

BEAM ASYMMETRY Σ FOR:

$\gamma p \rightarrow \eta p$

$\gamma p \rightarrow \eta' p$

WORK IN PROGRESS

η/η' DECAYING INTO:

$$\eta \rightarrow 2\gamma$$

$$\eta \rightarrow 3\pi^0 \rightarrow 6\gamma$$

$$\eta \rightarrow \pi^+ \pi^- \pi^0$$

$$\eta' \rightarrow 2\gamma$$

$$\eta' \rightarrow \pi^0 \pi^0 \eta \rightarrow 6\gamma$$

$$\eta' \rightarrow \pi^+ \pi^- \eta$$



Conception of a new technique for the simultaneous analysis of:

- several data taking periods
 - several meson decay channels with different efficiencies
- (also different polarization degrees)



Particle identification also in BGO-Barrel required for the analysis of charged multi-pion final states

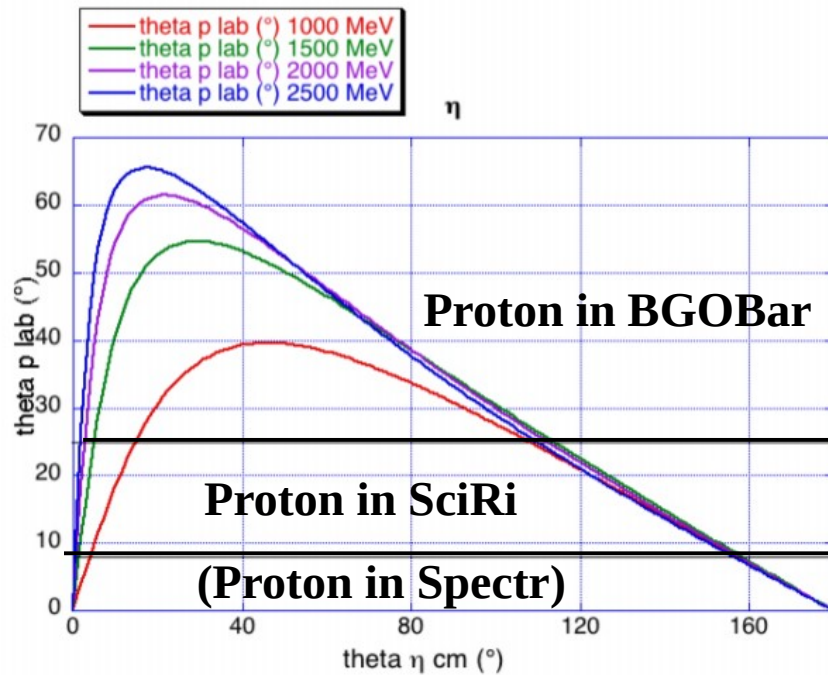
- barrel calibration
- proton/pion discrimination (dE/dx vs. E_{BGO})

Kinematical reconstruction technique based on energy and momentum conservation



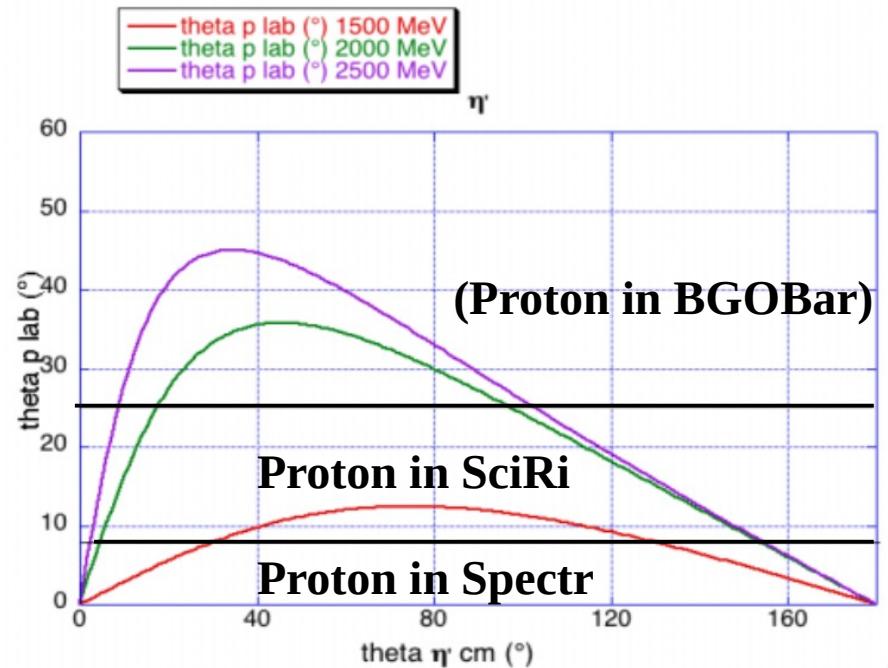
Extraction of the polarization degree of new data taking periods

η KINEMATICS



Central proton → BGOBar
 Interm proton → SciRi
 (Forward proton → Spectr)

η' KINEMATICS



Interm proton → SciRi
 Forward proton → Spectr

**TWO OR MORE DATA TAKING PERIODS P_i
WITH DIFFERENT POLARIZATION DEGREES AND EFFICIENCIES**

Ex. P_1 AND P_2 WITH

$$P_{Pol\pm}^{P_1} \neq P_{Pol\pm}^{P_2} \quad \text{AND} \quad \varepsilon^{P_1} \neq \varepsilon^{P_2}$$

IDEAL SOLUTION

- Extract the asymmetry from each single period
- Make the weighted average of the results

BUT

if the statistics of each period is not enough

=> Another strategy:

- Sum up the statistics of each period IN A PROPER WAY (see after)
- Extract the asymmetry

ONE DATA TAKING PERIOD P WITH:

$$P_{Pol+}^P = P_{Pol-}^P = P^P$$

$$N_{Pol\pm}^P = F_{Pol+}^P \left(\frac{d\sigma}{d\Omega} \right)_{UNP} \varepsilon^P(\varphi) N_{SC} (1 \pm P^P \Sigma \cdot \cos(2\varphi)) \quad (1a)$$

where: P = period

$$N_{UNP,norm}^P = \frac{N_{Pol+}^P}{F_{Pol+}^P} + \frac{N_{Pol-}^P}{F_{Pol-}^P} \propto 2\varepsilon(\varphi) \quad (1b)$$

$$\frac{\frac{N_{Pol+}^P}{F_{Pol+}^P}}{\frac{N_{Pol+}^P}{F_{Pol+}^P} + \frac{N_{Pol-}^P}{F_{Pol-}^P}} = \frac{\left(\frac{d\sigma}{d\Omega} \right)_{UNP} \varepsilon^P(\varphi) N_{SC} (1 + P^P \Sigma \cdot \cos(2\varphi))}{\left(\frac{d\sigma}{d\Omega} \right)_{UNP} \varepsilon^P(\varphi) N_{SC} 2}$$

NUMERATOR



DENOMINATOR



$$\frac{\frac{N_{Pol+}^P}{F_{Pol+}^P}}{N_{UNP,norm}^P} = \frac{1}{2} (1 + P^P \Sigma \cdot \cos(2\varphi)) \quad (2)$$

ONE DATA TAKING PERIOD P WITH:

$$P_{Pol+}^P \neq P_{Pol-}^P$$

NUMERATOR

$$\rightarrow \frac{N_{Pol+}^P}{F_{Pol+}^P}$$

where: P = period

DENOMINATOR

$$\rightarrow N_{UNP,norm}^P = \frac{2}{P_{Pol+}^P + P_{Pol-}^P} \left(P_{Pol-}^P \frac{N_{Pol+}^P}{F_{Pol+}^P} + P_{Pol+}^P \frac{N_{Pol-}^P}{F_{Pol-}^P} \right) \quad (3)$$

which has no dependence on $\cos(2\varphi)$. It only depends on φ through the efficiency ε :

$$N_{UNP,norm}^P \propto 2\varepsilon^P(\varphi) \quad (4)$$

NUMERATOR

$$\rightarrow \frac{N_{Pol+}^P}{F_{Pol+}^P}$$

DENOMINATOR

$$\frac{F_{Pol+}^P}{N_{UNP,norm}^P} = \frac{1}{2} \left(1 + P_{Pol+}^P \Sigma \cdot \cos(2\varphi) \right) \quad (5)$$

**TWO (OR MORE) DATA TAKING PERIODS P_i
WITH DIFFERENT POLARIZATION DEGREES AND EFFICIENCIES**

Ex. P_1 AND P_2 WITH

$$P_{Pol\pm}^{P_1} \neq P_{Pol\pm}^{P_2} \quad \text{AND} \quad \varepsilon^{P_1} \neq \varepsilon^{P_2}$$

NUMERATOR

$$\longrightarrow \frac{N_{Pol+}^{P_1}}{F_{Pol+}^{P_1}} + \frac{N_{Pol+}^{P_2}}{F_{Pol+}^{P_2}} \quad (6)$$

The (6) has the following behaviour in the azimuthal angle:

$$\frac{N_{Pol+}^{P_1}}{F_{Pol+}^{P_1}} + \frac{N_{Pol+}^{P_2}}{F_{Pol+}^{P_2}} \propto [\varepsilon^{P_1}(\varphi) + \varepsilon^{P_2}(\varphi)] [1 + P_{Pol+}^W \Sigma \cos(2\varphi)] \quad (7)$$

$$P_{Pol+}^W = \frac{\varepsilon^{P_1} P_{Pol+}^{P_1} + \varepsilon^{P_2} P_{Pol+}^{P_2}}{\varepsilon^{P_1} + \varepsilon^{P_2}} \quad (8)$$

DENOMINATOR

$$\longrightarrow N_{UNP,norm}^{P_1} + N_{UNP,norm}^{P_2} \propto 2[\varepsilon^{P_1}(\varphi) + \varepsilon^{P_2}(\varphi)] \quad (9)$$

where:
$$N_{UNP,norm}^{P_i} = \frac{2}{P_{Pol+}^{P_i} + P_{Pol-}^{P_i}} \left(P_{Pol-}^{P_i} \frac{N_{Pol+}^{P_i}}{F_{Pol+}^{P_i}} + P_{Pol+}^{P_i} \frac{N_{Pol-}^{P_i}}{F_{Pol-}^{P_i}} \right)$$

$$N_{UNP,norm}^{P_i} \propto 2\varepsilon^{P_i}(\varphi)$$

where: $P_i = i\text{-th period}$

NUMERATOR

DENOMINATOR

$$\frac{\frac{N_{Pol+}^{P1}}{F_{Pol+}^{P1}} + \frac{N_{Pol+}^{P2}}{F_{Pol+}^{P2}}}{\sum N_{UNP,norm}^{Pi}} = \frac{1}{2} \left(1 + P_{Pol+}^W \Sigma \cos(2\varphi) \right) \quad (10)$$

$$P_{Pol+}^W = \frac{\varepsilon^{P1} P_{Pol+}^{P1} + \varepsilon^{P2} P_{Pol+}^{P2}}{\varepsilon^{P1} + \varepsilon^{P2}} \quad (11)$$

Problem:

P_{Pol+}^W depends on the azimuthal angle φ through the efficiencies

Solution:

We can show that this dependency is quite smooth

→ We can average P_{Pol+}^W on all angles φ

The hypothesis is quite reasonable because:

- 1) we are simultaneously analysing several η decay channels which have different efficiencies (see after)
- 2) In the extraction of the weighted polarization only relative efficiencies are relevant

Followed procedure:

$$N_{UNP,norm}^{Pi}(E_\gamma, \theta_\eta^{C.M.}) = \frac{2}{P_{Pol+}^{Pi} + P_{Pol-}^{Pi}} \left(P_{Pol-}^{Pi} \frac{N_{Pol+}^{Pi}}{F_{Pol+}^{Pi}} + P_{Pol+}^{Pi} \frac{N_{Pol-}^{Pi}}{F_{Pol-}^{Pi}} \right)$$

$$N_{UNP,norm}^{Pi}(E_\gamma, \vartheta_\eta^{C.M.}) \propto \varepsilon^{Pi}(\varphi)$$

- in the formula

$$P_{Pol+}^W = \frac{\varepsilon^{P1} P_{Pol+}^{P1} + \varepsilon^{P2} P_{Pol+}^{P2}}{\varepsilon^{P1} + \varepsilon^{P2}} \quad (11)$$

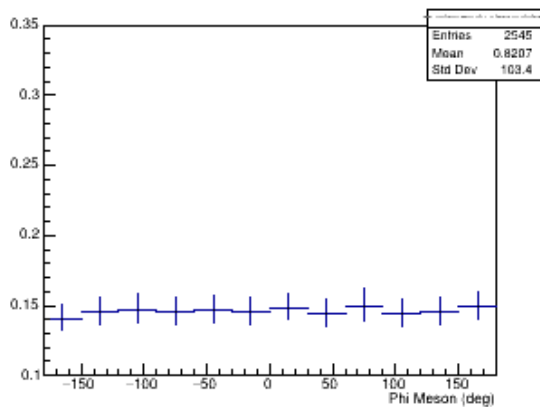
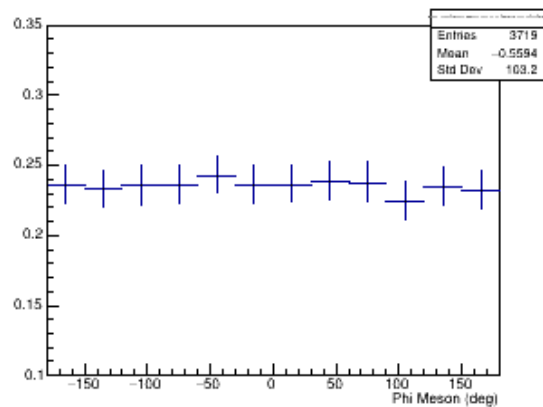
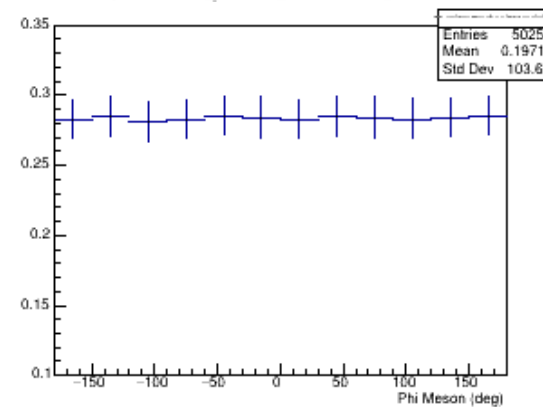
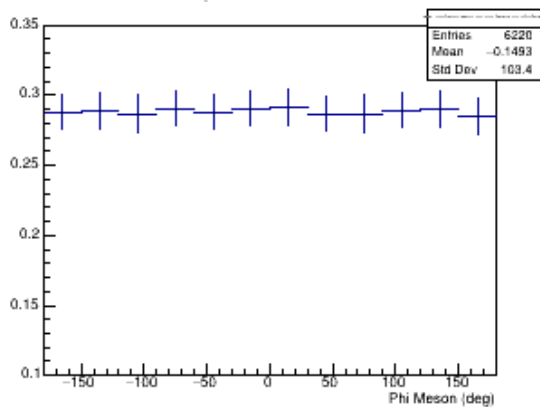
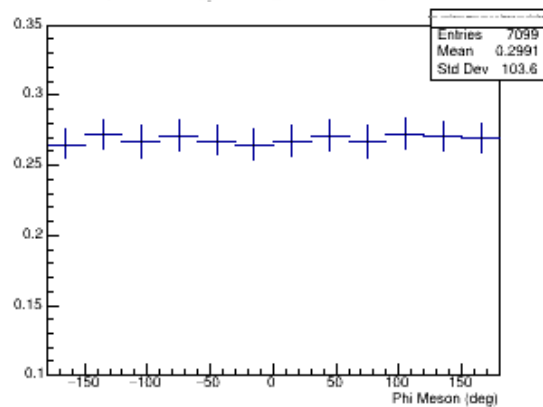
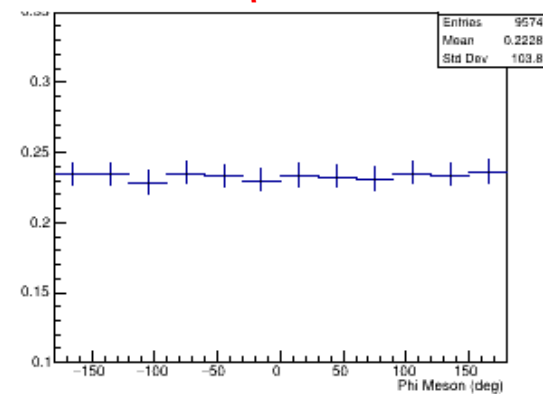
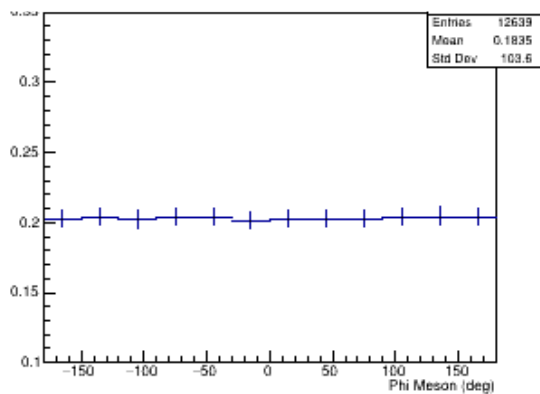
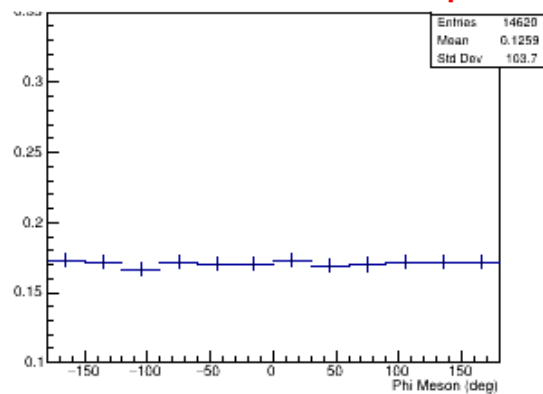
we replace:

$$\varepsilon^{Pi}(\varphi) \rightarrow N_{UNP,norm}^{Pi}(E_\gamma, \theta_\eta^{C.M.})$$

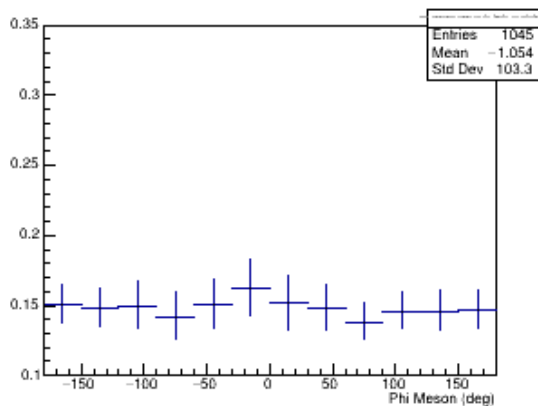
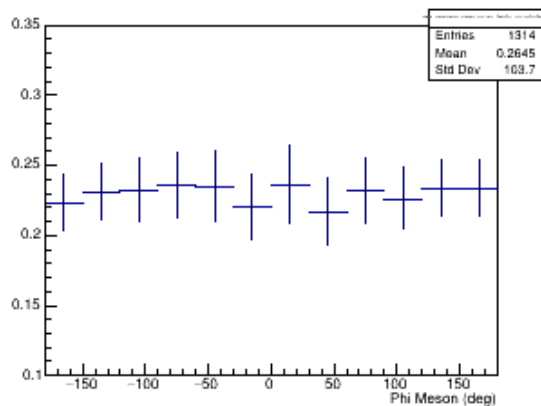
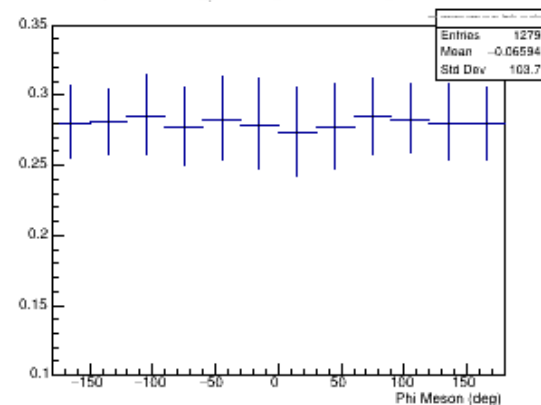
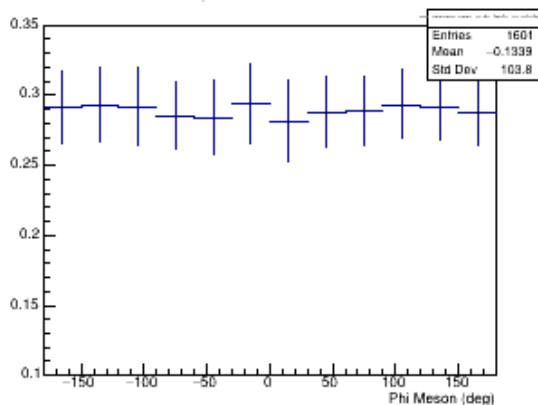
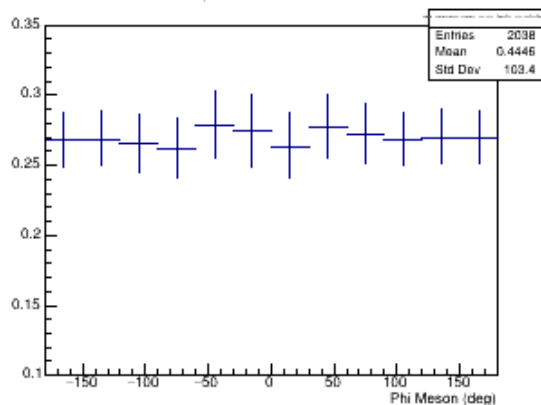
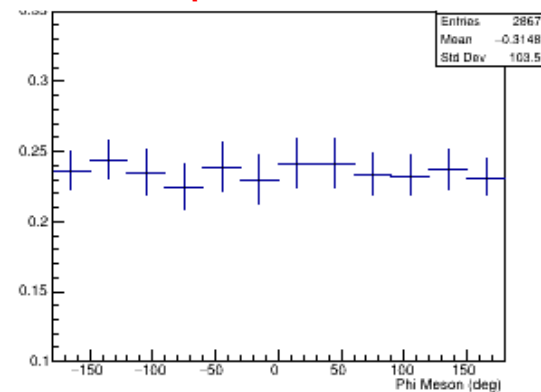
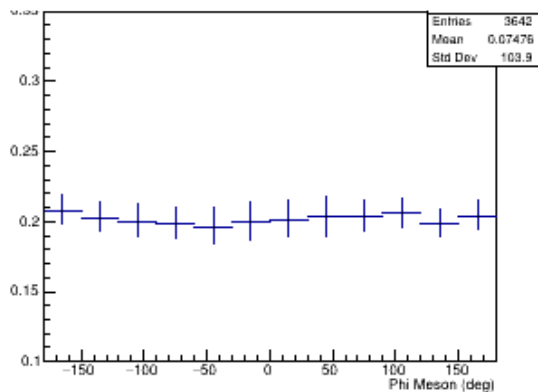
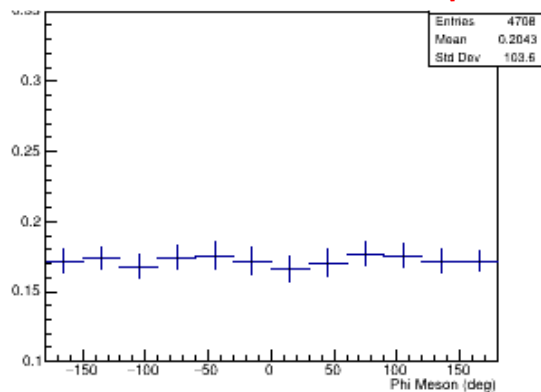
- We extract P_{Pol+}^W for every E_γ and η polar angle $\theta_\eta^{C.M.}$ in the c.m.
- make the average on all φ bins N_φ :

$$\overline{P_{Pol+}^W(E_\gamma)} = \frac{1}{N_\varphi} \sum_\varphi P_{Pol+}^W(E_\gamma, \varphi) \quad (21)$$

$P^W_{\text{pol}}(E_\gamma, \varphi) - \text{POLMINUS} - \text{CENTRAL PROTON} - \theta_\eta^{\text{CM}} = 68^\circ - 83^\circ$



$P_{\text{Pol}}^W(E_\gamma, \phi) - \text{POLMINUS} - \text{PROTON IN SciRi} - \theta_\eta^{\text{CM}} = 130^\circ - 140^\circ$



**ONE PERIOD P AND SEVERAL η DECAY CHANNELS $C1, C2, \dots$
WITH DIFFERENT EFFICIENCIES:**

$$\varepsilon^{P,C1} \neq \varepsilon^{P,C2}$$

POL. DEGREES AND FLUXES ARE THE SAME FOR ALL CHANNELS:

$$\begin{aligned} P_{Pol\pm}^{P,C1} &= P_{Pol\pm}^{P,C2} = P_{Pol\pm}^P & \text{BUT} & & P_{Pol+}^P &\neq P_{Pol-}^P \\ F_{Pol\pm}^{P,C1} &= F_{Pol\pm}^{P,C2} = F_{Pol\pm}^P & \backslash & & & \end{aligned}$$

Decay channels of one period can be summed up as follows:

$$\frac{\frac{N_{Pol+}^{P,C1}}{F_{Pol+}^P} + \frac{N_{Pol+}^{P,C2}}{F_{Pol+}^P}}{N_{UNP,norm}^{P,C1} + N_{UNP,norm}^{P,C2}} = \frac{(\varepsilon^{P,C1}(\varphi) + \varepsilon^{P,C2}(\varphi))(1 + P_{Pol+}^P \Sigma \cos(2\varphi))}{2(\varepsilon^{P,C1}(\varphi) + \varepsilon^{P,C2}(\varphi))} = \frac{1}{2}(1 + P_{Pol+}^P \Sigma \cos(2\varphi))$$

In general for more than two channels:

$$\frac{\sum_{Cj} \frac{N_{Pol+}^{P,Cj}}{F_{Pol+}^P}}{\sum_{Cj} N_{UNP,norm}^{P,Cj}} = \frac{1}{2}(1 + P_{Pol+}^P \Sigma \cos(2\varphi)) \quad (22)$$

**TWO OR MORE PERIODS P_i AND SEVERAL η DECAY CHANNELS $C_1, C_2, \dots$
WITH DIFFERENT EFFICIENCIES:**

$$\varepsilon^{P_i, C_1} \neq \varepsilon^{P_i, C_2}$$

$$\frac{\sum_{P_i} \sum_{C_j} \frac{N_{Pol+}^{P_i, C_j}}{F_{Pol+}^{P_i}}}{\sum_{P_i} \sum_{C_j} N_{UNP, norm}^{P_i, C_j}} = \frac{1}{2} \left(1 + P_{Pol+}^W \Sigma \cos(2\varphi) \right)$$

$$P_{Pol+}^W = \frac{\varepsilon^{P_1}(\varphi) P_{Pol+}^{P_1} + \varepsilon^{P_2}(\varphi) P_{Pol+}^{P_2}}{\varepsilon^{P_1}(\varphi) + \varepsilon^{P_2}(\varphi)}$$

$$\varepsilon^{P_1}(\varphi) = \varepsilon^{P_1, C_1}(\varphi) + \varepsilon^{P_1, C_2}(\varphi)$$

$$\varepsilon^{P_2}(\varphi) = \varepsilon^{P_2, C_1}(\varphi) + \varepsilon^{P_2, C_2}(\varphi)$$

Procedure for the extraction of the asymmetry: Summary

- for each period separately: **sum up** the contribution of all decay channels (**number of polarized events normalized to flux**)

$$\sum_{Cj} \frac{N_{Pol+}^{Pi, Cj}}{F_{Pol+}^P}$$

- for each period separately: extract the **average efficiency** of the period ε^{Pi} (proportional to the number of unpolarized and normalized events)

$$N_{UNP, norm}^{Pi}(E_\gamma, \vartheta_\eta^{C.M.}) \propto \varepsilon^{Pi}(\varphi)$$

- calculate the **weighted Polarization**

$$P_{Pol+}^W = \frac{\varepsilon^{P1}(\varphi) P_{Pol+}^{P1} + \varepsilon^{P2}(\varphi) P_{Pol+}^{P2}}{\varepsilon^{P1}(\varphi) + \varepsilon^{P2}(\varphi)}$$

- calculate the **average on the azimuthal angle φ of Weighted Polarization**

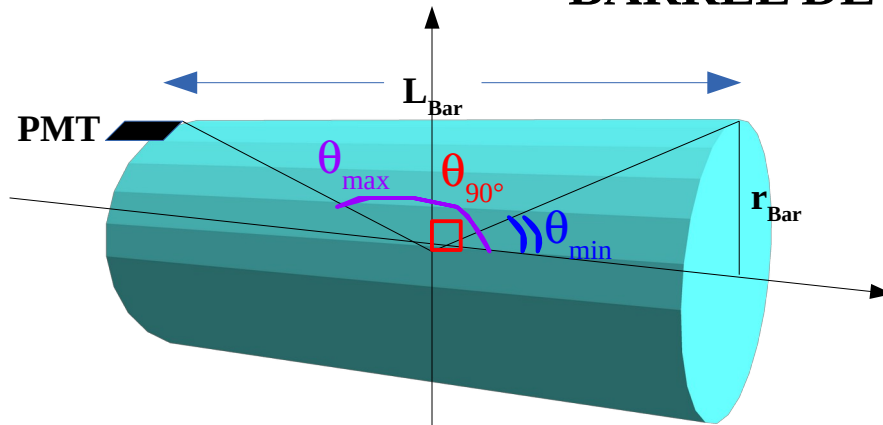
$$\overline{P_{Pol+}^W(E_\gamma)} = \frac{1}{N_\varphi} \sum_\varphi P_{Pol+}^W(E_\gamma, \varphi)$$

- extract the beam asymmetry from the fit:

$$\frac{\sum_{Pi} \frac{N_{Pol+}^{Pi}}{F_{Pol+}^{Pi}}}{\sum_{Pi} N_{UNP, norm}^{Pi}} = \frac{1}{2} \left(1 + \overline{P_{Pol+}^W} \Sigma \cos(2\varphi) \right) \quad (23)$$

BARREL CALIBRATION

BARREL DETAILS

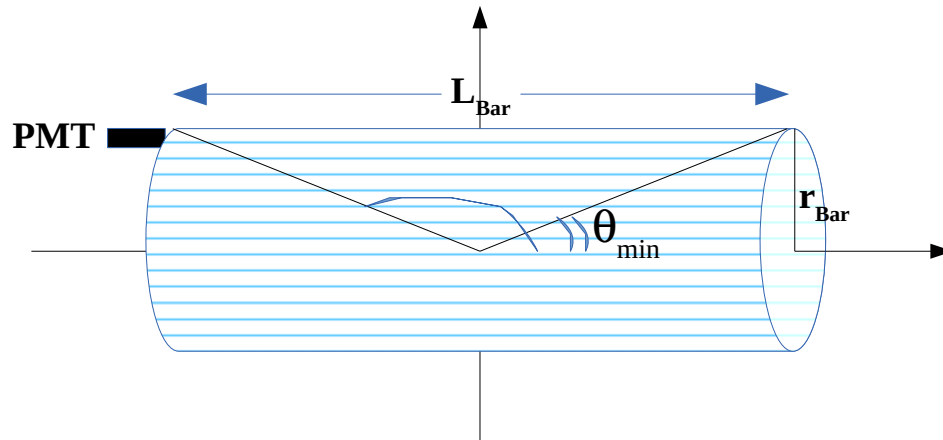


Polar Angle in LAB **Travelled distance To the PM**

$$\theta_{\min} = 25^\circ \Rightarrow x(\theta_{\min}) = L_{\text{Bar}}$$

$$\theta_{\max} = 155^\circ \Rightarrow x(\theta_{\max}) = 0$$

$$\theta_{90^\circ} = 90^\circ \Rightarrow x(\theta_{90^\circ}) = L_{\text{Bar}}/2$$

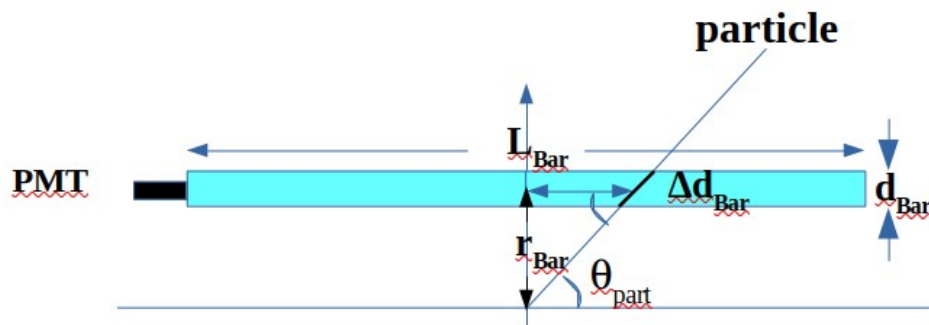


Barrel Length

$$L_{\text{bar}} = 43.4 \text{ cm}$$

Barrel Radius

$$r_{\text{Bar}} = L_{\text{Bar}}/2 * \text{tg}(\theta_{\min})$$



$$\Delta d_{\text{Bar}} = \frac{d_{\text{Bar}}}{\sin \theta_{\text{part}}}$$

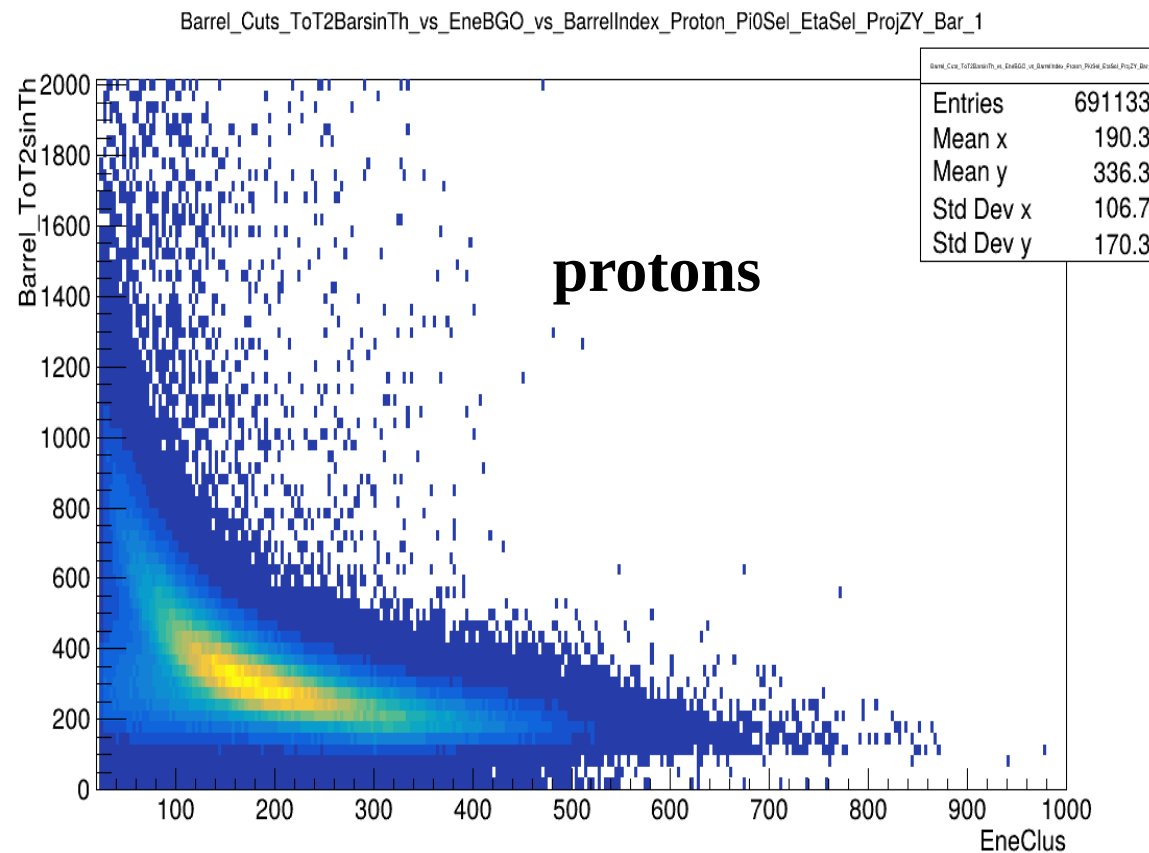
$$\text{ToT2}_{\text{Bar}} \propto \frac{1}{\sin \theta_{\text{part}}}$$

➡ If we multiply the ToT2_{Bar} by $\sin(\theta)$ we get free of this dependence

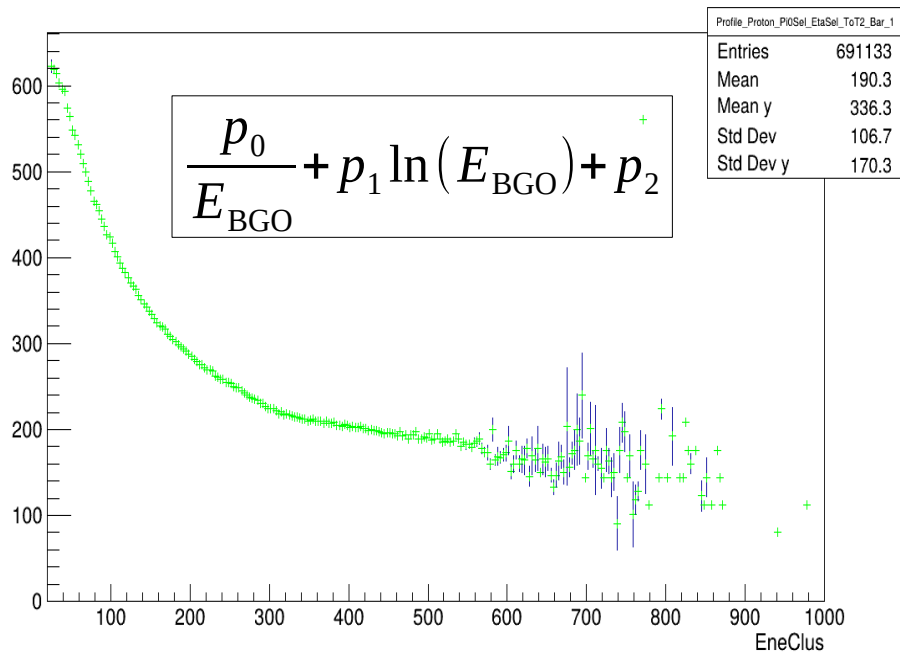
CUTS ON $ToT2_{Bar} * \sin(\theta)$ vs. E_{BGO}

FOR:

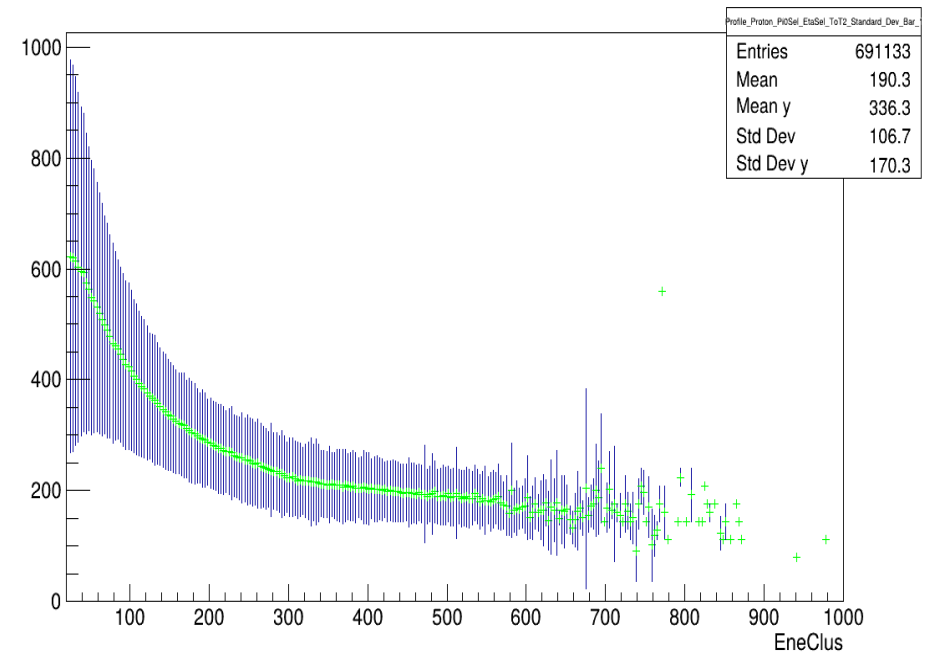
- **protons** from: $\gamma p \rightarrow \pi^0 p \rightarrow 2\gamma p$
- **pions** from: $\gamma p \rightarrow \eta p \rightarrow \pi^+ \pi^- \pi^0 p$



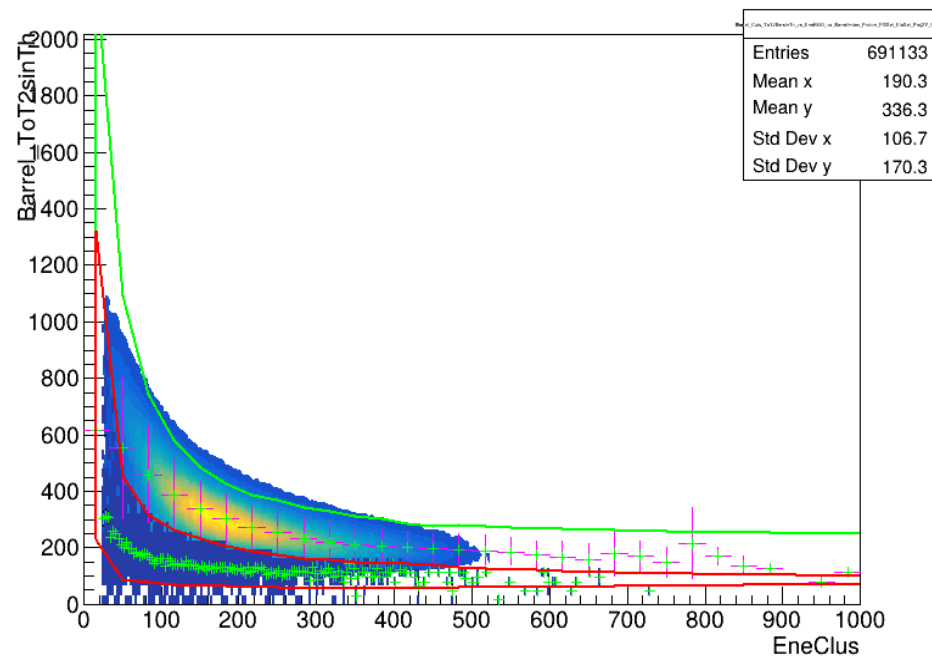
Barrel_Cuts_ToT2BarsinTh_vs_EneBGO_vs_BarrelIndex_Proton_Pi0Sel_EtaSel_ProjZY_Bar_1



Barrel_Cuts_ToT2BarsinTh_vs_EneBGO_vs_BarrelIndex_Proton_Pi0Sel_EtaSel_ProjZY_Bar_1



Barrel_Cuts_ToT2BarsinTh_vs_EneBGO_vs_BarrelIndex_Proton_Pi0Sel_EtaSel_ProjZY_Bar_1_1



CALIBRATION OF BARREL ADC

WITH:

- **pions@MIP** from: $\gamma p \rightarrow \pi^+ n$

BARREL CALIBRATION

$$\Delta E_{\text{Bar}}(i_{\text{Bar}}) = c.c.(i_{\text{Bar}}) * \text{ToT2}_{\text{Bar}}(i_{\text{Bar}})$$

Technique

Choose a reaction for which energy deposited by a charged particle in the BGO and in the barrel (plastic EJ248) is fixed

$$\gamma p \rightarrow \pi^+ n$$

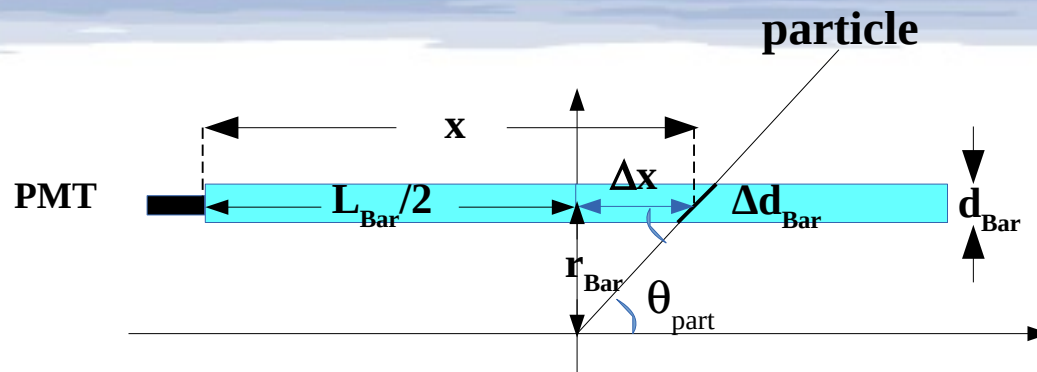
in energy range where the pion is a MIP particle both in BGO and in barrel.

BGO

$$E_{\text{BGO}}^{\text{MIP}} = \left(\frac{dE}{dx} \right)_{\text{BGO}}^{\text{MIP}} L_{\text{BGO}} = 213.2 \text{ MeV}$$

BARREL

$$\left(\frac{dE}{dx} \right)_{\text{Bar}}^{\text{MIP}} = 2.0 \frac{\text{MeV}}{\text{g}} \text{ cm}^2 * \rho_{\text{Bar}} = 2.0 \frac{\text{MeV}}{\text{g}} \text{ cm}^2 * 1.023 \frac{\text{g}}{\text{cm}^3} = 2.046 \frac{\text{MeV}}{\text{cm}}$$



- Energy loss of a particle emitted at a polar angle θ_{part} depends on crossed material thickness Δd_{Bar} , i.e. on θ_{part} :

$$(\Delta E)_{\text{Bar}}^{\text{MIP}} = 2.046 \frac{\text{MeV}}{\text{cm}} * \Delta d_{\text{Bar}}(\theta_{\text{part}})$$

$$\Delta d_{\text{Bar}} = \frac{d_{\text{Bar}}}{\sin \theta_{\text{part}}}$$

- Distance x travelled by light before collection in PMT causes light exponential attenuation and can be expressed in terms of the polar angle θ_{part} (see picture):

$$x = \frac{L_{\text{Bar}}}{2} + \Delta x = \frac{L_{\text{Bar}}}{2} + r_{\text{Bar}} \frac{\cos \theta_{\text{part}}}{\sin \theta_{\text{part}}}$$

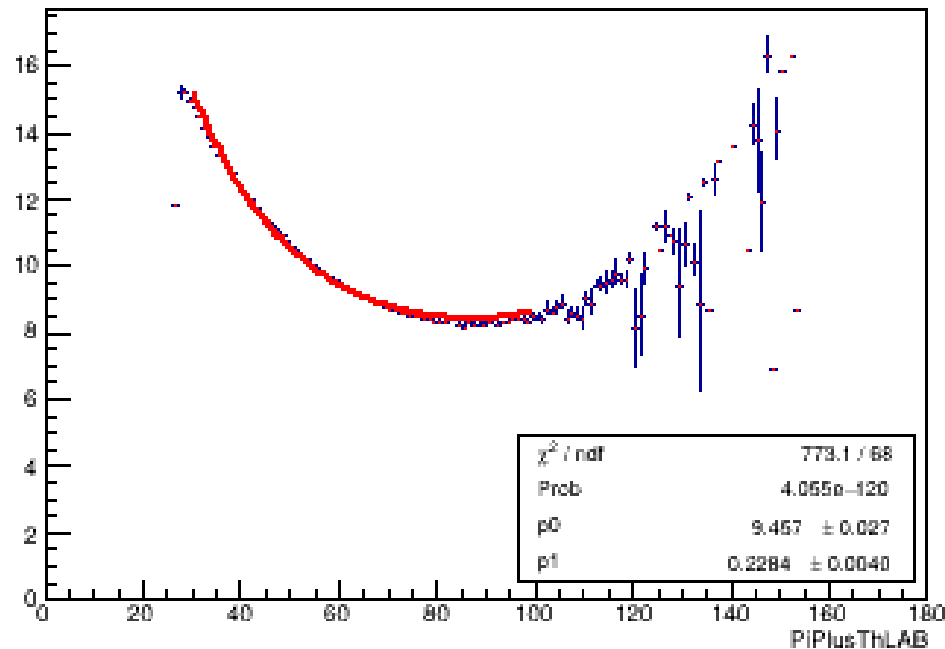
➡ Barrel ToT2 can be parametrized as follows:

$$ToT2_{\text{Bar},\pi} = \frac{A}{\sin \theta_{\pi}} \exp\left(-\frac{B}{L_{\text{Bar}}} x\right) \quad \rightarrow \quad ToT2_{\text{Bar},\pi} = \frac{A}{\sin \theta_{\pi}} \exp\left[-\frac{B}{L_{\text{Bar}}} \left(\frac{L_{\text{Bar}}}{2} + r_{\text{Bar}} \frac{\cos \theta_{\pi}}{\sin \theta_{\pi}}\right)\right]$$

Parameters **A** and **B** can be derived from data.

$$ToT2_{Bar,\pi} = \frac{A}{\sin \theta_{\pi}} \exp \left[- \frac{B}{L_{Bar}} \left(\frac{L_{Bar}}{2} + r_{Bar} \frac{\cos \theta_{part}}{\sin \theta_{part}} \right) \right]$$

ToT2

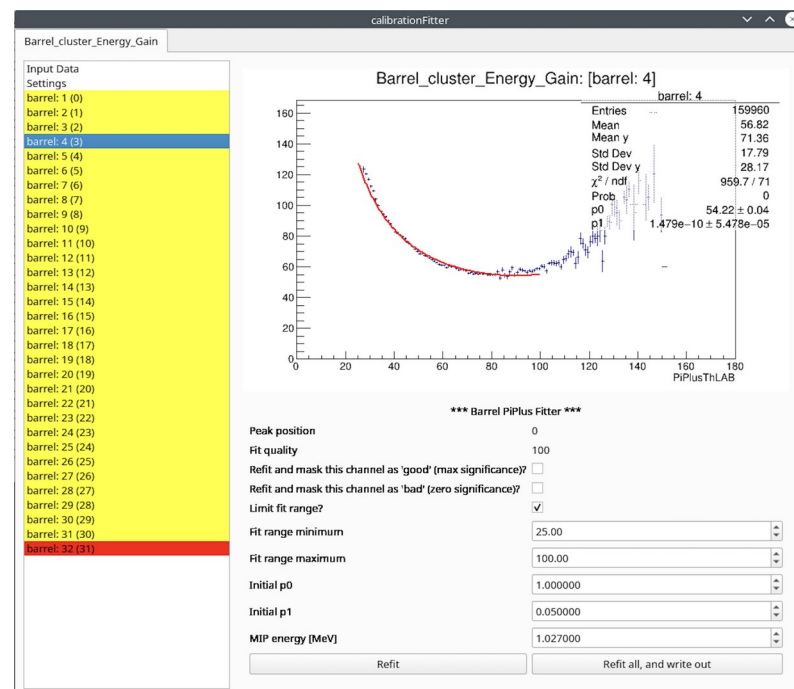
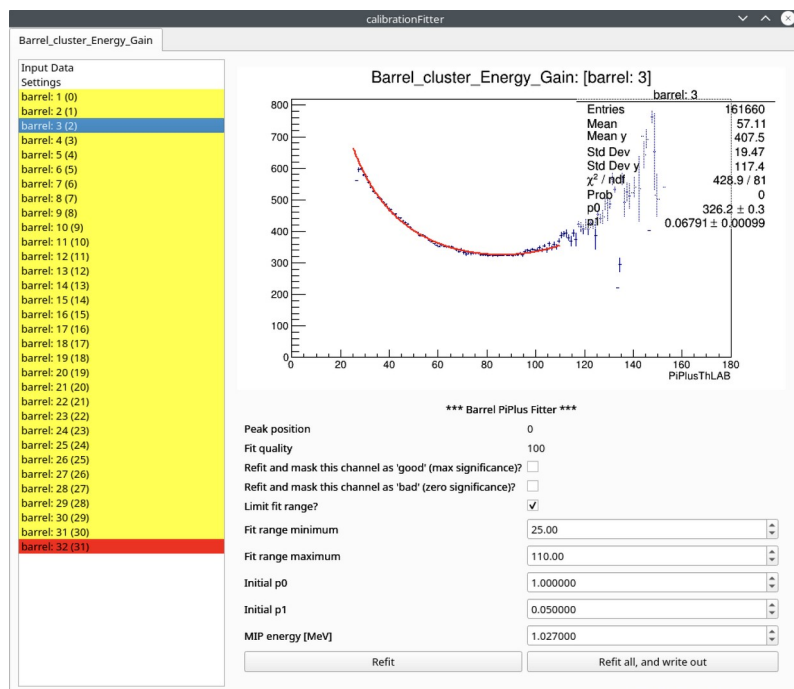
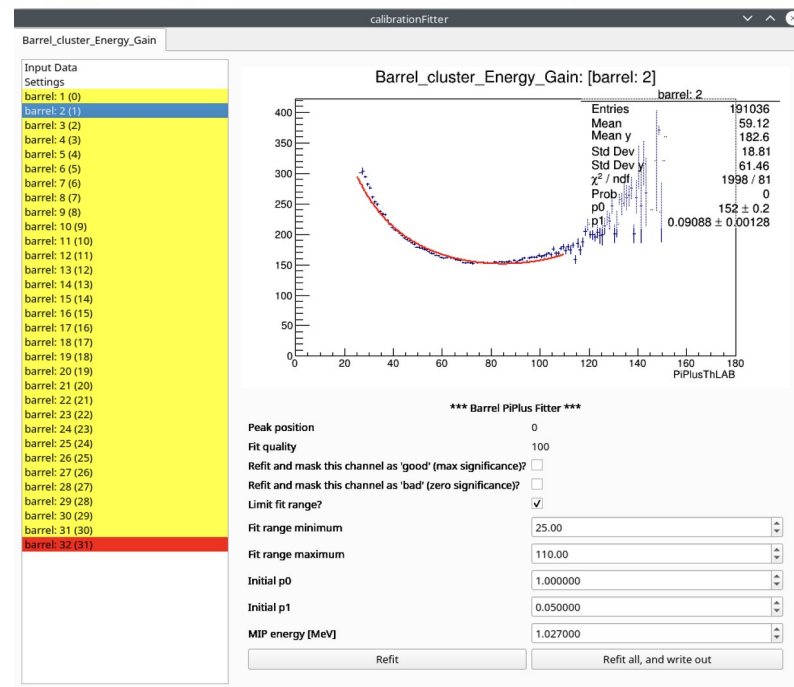
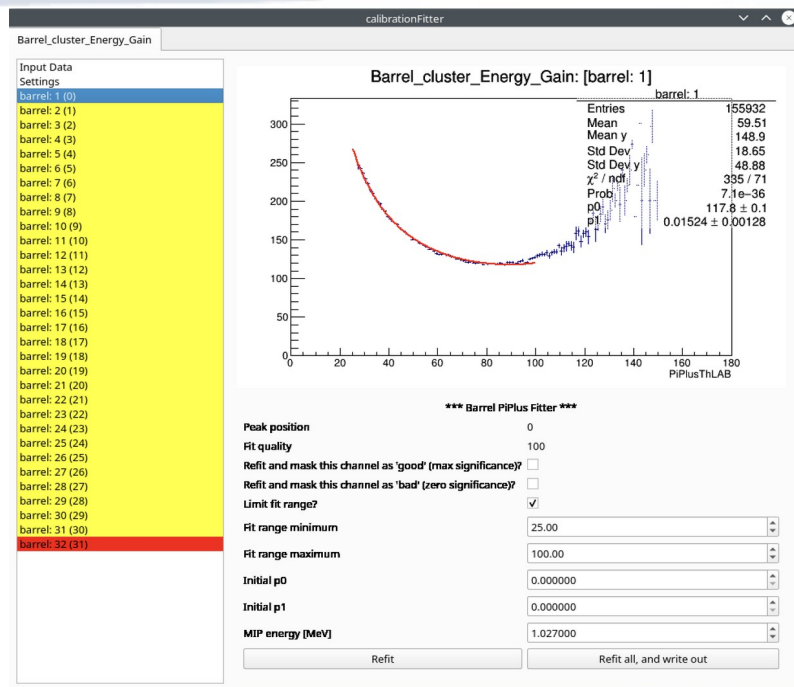


- Fit procedure has been implemented inside a **calibrationFitter program** (similar to the BGO calibration one) → production of the calibration file necessary for the database
- Analysis outputs are provided to this program in the form of 3D histos

$ToT2_{Bar}$ vs. $\theta_{\pi+,LAB}$ vs. Barrel Index

BUT

We still have to find how to read in explora barrel calibration constants from database



BEAM ASYMMETRY Σ FOR:

$$\gamma p \rightarrow \eta p$$

- Analysis of 3 data taking periods
(OctNov2015, AprMay2017, NovDec2018)

- Analysis of 3 decay channels

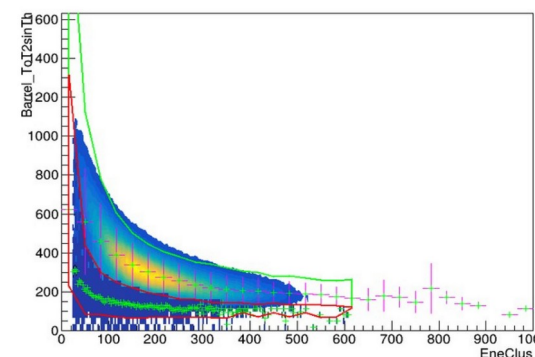
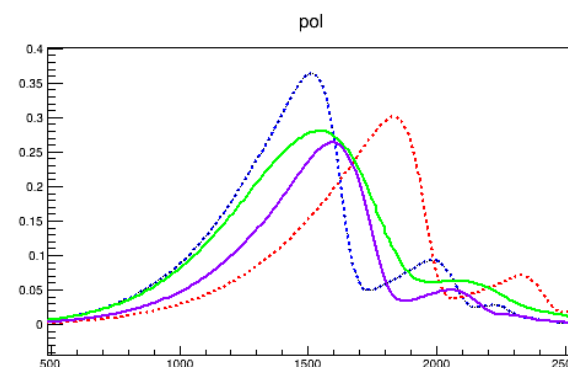
$$\eta \rightarrow 2\gamma \quad \eta \rightarrow 3\pi^0 \rightarrow 6\gamma \quad \eta \rightarrow \pi^+ \pi^- \pi^0$$

With proton in BGO/Bar and in SciRi

- Calibration constants not yet applied

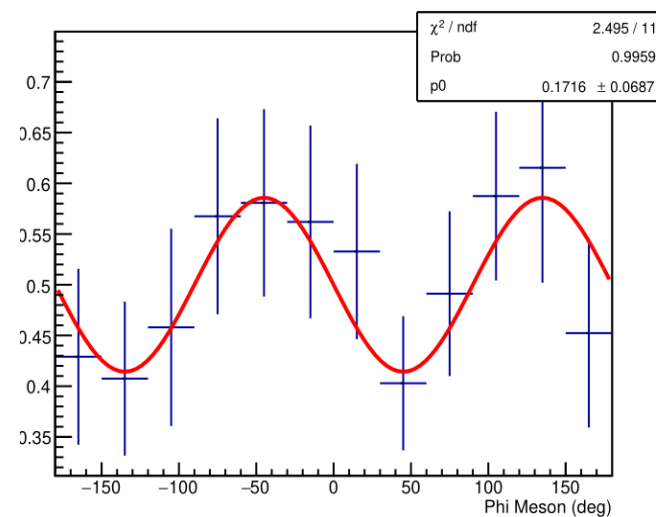
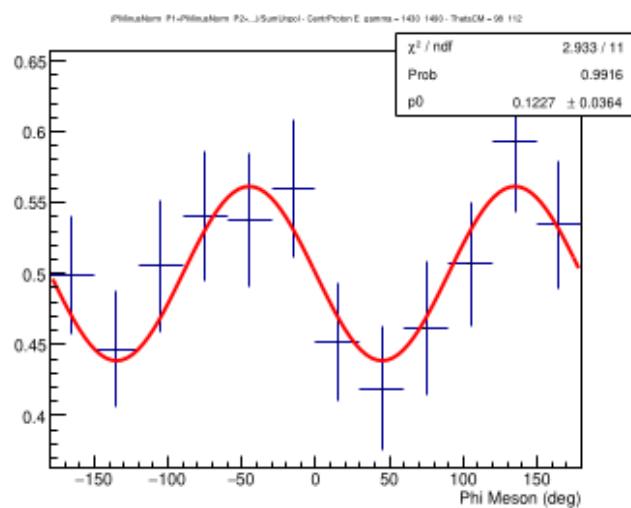
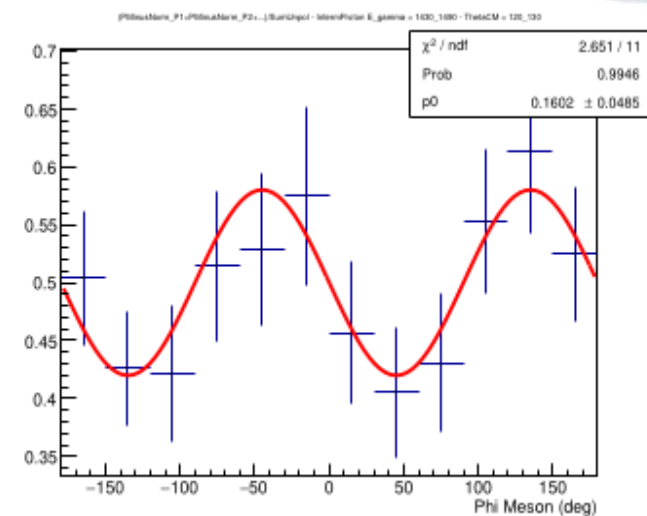
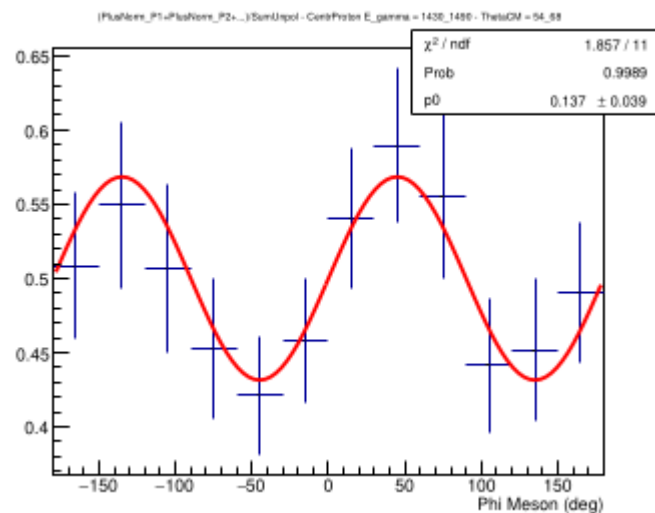
→ Only cuts $ToT2 \cdot \sin(\theta_{CM})$ vs. Ene_{BGO}

- Different efficiencies
- Different polarization degrees



COMPARISON EFFICIENCY DATA TAKING PERIODS - H2

	Jun15	Oct15	AprMay17	NovDec18
	CENTRAL	PROTON	AND $\eta \rightarrow 2\gamma$	
$E_\gamma > 0.5 \text{ GeV}$				
# η /Flux	$0.3 \cdot 10^{-8}$	$0.4 \cdot 10^{-8}$	$1.6 \cdot 10^{-8}$	$2.3 \cdot 10^{-8}$
$E_\gamma > 1.2 \text{ GeV}$				
# η /Flux	$0.6 \cdot 10^{-8}$	$0.6 \cdot 10^{-8}$	$2.2 \cdot 10^{-8}$	$3. \cdot 10^{-8}$
	INTERM.	PROTON	AND $\eta \rightarrow 2\gamma$	
$E_\gamma > 0.5 \text{ GeV}$				
# η /Flux	$0.5 \cdot 10^{-8}$	$0.6 \cdot 10^{-8}$	$2.3 \cdot 10^{-8}$	$2.7 \cdot 10^{-8}$
$E_\gamma > 1.2 \text{ GeV}$				
# η /Flux	$0.2 \cdot 10^{-8}$	$0.2 \cdot 10^{-8}$	$0.8 \cdot 10^{-8}$	$0.7 \cdot 10^{-8}$

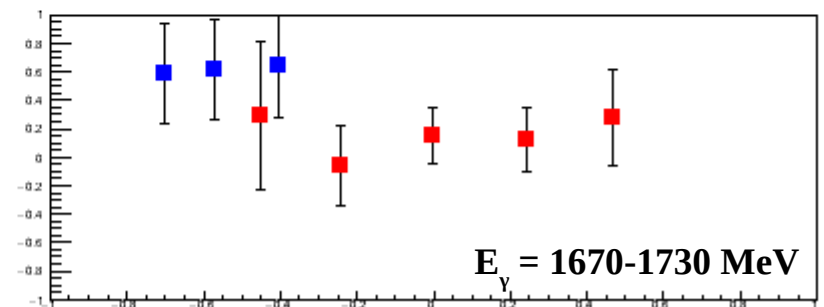
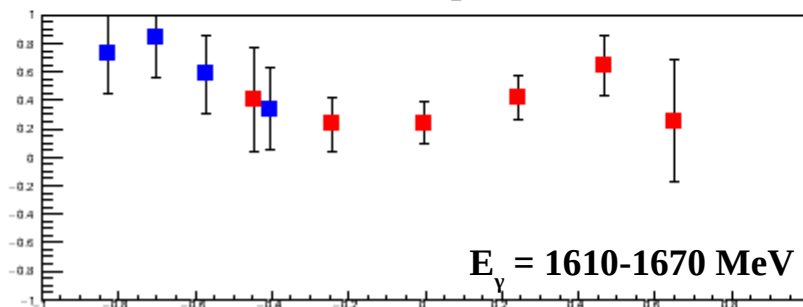
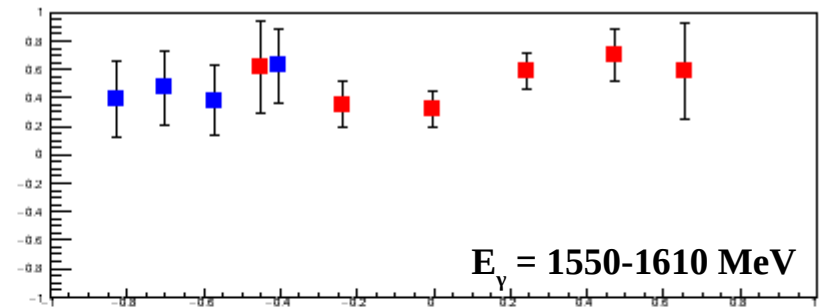
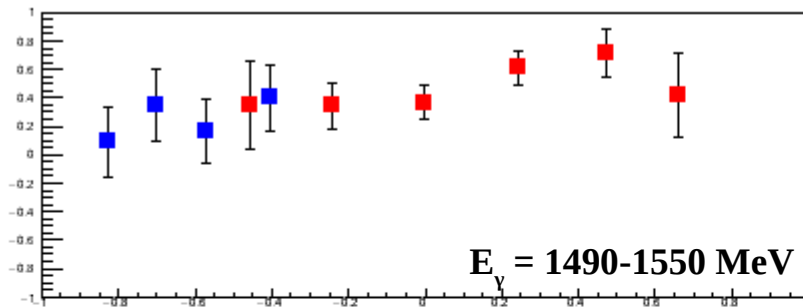
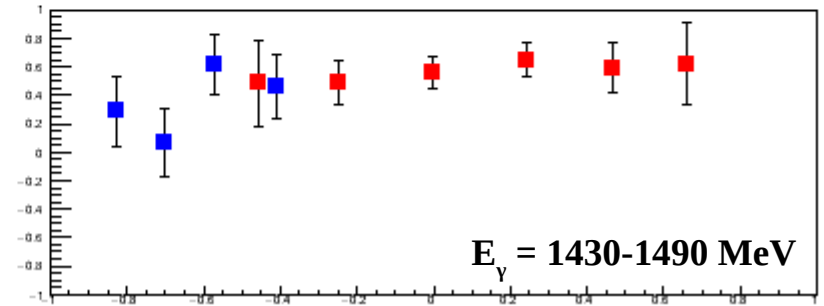
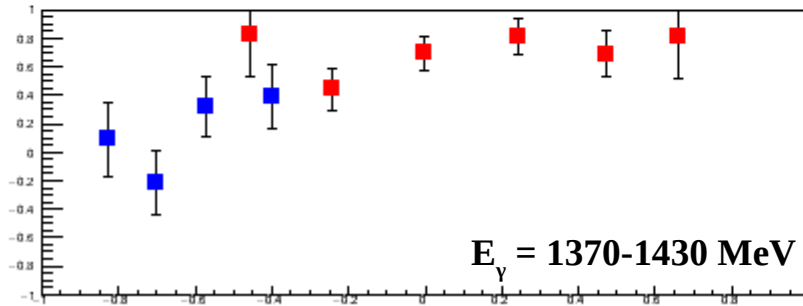
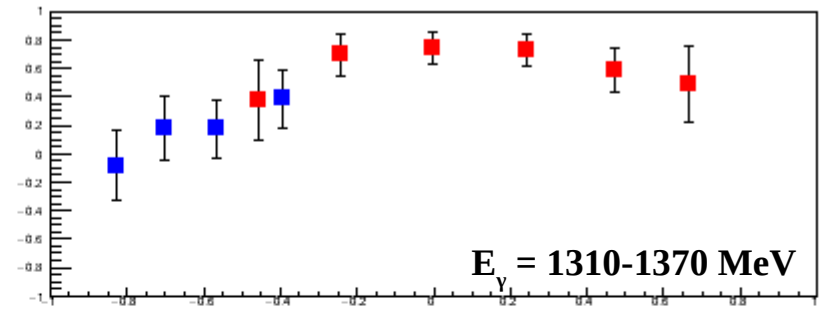
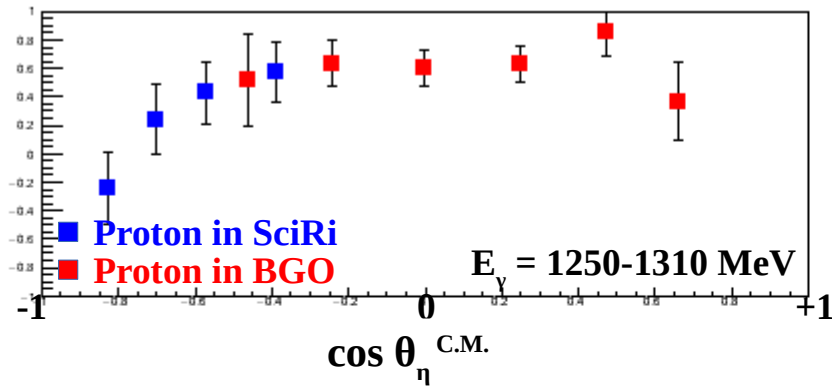


Proton in BGO/Bar

Proton in SciRi

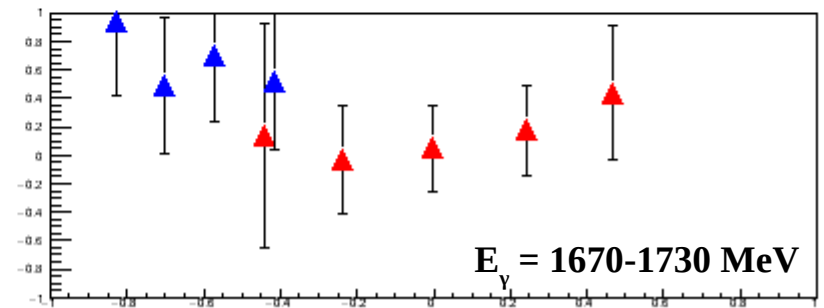
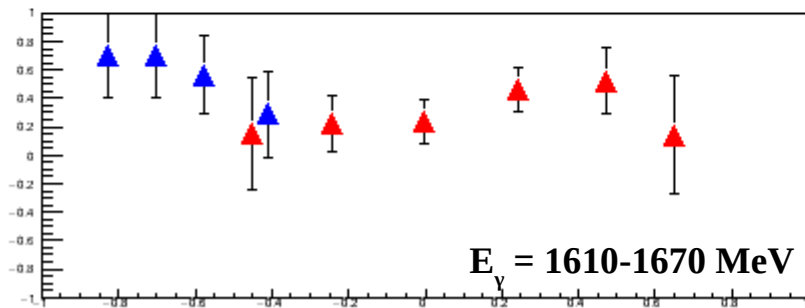
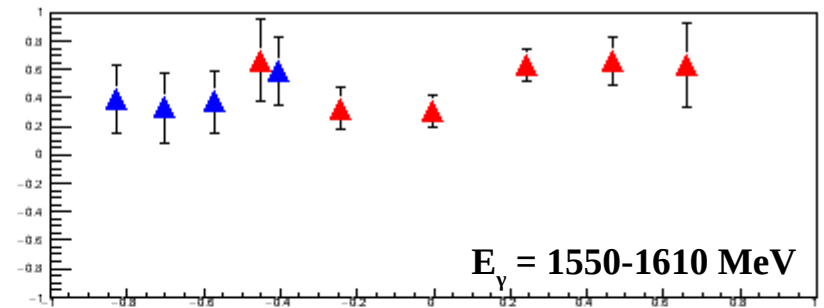
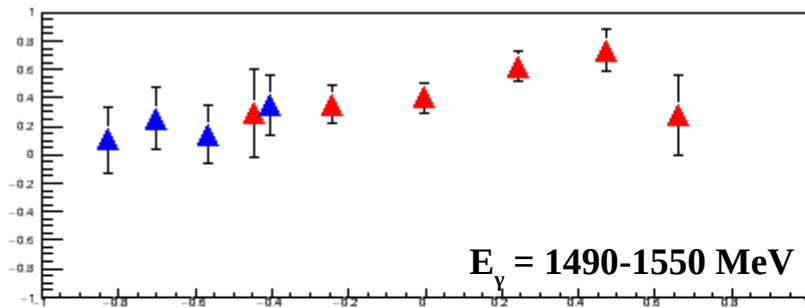
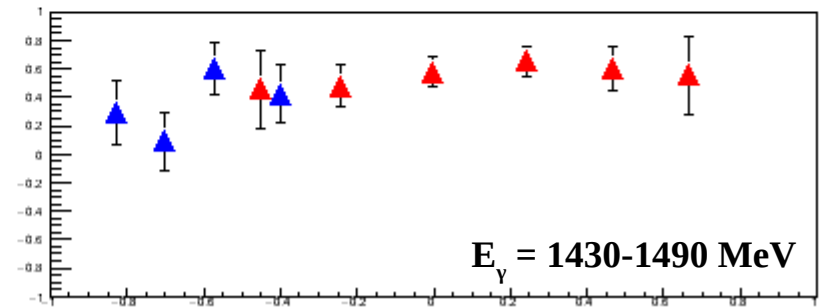
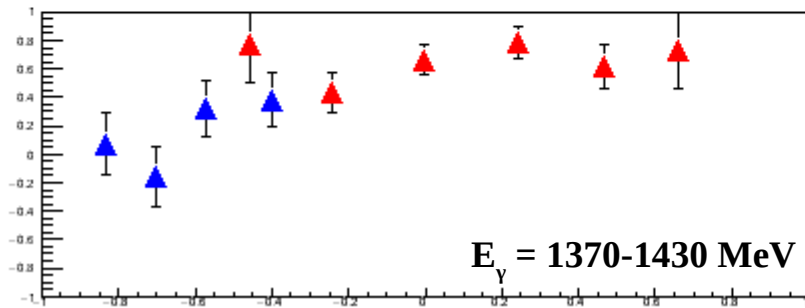
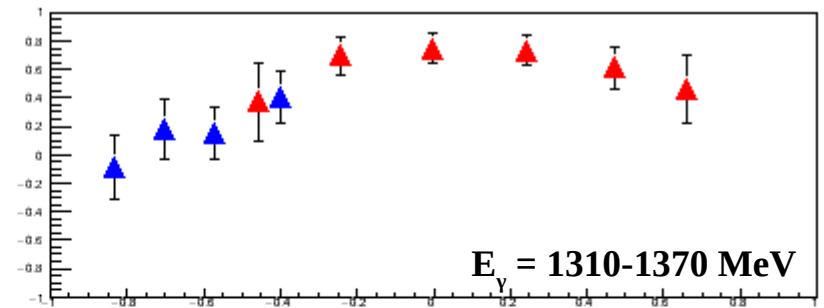
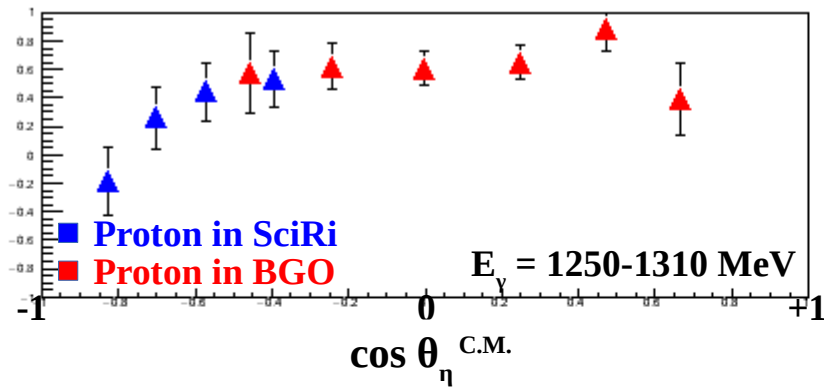
PolPlus

Σ



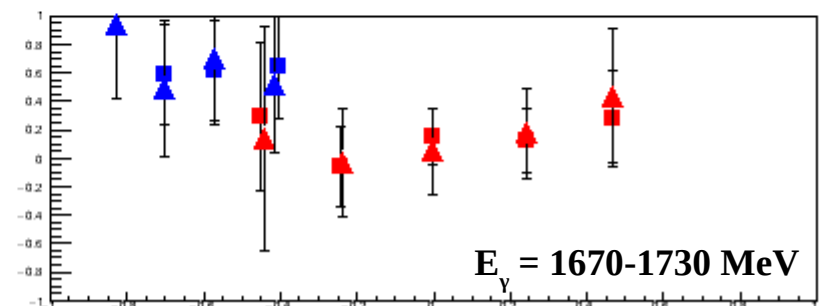
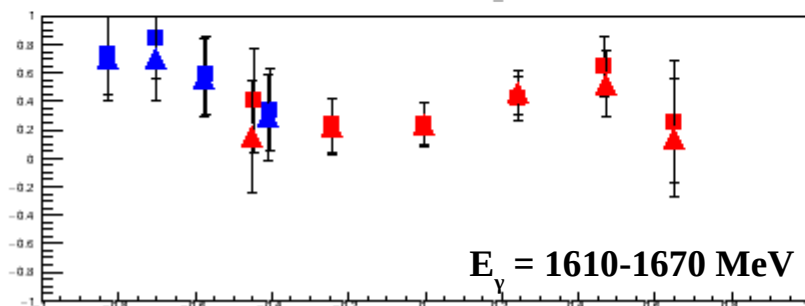
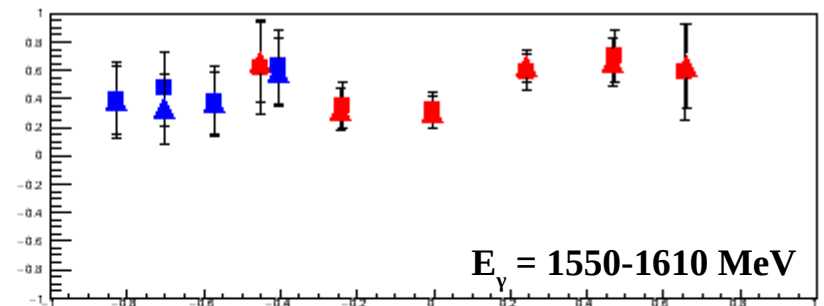
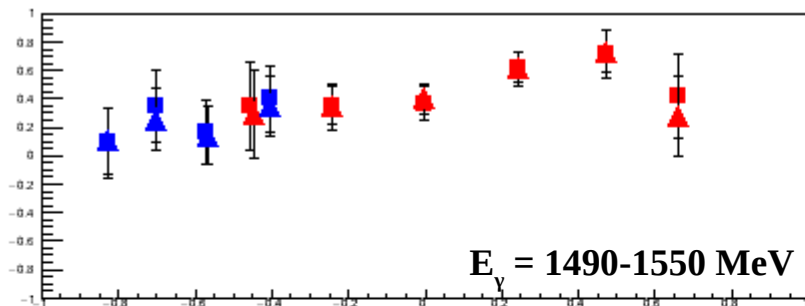
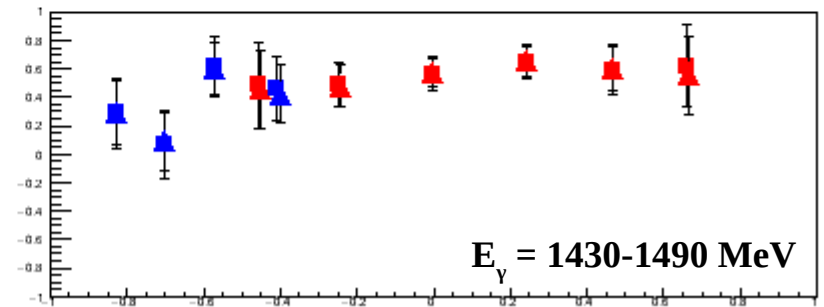
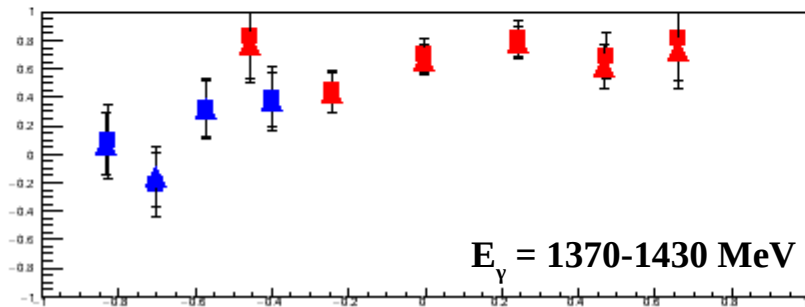
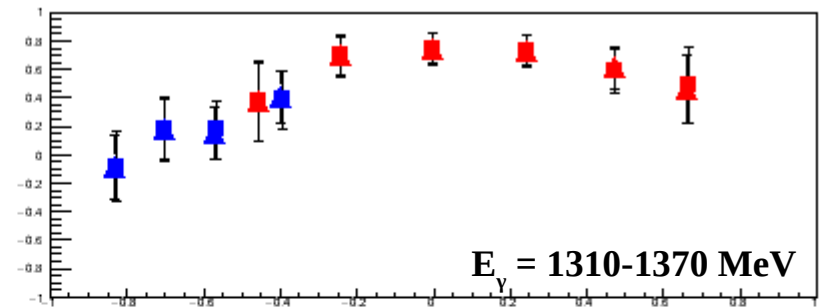
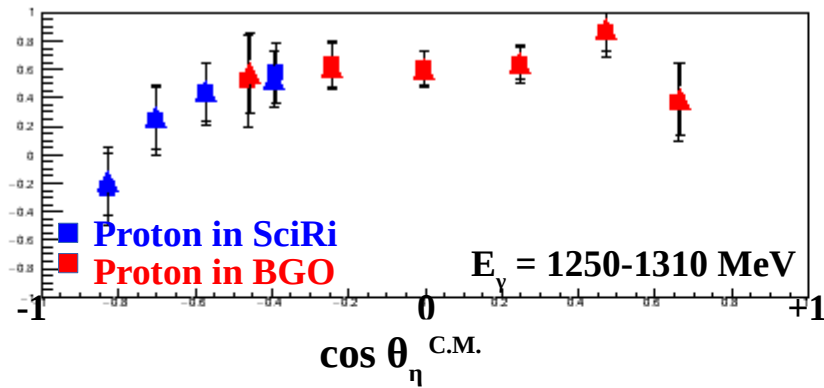
PolMinus

Σ



PolPlus/PolMinus

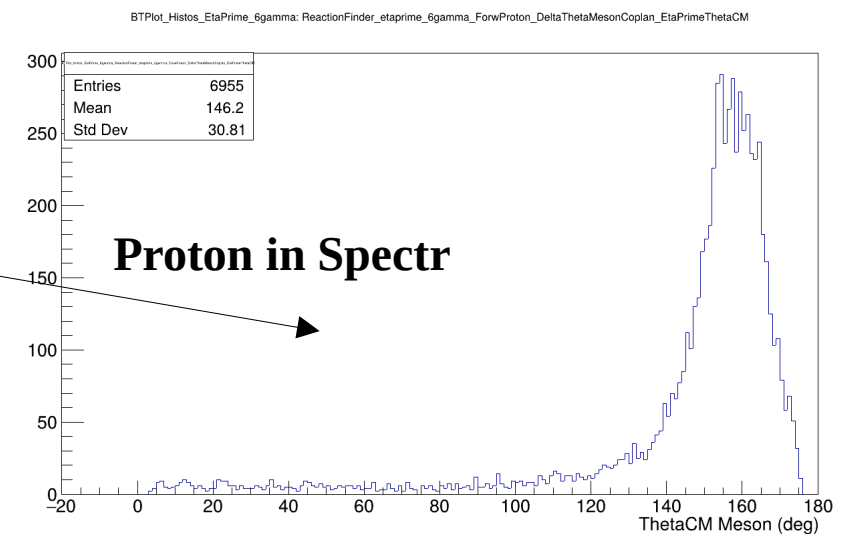
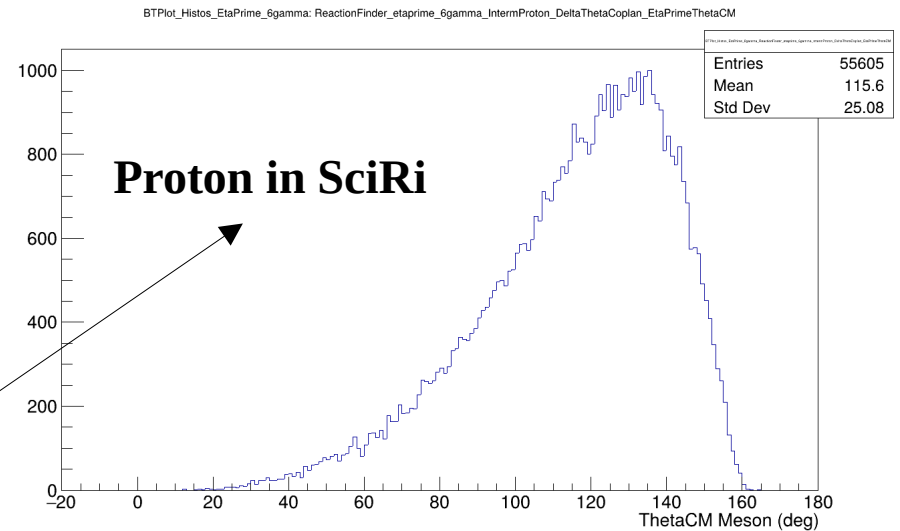
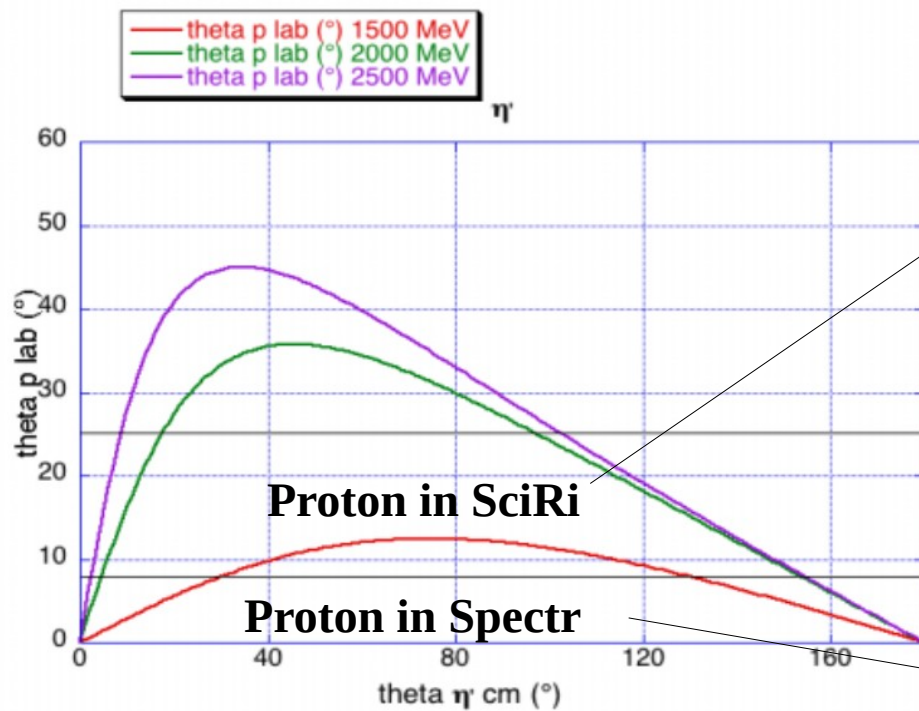
Σ



η' ANALYSIS

DATA

η' KINEMATICS



BEAM ASYMMETRY Σ FOR:

$$\gamma p \rightarrow \eta' p$$

- Analysis of 3 data taking periods
(OctNov2015, AprMay2017, NovDec2018)

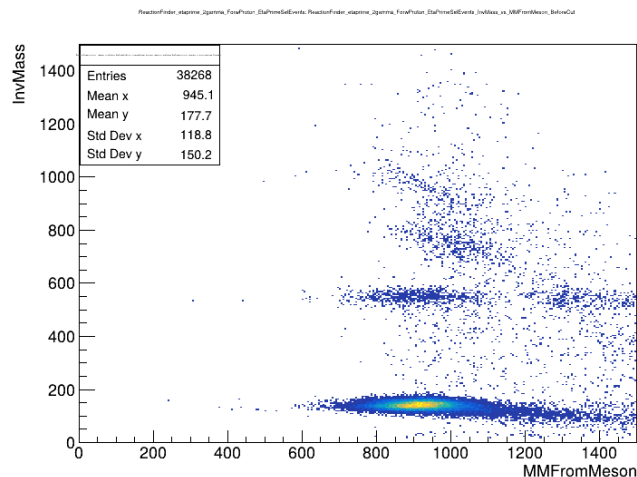
- Analysis of:
Proton in Spectr $\rightarrow$ 3 decay channels

$$\eta' \rightarrow 2\gamma \quad \eta' \rightarrow \pi^0 \pi^0 \eta \rightarrow 6\gamma \quad \eta' \rightarrow \pi^+ \pi^- \eta$$

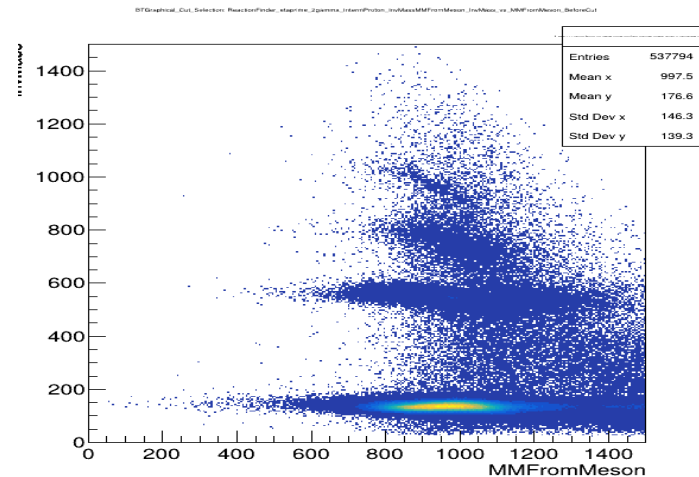
Proton in SciRi $\rightarrow$ 2 decay channels

$$\eta' \rightarrow 2\gamma \quad \eta' \rightarrow \pi^0 \pi^0 \eta \rightarrow 6\gamma$$

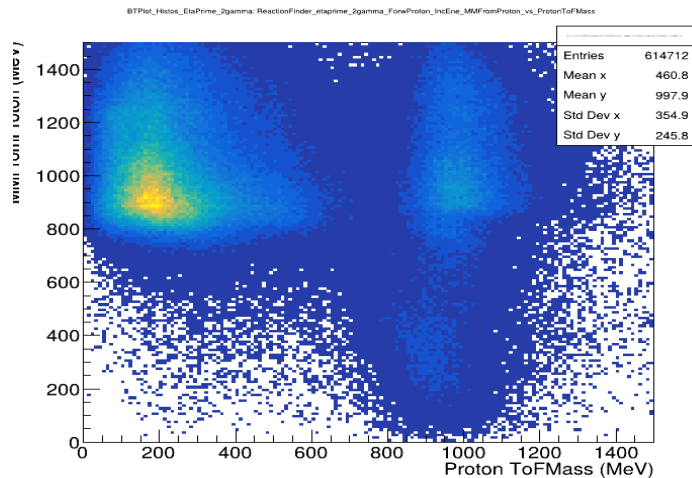
Results are promising BUT not yet final



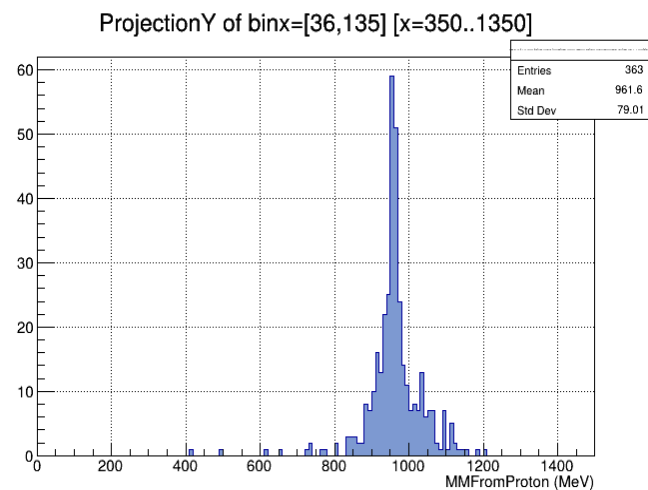
$\eta' \rightarrow 2\gamma$ – Proton in Spectr
 2γ InvMass vs MMFromMeson



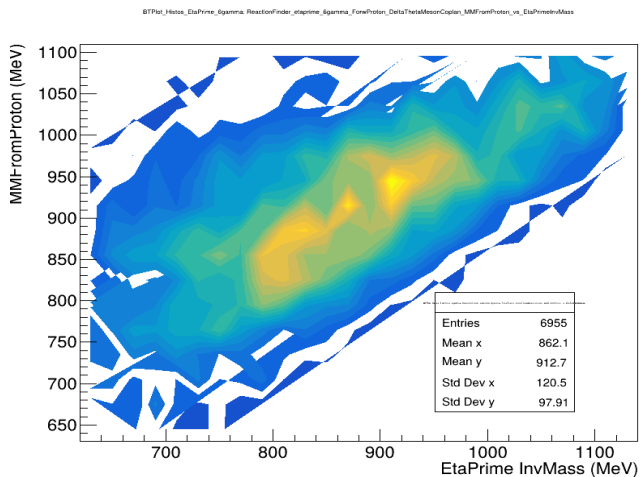
$\eta' \rightarrow 2\gamma$ – Proton in SciRi
 2γ InvMass vs MMFromMeson



$\eta' \rightarrow 2\gamma$ – Proton in Spectr
MMFromProton vs ToFMass

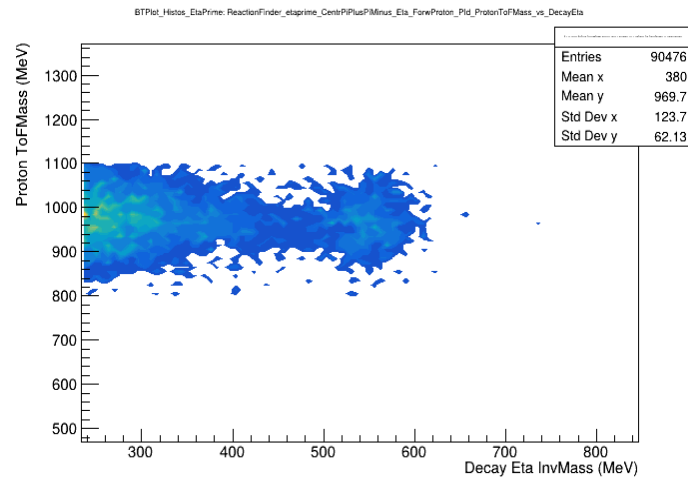


$\eta' \rightarrow 2\gamma$ – Proton in Spectr
MMFromProton After Cuts



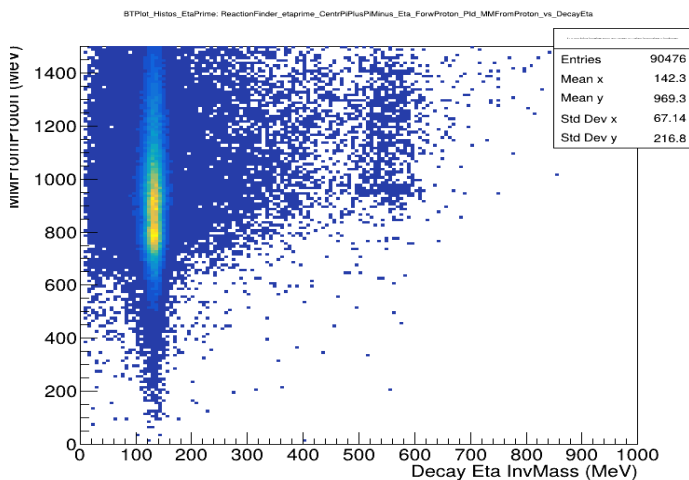
$\eta' \rightarrow 6\gamma$ – Proton in Spectr

MMFromProton vs 6γ InvMass



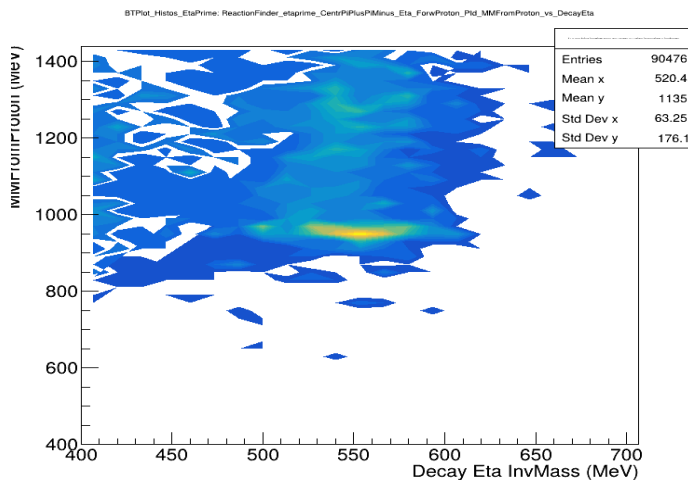
$\eta' \rightarrow \pi^+\pi^-\eta$ – Proton in Spectr

ToFMass vs η DecayInvMass



$\eta' \rightarrow \pi^+\pi^-\eta$ – Proton in Spectr

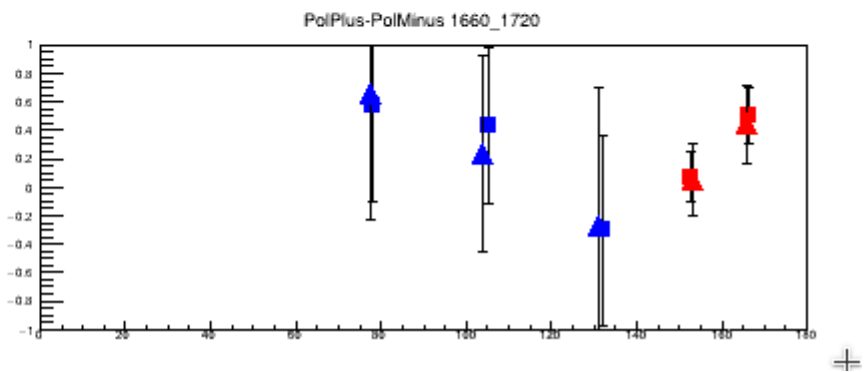
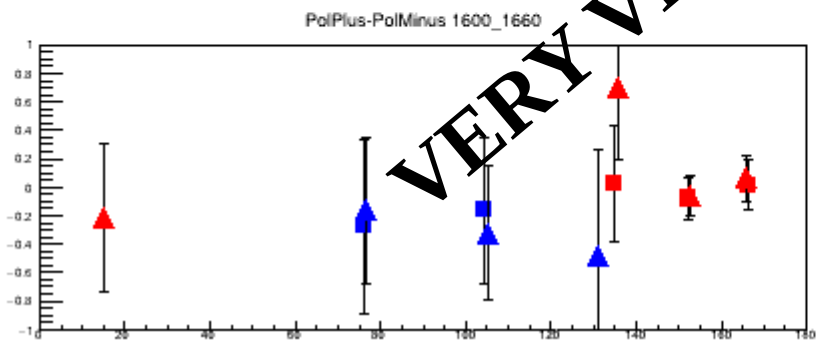
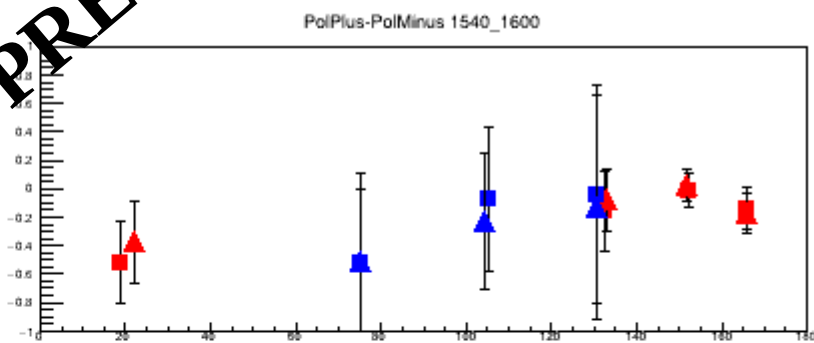
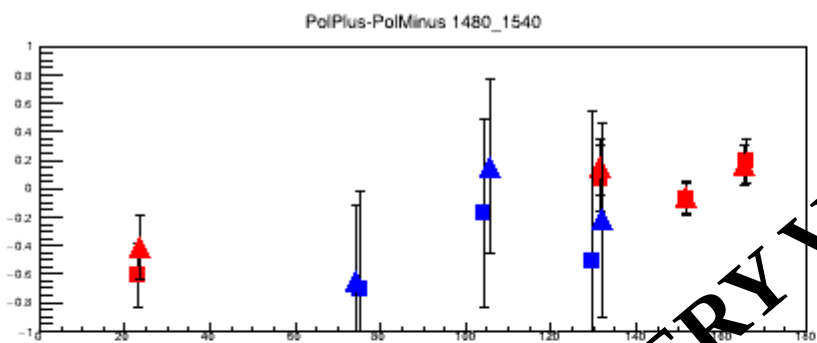
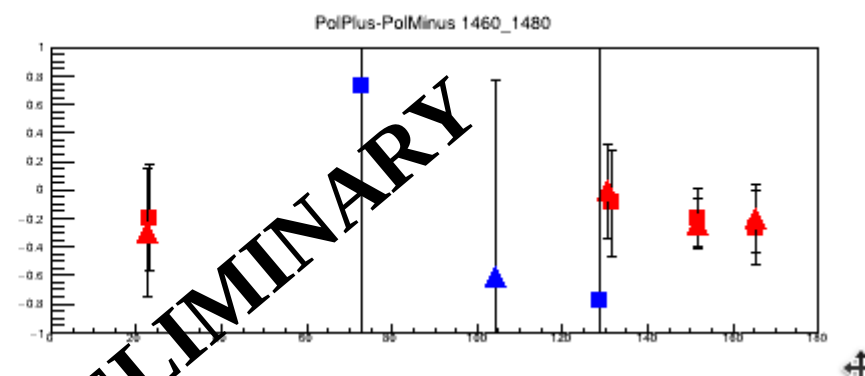
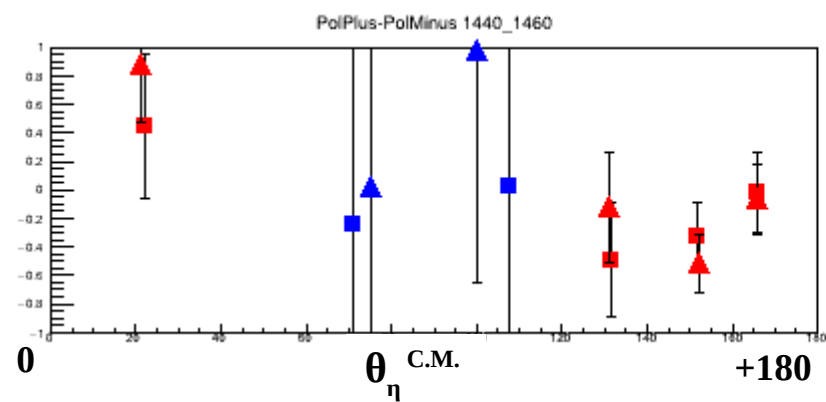
MMFromProton vs η DecayInvMass



$\eta' \rightarrow \pi^+\pi^-\eta$ – Proton in Spectr

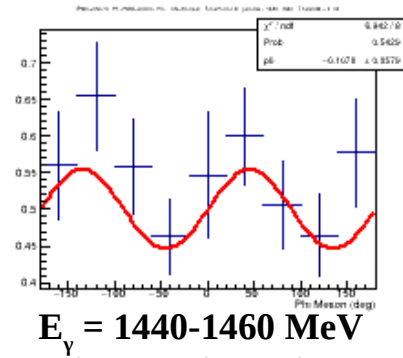
MMFromProton vs η DecayInvMass

PolPlus/PolMinus

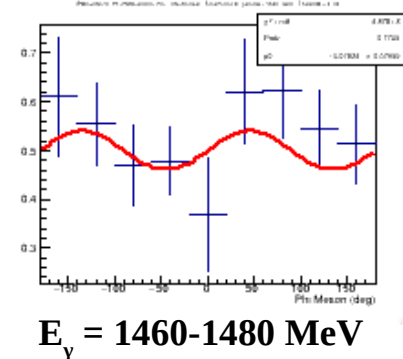


VERY VERY PRELIMINARY

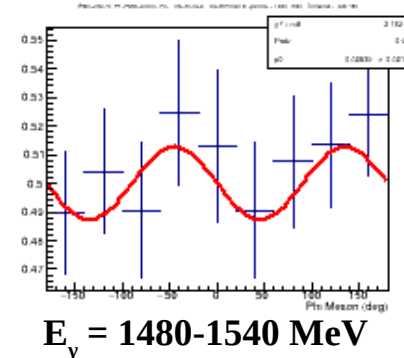
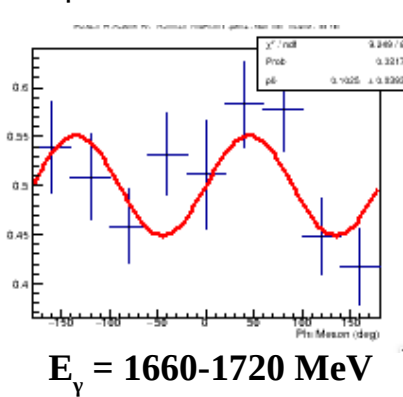
**Proton in Spectr
First Bin in ThetaCM**



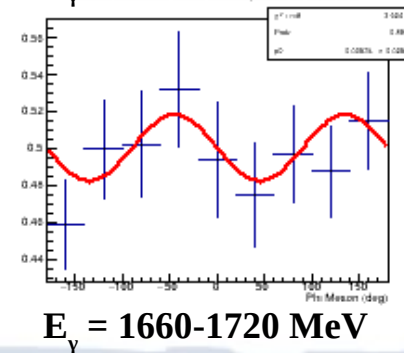
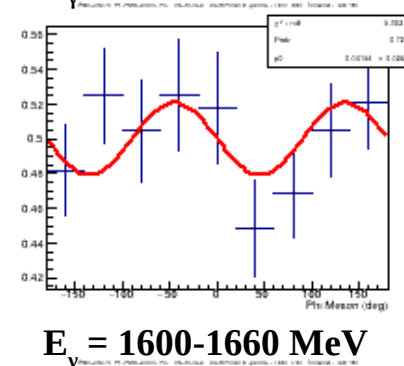
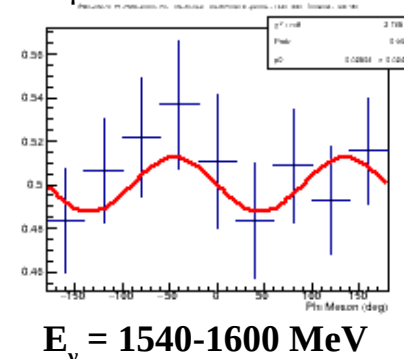
**Proton in Spectr
First Bin in ThetaCM**



**Proton in Spectr
Last Bin in ThetaCM**



**Proton in SciRi
First Bin in ThetaCM**



TO DO LIST

- We analysed only 3 data taking periods on H2 (OctNov2015, AprMay2017, NovDec2018)
 - ➡ We need to implement a reliable procedure for the extraction of the pol. degree of new periods (not trivial because pol. changes also during a d.t. period)
- In η analysis we checked that asymmetry results obtained with a Proton in BGO/Bar or in SciRi are compatible
 - ➡ We can sum up the statistics BEFORE the extraction of the asymmetry in bins with same ThetaCM
 - ➡ This can be applied in the η' analysis to the statistics of Protons in Spectr and in SciRi in order to increase the statistics
- In η analysis we showed that the procedure of kinematical reconstruction applied to $\eta \rightarrow \pi^+ \pi^- \pi^0$ is valid and allows the correct identification of this decay channel when Proton is in SciRi
 - ➡ We can apply this procedure in the analysis of $\eta' \rightarrow \pi^+ \pi^- \eta$ with Proton in SciRi in order to increase the statistics

THANK YOU!

$\gamma p \rightarrow \eta p \rightarrow \pi^+ \pi^- \pi^0 + p$
with: 2 γ in BGO (π^0) + 1 proton in Centr/Interm Det
 $\pi^+ \pi^-$ in BGO

We miss the energies of pions (and also proton when it is in Centr or Interm Det)

We apply momentum conservation along the three directions and we get a system of three equations:

$$0 = p_p \sin \theta_p \cos \varphi_p + E_{\gamma 1} \sin \theta_{\gamma 1} \cos \varphi_{\gamma 1} + E_{\gamma 2} \sin \theta_{\gamma 2} \cos \varphi_{\gamma 2} + p_{\pi 1} \sin \theta_{\pi 1} \cos \varphi_{\pi 1} + p_{\pi 2} \sin \theta_{\pi 2} \cos \varphi_{\pi 2}$$

$$0 = p_p \sin \theta_p \sin \varphi_p + E_{\gamma 1} \sin \theta_{\gamma 1} \sin \varphi_{\gamma 1} + E_{\gamma 2} \sin \theta_{\gamma 2} \sin \varphi_{\gamma 2} + p_{\pi 1} \sin \theta_{\pi 1} \sin \varphi_{\pi 1} + p_{\pi 2} \sin \theta_{\pi 2} \sin \varphi_{\pi 2}$$

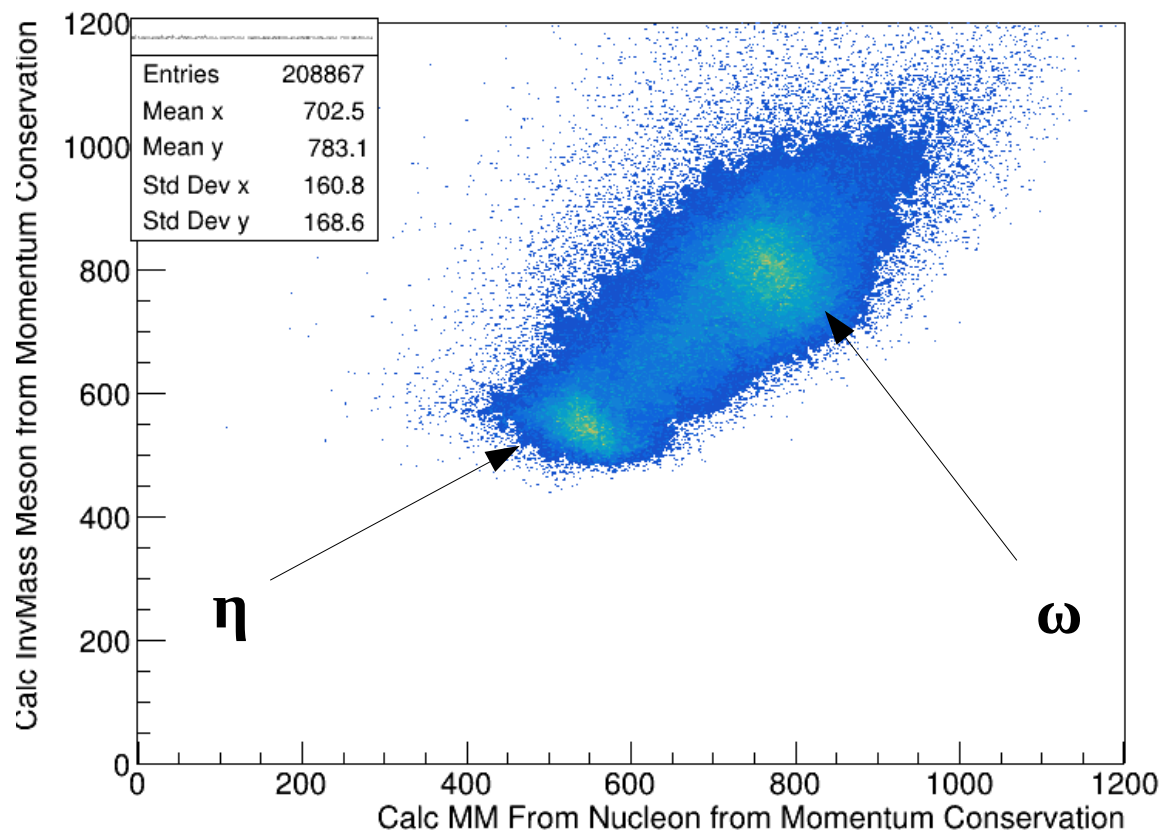
$$E_\gamma = p_p \cos \theta_p + E_{\gamma 1} \cos \theta_{\gamma 1} + E_{\gamma 2} \cos \theta_{\gamma 2} + p_{\pi 1} \cos \theta_{\pi 1} + p_{\pi 2} \cos \theta_{\pi 2}$$

Where:

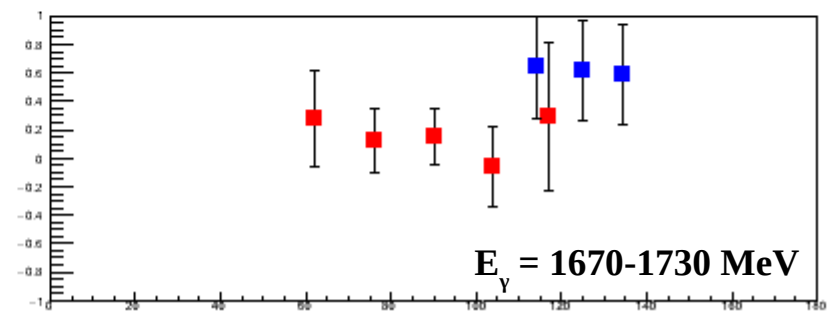
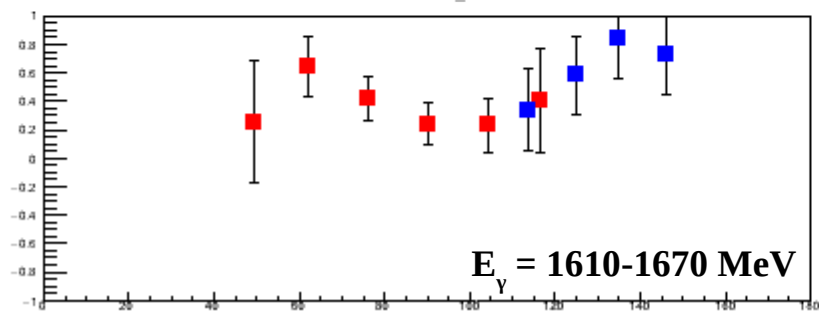
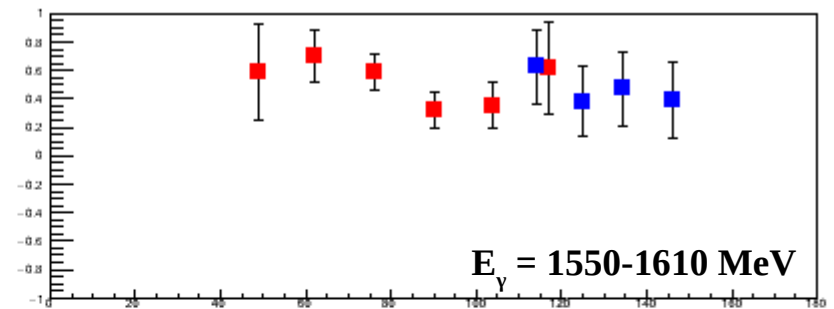
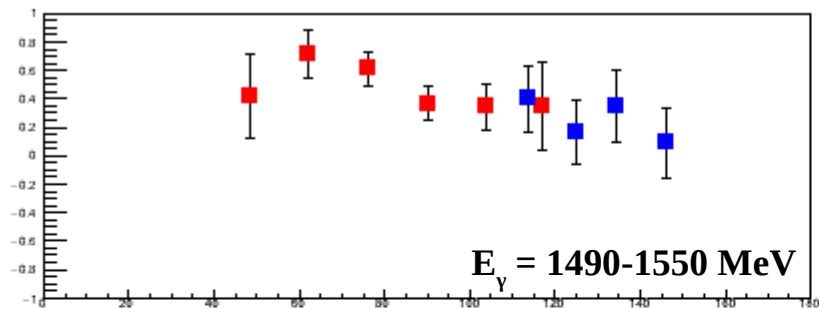
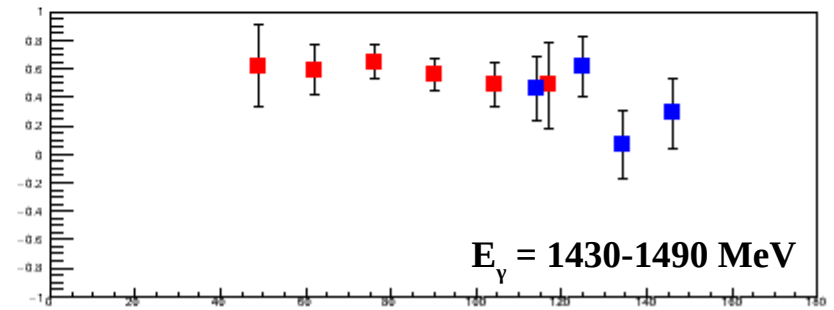
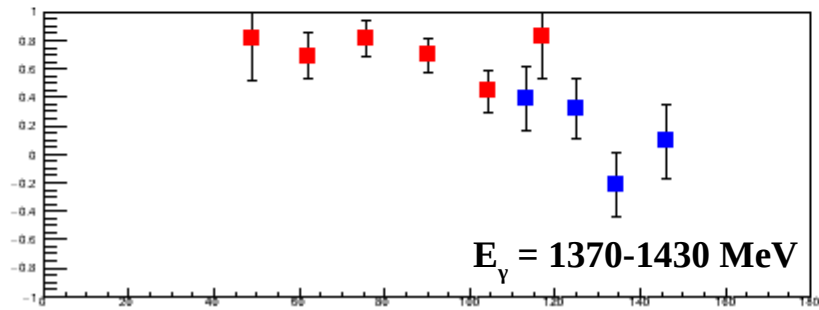
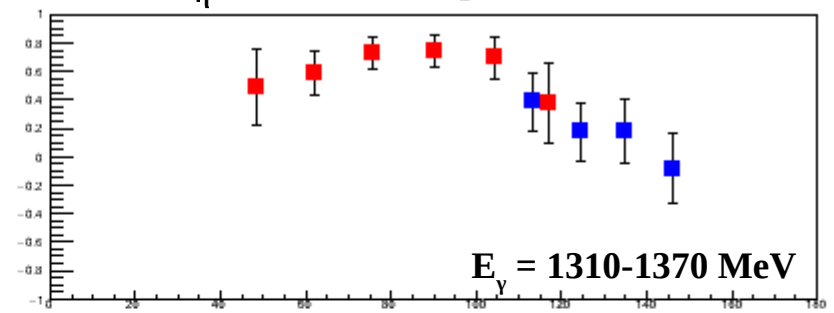
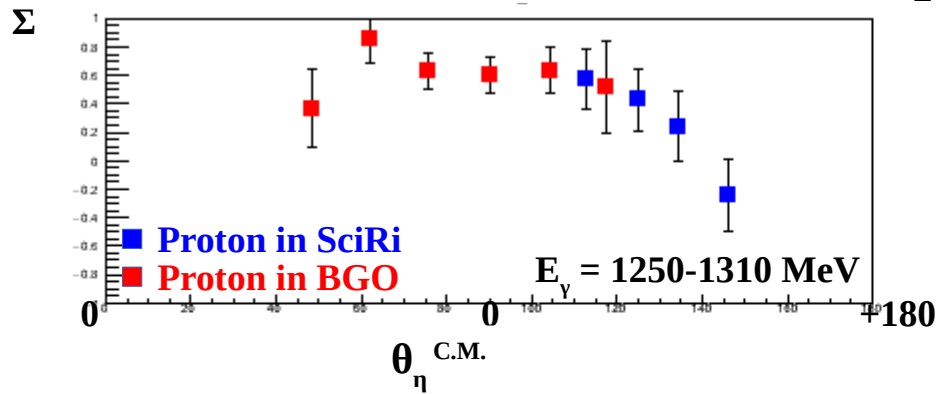
- **proton and charged pions angles** (θ_p, φ_p) $(\theta_{\pi 1}, \varphi_{\pi 1})$ $(\theta_{\pi 2}, \varphi_{\pi 2})$ are measured
- **photons energies and angles** $(\theta_{\gamma 1}, \varphi_{\gamma 1})$ $(\theta_{\gamma 2}, \varphi_{\gamma 2})$ are measured
- **proton and charged pions momenta are unknown**
- **NO HYPOTHESIS has been done on the nature of the meson decaying into $\pi^+ \pi^- \pi^0$**

These means that we can build two independent variables:

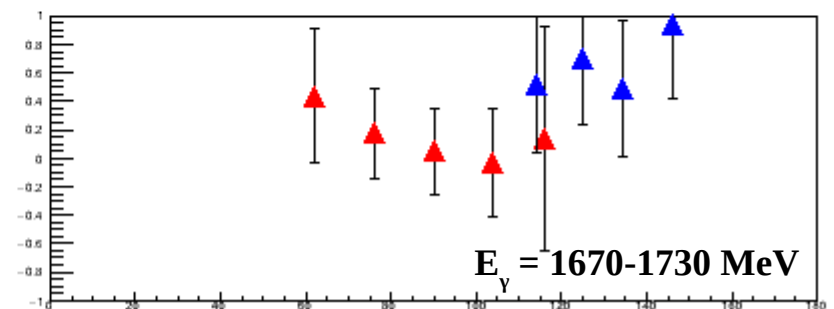
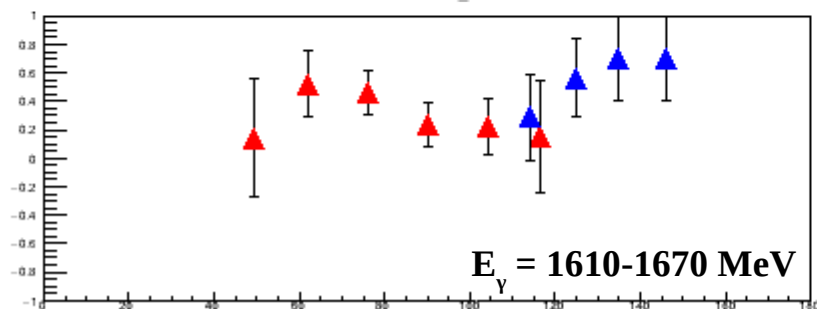
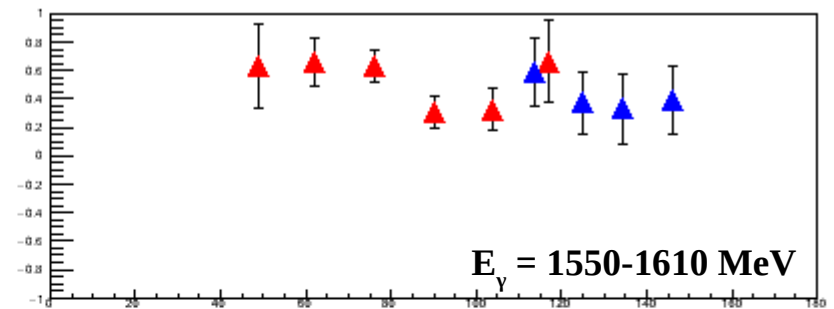
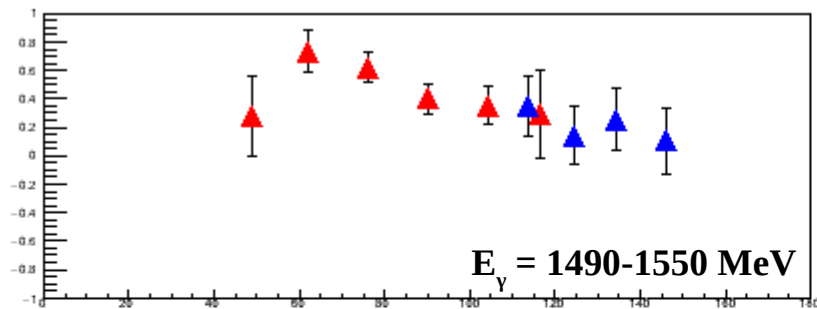
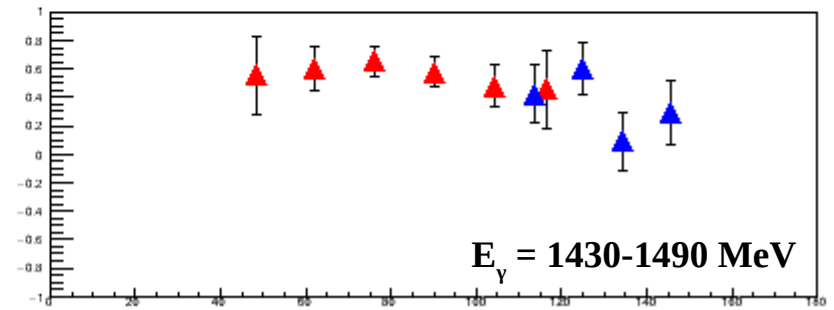
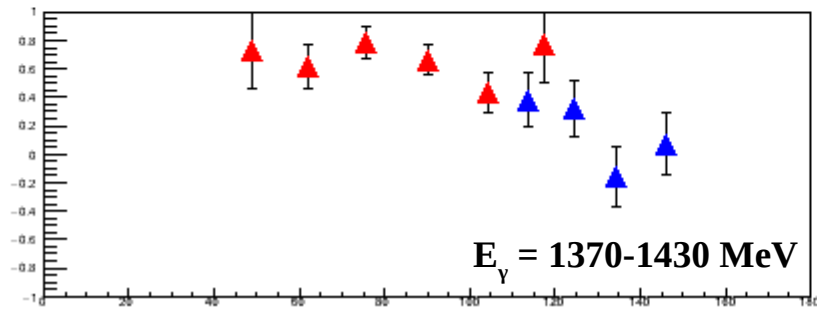
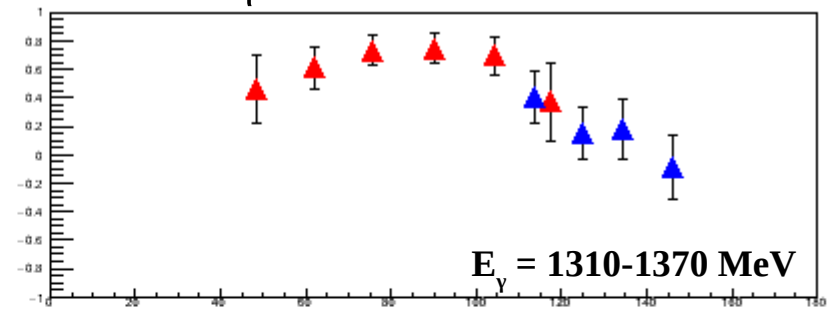
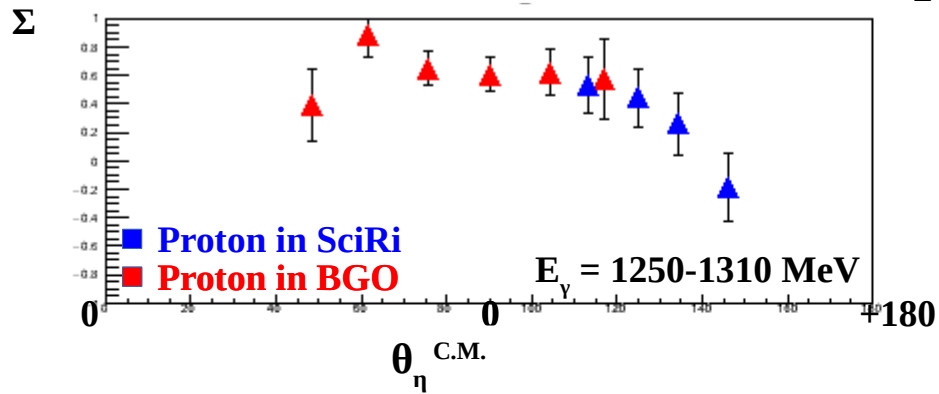
- **Calculated Invariant Mass of ($\pi^0 \pi^+ \pi^-$)** which involves:
Reconstructed $\pi^+ \pi^-$ momenta
Detected π^0
- **Calculated Missing Mass from Proton ($\gamma p \rightarrow p X$)** which involves:
Reconstructed proton momentum
Incident Beam Energy



PolPlus – Σ vs. $\theta_{\eta}^{\text{CM}}$

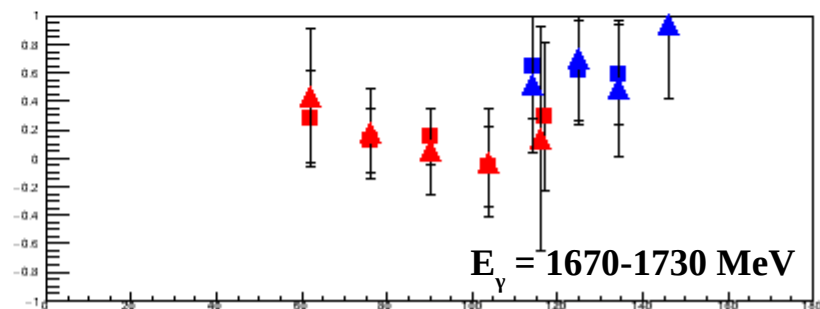
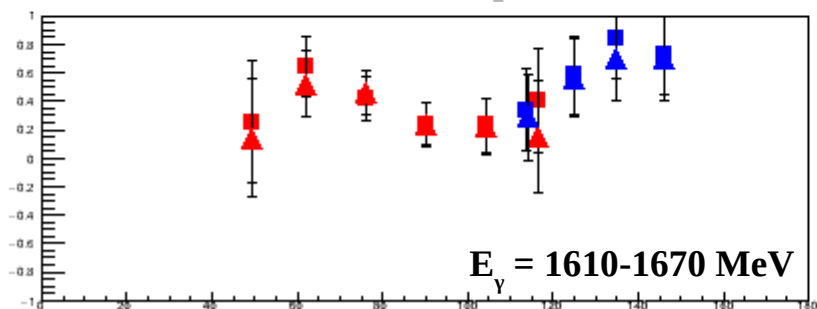
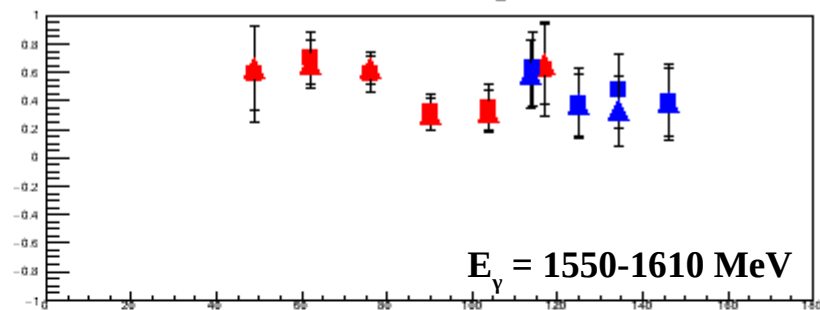
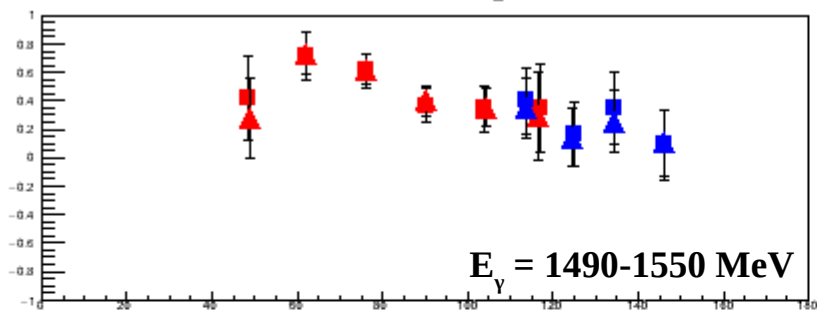
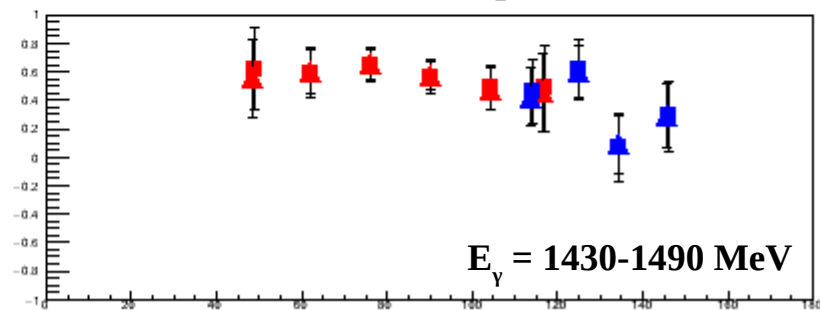
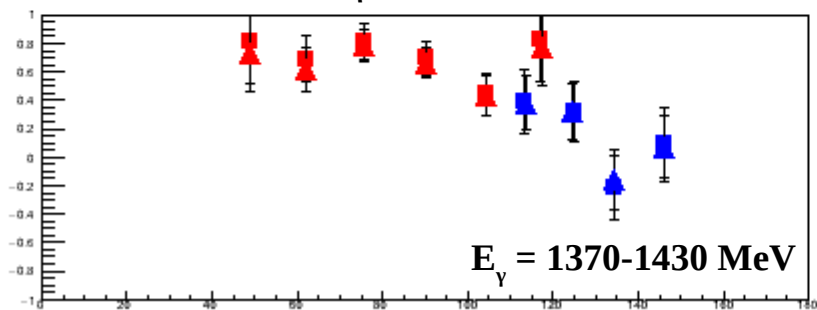
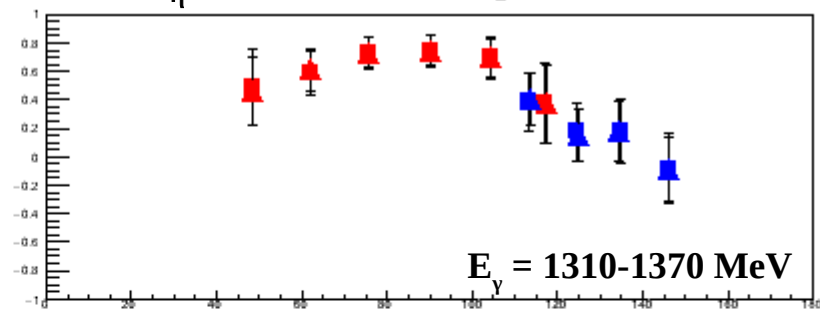
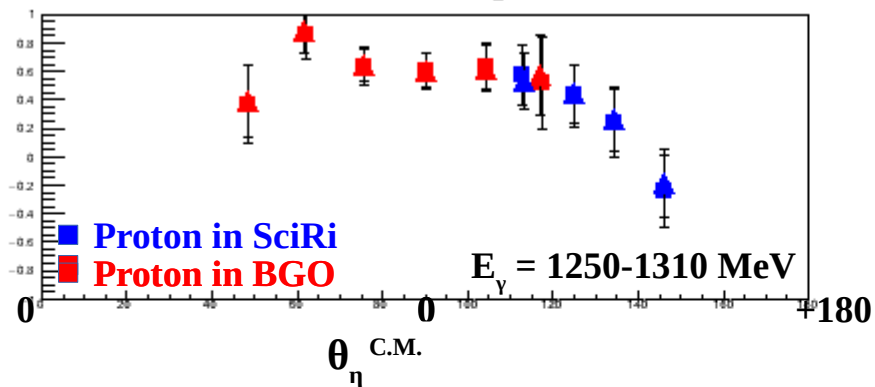


PolMinus - Σ vs. θ_η^{CM}

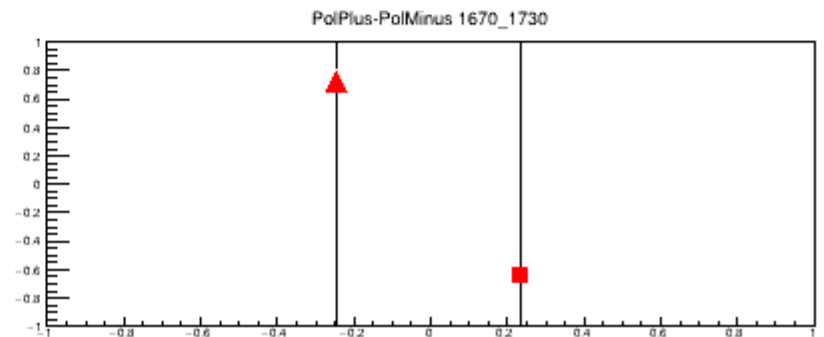
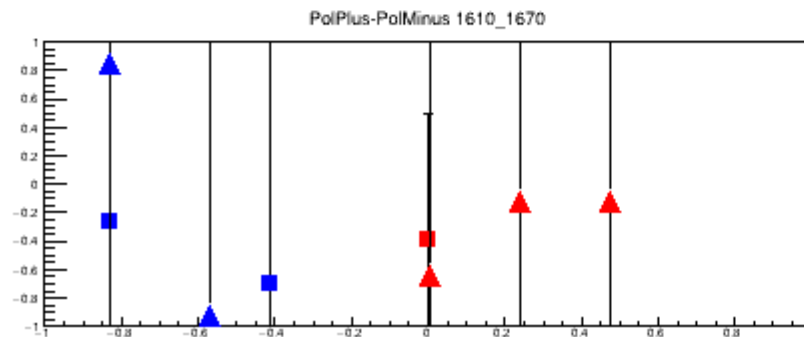
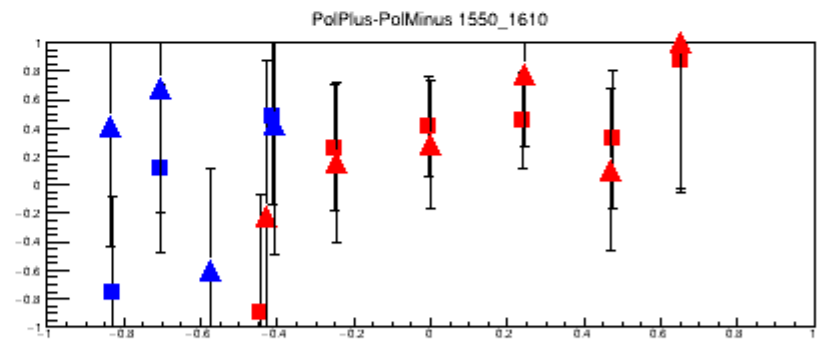
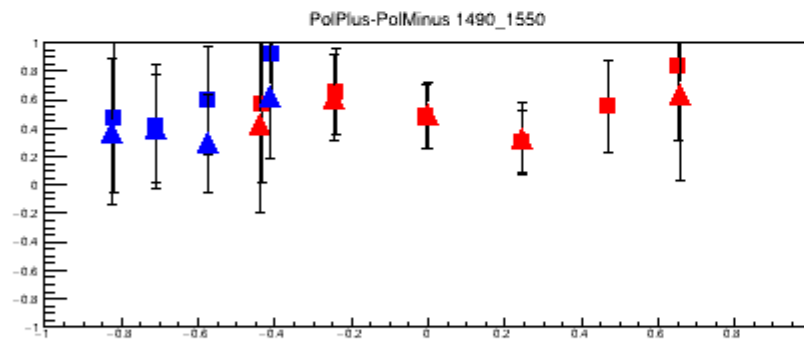
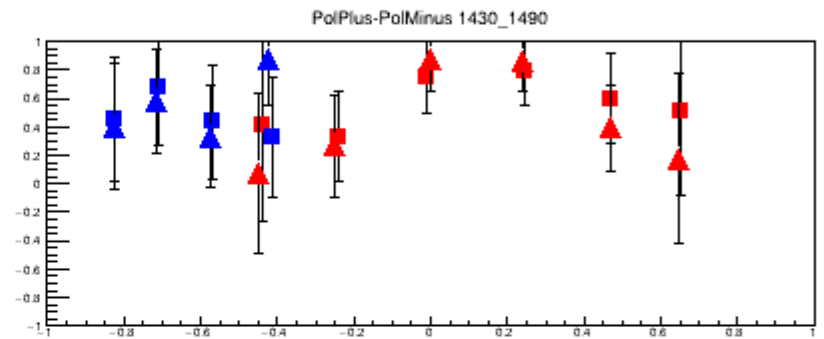
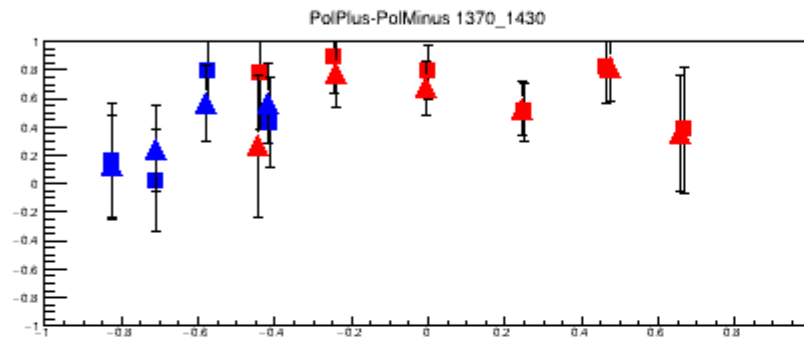
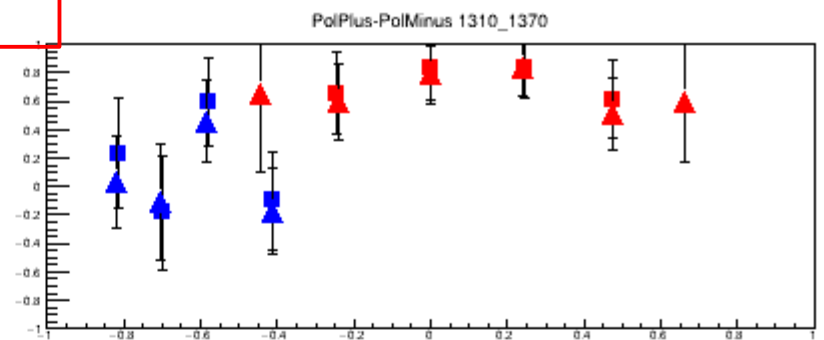
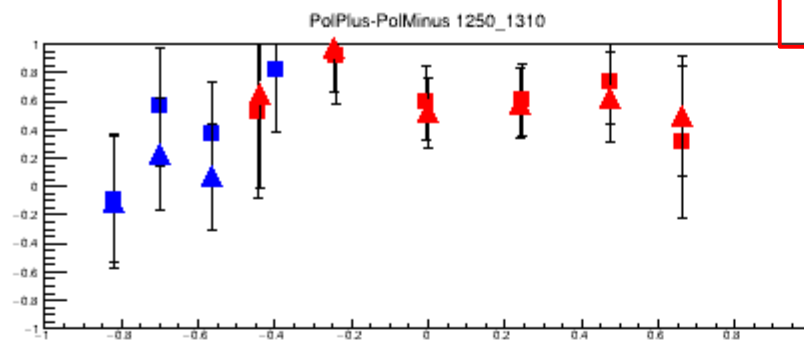


PolPlus-PolMinus – Σ vs. θ_η^{CM}

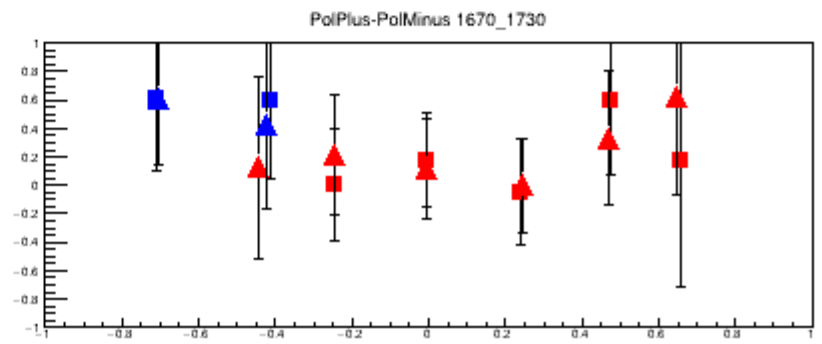
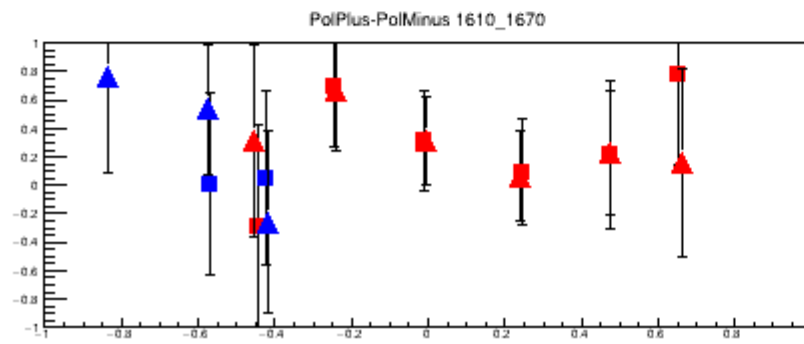
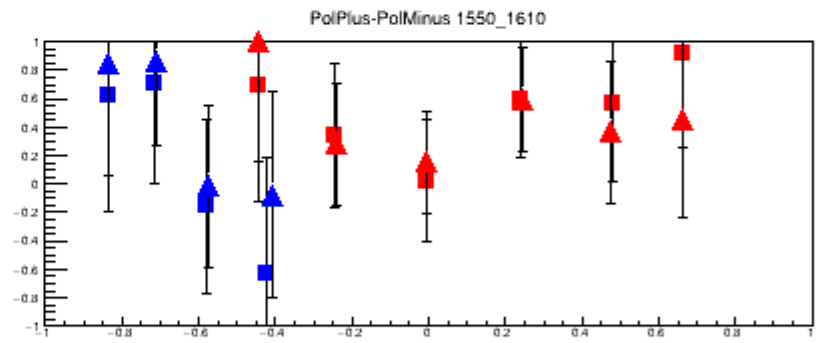
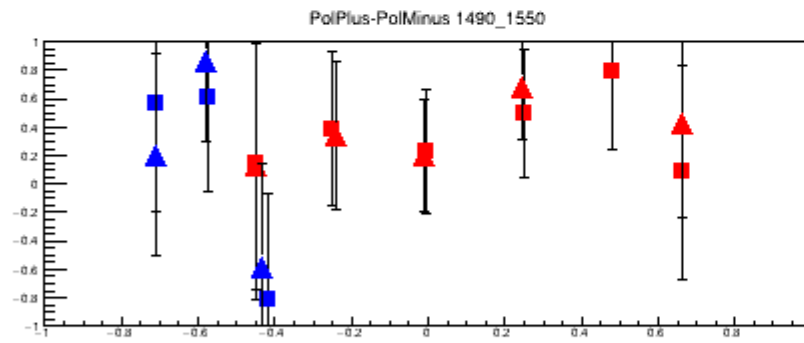
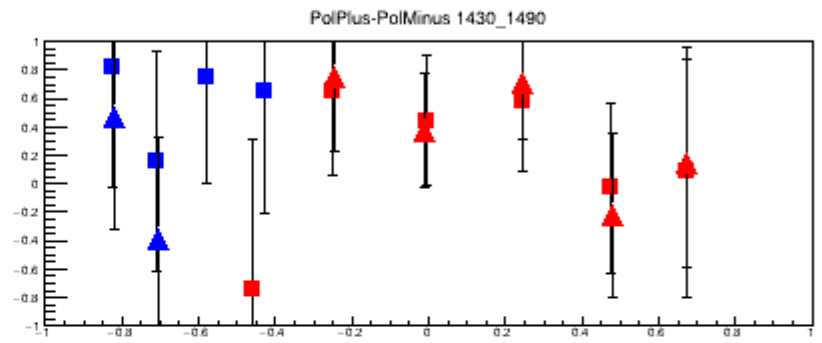
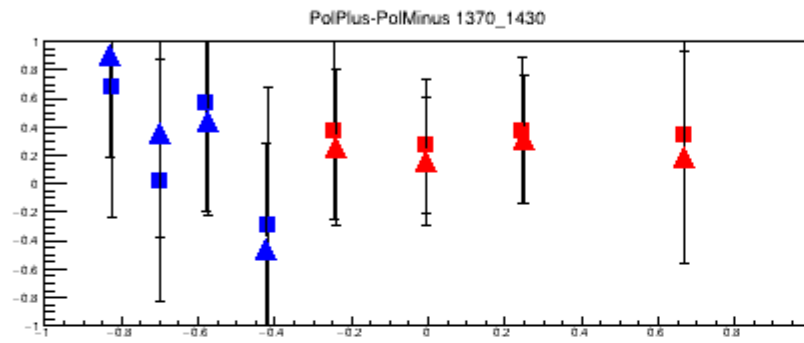
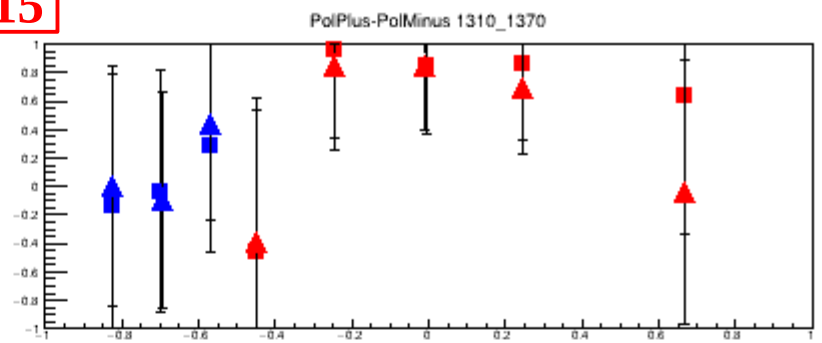
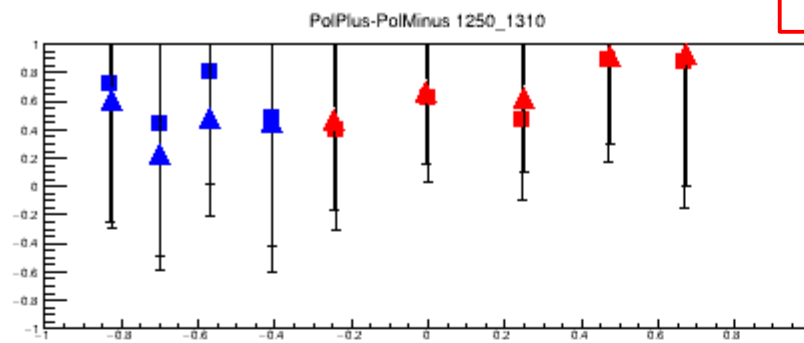
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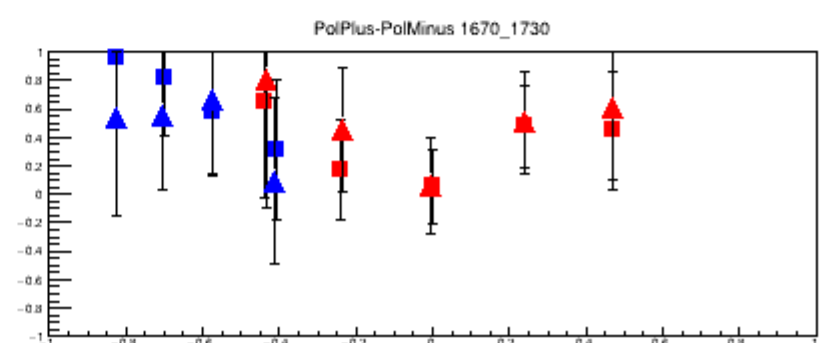
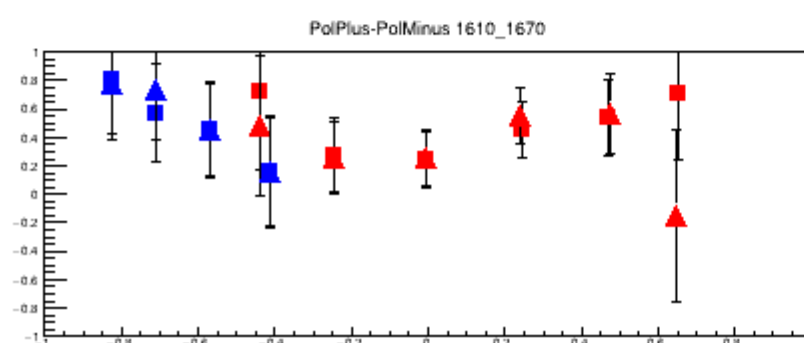
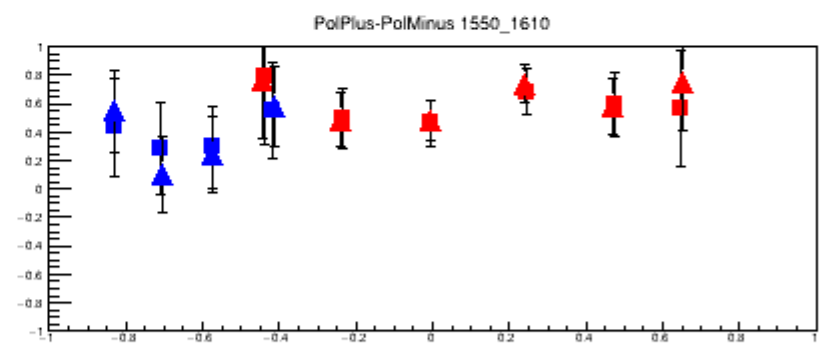
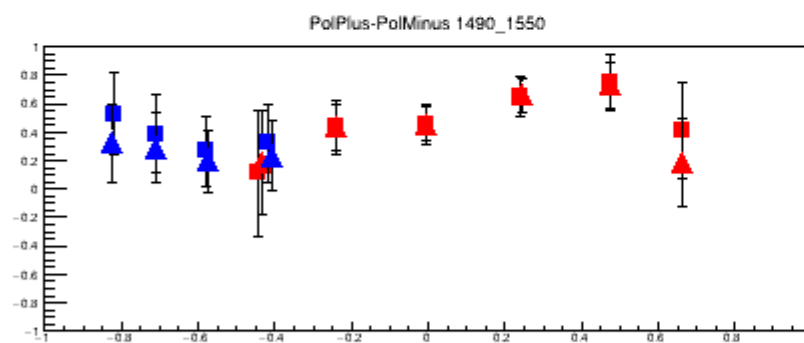
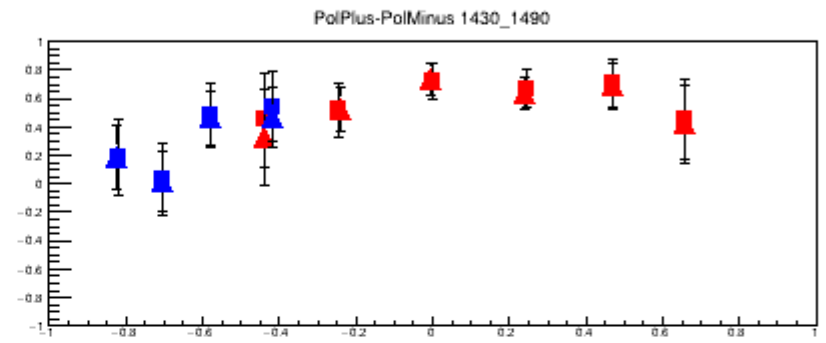
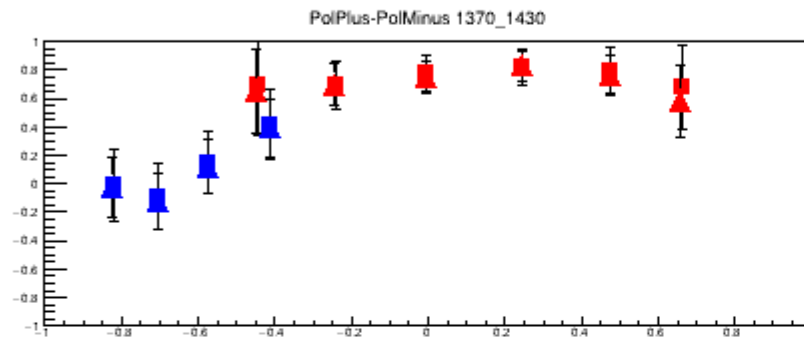
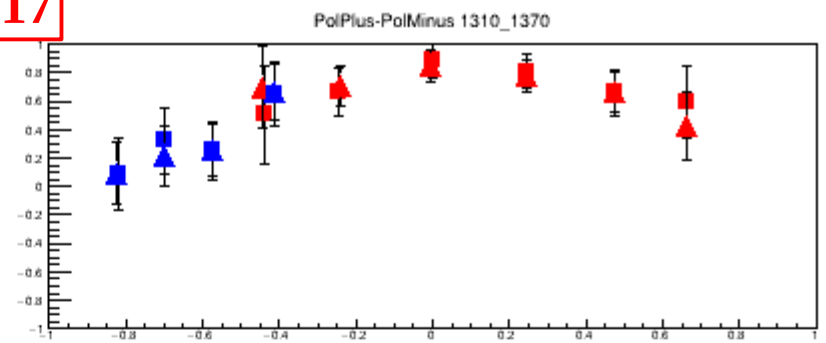
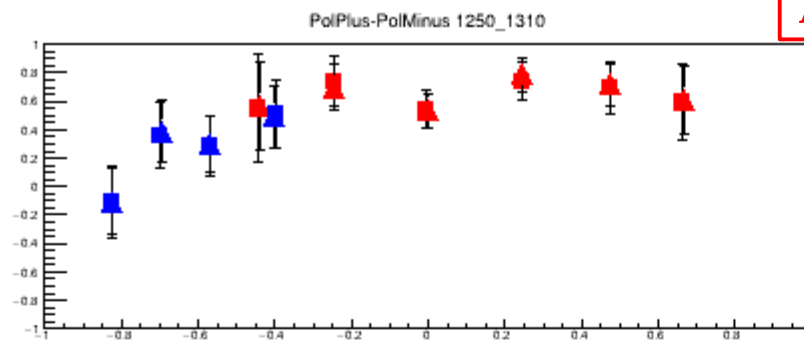
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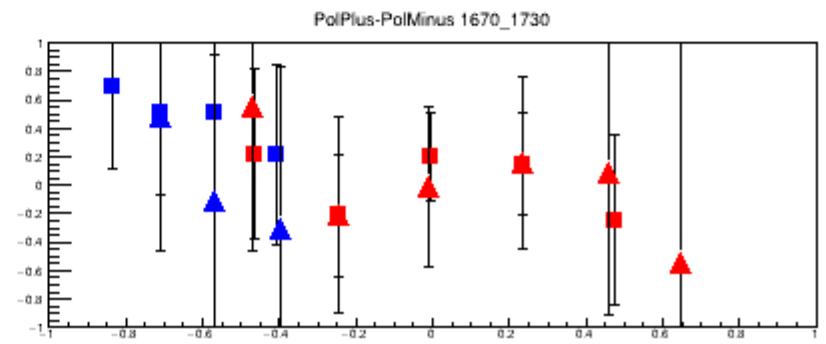
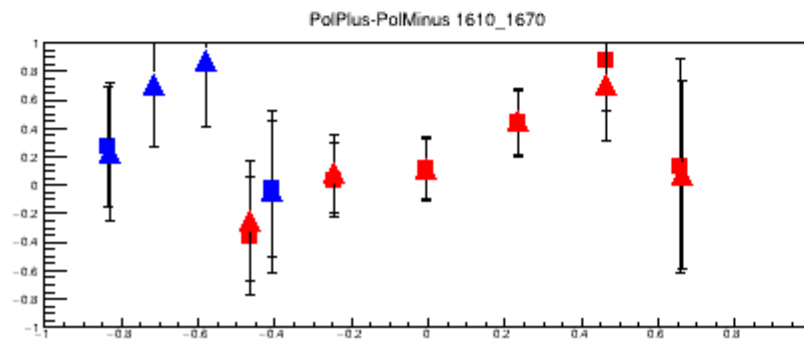
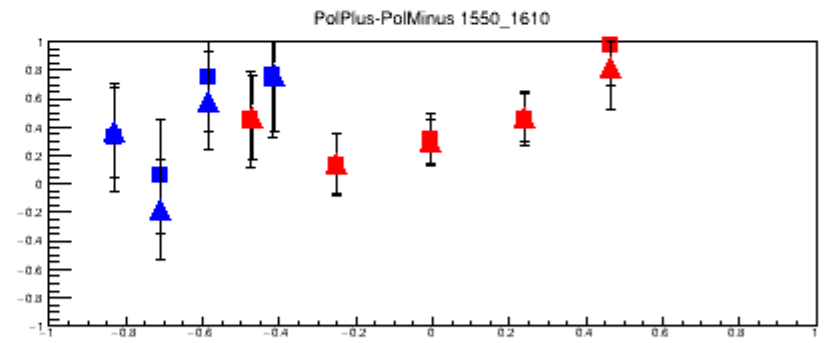
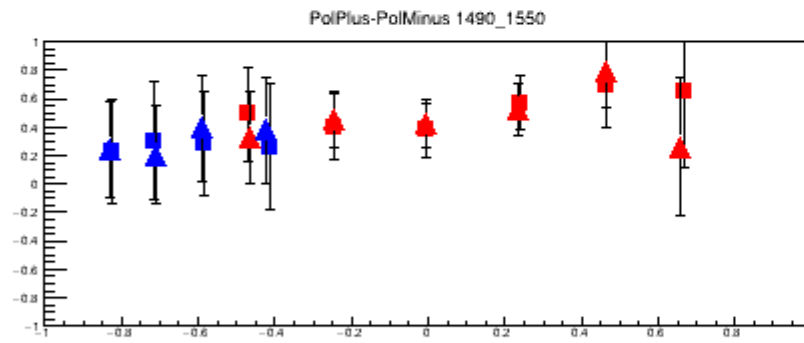
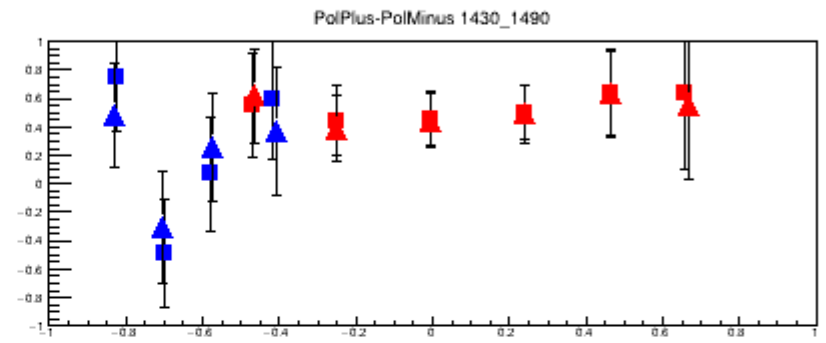
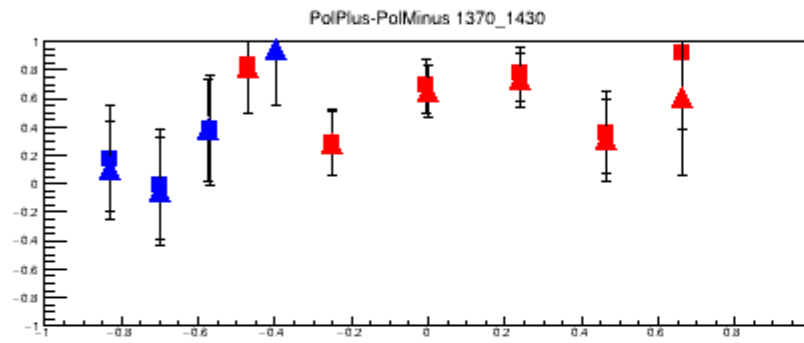
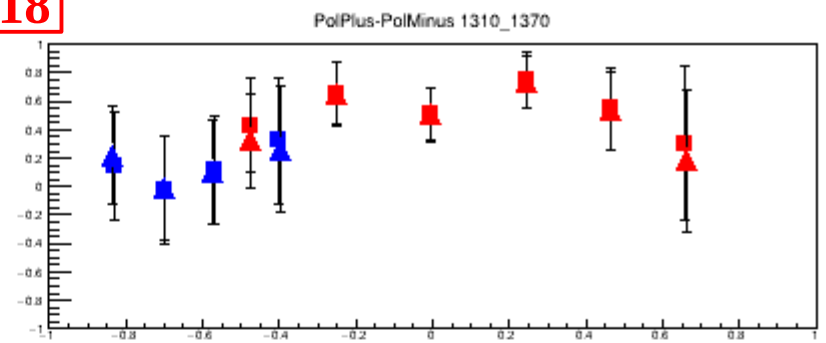
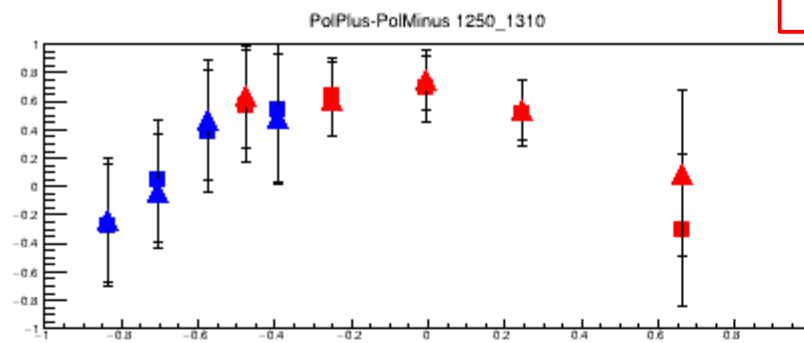
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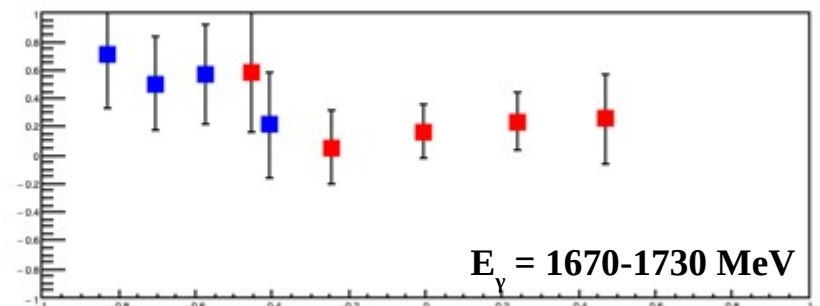
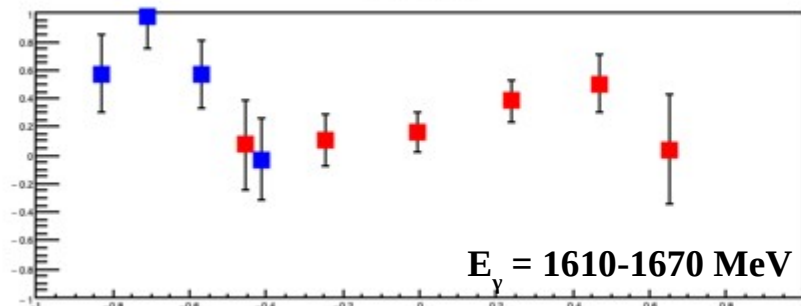
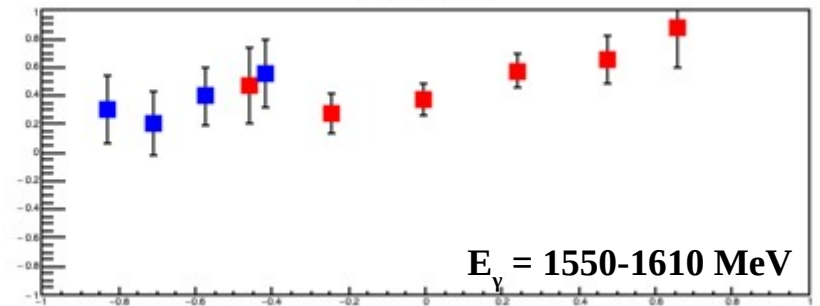
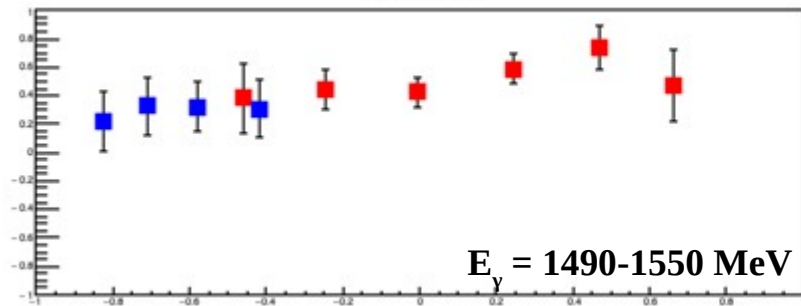
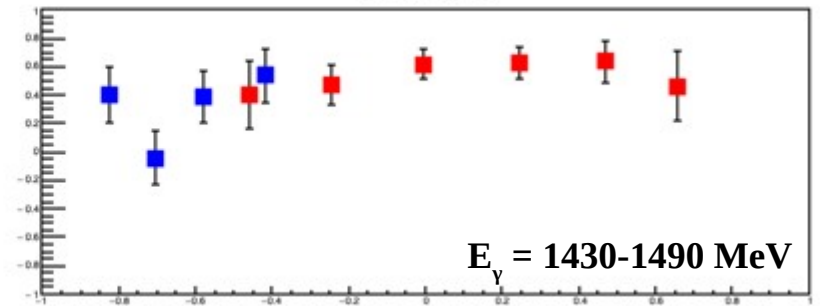
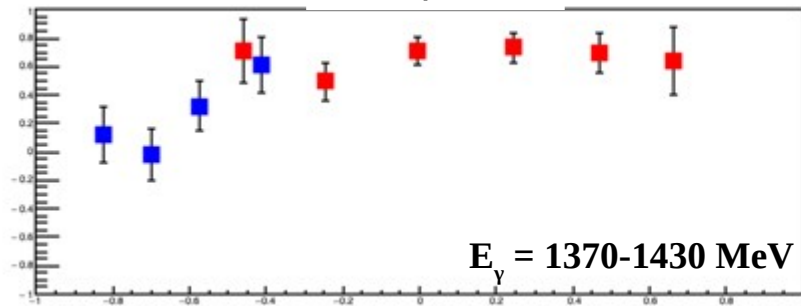
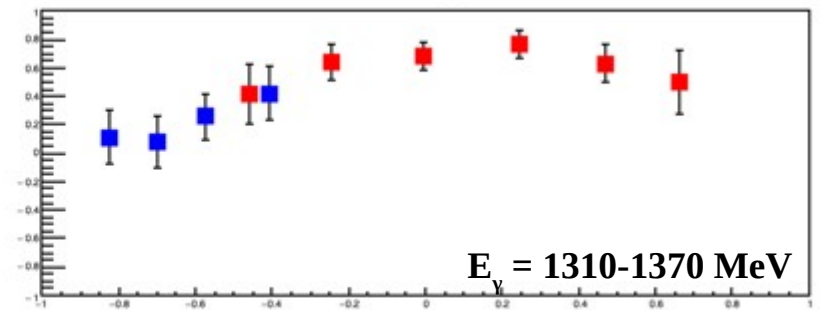
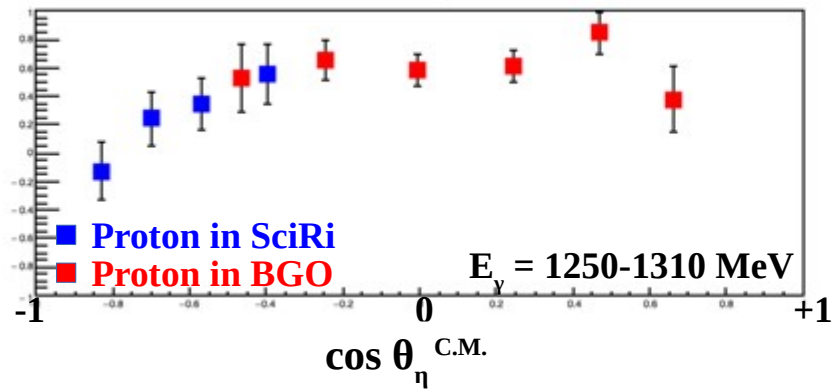


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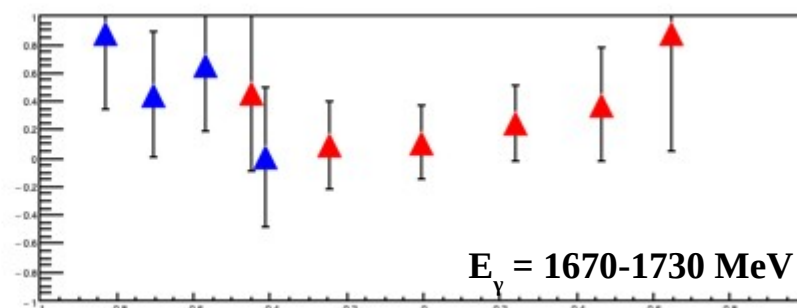
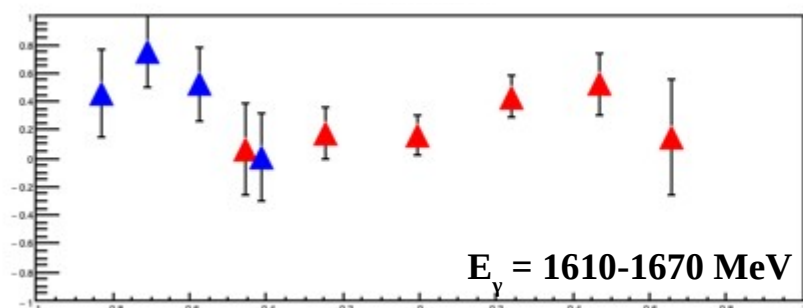
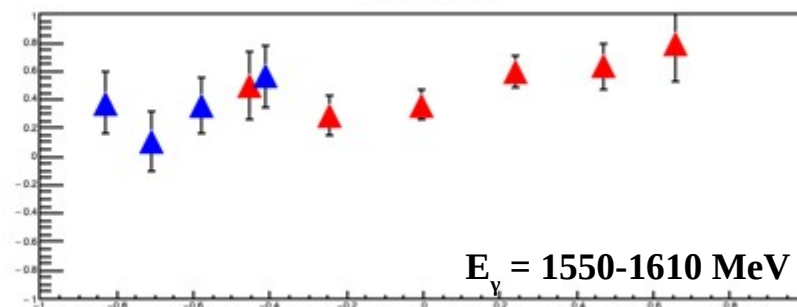
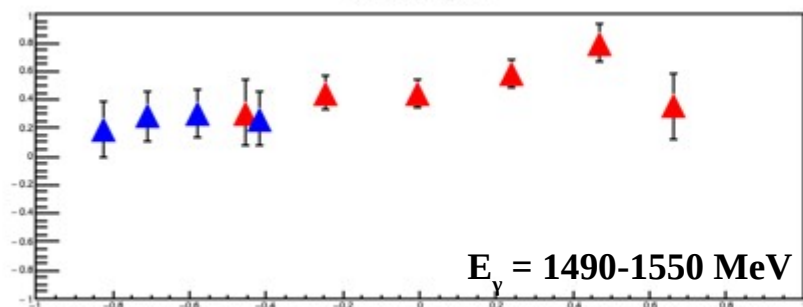
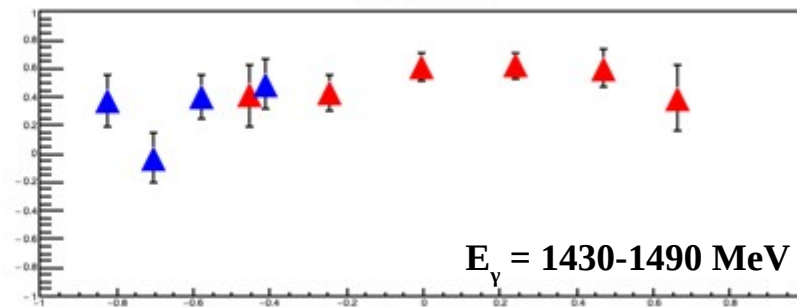
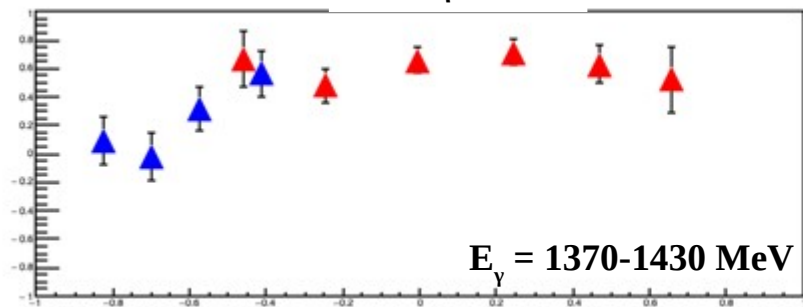
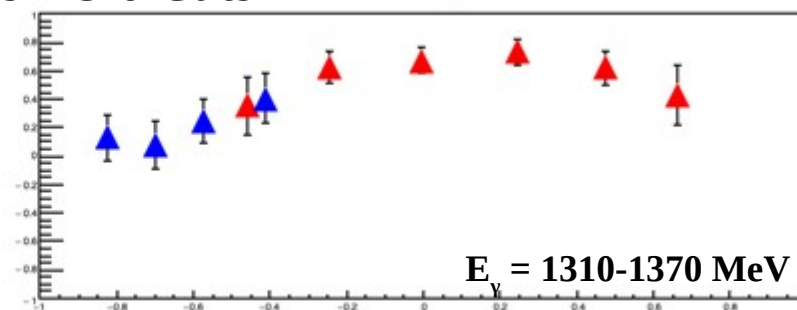
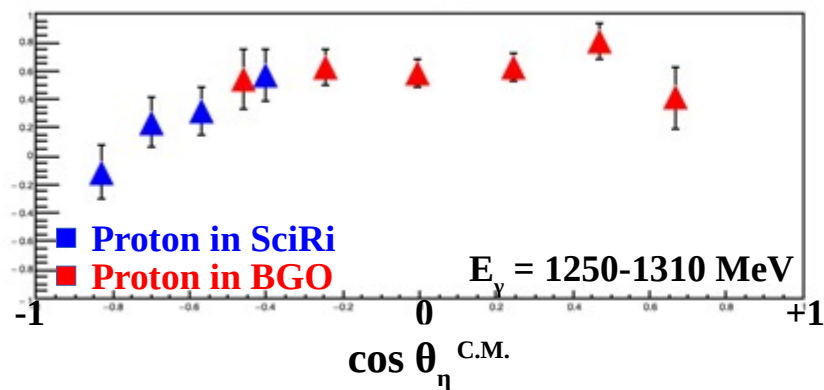
PolPlus – Old Cuts

Σ



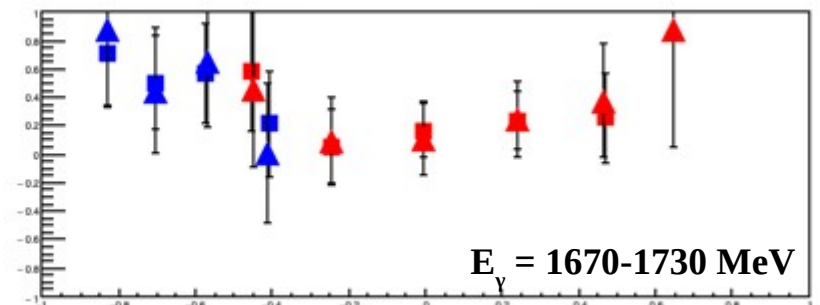
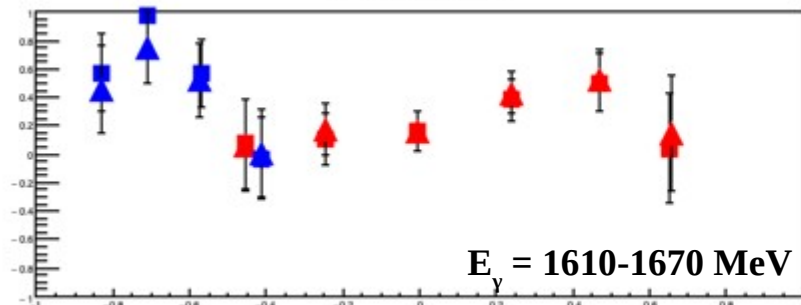
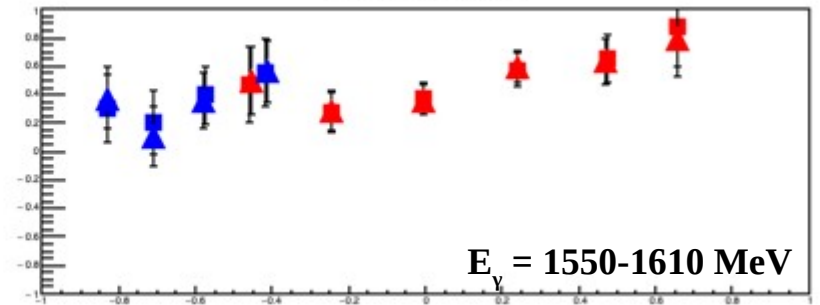
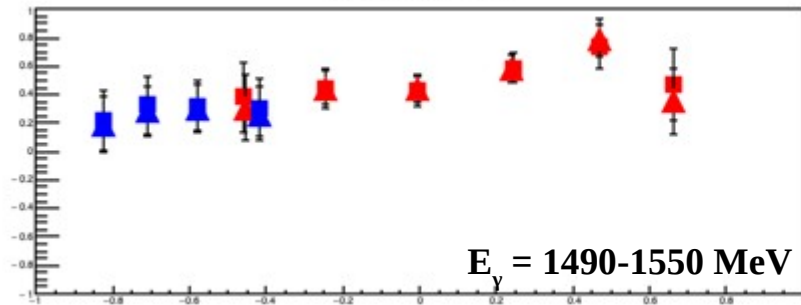
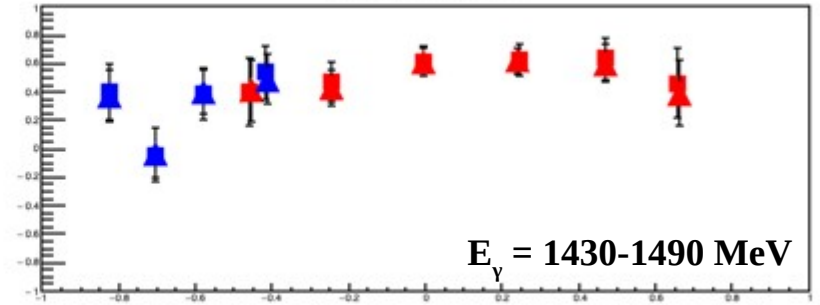
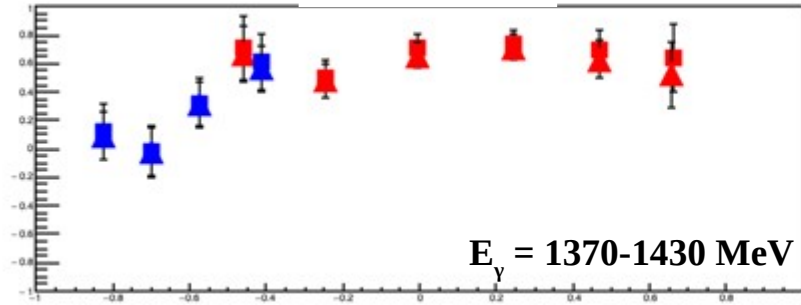
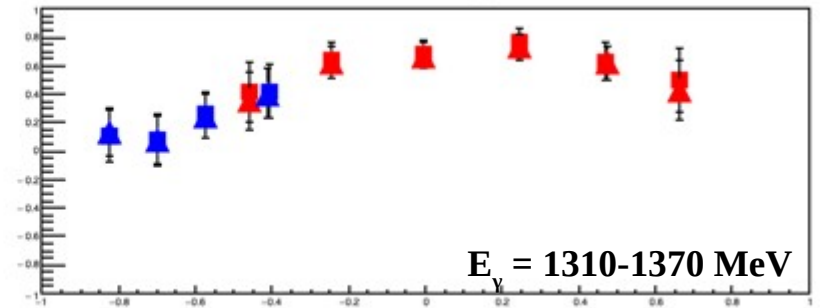
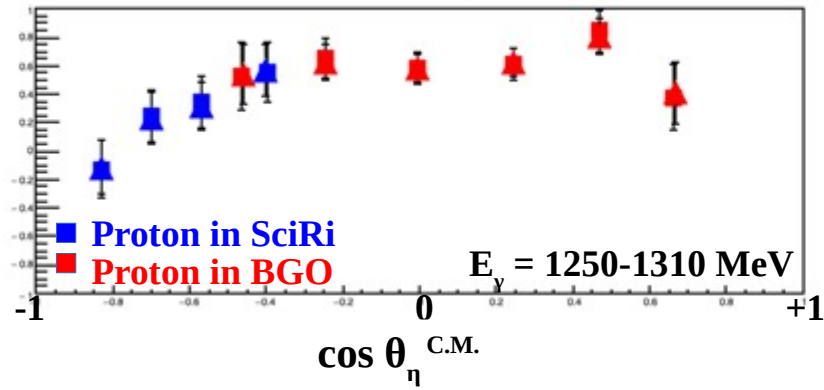
PolMinus – Old Cuts

Σ



PolPlus/PolMinus – Old Cuts

Σ



1) PREVIOUSLY USED TECHNIQUE

We first sum up the $Pol+$ number of events of the two periods and then normalize to the sum of the fluxes of the two periods:

$$\frac{N_{Pol+}^{P1} + N_{Pol+}^{P2}}{F_{Pol+}^{P1} + F_{Pol+}^{P2}} \quad (9)$$

The behaviour of the average of normalized and unpolarized number of events on the number of periods N_{PER} ($N_{PER}=2$ in this case) is in general:

$$\langle N_{UNP, norm} \rangle = \frac{1}{N_{PER}} \sum 2 \varepsilon^{Pi}(\varphi) \quad (10)$$

The behaviour of (9) (with some mathematical passages) becomes:

$$\frac{N_{Pol+}^{P1} + N_{Pol+}^{P2}}{F_{Pol+}^{P1} + F_{Pol+}^{P2}} \propto \frac{\varepsilon^{P1}(\varphi) F_{Pol+}^{P1} (1 + P_{Pol+}^{P1} \Sigma \cos(2\varphi)) + \varepsilon^{P2}(\varphi) F_{Pol+}^{P2} (1 + P_{Pol+}^{P2} \Sigma \cos(2\varphi))}{F_{Pol+}^{P1} + F_{Pol+}^{P2}}$$

ONLY IF the detection/reconstruction efficiency of the two periods is the same:

$$\varepsilon^{P_i}(\varphi) = \varepsilon(\varphi) \quad \forall P_i \quad i=1, \dots, N_{PER}$$

the behaviour of (9) becomes:

$$\frac{N_{Pol+}^{P1} + N_{Pol+}^{P2}}{F_{Pol+}^{P1} + F_{Pol+}^{P2}} = \varepsilon(\varphi) \left[1 + \langle P_{Pol+} \rangle \Sigma \cos(2\varphi) \right] \quad (9)$$

where:

$$\langle P_{Pol+} \rangle = \frac{P_{Pol+}^{P1} F_{Pol+}^{P1} + P_{Pol+}^{P2} F_{Pol+}^{P2}}{F_{Pol+}^{P1} + F_{Pol+}^{P2}} \quad (11)$$

And the efficiency can be simplified in the ratio (9)/(10):

$$\frac{\frac{N_{Pol+}^{P1} + N_{Pol+}^{P2}}{F_{Pol+}^{P1} + F_{Pol+}^{P2}}}{\langle N_{UNP,norm} \rangle} = \frac{\varepsilon(\varphi) \left[1 + \langle P_{Pol+} \rangle \Sigma \cos(2\varphi) \right]}{2\varepsilon(\varphi)} = \frac{1}{2} \left(1 + \langle P_{Pol+} \rangle \Sigma \cdot \cos(2\varphi) \right) \quad (12)$$

This is no more possible if the efficiencies are not the same. For instance for two periods the ratio (9)/(10) becomes:

$$\frac{\frac{N_{Pol+}^{P1} + N_{Pol+}^{P2}}{F_{Pol+}^{P1} + F_{Pol+}^{P2}}}{\langle N_{UNP,norm} \rangle} = \frac{\frac{[\varepsilon^{P1}(\varphi) F_{Pol+}^{P1} + \varepsilon^{P2}(\varphi) F_{Pol+}^{P2}]}{F_{Pol+}^{P1} + F_{Pol+}^{P2}} [1 + \langle P_{Pol+} \rangle \Sigma \cos(2\varphi)]}{\frac{1}{2} (2\varepsilon^{P1}(\varphi) + 2\varepsilon^{P2}(\varphi))} \quad (13)$$



With this technique **it is not possible** to extract the asymmetries from two or more data taking periods **unless the efficiencies are the same**.