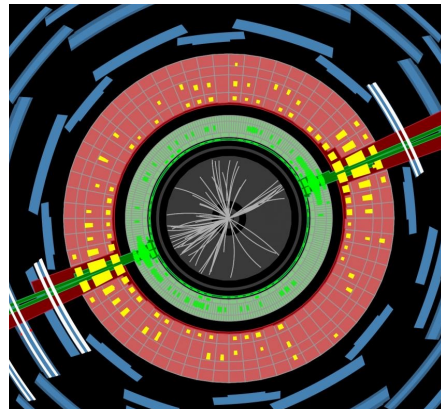


Pursuing Physics Beyond the Standard Model



...using model independent measurements.



Herbifest, Bonn,



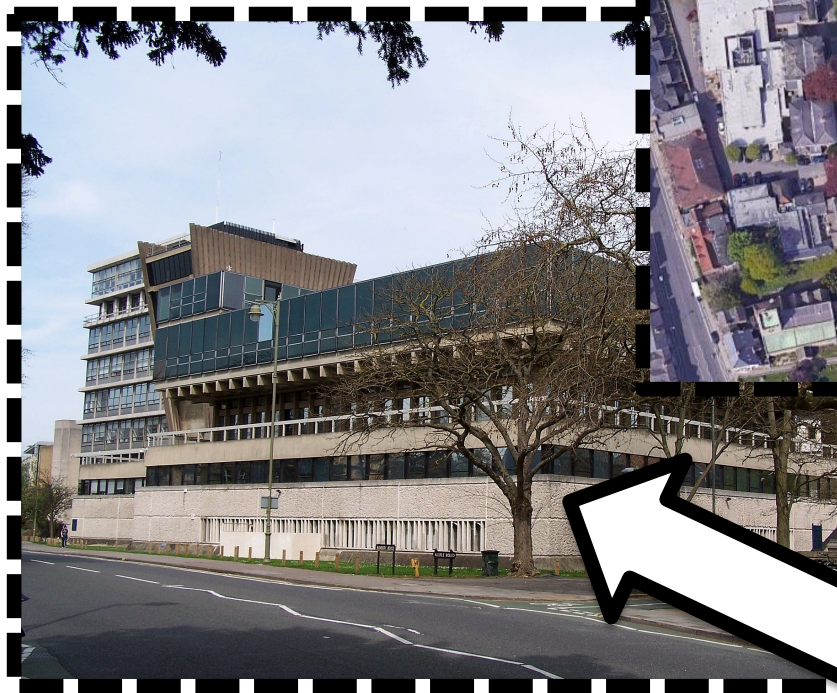
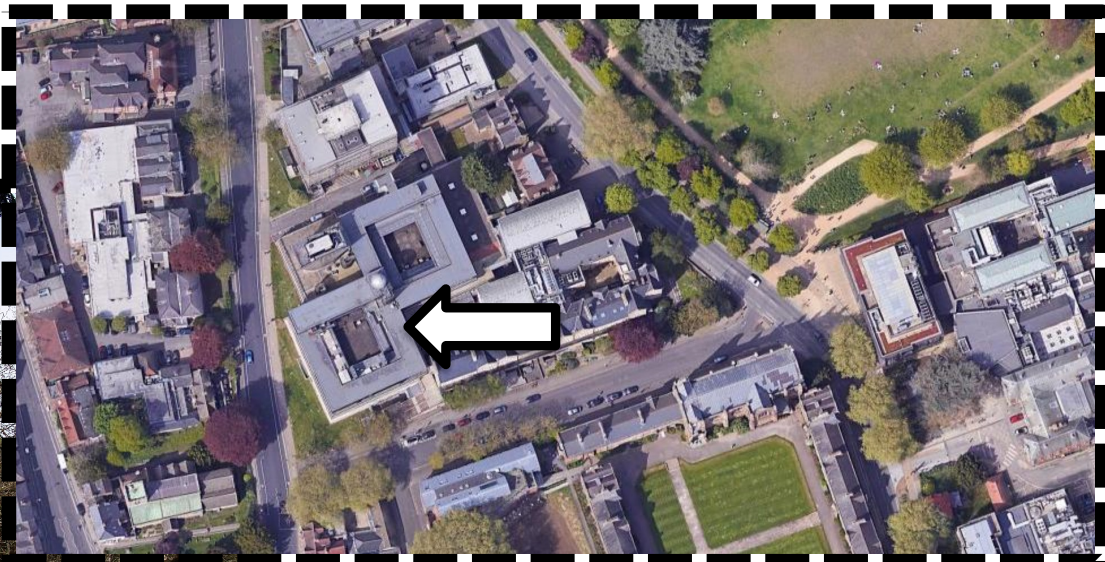
28/3/2023

Jon Butterworth



Physics & Astronomy

Oxford, 1990



Herbi goes to Monte Carlo



Finding Supersymmetry via R-parity Violation
 J. Butterworth, H. Dreiner
 Department of Physics, University of Oxford,
 1 Keble Rd, Oxford OX1 2NP

1 Introduction

Recent theoretical developments [1, 2] show that supersymmetry with broken R-parity (\tilde{R}_p) should be considered on equal footing with the "Minimal supersymmetric standard model" (MSSM) [3]. It is known that the phenomenology of the two models is significantly different [2, 4]. Thus if one wants to find supersymmetry (SUSY) it is imperative to also consider the new signals of \tilde{R}_p . It is the purpose of this report to study a class of \tilde{R}_p signals at HERA. As we shall see below, HERA is ideally suited to study this class and there has a unique chance of discovering supersymmetry. (For a previous study see Ref. [5].)

R-parity (\tilde{R}_p) is a multiplicative discrete symmetry. For a previous study see Ref. [5]. \tilde{R}_p and \tilde{B} are the spin, the lepton- and the baryon-number, respectively. All particles of the standard model (SM) have $\tilde{R}_p = +1$; their respective superpartners have $\tilde{R}_p = -1$. \tilde{R}_p conservation is imposed in the MSSM and their scalar superpartners. When relaxing this symmetry, but still restricting the model to the minimal particle content, the superpotential in general also contains the following gauge and supersymmetric invariant terms which violate \tilde{R}_p .

$$W_{\tilde{R}_p} = (\lambda_{ijk} L_i L_j E_k + \lambda'_{ijk} L_i Q_j \tilde{D}_k + \lambda''_{ijk} \tilde{L}_i \tilde{D}_j \tilde{D}_k) \nu_i \quad (1)$$

Here L and \tilde{L} (Q and \tilde{Q} , \tilde{D}) are the (left-handed) lepton doublet and the antilepton singlet (quark doublet and antiquark singlets) chiral superfields respectively. $\lambda, \lambda', \lambda''$ are dimensionless Yukawa couplings and i, j, k are generation indices. Thus there are in general 45 extra terms. The first two sets of terms in Eq.(1) violate lepton-number, the third set, thus one usually considers either only the lepton- or only the baryon number violating terms. In addition, we make the simplifying assumption that, in turn, only one operator of Eq.(1) is non-zero, leading to 45 different models. We consider this a likely scenario, since in the SM the Yukawa couplings show a strong hierarchy.

2 \tilde{R}_p at HERA

The phenomenology of \tilde{R}_p differs in two important aspects from that of the MSSM. First, the lightest supersymmetric particle (LSP) is no longer protected by a symmetry and can thus decay. Throughout we shall make the simplifying assumption that the LSP is the photino, $\tilde{\gamma}$. Furthermore we shall assume that the mass of the photino $m_{\tilde{\gamma}} = 50 \text{ GeV}$. In

the scenario that we consider, light photinos can most likely be experimentally produced. A typical decay through the operator $L_i Q_j \tilde{D}_k$, for example, is

$$\tilde{\gamma} \rightarrow e^+ \bar{\nu} \rightarrow e^+ \bar{d} \quad (2)$$

Therefore the missing p_T signals of the MSSM are either diluted or completely absent. Second, in \tilde{R}_p beyond the standard pair production mechanisms of the MSSM, it is also possible to produce single supersymmetric particles through the operators (1), the pair (or singly) produced SUSY particles subsequently decay to \tilde{R}_p even final states. Combining these results in three sets of promising signals at HERA:

- 1: The most viable MSSM pair production mechanism at HERA is $e^+ q \rightarrow \tilde{e}^+ q$. In \tilde{R}_p , the selectron and the squark will dominantly cascade decay to the LSP $\tilde{\gamma}$, the most visible signal is obtained in the case of a dominant operator $L_i L_j E_k$, and the charges can vary due to the Majorana nature of the photino. For the complete process we have

$$e^+ q \rightarrow \tilde{e}^+ q \rightarrow e^+ q + 2 \tilde{\gamma} \rightarrow e^+ q + 2 e^+ \bar{\nu} \quad (3)$$

and the photinos do not have to decay identically. This process can be analyzed by using the existing MSSM Monte Carlo but now adding the \tilde{R}_p photino decay. For \tilde{R}_p the discovery region is larger than in the MSSM case since the $m_{\tilde{e}} - m_{\tilde{q}}$ plane is experimentally less restricted. We shall not discuss these signals here.

- 2: At HERA, to leading and next to leading order, all \tilde{R}_p single SUSY production mechanisms proceed through the operator $L_i Q_j \tilde{D}_k$, and they all involve the s-channel resonance production of a squark or an anti-squark $e^+ q \rightarrow \tilde{q}^*$. Therefore, in principle, squarks with mass up to the total CM-energy can be produced. This makes HERA an ideal machine to look for supersymmetry, in the case of a dominant operator $L_i Q_j \tilde{D}_k$. The produced squarks can decay in several ways: $\tilde{q} \rightarrow (e \bar{\nu}, q \bar{q}', q' \bar{q})$, where \bar{q} denotes an arbitrary MSSM gaugino other than the photino, and \tilde{R}_p operator and below discuss the decay to the photino. We shall not further consider the other cases.

$$e^+ u_i \rightarrow \tilde{d}_j \rightarrow e^+ + u_i \quad (4)$$

$$e^+ d_j \rightarrow \tilde{u}_i \rightarrow e^+ + d_j \quad (5)$$

In the case of the \tilde{R}_p decay of the squark, the inclusive processes are given by and will result in two resonant bumps at $x = m_{\tilde{q}}^2/s$ and at $x = m_{\tilde{q}'}^2/s$ in the cross section $d\sigma/dx(ep \rightarrow e^+ X)$. These signals correspond exactly to leptoquark production [5, 6], and the analysis need not be repeated. The first process (4) corresponds, in the notation of Ref. [6], to the production of a S_1 leptoquark. The process (5) corresponds to the production of a R_2 leptoquark. The quantum numbers F of Table 1 in Ref. [6] do not agree with those of the produced squarks, since in determining F it was assumed that lepton- and baryon-number are conserved, this is explicitly not the case in \tilde{R}_p . We now address the question of how one could experimentally distinguish leptoquark- and baryon-number conservation, as seen in Eqs.(4,5), an \tilde{R}_p operator will lead to two resonance bumps, provided the squarks are not mass degenerate. However, theoretically

$M_{\tilde{e}} = 200 \text{ GeV}/c^2$
 $M_{\tilde{q}} = 150 \text{ GeV}/c^2$
 $M_{\tilde{q}'} = 250 \text{ GeV}/c^2$
 $M_{\tilde{\gamma}} = 50 \text{ GeV}/c^2$
 0.01

Which led to...



INSPIRE HEP

literature find author H. K. Dreiner and author J. Butterworth

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R-parity violation at HERA #1
J. Butterworth (Oxford U.), Herbert K. Dreiner (Oxford U.) (Nov 20, 1992)
Published in: *Nucl.Phys.B* 397 (1993) 3-34 • e-Print: [hep-ph/9211204](#) [hep-ph]
pdf DOI cite claim reference search 170 citations

Finding supersymmetry at HERA via R-parity violation #2
J. Butterworth (Oxford U.), Herbert K. Dreiner (Oxford U.) (Oct, 1991)
Contribution to: *Workshop on Physics at HERA*, 1079-1087
pdf cite claim reference search 1 citation

Date of paper
1991 1992

Number of authors
 10 authors or less 2

Exclude RPP
 Exclude Review of Particle Physics 2

Document Type

But actually...

Which led to...



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World ▶ Europe US Americas Asia Australia Middle East Africa Inequality Global development

Life and Physics
Science

• This article is more

Can physics refugees...

Herbi Dreiner

Tue 19 Jan 2016 17:24 GMT

f t e 54

Since nature is u...
can reach out to...
practice? And if...

Life and Physics
Science

SLAM POW BANG: On Tour in England with the Bonn Physics Show

My old friend Herbi Dreiner has won awards for his classic(al) physics shows. Next week he brings a new one, with added quantum, to the UK. Here he describes the experience of trying to put the Higgs boson on stage. Plus! DEMO: The Movie

Jon Butterworth
@jonmbutterworth
Thu 13 Mar 2014 07:00 GMT

f t e 0

And also, indirectly, the rest of this talk...

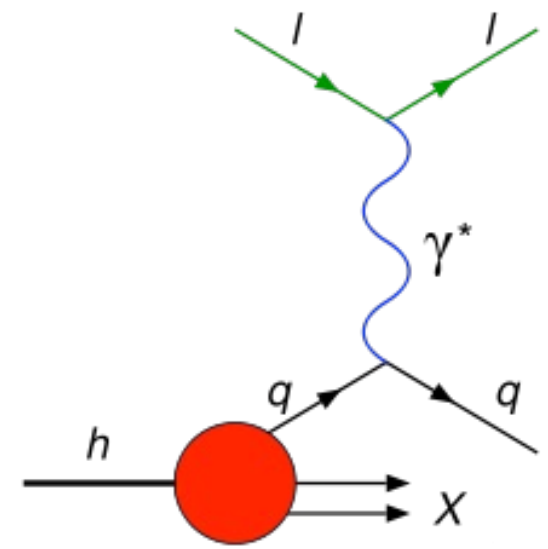
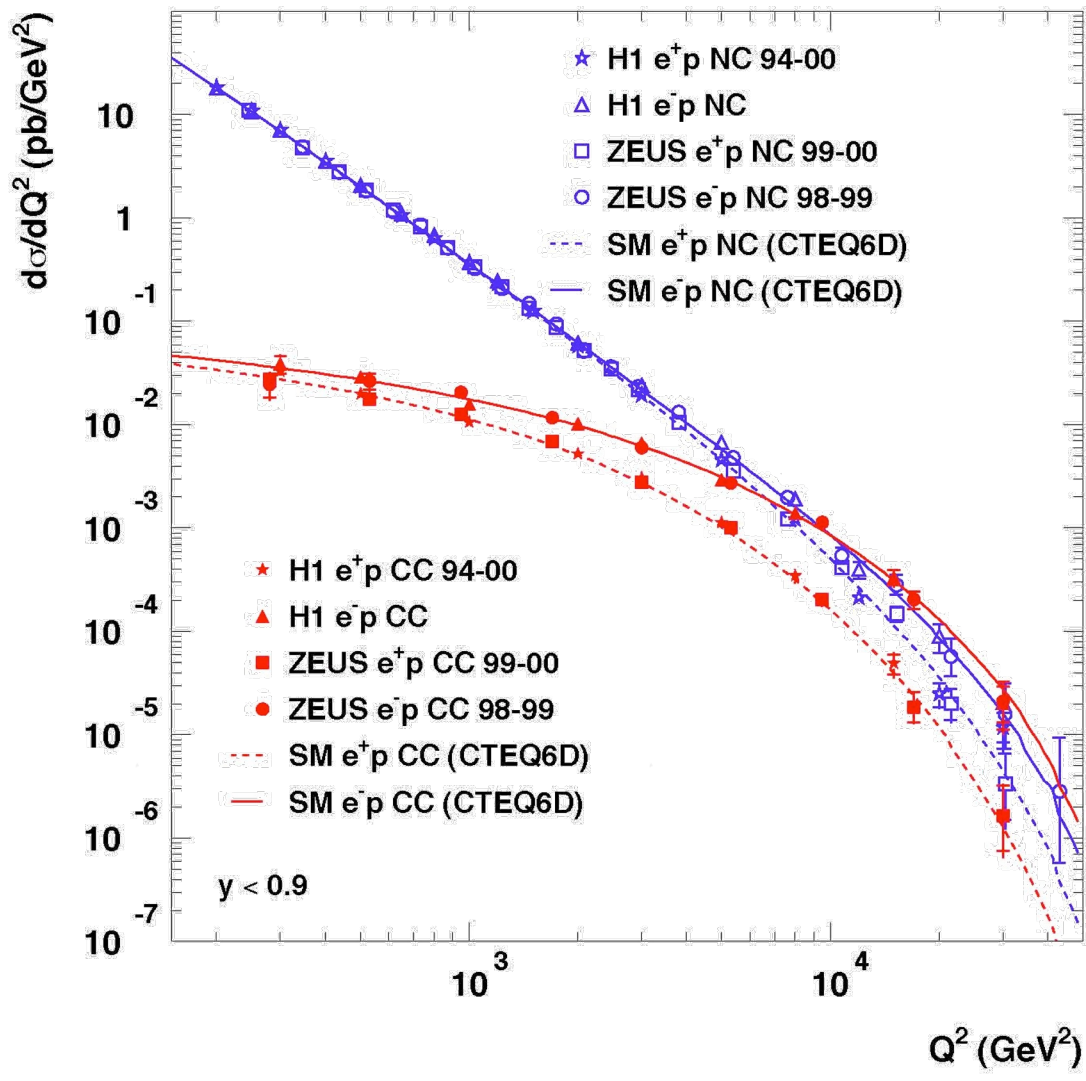
A MAP OF THE
INVISIBLE
GREAT TRAIN
JOURNEYS

Particle physics in the 90's

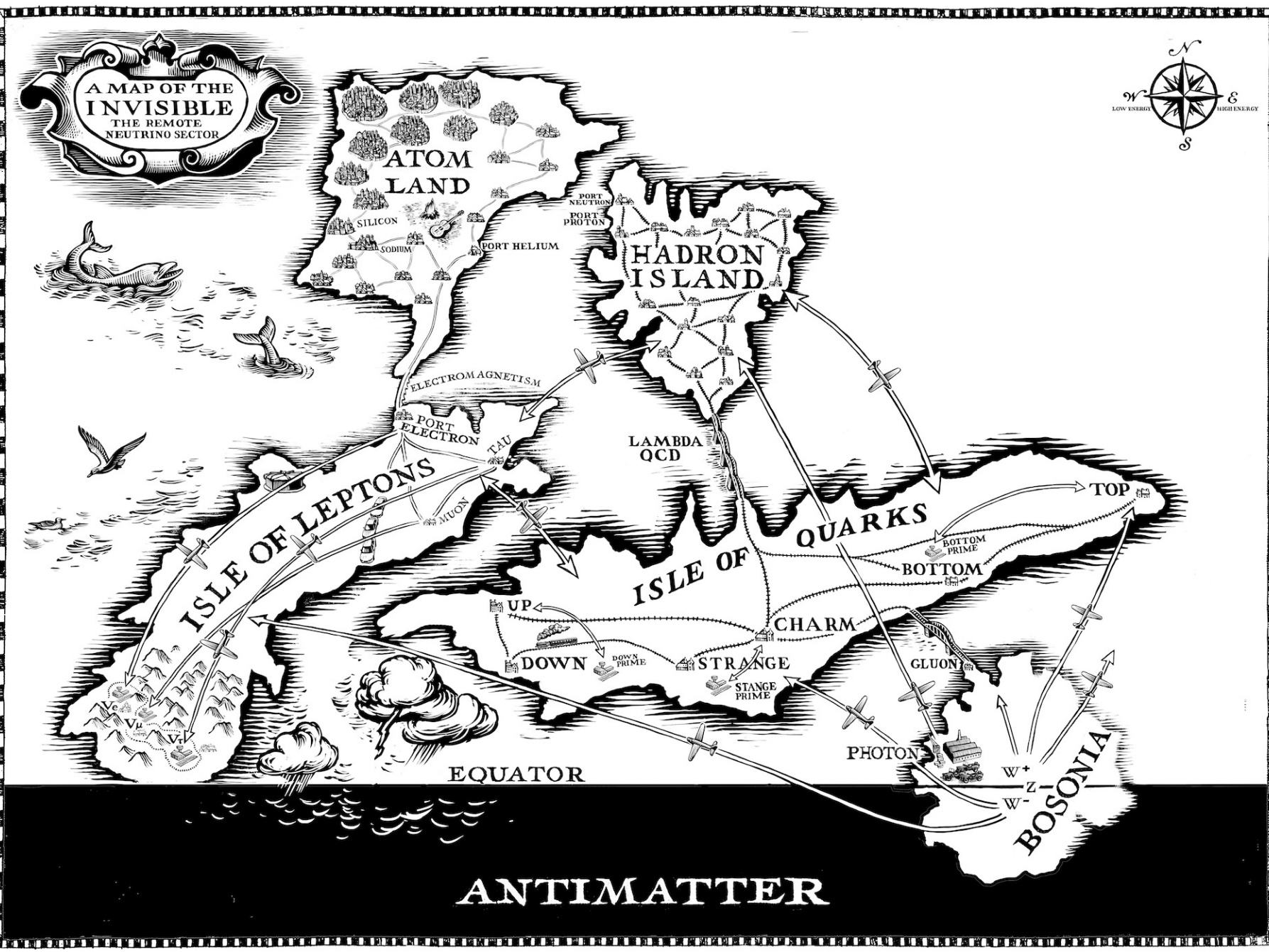


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HERA



A MAP OF THE INVISIBLE
THE REMOTE NEUTRINO SECTOR



ANTIMATTER

ISLE OF LEPTONS

ISLE OF QUARKS

QUARKS

BOSONIA

ATOM LAND

HADRON ISLAND

EQUATOR

SILICON

SODIUM

HELIUM

PORT NEUTRON

PORT PROTON

ELECTROMAGNETISM

PORT ELECTRON

TAU

MUON

LAMBDA QCD

UP

DOWN

DOWN PRIME

STRANGE

STRANGE PRIME

CHARM

BOTTOM

BOTTOM PRIME

TOP

GLUON

PHOTON

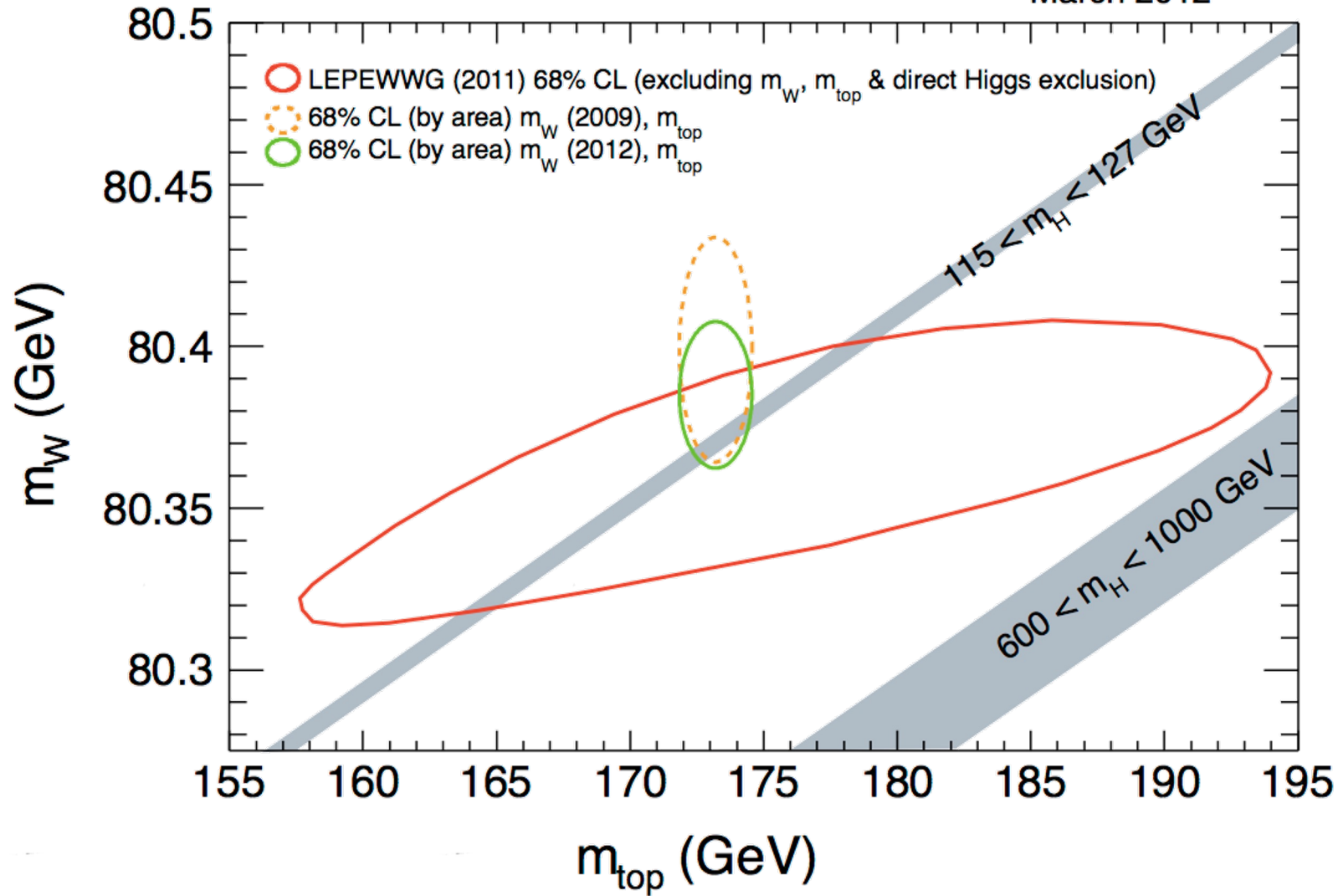
W+

W-

Z

Before the LHC

March 2012





SUISSE
FRANCE

CMS

LHCb

CERN
Prévessin

ATLAS

CERN
Meyrin

SPS 7 km

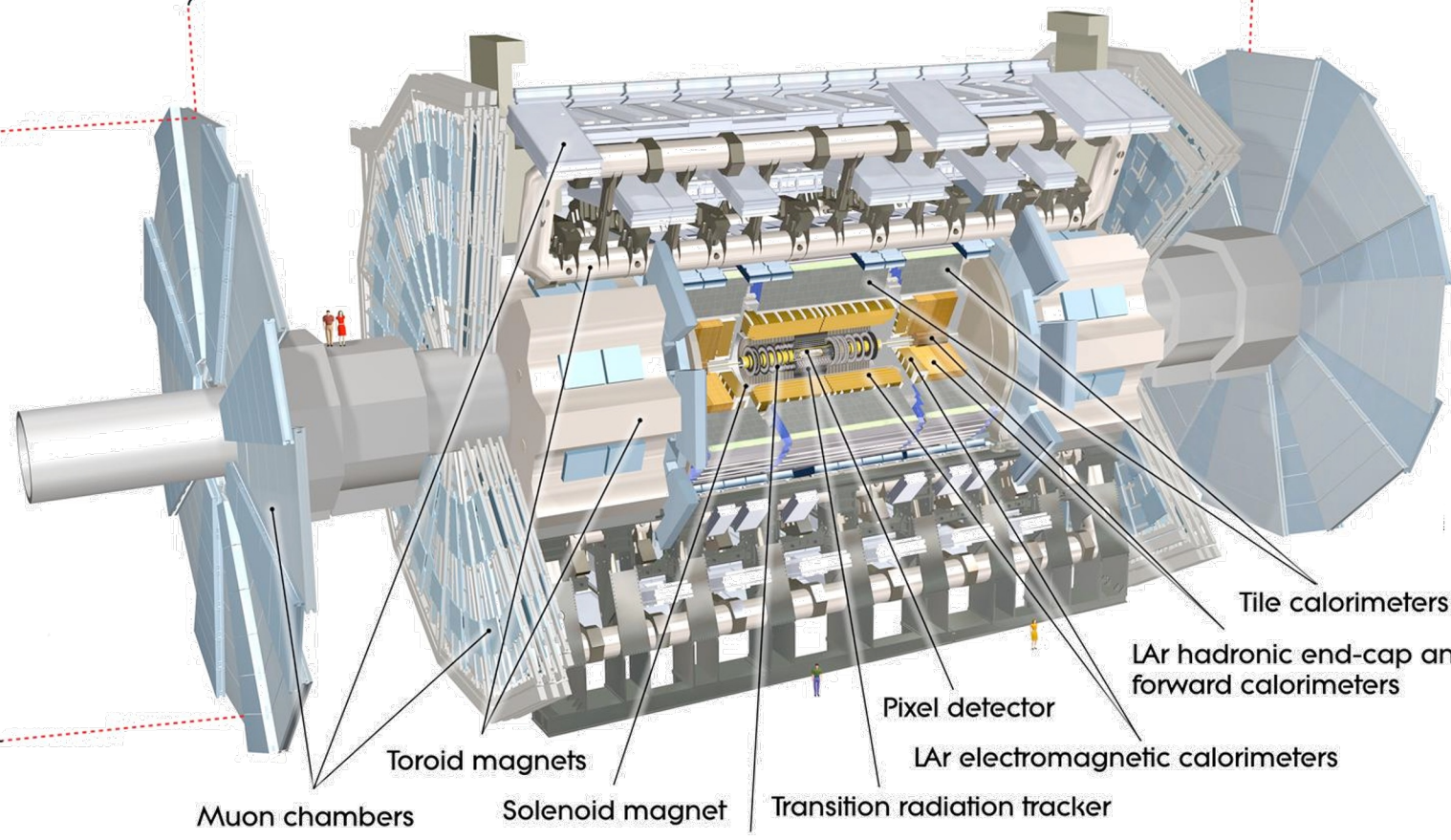
ALICE

LHC 27 km



44m

25m



Muon chambers

Toroid magnets

Solenoid magnet

Semiconductor tracker

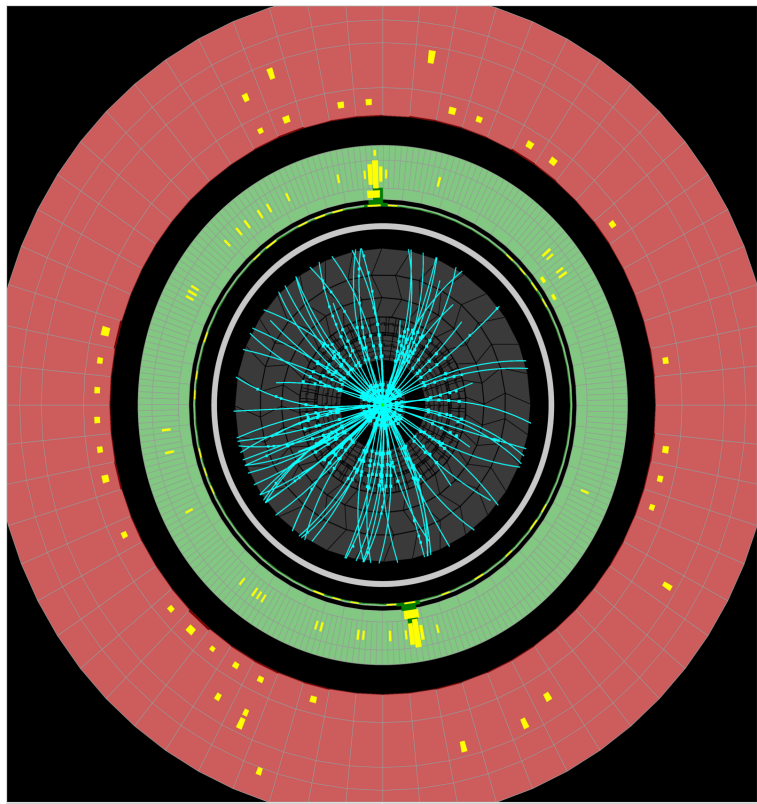
Transition radiation tracker

Pixel detector

LAr electromagnetic calorimeters

LAr hadronic end-cap and forward calorimeters

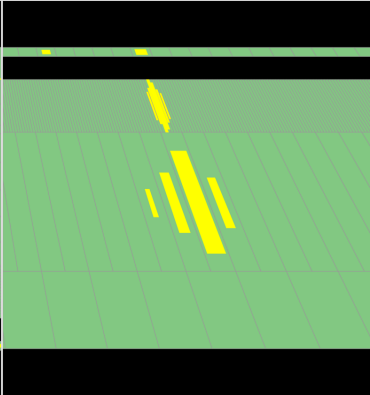
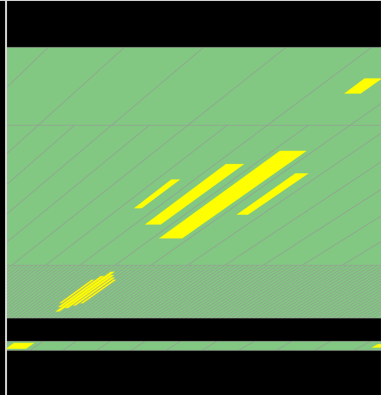
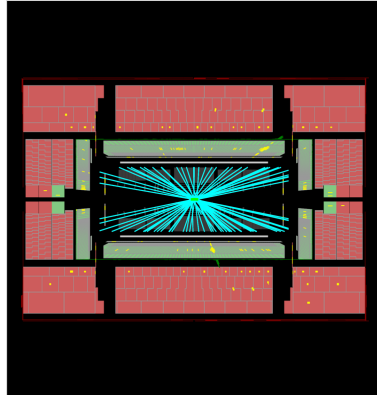
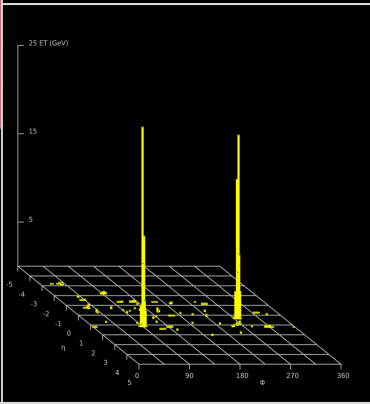
Tile calorimeters



ATLAS EXPERIMENT

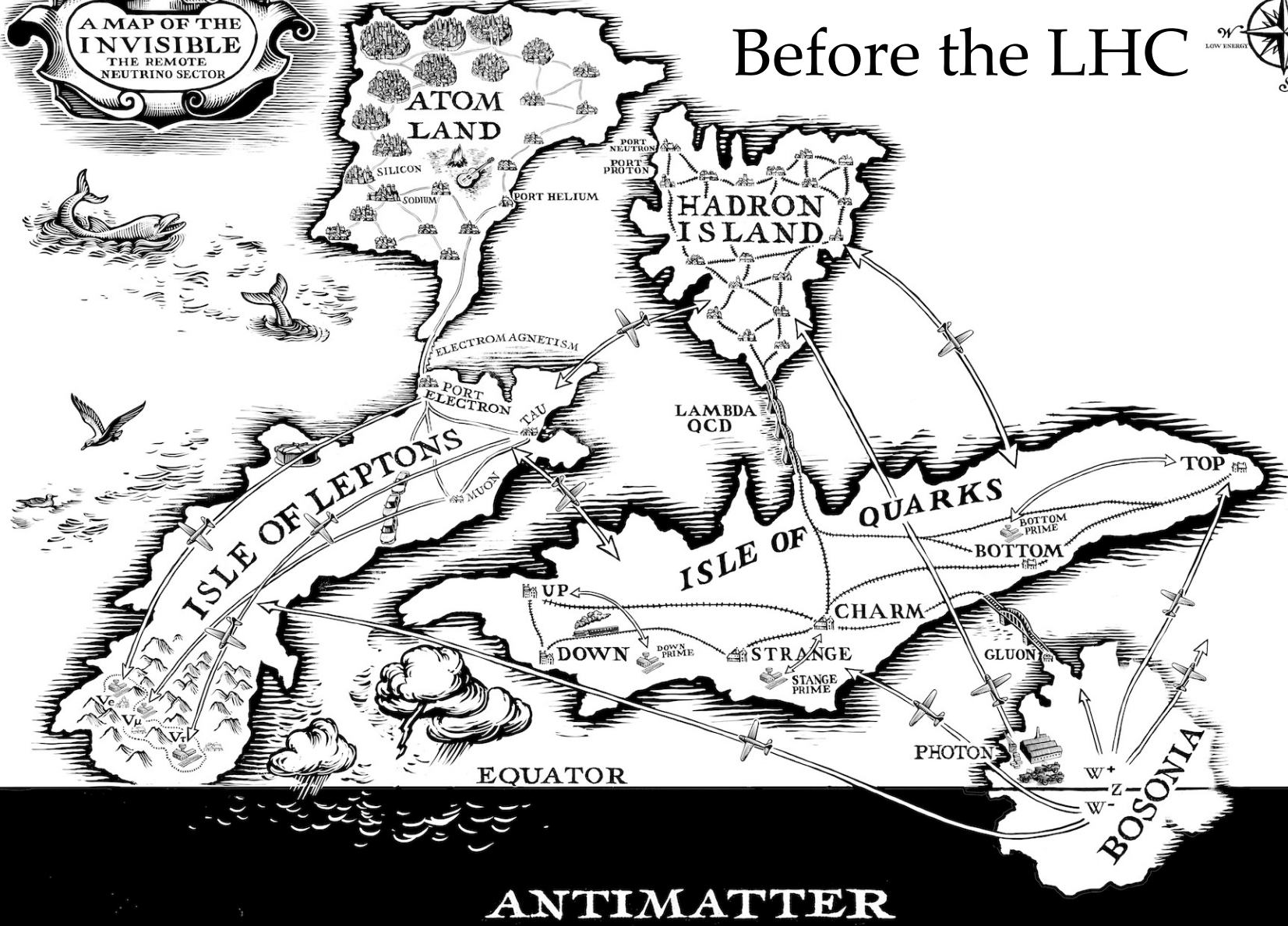
Run Number: 203779, Event Number: 56662314

Date: 2012-05-23 22:19:29 CEST



A MAP OF THE
INVISIBLE
THE REMOTE
NEUTRINO SECTOR

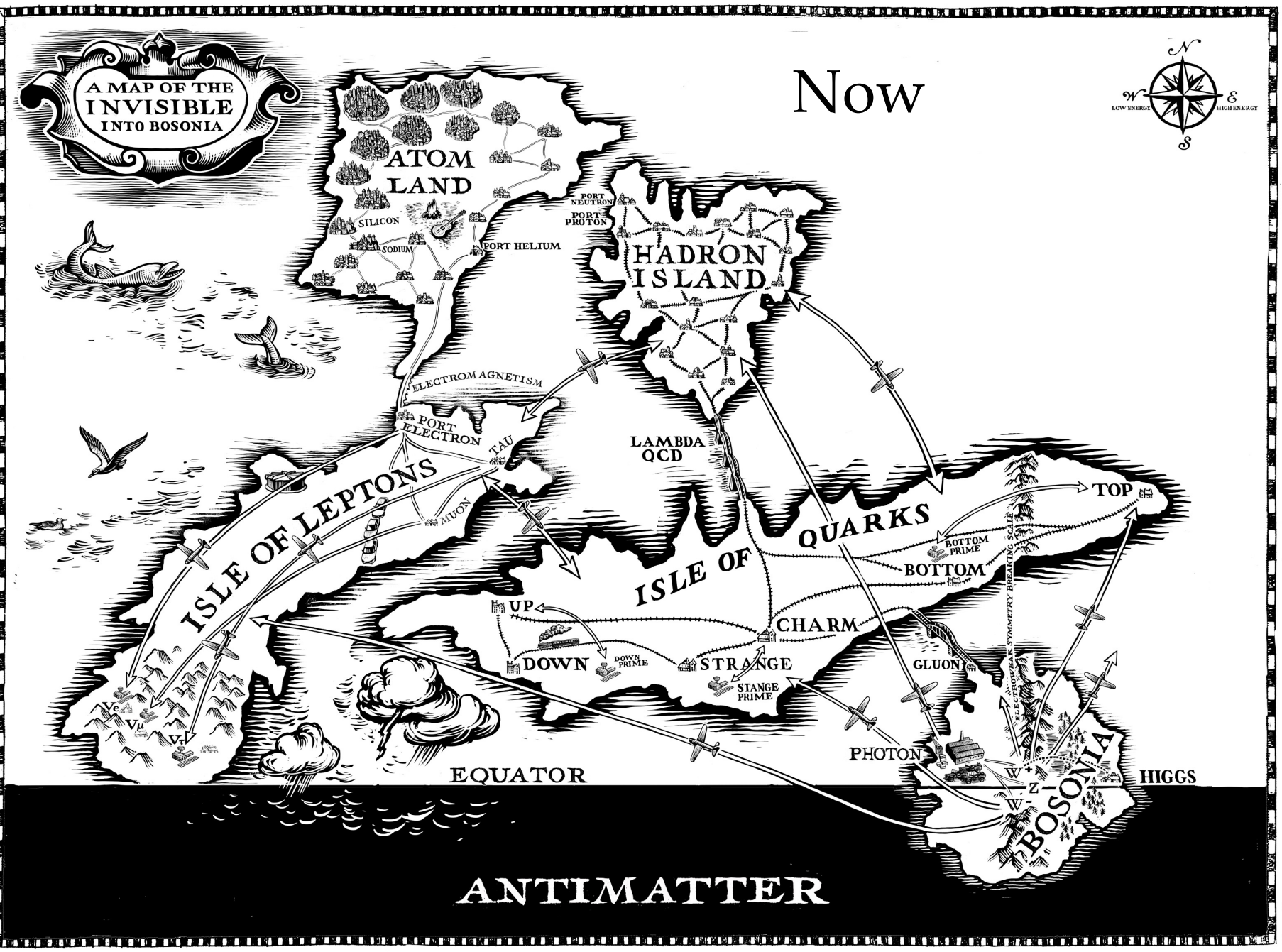
Before the LHC



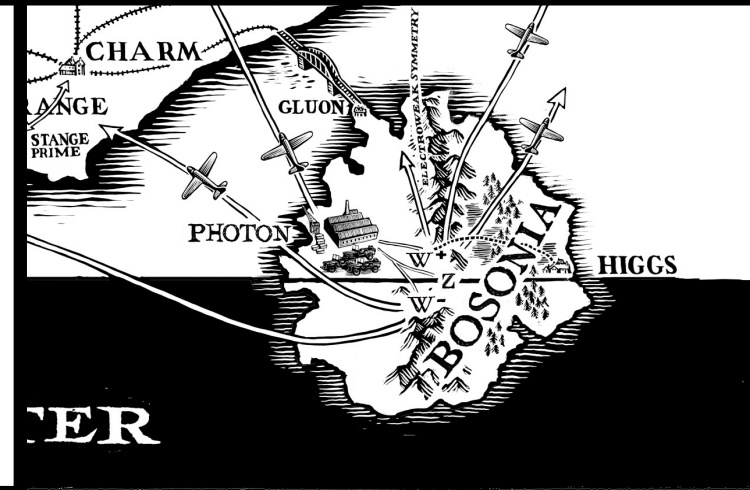
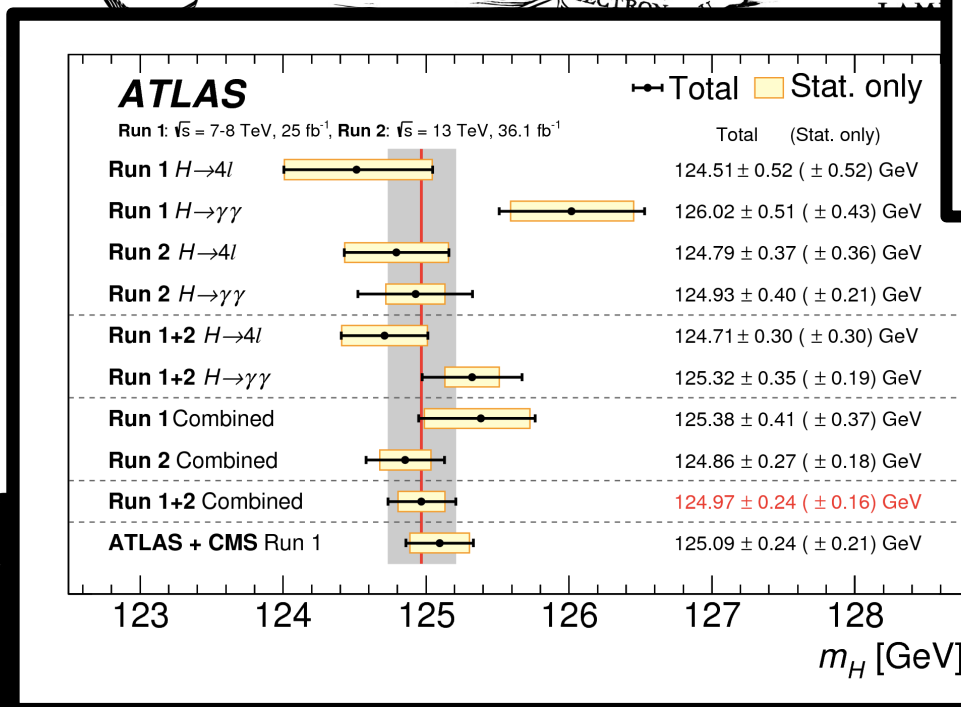
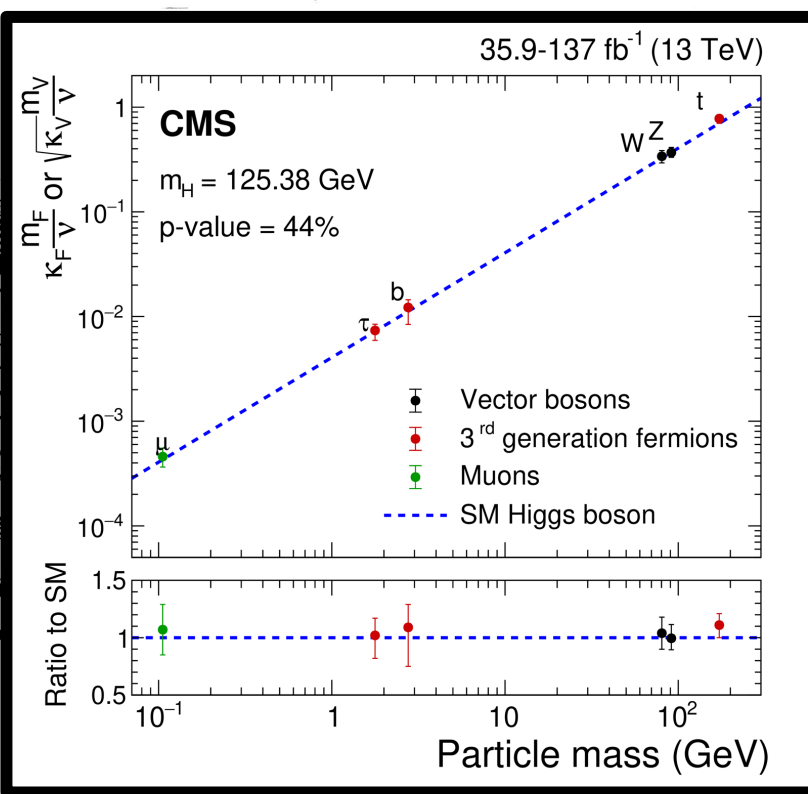
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A MAP OF THE INVISIBLE INTO BOSONIA

Now



A MAP OF THE INVISIBLE INTO BOSONIA

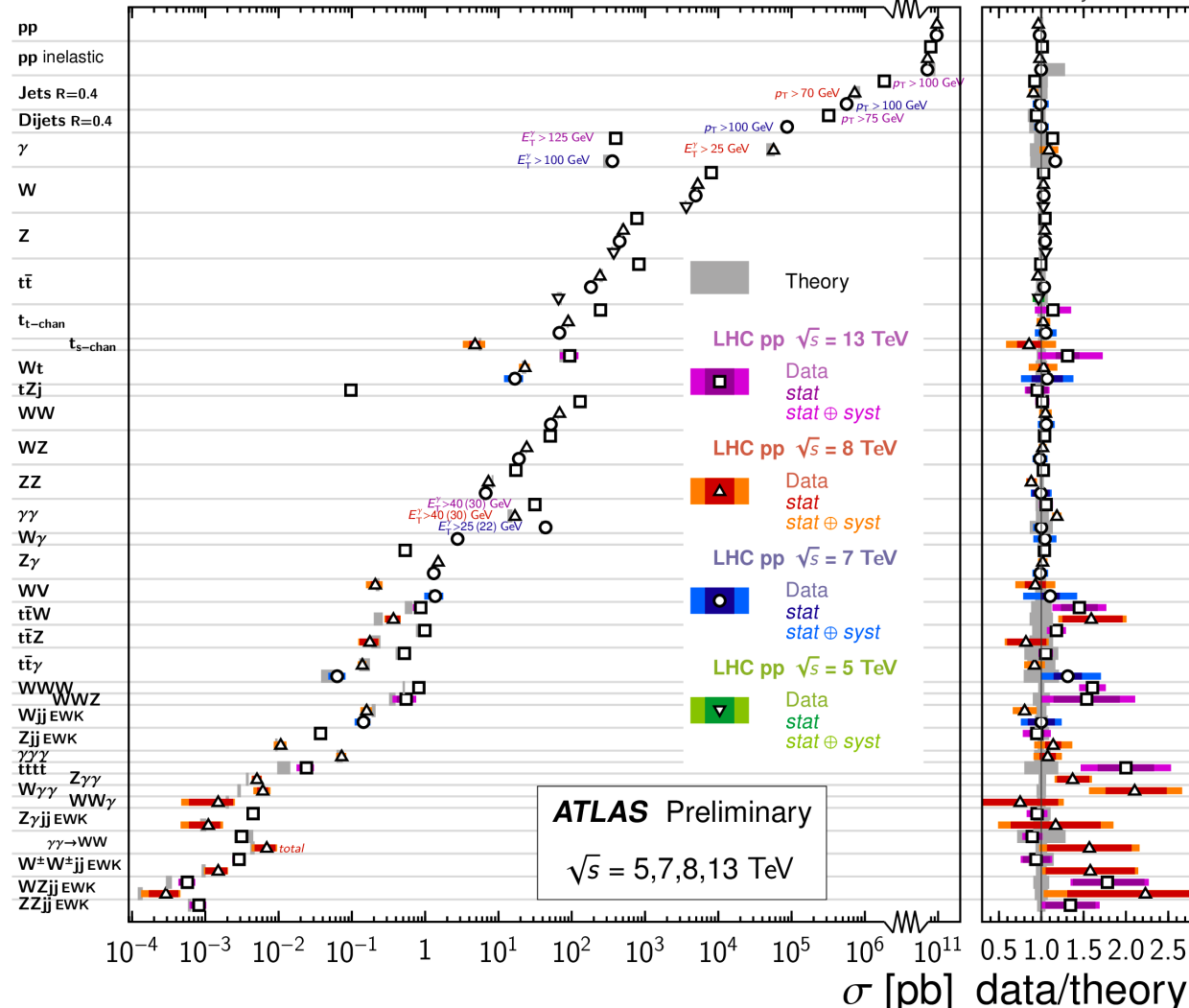


ER



Standard Model Production Cross Section Measurements

Status:
February 2022

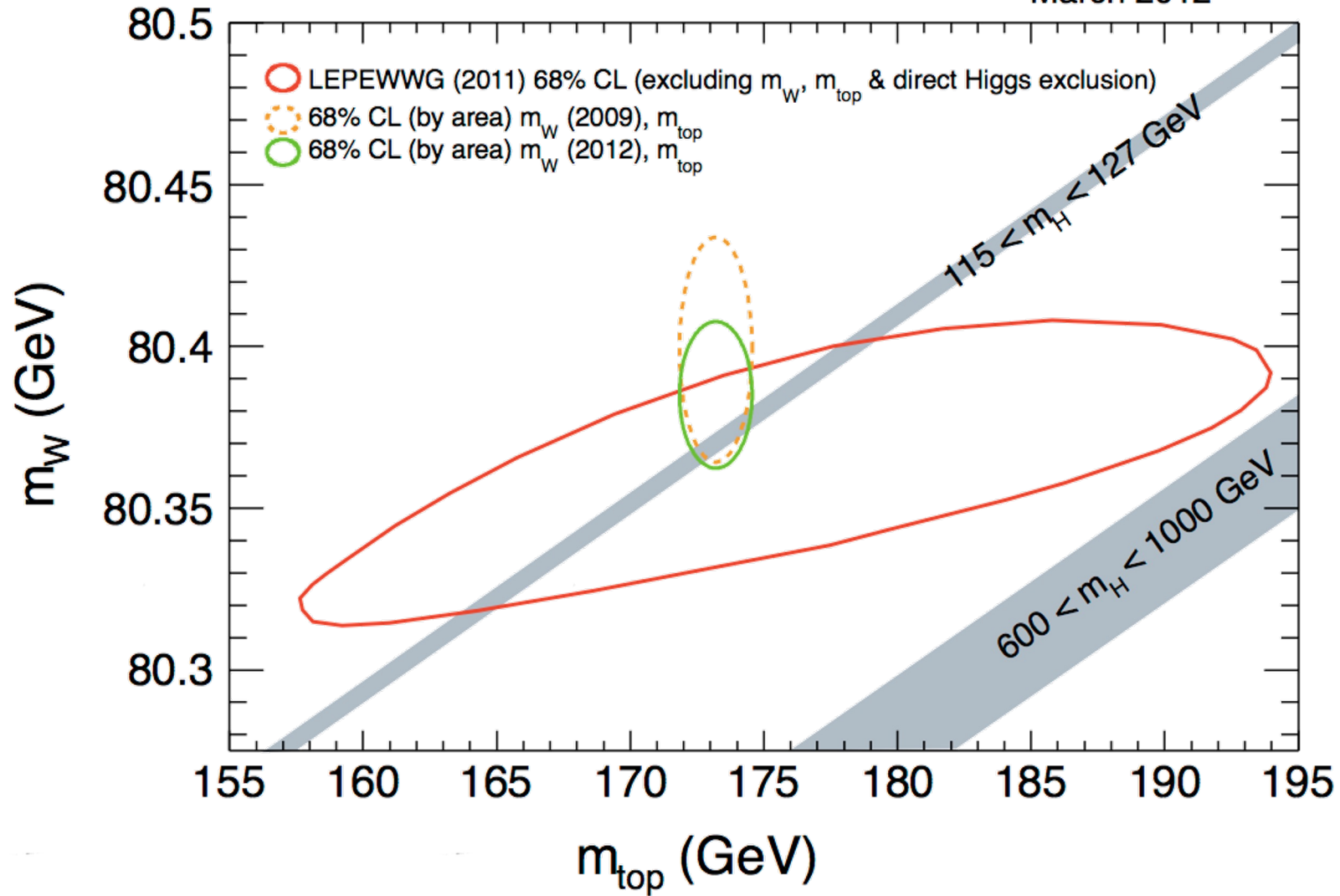


$\int \mathcal{L} dt$ [fb ⁻¹]	Reference
50×10^{-8}	PLB 761 (2016) 158
20×10^{-8}	Nucl. Phys. B, 486-548 (2014)
8×10^{-8}	PRL 117, 182002 (2016)
50×10^{-8}	PLB 761 (2016) 158
20×10^{-8}	Nucl. Phys. B, 486-548 (2014)
8×10^{-8}	JHEP 05 (2018) 195
20×10^{-8}	JHEP 09 (2017) 020
4×10^{-8}	JHEP 02, 153 (2015)
20×10^{-8}	JHEP 05 (2018) 195
20×10^{-8}	JHEP 05, 059 (2014)
20×10^{-8}	PLB 2017 04 072
20×10^{-8}	JHEP 06 (2016) 005
4×10^{-8}	PRD 89, 052004 (2014)
0×10^{-8}	PLB 759 (2016) 601
20×10^{-8}	EPJC 79 (2019) 760
4×10^{-8}	EPJC 77 (2017) 367
20×10^{-8}	EPJC 79 (2019) 124
0×10^{-8}	JHEP 02 (2013) 128
20×10^{-8}	JHEP 02 (2017) 117
4×10^{-8}	JHEP 02 (2017) 117
20×10^{-8}	EPJC 80 (2020) 528
20×10^{-8}	EPJC 74 (2014) 3109
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20×10^{-8}	JHEP 01 (2018) 63
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Before the LHC

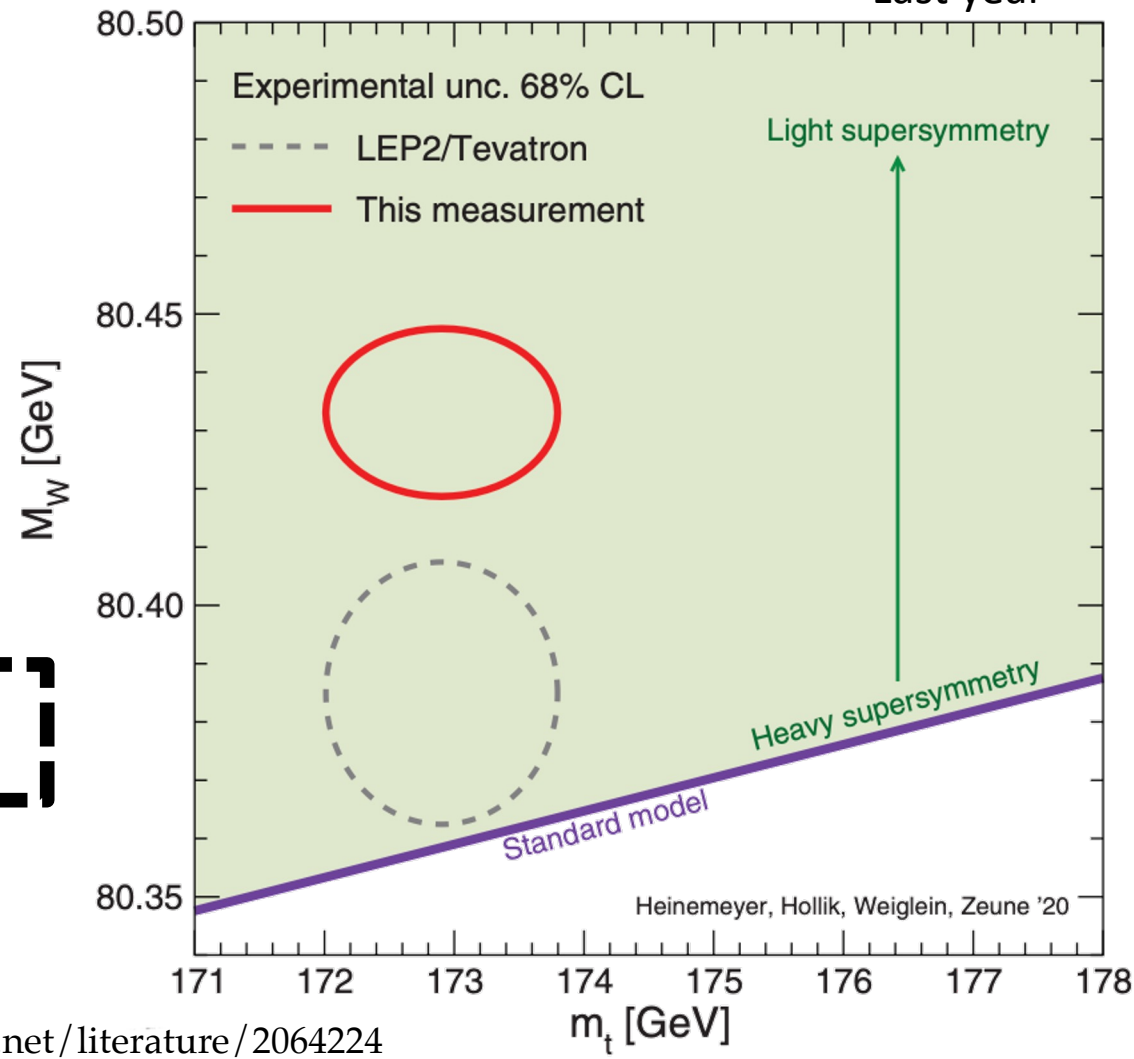
March 2012



Update from CDF



Last year

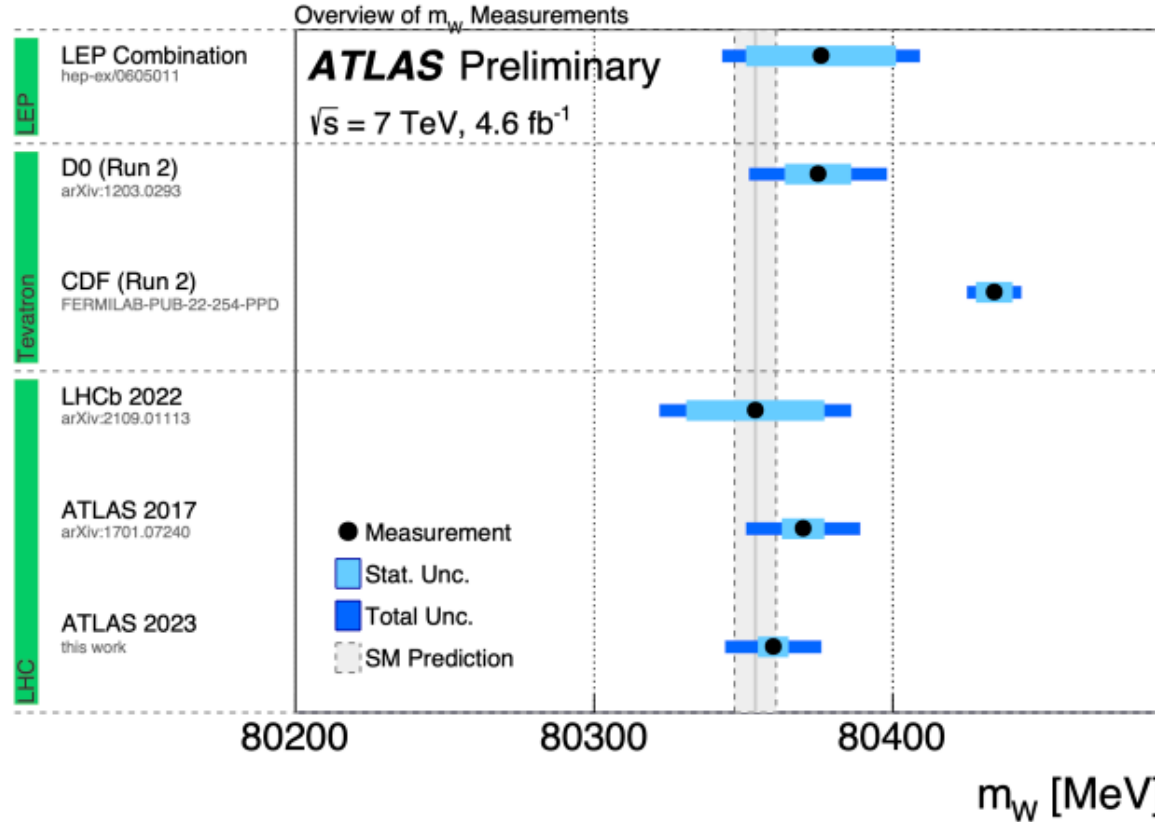


Plot ©GW et al,
caveats noted

Update from ATLAS

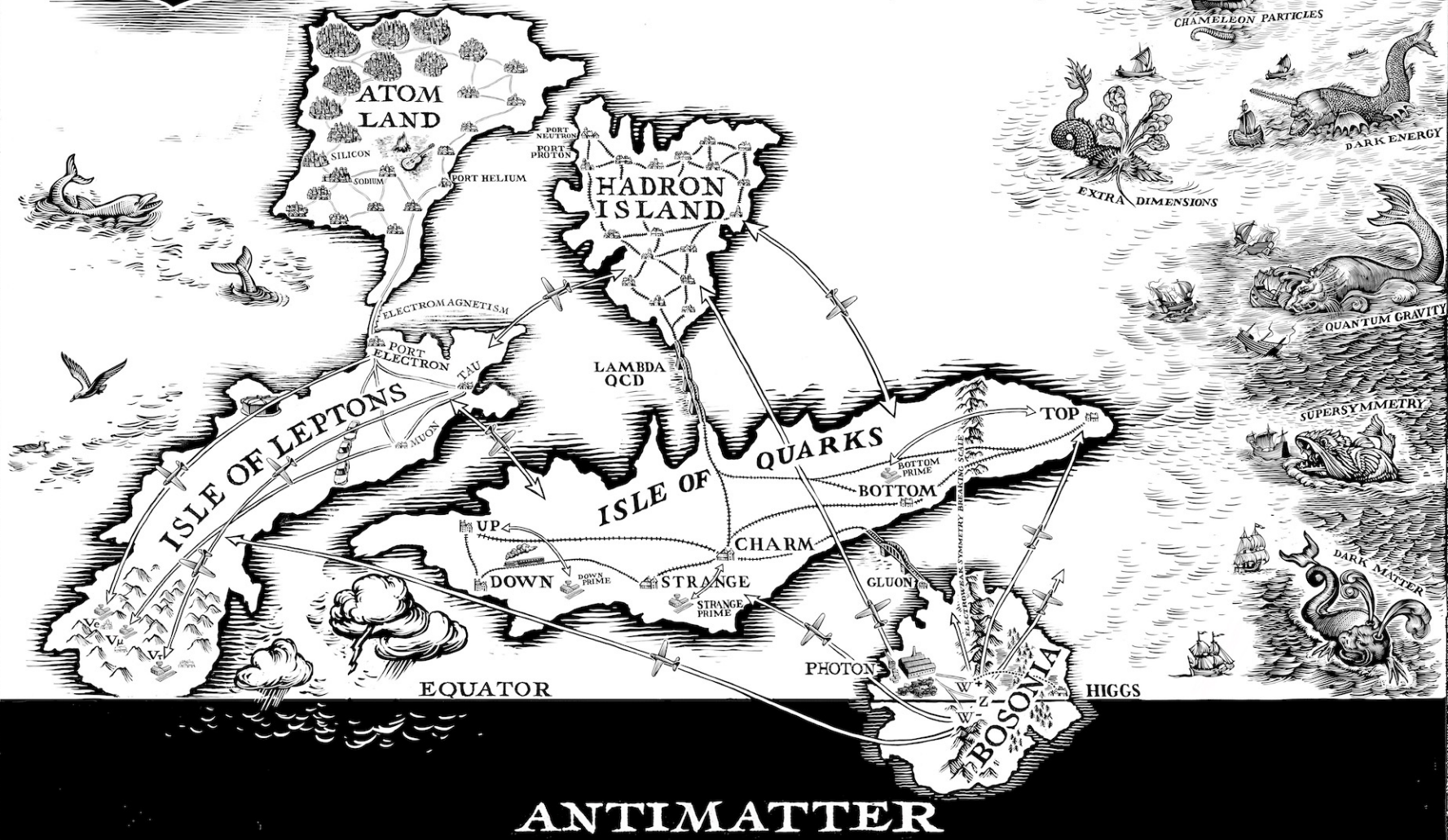


Last week



A MAP OF THE
INVISIBLE
FAR EAST

Off the map...?



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A MAP OF THE
INVISIBLE
FAR EAST

Measurements at the Energy Frontier

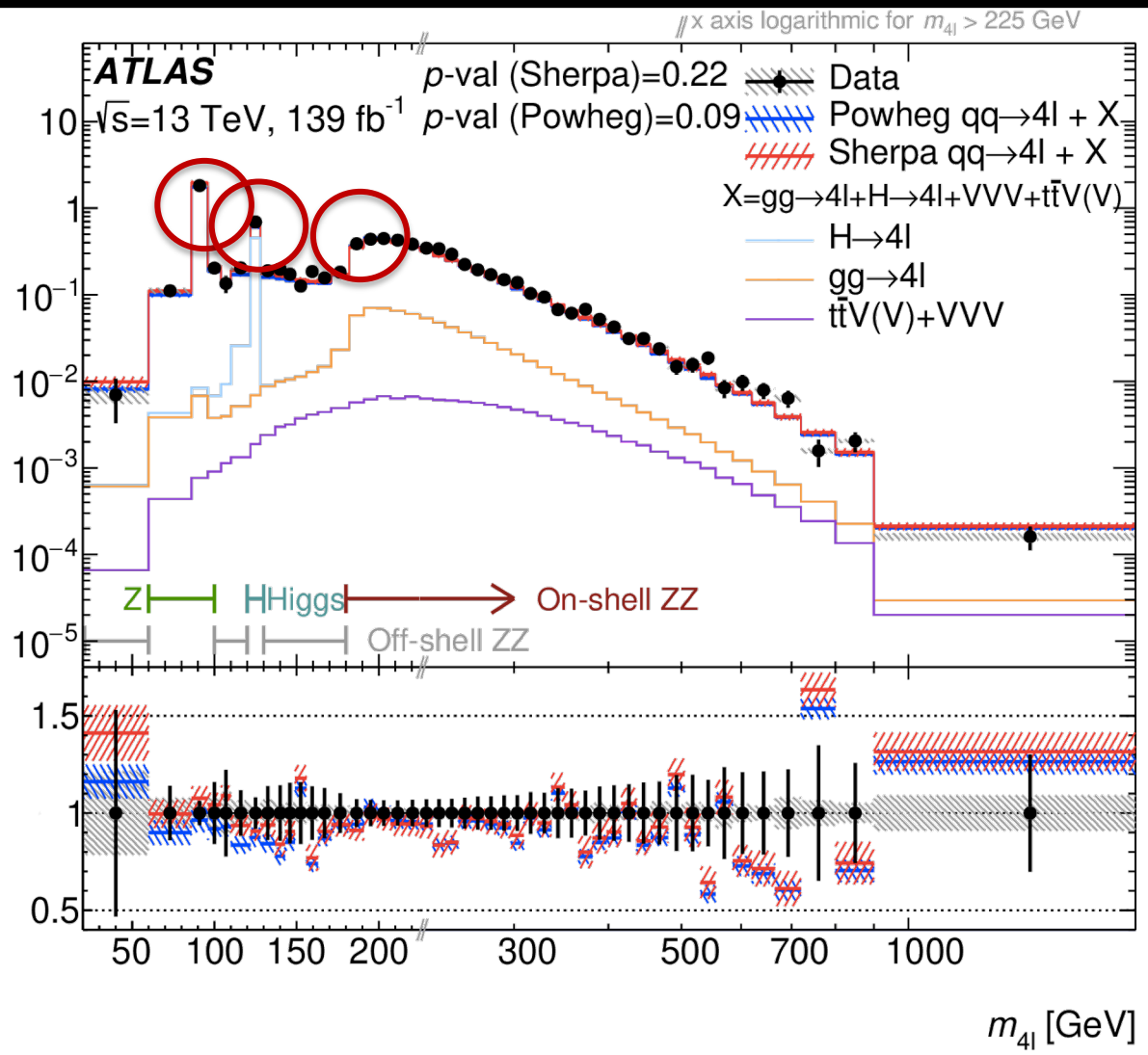


Run 3,
HL-LHC

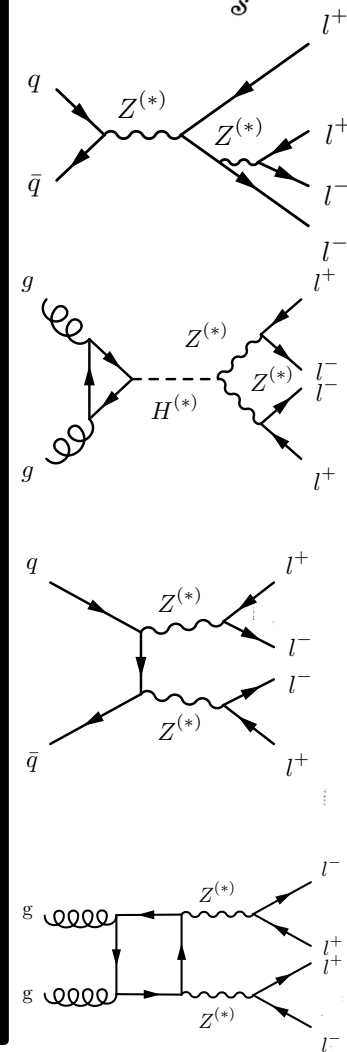
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$d\sigma/dm_{4l}$ [fb/GeV]

Prediction/Data



arXiv:2103.01918 (JHEP)

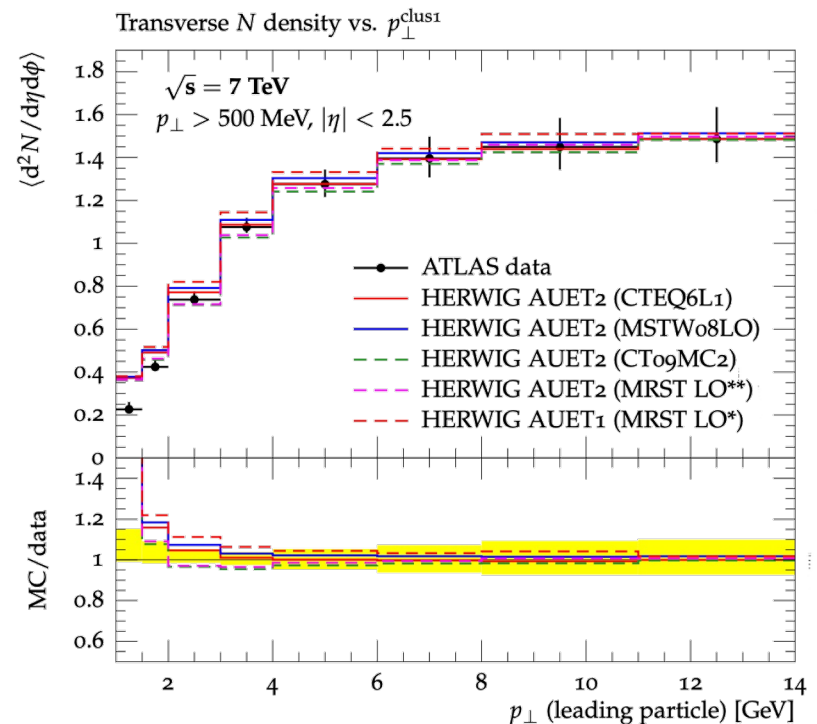


Introducing Rivet



“Robust Independent Validation of Experiment and Theory”

- Direct legacy from HERA (1990s, HZTOOL)
- Developed by MCnet for tuning and validation of new MC event generators
 - e.g. What does the underlying event look like in 7 TeV pp collisions?
- Vast library of measurements of final state particles produced in collisions, and variables derived from them



Buckley et al, Bierlich et al *arXiv:1003.0694* (CPC),
arXiv:1912.05451 (SciPost)

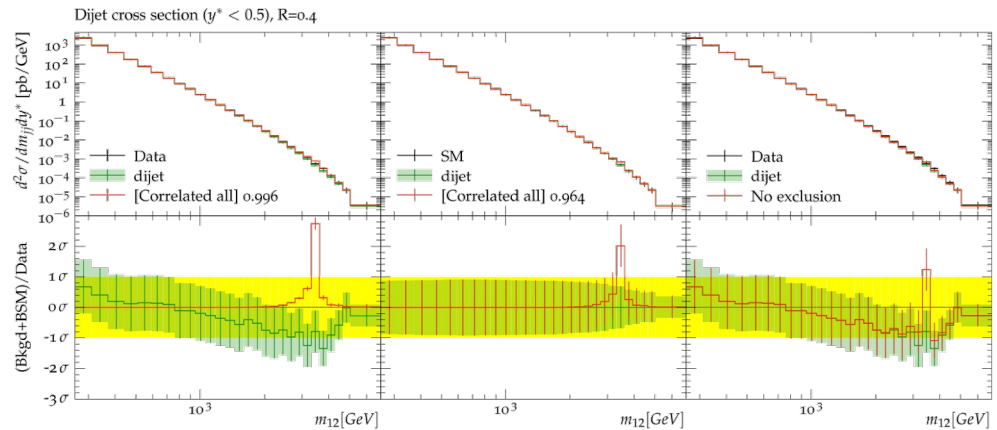
From ATL-PHYS-PUB-2011-008

Introducing **Contur**



“Constraints On New Theories Using Rivet”

- Extend the power of Rivet beyond the Standard Model
- Signal-injection of final-state particles from Beyond-the-SM physics events on to the measured cross sections in Rivet



From Altak
arXiv:2111.

Importance of of
reinterpretability. See
also **CheckMate** for
searches

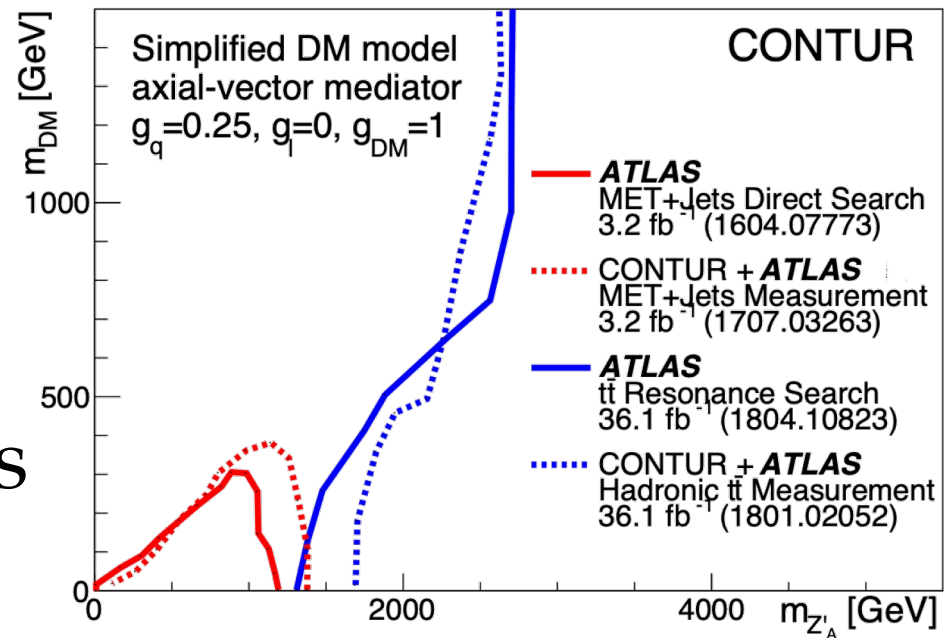
- Increasingly precise measurements also extend to calculations **together**

JMB, Grellscheid, Krämer, Sarrazin, Yallup; Buckley et al
arXiv:1606.05296 (JHEP), arXiv:2102.04377 (SciPost)

Unleashing the power of high luminosity LHC data (*example case studies*)



- Composite Dark Matter
- Back to the W mass
- Vector-like Quarks
- Back to SUSY



Composite Dark Matter Models

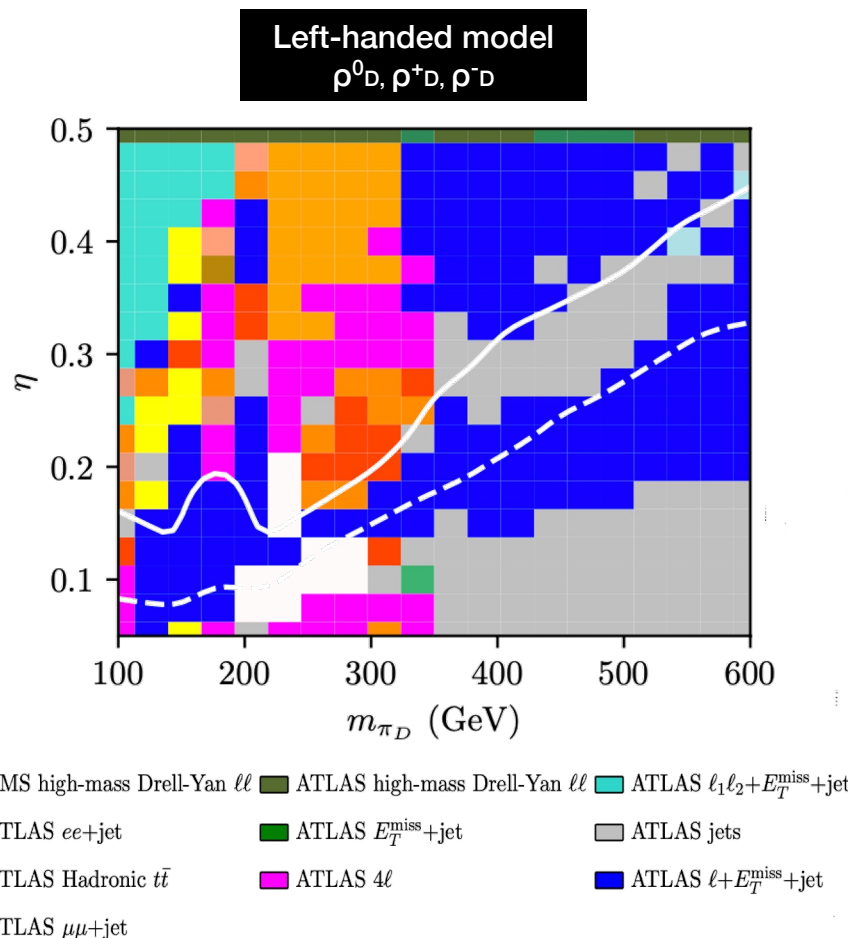


- What if Dark Matter is a composite particle arising from e.g. an $SU(4)$ symmetry which confines at some scale Λ_{dark} ?
- Lead to bound states "dark" mesons and baryons.
 - Kribs et al. [arXiv:1809.10183](https://arxiv.org/abs/1809.10183) (JHEP)
- Dark fermions transform under electroweak part of the Standard Model: communication with SM
- There are **no direct searches** for this model by ATLAS or CMS:
instead to constrain this model using the bank of existing LHC measurements using Contur
- Dynamics of the theory depend a lot on $\eta = m(\pi_D) / m(\rho_D)$

Composite Dark Matter Models



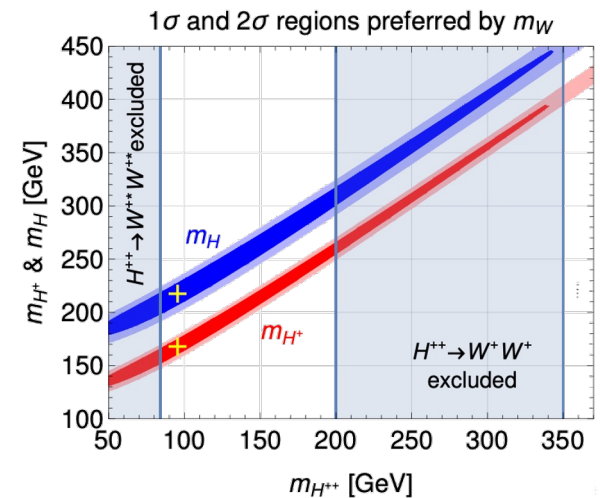
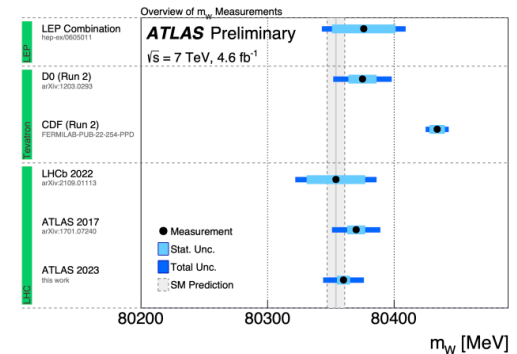
- Large areas excluded:
 - When pion mass is close to Higgs mass, $H \rightarrow gg$ analysis contributes
 - Boosted hadron "top" measurements contribute when pion mass ~ 200 GeV: Pions decay to tb and are boost from heavy r .
 - Other sensitivity from Z -pole dileptons, and lepton+missing energy (Z , top, W production in decay chains)



Back to the W mass



- Numerous possible explanations (not all BSM!)
- *Type-II seesaw* model is well-motivated by other physics (neutrino masses)
 - New scalars $\Delta^{\pm\pm}, \Delta^{\pm}, \Delta^0, \xi^0$
 - Predicts shifts in W/Z mass ratio
 - Shift can be negative or positive depending on scalar mass splittings
 - Region around $100 < M_{\Delta} < 300$ GeV explains M_W , and evades LFV constraints and LHC searches
- Does it survive our “SM” measurements?

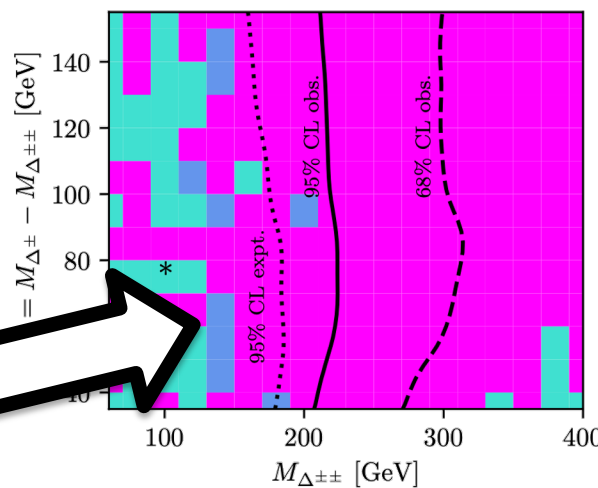
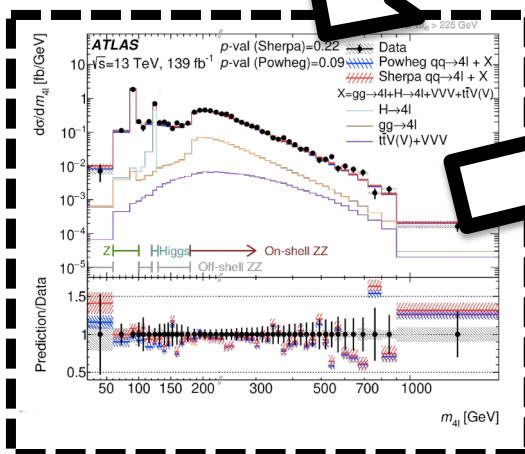
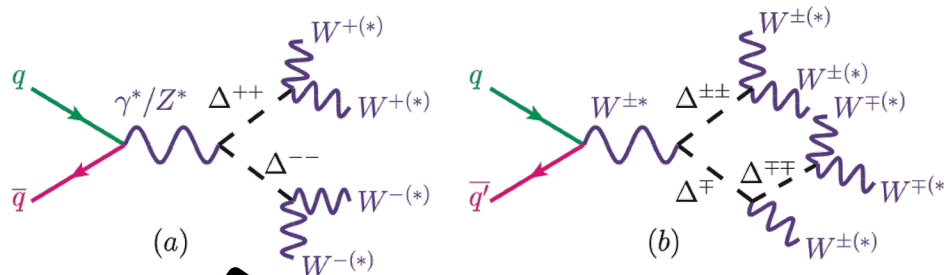


J. Heeck, [2204.10274](#), PRD

Does it survive our “SM” measurements?



- In a word, **no**.
- $\Delta^{\pm\pm}$ production leads to multilepton production
- Would have been visible in several ATLAS measurements, most notably inclusive four-leptons



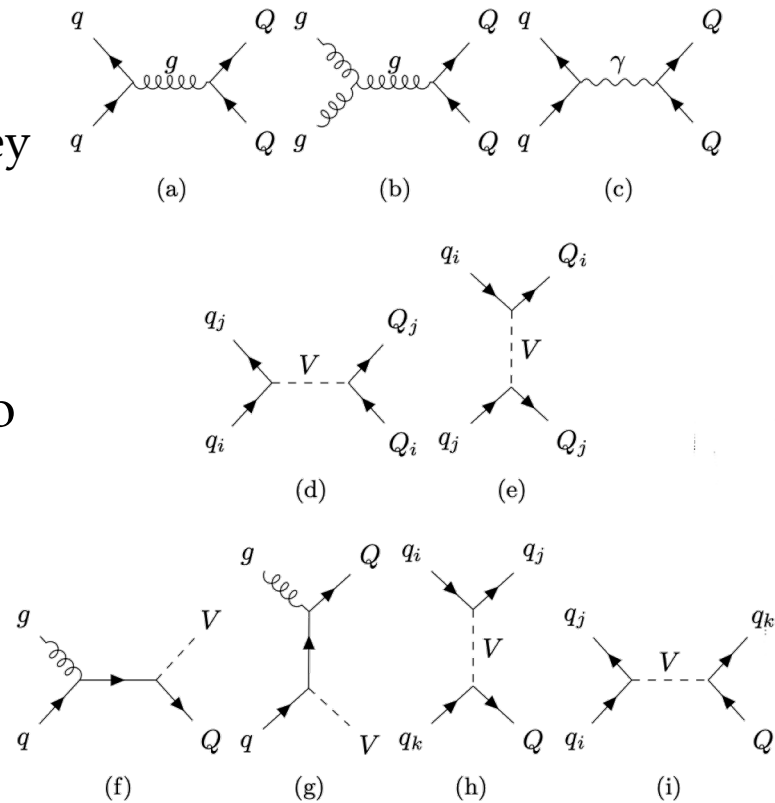
■ ATLAS 4ℓ JHEP,07:005, 2021
■ ATLAS $\ell_1\ell_2+E_T^{\text{miss}}$ EPJC,77(3):141,2017.
■ ATLAS $\ell_1\ell_2+E_T^{\text{miss}}+\text{jet}$ JHEP,06:003, 2021.

JMB, J. Heck, S. Jeon, O. Matteler,
 R. Ruiz [2210.13496](https://arxiv.org/abs/2210.13496) (PRD,
 accepted)

Vector-like Quarks



- Very common extension to SM, general model by *Buchkremer et al* [arXiv:1305.4172 \(NPB\)](https://arxiv.org/abs/1305.4172). Introduces up to four quark partners, B, T, X, Y.
 - Usual strong couplings to SM
 - Evade bounds from Higgs because they are vectors
 - B, T interact with with W, Z, H with modified weak couplings
 - X, Y interact with W (only) similarly
- Three sets of parameters (in additon to masses)
 - κ : absolute coupling of VLQs to SM quarks
 - ζ_i : relative coupling of VLQs to i^{th} generation
 - ξ_v : relative coupling of B,T to V in {W, H, Z}




Vector-like Quarks




- Compare to (quite limited) direct searches: ATLAS limits from *arXiv:1808.02343* (PRL)
- Assumes 3rd generation coupling only, and X, Y are decoupled.
- Only include pair production

EUROPEAN ORGANISATION FOR NUCLEAR RESEARCH (CERN)



Phys. Rev. Lett. 121 (2018) 211801
DOI: 10.1103/PhysRevLett.121.211801



CERN-EP-2018-205
November 26, 2018

Combination of the searches for pair-produced vector-like partners of the third-generation quarks at $\sqrt{s} = 13$ TeV with the ATLAS detector

The ATLAS Collaboration

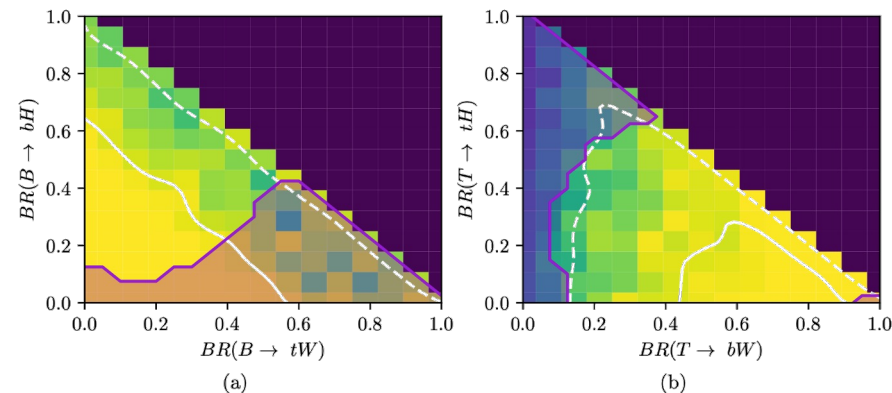
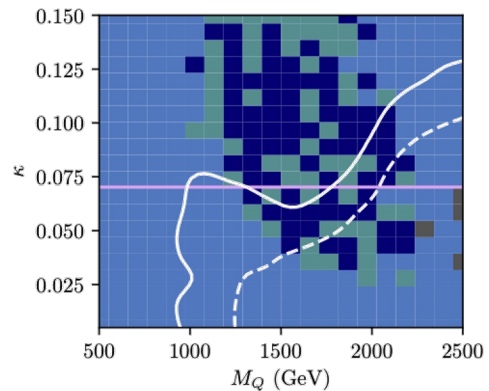


Figure 5: Sensitivity of LHC measurements to (a) B -production for $M_B = 1200$ GeV and (b) T -production for $M_T = 1350$ GeV. The CONTUR exclusion is shown in the bins in which it is evaluated, graduated from yellow through green to black on a linear scale, with the 95% CL (solid white) and 68% CL (dashed white) exclusion contours superimposed. The mauve region is excluded at 95% CL by the ATLAS combination [16].

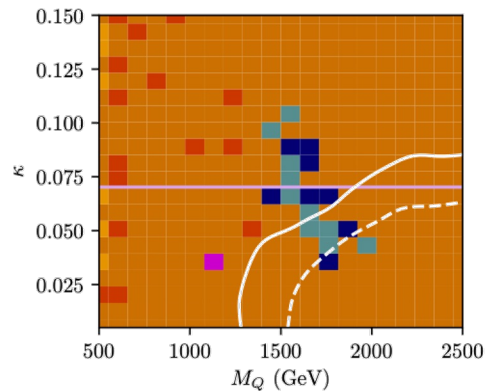
Vector-like Quarks



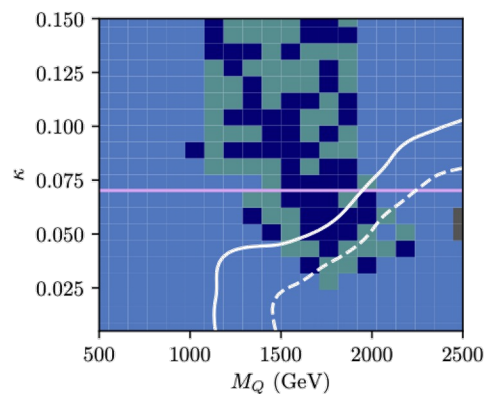
- Coupling to 1st generation.
- Region above line excluded by non-collider constraints
- No LHC search analyses exist
- Measurements exclude most of the plane.
- Single VLQ production very important at highest masses



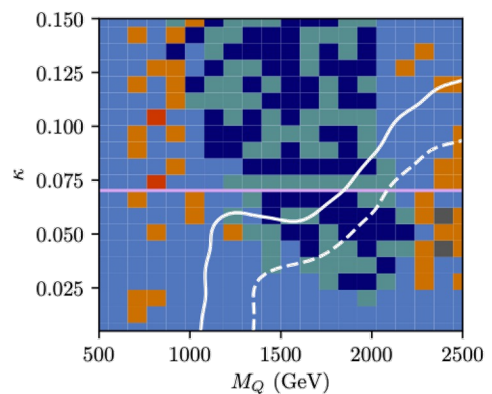
(a)



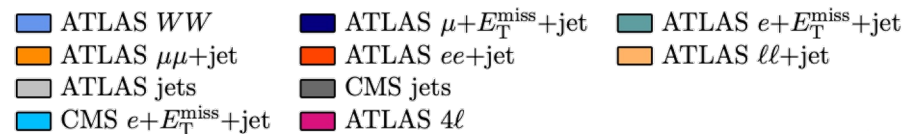
(b)



(c)



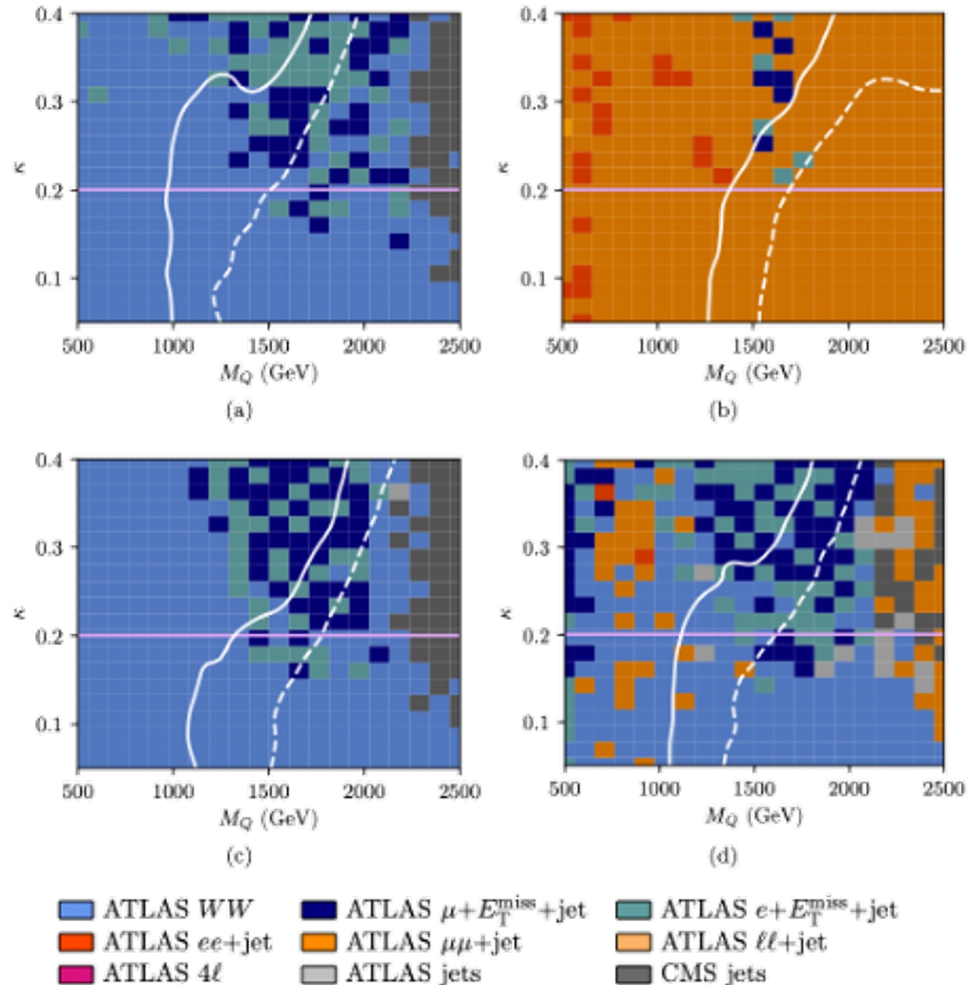
(d)



Vector-like Quarks



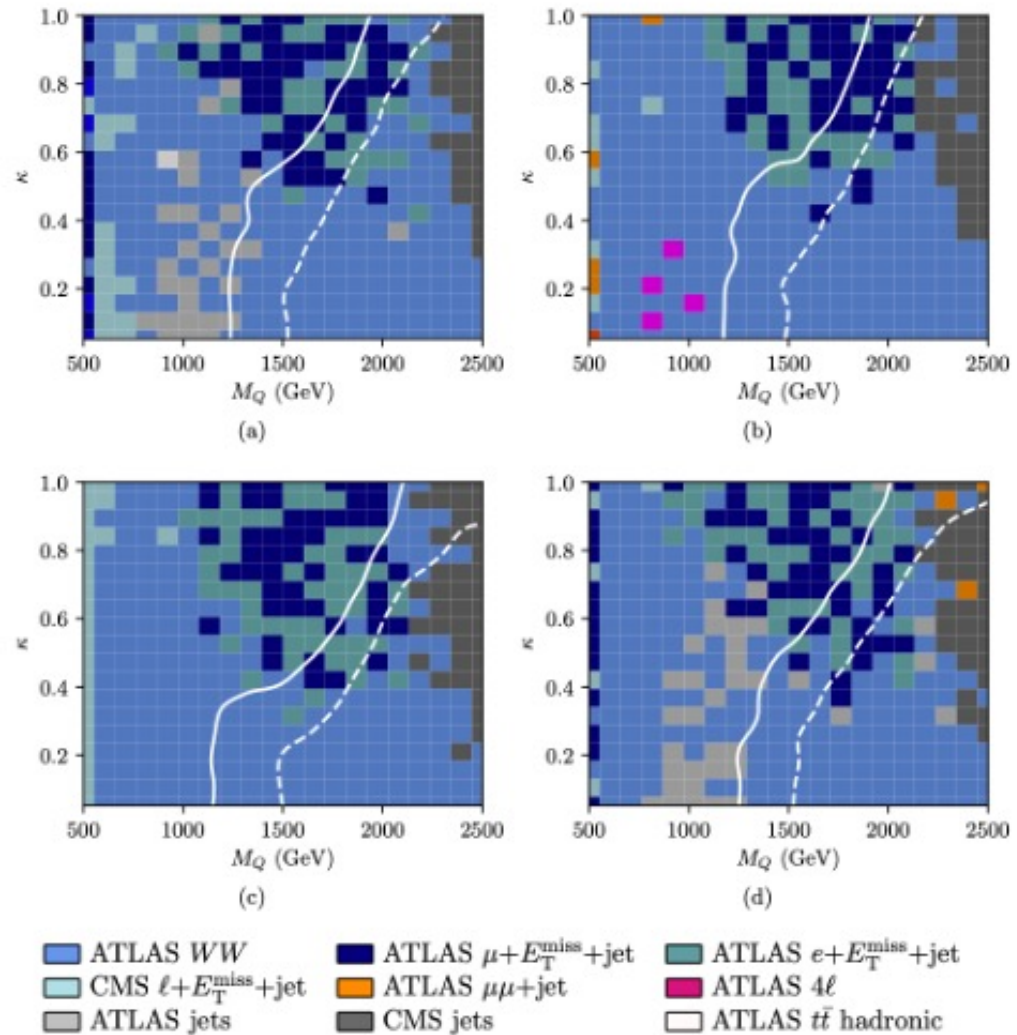
- Coupling to 2nd generation.
- Region above line excluded by non-collider constraints
- No LHC search analyses exist
- Measurements exclude significant part of the plane.
- Single VLQ production again very important at highest masses



Vector-like Quarks



- Coupling to 3rd generation.
- No exclusion from non-collider, but there are several LHC searches
- Measurements also exclude significant part of the plane.
- Single VLQ production still significant at highest masses

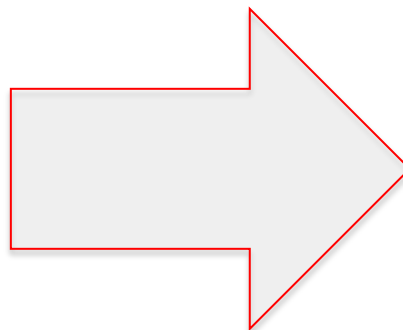


Vector-like Quarks



- Addendum: During journal review for this paper, it was pointed out that we'd missed some of the most compelling scenarios, and should instead consider:
 - B, T singlets
 - BT, XT, TY doublets
 - BYX, BTY triplets
- ... for each generational coupling scenario and for four different decay branching benchmarks to W, Z, H.
- i.e. $7 \times 3 \times 4$ two dimensional parameter scans
- Hmm. A challenge for Contur?

Vector-like Quarks



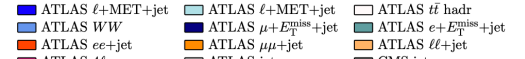
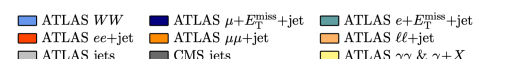
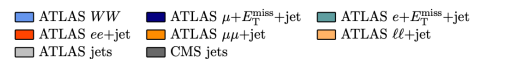
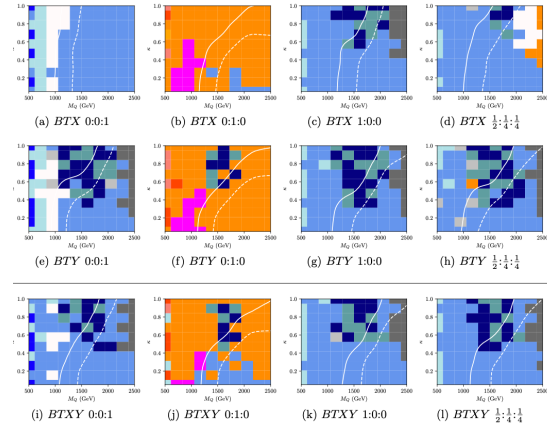
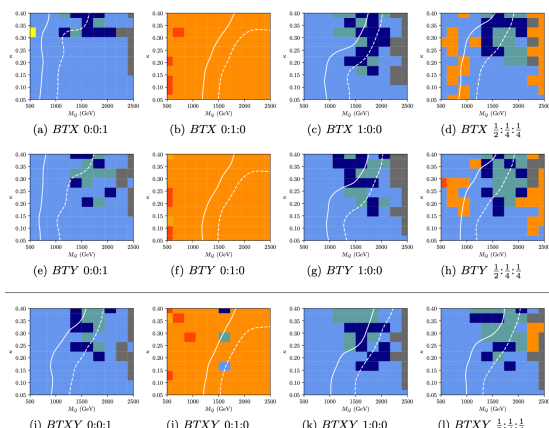
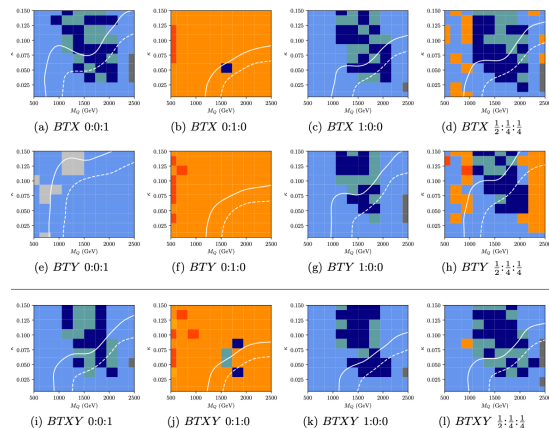
Vector-like Quarks



1st Generation

2nd Generation

3rd Generation



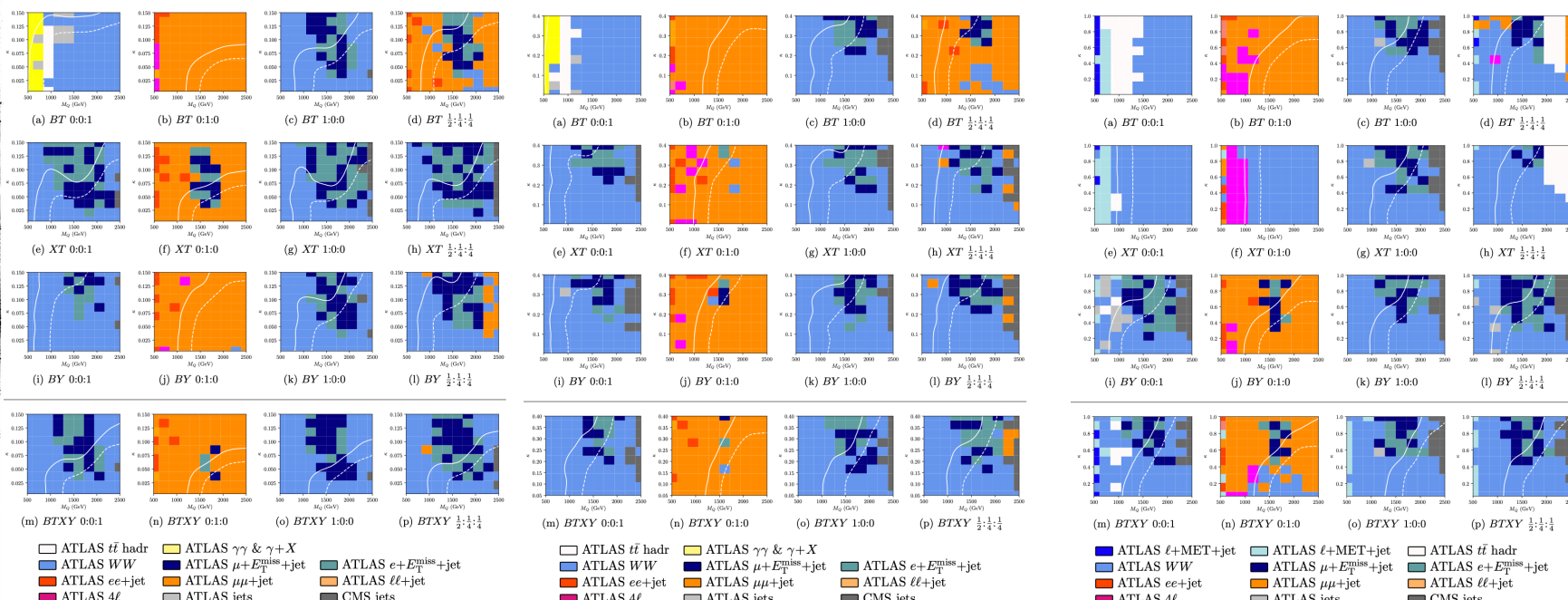
Vector-like Quarks



1st Generation

2nd Generation

3rd Generation



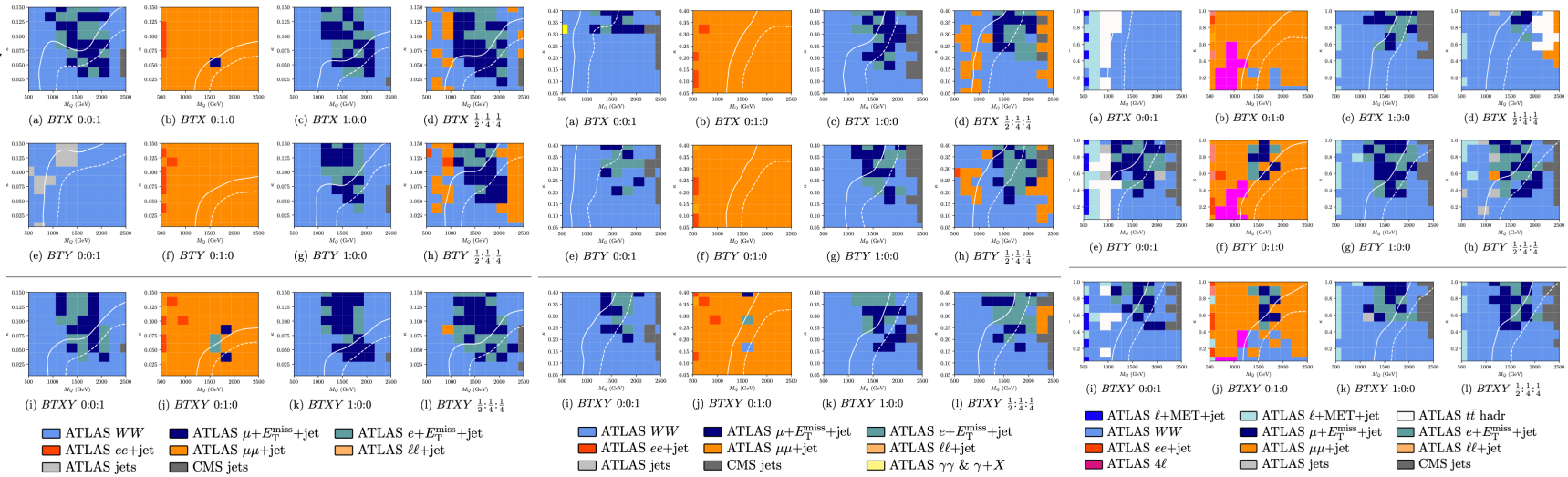
Vector-like Quarks



1st Generation

2nd Generation

3rd Generation



Back to SUSY?



TTP23-009, KCL-PH-TH/2023-21, gambit-physics-23, MCnet-23-05, ADP-23-08/T1217, CERN-TH-2023-043

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Received: date / Accepted: date

Abstract Using the GAMBIT global fitting framework, we constrain the MSSM with an eV-scale gravitino as the lightest supersymmetric particle, and the six electroweakinos (neutralinos and charginos) as the only other light new states. We combine 15 ATLAS and 12

CMS searches at 13 TeV, along with a large collection of ATLAS and CMS measurements of Standard Model signatures. This model, which we refer to as the \tilde{G} -EWMSSM, exhibits quite varied collider phenomenology due to its many permitted electroweakino production processes and decay modes. Characteristic \tilde{G} -EWMSSM signal events have two or more Standard Model bosons and missing energy due to the escaping gravitinos. While much of the \tilde{G} -EWMSSM parameter space is excluded,

we find several viable parameter regions that predict phenomenologically rich scenarios with multiple neutralinos and charginos within the kinematic reach of the LHC during Run 3, or the High Luminosity LHC. In particular, we identify scenarios with Higgsino-dominated electroweakinos as light as 140 GeV that are consistent with our combined set of collider searches and measurements. The full set of \tilde{G} -EWMSSM parameter samples and GAMBIT input files generated for this work is available via Zenodo.

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3	Collider likelihoods	5
	3.1 LHC searches	5

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So where are we now?



- *No agreed extensions to the Standard Model*
 - *But many ideas, connecting various anomalous phenomena*
- *Change of approach required*
 - *This is about **exploration** of new physics territory*
 - ***No guarantee** that Dark Matter, Supersymmetry, or indeed anything else beyond the Standard Model will be within reach*
- *Need precise, theory-independent measurements, and comparable calculations, in Standard Model & beyond. (As well as looking for “outliers”.)*

ANTIMATTER

So where are we now?



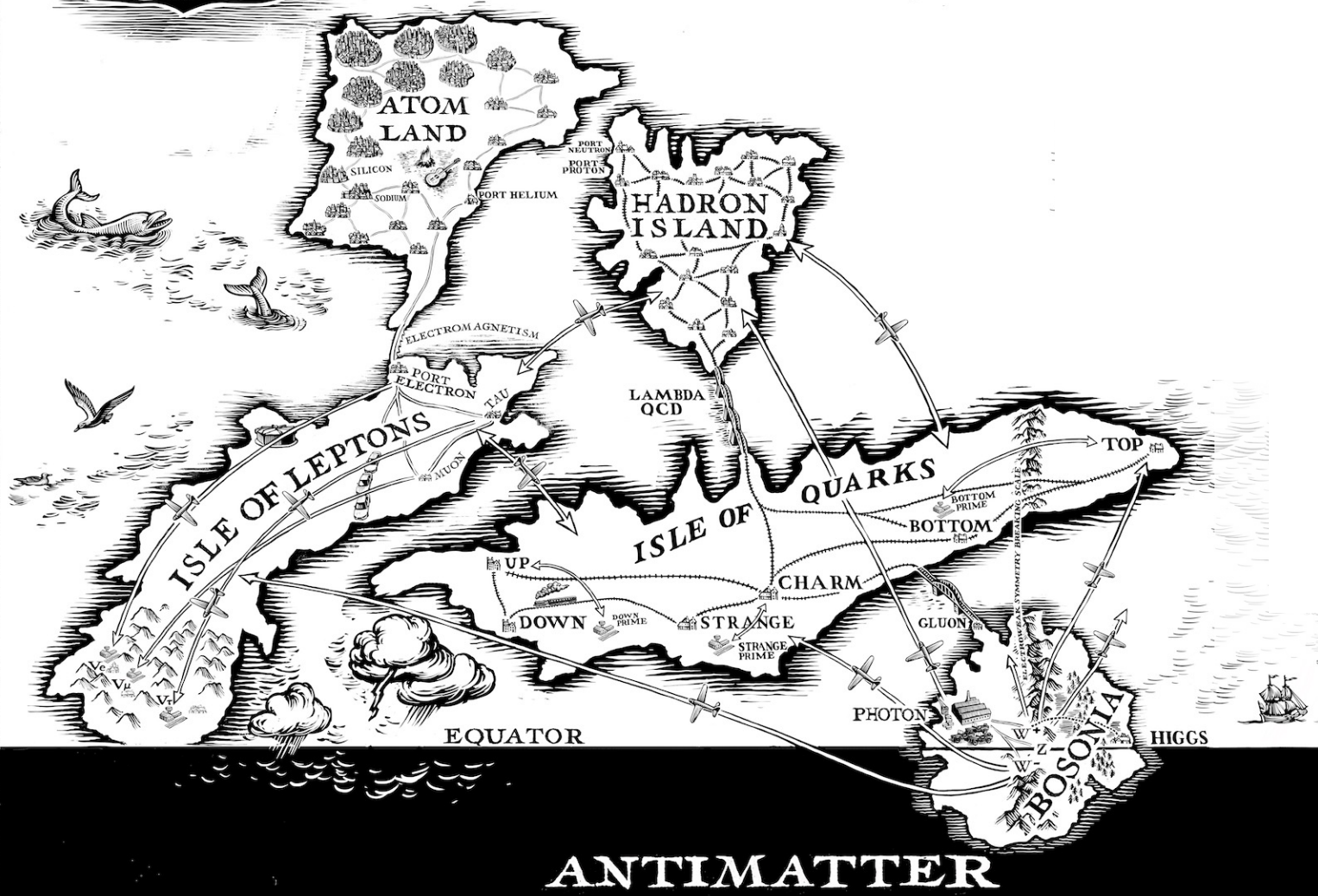
The *guarantee* is we will find out

- Whether the Higgs-self coupling is as the Standard Model predicts
- Whether or not the Standard Model continues to apply, well beyond the region in which it was developed, and to what precision
- We will also push some *amazing technologies*, with likely benefits elsewhere
- *And there may be big surprises!*

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A MAP OF THE
INVISIBLE
FAR EAST

Is the Standard Model Isolated?



A MAP OF THE
INVISIBLE
FAR EAST

Is the Standard Model Isolated?



*We'll never know
if we don't look...*

Maps © Chris Wormell from "A Map of the
Invisible" (Penguin)

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Happy New Decade, Herbi!



'A magnificent, compelling and insightful voyage'
BRIAN COX

JON BUTTERWORTH

JON BUTTERWORTH

Maps by Chris Wormell

QUATOR

QUARKS

ARM

LU