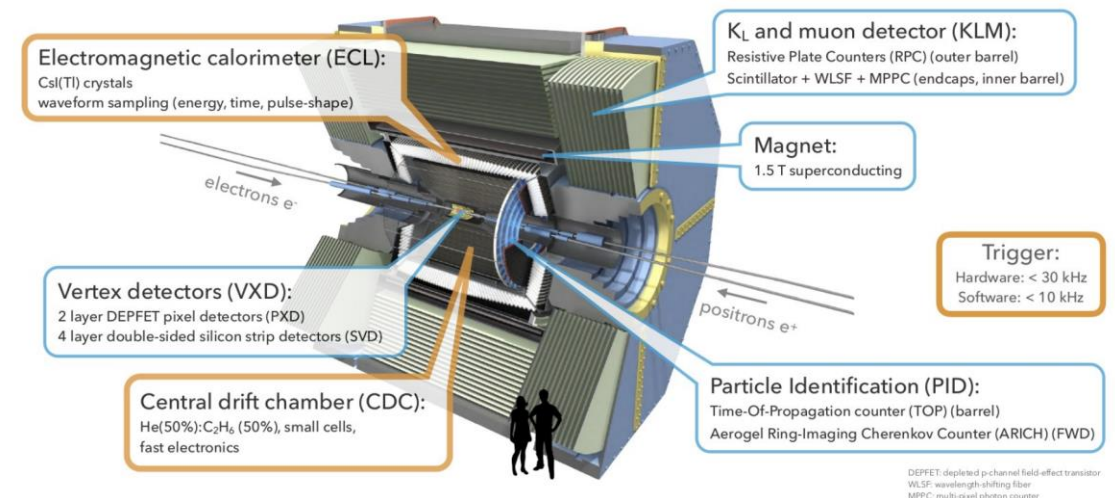
EINN2023



Belle 2

- Belle 2 is an **asymmetric** collider, electron and positron have energies 7 GeV and 4 GeV. This requires a **boost** with a relative velocity $\beta = 0.27$ to the CMS frame
- For **2018** analysis we require **three resolved photons** with energies higher than 0.65 GeV in the center-of-momentum frame as a criteria
- These requirements are slightly different from those used in the Belle II 2018 report, where the selection of photons with energies above 0.65 GeV for $m_a > 4$ GeV and 1 GeV for $m_a \leq 4$ GeV in the lab frame was performed
- The difference is negligible since $g_{a\gamma\gamma}$ is sensitive to $\sigma_{QED}^{-1/4}$

- The region $37.3^\circ < \theta < 123.7^\circ$ provides the best energy resolution, avoiding regions close to detector gaps, and offers the lowest beam background levels – used for both 2018 and projected analyses
- Following the work of M. J. Dolan and others, we set the photon energy selection threshold of 0.25 GeV in the CMS frame for projected curves



Belle 2

