

# Period Geometry of Calabi-Yau $n$ -folds for Feynman integrals

Bethe Forum: Geometries and Special functions for Physics  
and Mathematics

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[1]=arXiv:1912.06201v2, [2]=arXiv:2008.10574v1,

[3]=arXiv:2108.05310, published JHEP

[4]=arXiv:2209.05291 in PRL and [5]=arXiv: 2212.09550 in JHEP,  
in progress

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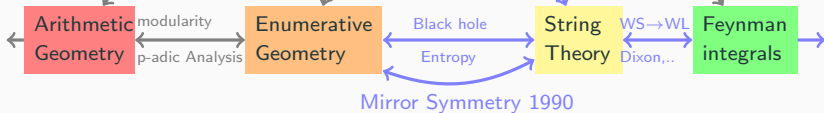
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The latter condition is equivalent to

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- 2) given a Kähler class,  $\exists$  metric  $g$  with  $R_{i\bar{j}}(g) = 0$ ,
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- CY 1-fold is an elliptic curve, say  $y^2 = x(x - 1)(x - z)$  with  $\Omega$  given by  $\frac{dx}{y}$  and  $\omega = \frac{dx}{y} \wedge \frac{d\bar{x}}{\bar{y}}$  is its volume form.

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- We use  **$SU(n)$**  rather than  $\subset SU(n)$  to avoid trivial products of lower CY  $n$ -folds in the generalisation.



# Construction of Calabi-Yau n-folds hypersurface in projective spaces

Let  $M$  be a degree  $\mathcal{N} = dH$  embedding of  $M$  into  $H \subset \mathbb{P}^{n+1}$  .

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$$\frac{(1 + H)^{n+2}}{1 + dH} = 1 + \underbrace{[(n + 2) - d]H}_{c_1(TM)=0!} + \underbrace{[(1 - d)^2 + \frac{1}{2}n(n + 3 - 2d)]H^2}_{c_2(TM)=c_2H^2} + \dots$$

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- $\Rightarrow$  3) quintic in  $\mathbb{P}^4$  is a CY 3-fold with **101 complex moduli**.

## More on constructions of Calabi-Yau n-folds

Number of complex moduli  $\#mon - |Aut(\mathbb{P}^*)|$ :

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- 2) By  $c_2(TM) = 6H^2 \Rightarrow \chi = 24$ . HRR for arithmetic genus of surface  $\chi_0 = \sum_{i=0}^2 (-1)^i h^{0,i} = \frac{1}{12} \int_{M_2} (c_1^2 + c_2)$ . Now by definition  $h^{00} = h^{02} = 1, h^{01} = 0$  because of  $SU(2)$  hol, i.e.  $\chi_0(M_2) = 2$  and since  $c_1 = 0 \Rightarrow \chi(M_2) = 24$  and we have only **one topological type the K3 surface**
- 3) By  $c_3(TM) = -40H^3 \Rightarrow \chi = -200$ . Hirzebruch Riemann Roch  $\chi_0 = \frac{1}{24} \int_{M_3} c_1 c_2 = 1 - 0 + 0 - 1 \checkmark$ ,  
 $\chi_1 = -h^{11} + h^{21} = \frac{1}{24} \int_{M_3} c_1 c_2 - 12c_3 \Rightarrow \chi = 2(h^{11} - h^{21}) \checkmark$

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**Theorem (C.T.C Wall):** The topological type of a Calabi-Yau 3-fold  $M$  is fixed by their Hodge numbers  $(h_{2,1}, h_{1,1})$ , their triple intersection  $D_i \cap D_j \cap D_k \in \mathbb{N}$  and  $c_2(TM) \cdot \dots \cdot D_k, D_k \in H_4(M, \mathbb{Z})$ .

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CICYs: Complete intersections: The vanishing locus of  $r$  polynomials  $P_k = 0, k = 1, \dots, r$  in  $\mathbb{P} = \otimes_{l=1}^m \mathbb{P}^{n_l}$  define a CY  $(\sum_{l=1}^m n_l - r)$ -fold if  $\sum_{k=1}^r d_{kl} = n_l + 1, \forall l = 1, \dots, m$ , with  $d_{kl}$  are degrees of the  $k$ 'th polynomial in the  $l$ 'th factor: **2d n-1 loop bananas**.

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BCs: Branched covers: Let  $\mathbb{P}$  be a  $n$ -dimensional Fano variety with positive canonical class  $K(\mathbb{P}) = c_1(\mathbb{P}) > 0$  then a  $b$ -fold cover that is branched at  $bK(\mathbb{P})$  is a non necessarily smooth CY  $n$ -fold: **2d n loop fishnets** .

# Hypersurface in toric ambient spaces

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Huge classified construction: CY 3folds as hypersurfaces in toric ambient spaces  $\mathbb{P}_\Delta, \mathbb{P}_{\hat{\Delta}}$  defined by reflexive pairs  $(\Delta, \hat{\Delta})$ :

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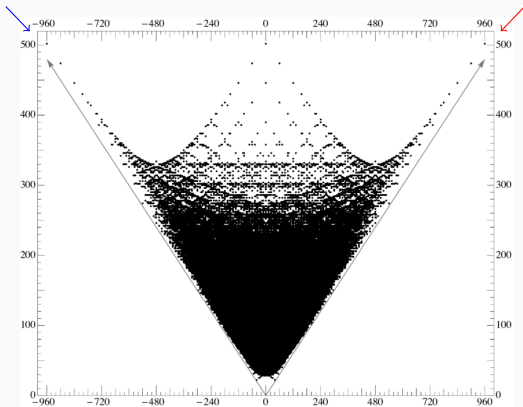
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$$M = \{[\rho_{\hat{\Delta}} = 0] = [\sum_i H_i] \subset \mathbb{P}_{\hat{\Delta}}\}$$

Batyrev:  
(M, W) mir-  
ror pairs

$$W = \{[\rho_{\Delta} = 0] = [\sum_i \hat{H}_i] \subset \mathbb{P}_{\Delta}\}$$





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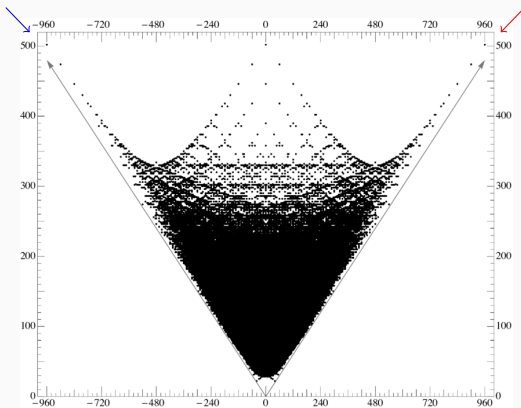
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Geom.	$\#(\Delta_k, \hat{\Delta}_k)$
2 pts	1
ell crv	16
K3	$4319^1$
CY 3-flds	$473800776^1$
CY 4-flds	$\mathcal{O}(10^{22})?$
$\vdots$	$\vdots$

<sup>1</sup> Kreuzer & Skarke

'02,  $k = 3, 4$



## General properties of Calabi-Yau $n$ -fold families

**Theorem Tian/Todorov:** The complex moduli space  $\mathcal{M}_{cs}(M)$  of a CY  $n$ -fold  $M$  is parametrized for by  $h^{n-1,1} = \dim_{\mathbb{C}}(H^{n-1,1}(M))$  globally unobstructed complex deformation parameters  $z$ , i.e. is a manifold of complex dimension  $h^{n-1,1} =: r$  ( $E$  and  $K_3$  are special).

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**Example:** We counted  $101 = h^{2,1}$  complex deformation parameters for the quintic in  $\mathbb{P}^4$  and by the Lefschetz hyperplane theorem  $h^{1,1} = 1$  (inherited from  $\mathbb{P}^4$ ), hence  $\chi = 2(h^{1,1} - h^{2,1}) = -200$  in accordance with Gauss-Bonnet.

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**Application:** The complex moduli dependent **period integrals** on CY n-fold families **generalize elliptic functions**. They are identified for important examples with the maximal cut Feynman  $n - 1$ -loop integrals, where the complex moduli  $z$  are identified with the scale invariant physical parameters  $z_i = p^2/m_i^2, \dots$

# Periods on Calabi-Yau n-folds

## Periods integrals

$$\Pi_{ij}(z) = \int_{\Gamma_i} \gamma^j(z)$$

define a non-degenerate pairing between (middle) homology and (middle) cohomology well defined by the theorem of Stokes:

$$\Pi : H_n(M_n, \mathbb{K}) \times H^n(M_n, \mathbb{C}) \rightarrow \mathbb{C}.$$

It is possible and natural to have  $\mathbb{K}$  to be  $\mathbb{Z}$ . There is an intersection pairing

$$\Sigma : H_n(M_n, \mathbb{K}) \times H_n(M_n, \mathbb{K}) \rightarrow \mathbb{K},$$

that can be made in particular integral. If  $n$  is odd  $\Sigma$  is antisymmetric and can be made symplectic. If  $n$  is even  $\Sigma$  is a symmetric on the even self dual lattice  $H_n(M_n, \mathbb{K})$ . E.g. for  $K3$   $b_2 = 22$  and

$$\sigma = b_2^+ - b_2^- = \frac{1}{3} \int_{M_2} c_1^2 - 2c_2 = -16 \text{ hence } b_2 \text{ has signature } (3, 19) \text{ and is } E_8(-1)^{\oplus 2} \oplus \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}^{\oplus 3}.$$

If  $n$  is odd we fix an integral symplectic basis  $\underline{\Gamma} = \{A_I, B^I\}$ ,  $I = 0, \dots, r$  with  $\text{Span}_{\mathbb{Z}}(\underline{\Gamma}) = H_n(W, \mathbb{Z})$  and

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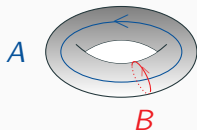
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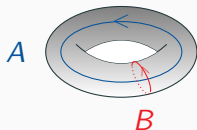


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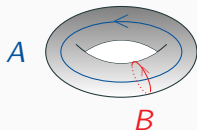


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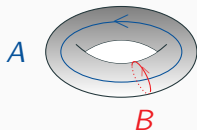
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Well studied in part because they solve Keplers problem

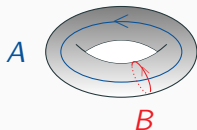
Periods annihilated by Picard-Fuchs (1881) 2cd order linear operator  $L^{(2)}$ .

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Periods annihilated by Picard-Fuchs (1881) 2cd order linear operator  $L^{(2)}$ .

$$L^{(2)} \int_{\Gamma} \Omega = \left[ (1-z)\partial_z^2 + (1-2z)\partial_z - \frac{1}{4} \right] \int_{\Gamma} \Omega = 0.$$

## The Picard-Fuchs differential ideal

We can always expand  $\Omega = \sum_{i=1}^{b_3(W)} \Pi_i(z) \gamma_i$  in terms of periods  
 $\Pi_i(z) = \int_{\Gamma_i} \Omega(z)$ .

The  $b_n(M_n)$  periods span a vector space that is identified with the solutions space of linear Picard-Fuchs differential ideal  $\mathcal{L}\Pi_i(z) = 0$ .

For one parameter families  $\mathcal{L}$  is generated by a  $b_n(M_n) + 1$  order Picard-Fuchs operator  $L^{(b_n(M_n)+1)}$ , while for multiparameter families  $\mathcal{L} = \{L_i^{(k)}, i = 1, \dots, |\mathcal{L}|, k = 2, \dots, b_n(M_n) + 1\}$  with several  $L_i^{(k)}$ .

The latter can be derived using the Griffiths reduction method and for CY embedded in toric ambient space also as a reduction of a Gelfand Kapranov Zelevinskii system.

## Finding and integral basis

The Feynman integrals correspond i.a. to periods over integral cycles, e.g.  $\underline{\Gamma} = \{A_I, B^I\}$ . Such are not specified by  $\mathcal{L}$  alone.

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If  $\overline{\mathcal{M}_{cs}(M_n)}$  has a point of maximal unipotent monodromy (MUM) with a known mirror  $W_n$  one can calculate an integral period vector using the  $\hat{\Gamma}(TW_n)$ -class.

## One Parameter CY 3 fold operators

**Examples:** There are 14 hyper geometric  ${}_3F_4$  CY 3-fold operators given by

$$L^{(4)} = \theta^4 - \mu^{-1}z \prod_{k=1}^4 (\theta + a_k) ,$$

where  $\theta = z \frac{d}{dz}$  and  $z$  parametrizes  $\mathcal{M}_{cs}(M_3) = \mathbb{P}^1 \setminus \{0, \mu, \infty\}$



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#	$W$	$\kappa$	$c_2 \cdot D$	$\chi(W)$	$a_1, a_2, a_3, a_4$	$\mu^{-1}$	$dT_\infty$
1	$X_5(1^5)$	5	50	-200	$\frac{1}{5}, \frac{2}{5}, \frac{3}{5}, \frac{4}{5}$	$5^5$	$O_5^{\text{DG}}$
2	$X_{4,2}(1^6)$	8	56	-176	$\frac{1}{4}, \frac{1}{2}, \frac{1}{2}, \frac{3}{4}$	$2^{10}$	$C_4$
3	$X_{3,3}(1^6)$	9	54	-144	$\frac{1}{3}, \frac{1}{3}, \frac{2}{3}, \frac{2}{3}$	$3^6$	$K_3$
4	$X_{2,2,2,2}(1^8)$	16	64	-128	$\frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}$	$2^8$	$M_2$
$\vdots$	$\vdots$	$\vdots$	$\vdots$	$\vdots$	$\vdots$	$\vdots$	$\vdots$

# One Parameter CY 3 fold operators

Their Riemann symbols are

$$\mathcal{P} \left\{ \begin{array}{c} 0 \quad \mu \quad \infty \\ \hline 0 \quad 0 \quad a_1 \\ 0 \quad 1 \quad a_2 \\ 0 \quad 1 \quad a_3 \\ 0 \quad 2 \quad a_4 \end{array} \right\} .$$

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At  $z = 0$  the local exponents are completely degenerate and we have a MUM point. A Frobenius  $\mathbb{C}$ -basis of solutions is

$$\vec{\Pi}_0(z) = \begin{pmatrix} f_0(z) \\ f_0(z) \log(z) + f_1(z) \\ \frac{1}{2} f_0(z) \log^2(z) + f_1(z) \log(z) + f_2(z) \\ \frac{1}{6} f_0(z) \log^3(z) + \frac{1}{2} f_1(z) \log^2(z) + f_2(z) \log(z) + f_3(z) \end{pmatrix}$$

for power series normalized by  $f_0(0) = 1$  and  $f_1(0) = f_2(0) = f_3(0) = 0$ .

# One Parameter CY 3 fold operators

The  $\hat{\Gamma}(TW)$  class determines an integral basis at  $z = 0$

$$\vec{\Pi} = \begin{pmatrix} \int_{B^0} \Omega \\ \int_{B^1} \Omega \\ \int_{A_0} \Omega \\ \int_{A_1} \Omega \end{pmatrix} = (2\pi i)^3 \begin{pmatrix} \frac{\zeta(3)\chi(M)}{(2\pi i)^3} & \frac{c_2 \cdot D}{24 \cdot 2\pi i} & 0 & \frac{\kappa}{(2\pi i)^3} \\ \frac{c_2 \cdot D}{24} & \frac{\sigma}{2\pi i} & -\frac{\kappa}{(2\pi i)^2} & 0 \\ 1 & 0 & 0 & 0 \\ 0 & \frac{1}{2\pi i} & 0 & 0 \end{pmatrix} \Pi_0. \quad (1)$$

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in terms of the C.T.C Wall data. The monodromies in  $SP(4, \mathbb{Z}) = O(\Sigma, \mathbb{Z})$  are generated by

$$M_0 = \begin{pmatrix} 1 & -1 & \frac{\kappa}{6} + \frac{c_2 \cdot D}{12} & \frac{\kappa}{2} + \sigma \\ 0 & 1 & \sigma - \frac{\kappa}{2} & -\kappa \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 1 & 1 \end{pmatrix}, \quad M_\mu = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ -1 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

with  $M_\infty = (M_0 M_\mu)^{-1}$ . Note that by HRR  $\frac{\kappa}{6} + \frac{c_2 \cdot D}{12} = \chi(\mathcal{O}_D) + 1 \in \mathbb{Z}$

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One parameter CY 3-fold differential operators  $L^{(4)} = \sum_{i=0}^4 c_i(z) \partial_z^i$  have been classified by [Almkvist, Enckevort, van Straten and Zudilin \(AESZ list\)](#) at least to finite order in  $c_i(z)$  in  $z$ . E.g. the AESZ34 operator

$$L^{(4)} = 1 - 5z - (4 - 28z)\theta + (6 - 63z + 26z^2 - 225z^3) \theta^2 - (4 - 70z + 450z^3) \theta^3 + (1 - z)(1 - 9z)(1 - 25z)\theta^4$$

with Riemann symbol

$$\mathcal{P}_4 \left\{ \begin{array}{ccccc} 0 & \frac{1}{25} & \frac{1}{9} & 1 & \infty \\ 1 & 0 & 0 & 0 & 0 \\ 1 & 1 & 1 & 1 & 0 \\ 1 & 1 & 1 & 1 & 1 \\ 1 & 2 & 2 & 2 & 1 \end{array} \right\}$$

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corresponds to the 4 loop equal mass Banana maximal cut integral with  $z = m^2/p^2$ . Which itself is the diagonal specialisation of the five parameter GKZ system of the complete inter section of two degree  $d_{1,k} = (1, 1, 1, 1, 1)$ ,  $d_{2,k} = (1, 1, 1, 1, 1)$  constraints in  $(\mathbb{P}^1)^5$  describing the general mass case  $z_i = m_i^2/p^2$ .

## Period geometry on CY n-fold

The main constraints which govern the period geometry of CY-folds are **the Riemann bilinear** relations

$$e^{-K} = i^{n^2} \int_{M_n} \Omega \wedge \bar{\Omega} > 0 \quad (2)$$

defining the real positive exponential of the **Kähler potential**  $K(z)$  for the **Weil-Peterssen metric**  $G_{i\bar{j}} = \partial_{z_i} \bar{\partial}_{\bar{z}_j} K(z)$  on  $\mathcal{M}_{CS}(M_n)$ .



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$$\int_{M_n} \Omega \wedge \underline{\partial}_{I_k}^k \Omega = \begin{cases} 0 & \text{if } k < n \\ C_{I_n}(z) & \text{if } k = n . \end{cases} \quad (3)$$

Here  $\underline{\partial}_{I_k}^k \Omega = \partial_{z_{I_1}} \dots \partial_{z_{I_k}} \Omega \in F^{n-k} := \bigoplus_{p=0}^k H^{n-p,p}(W)$  are arbitrary combinations of derivatives w.r.t. to the  $z_i$ ,  $i = 1, \dots, r$ .

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The  $C_{I_n}(z)$  are rational functions labelled by  $I_n$  up to permutations.  
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**Remark 1:** W.r.t the Hodge decomposition the pairings (2) and (3) have the property that if  $\alpha_{m,n} \in H^{m,n}(M_n)$  and  $\beta_{p,q} \in H^{r,s}(M_n)$  then  $\int_W \alpha_{m,n} \wedge \beta_{p,q} = 0$  unless  $m + p = n + q = 3$ .

**Remark 2:** In terms of a basis of periods compatible with  $\Sigma$  they can be written as

$$\int_{M_n} \Omega \wedge \bar{\Omega} = \vec{\Pi}^\dagger \Sigma \vec{\Pi}, \quad \int_{M_n} \Omega \wedge \underline{\partial}_{I_k}^k \Omega = -\vec{\Pi}^T \Sigma \underline{\partial}_{I_k}^k \vec{\Pi} ,$$

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We will focus here on one aspect of the latter that is relevant for Feynman integrals for the following reason: While the maximal integrals are periods and as such solutions of the homogeneous differential equations  $\mathcal{L}\Pi = 0$  the actual Feynman integral is a solution of an inhomogeneous extension  $\mathcal{L}\Pi = g(z)$ .

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When determining the inhomogeneous solutions by the variation of constants method one considers the Wronskian  $[W(z)]_{i,j} = \partial_z^i \Pi_j$ ,  $i, j = 0, \dots, r$  and in particular its inverse.

## A simple consequence of Griffiths transversality

Let us define the skew symmetric matrix

$$Z = W\Sigma W^T, \quad \text{i.e. } [Z(z)]_{ij} = \partial_z^i \Pi^T \Sigma \partial_z^j \Pi, \quad \text{for } i, j = 0, \dots, r$$

Then (3) implies that  $Z$  is rational and its entries are calculated recursively from derivatives of (3) using  $\Pi^T \Sigma \partial_{I_k}^k \mathcal{L} \Pi = 0$ .



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E.g. for the one parameter case with  $r = 3$ , with the abbreviations  $C = C_{111}$ ,  $C' = \partial_z C$  one finds

$$Z^{-1} = \frac{(2\pi i)^3}{C} \begin{pmatrix} 0 & \frac{C''}{C} - 2\frac{C'}{C} + \frac{\varrho_2}{c_4} & -\frac{C'}{C} & 1 \\ 2\frac{C'}{C} - \frac{C''}{C} - \frac{\varrho_2}{c_4} & 0 & -1 & 0 \\ \frac{C'}{C} & 1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix}$$

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$$\Rightarrow W^{-1} = \Sigma W^T Z^{-1} .$$

depends up to rational functions linear on the periods and its derivatives and the inhomogeneous solution becomes an iterated integral.

## Special coordinates in special geometry

The local **Torelli theorem** states that a sufficiently small domain  $U_* \subset \mathcal{M}_{cs}(W)$  can be identified with a chart in  $\mathbb{P}^r$  using the period map  $z \mapsto (X_*^0(z) : \dots : X_*^r(z)) \in \mathbb{P}^r$  and parametrized by inhomogeneous coordinates  $t_*^i(z) = X_*^i / X_*^0$ . Clearly the  $P_I^* = \int_{B_I} \Omega$  are then **homogeneous functions** of the  $X_*^I = \int_{A^I} \Omega$ .

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$$P_i^*(X_*) = \frac{\partial}{\partial X_*^I} F^*(X_*), \quad 2F^*(X_*) = X_*^I P_i^*(X_*),$$

where the Newton equation with Einstein sum conventions implies that  $F(X_*)$  is of degree two in the  $X_*^I$ .

## Special coordinates in special geometry

Writing  $F_*(X_*) = \mathcal{F}_*(t_*)(X_*^0)^2 \vec{\Pi} = (P_l^*, X_*^l)^t$  becomes  $\vec{\Pi} = X^0(2\mathcal{F}_* - t^i \partial_{t_*^i} \mathcal{F}_*, \partial_{t_*^i} \mathcal{F}_*, 1, t_*^i)^t$ , and inserting this into (??), changing variables from  $z_k$  to  $t^k$  and using the transformations properties of  $C_{z_i z_j z_k}$  one establishes  $C_{t_*^i t_*^j t_*^k} = \partial_{t_*^i} \partial_{t_*^j} \partial_{t_*^k} \mathcal{F}_*(t_*)$

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$$\partial_{t_*^i} \begin{pmatrix} \mathcal{V}_0 \\ \mathcal{V}_j \\ \mathcal{V}^j \\ \mathcal{V}^0 \end{pmatrix} = \begin{pmatrix} 0 & \delta_{ik} & 0 & 0 \\ 0 & 0 & C_{ijk} & 0 \\ 0 & 0 & 0 & \delta_i^j \\ 0 & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} \mathcal{V}_0 \\ \mathcal{V}_k \\ \mathcal{V}^k \\ \mathcal{V}^0 \end{pmatrix}.$$

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This is the **Gauss-Manin connection** in projective flat coordinates and in **special Kähler gauge**. These formulas allow simpler iterated integrals and generalise to all  $n$  provided one knows  $\Sigma$ .

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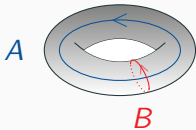
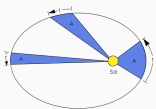
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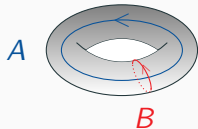
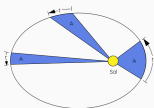
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So far we explored mainly the holomorphic story following from (3) and not the combinations of (3) with (2) and *Remark 1*. This leads to an **integrable structure**,  **$tt^*$ -equations**, **Kodaira Spencer gravity**, **topological string theory** and related topics ....

# Conclusion and Outlook

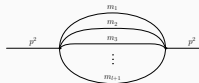


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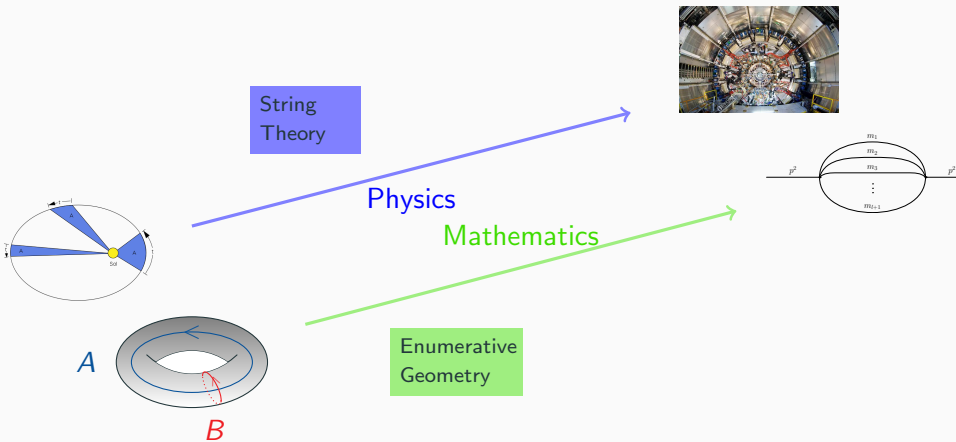


Physics

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