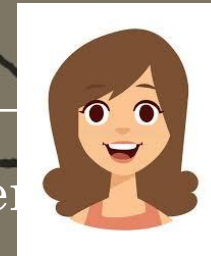
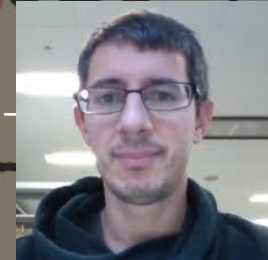
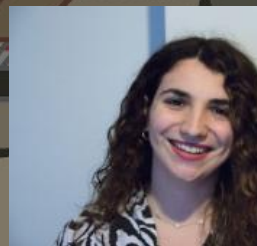


Probing hadronic CP violation with Electric Dipole Moments of paramagnetic molecules

Nikhef



Collaborators: Jordy de Vries, Rob Timmermans, Lemonia Gialidi, Wouter Dekens, Javier Menéndez, Beatriz Romeo

Contents

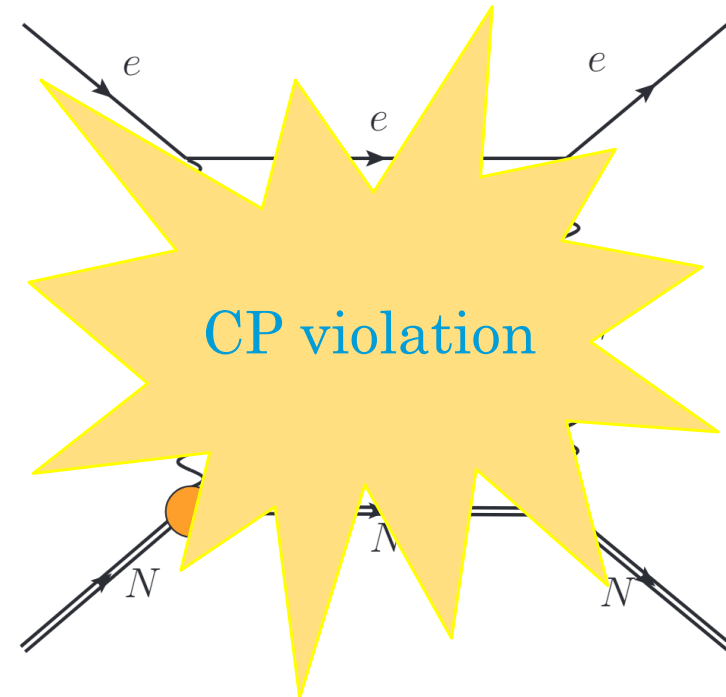
Part 1: hadronic CPV in **para**magnetic EDMs

- Intro on **E**lectric **D**ipole **M**oments
- Sources of CP violation
- Precision development

Part 2: e - ~~N~~ interactions

- Particle & nuclear physics
- Results: bound & disentangle

Nucleus

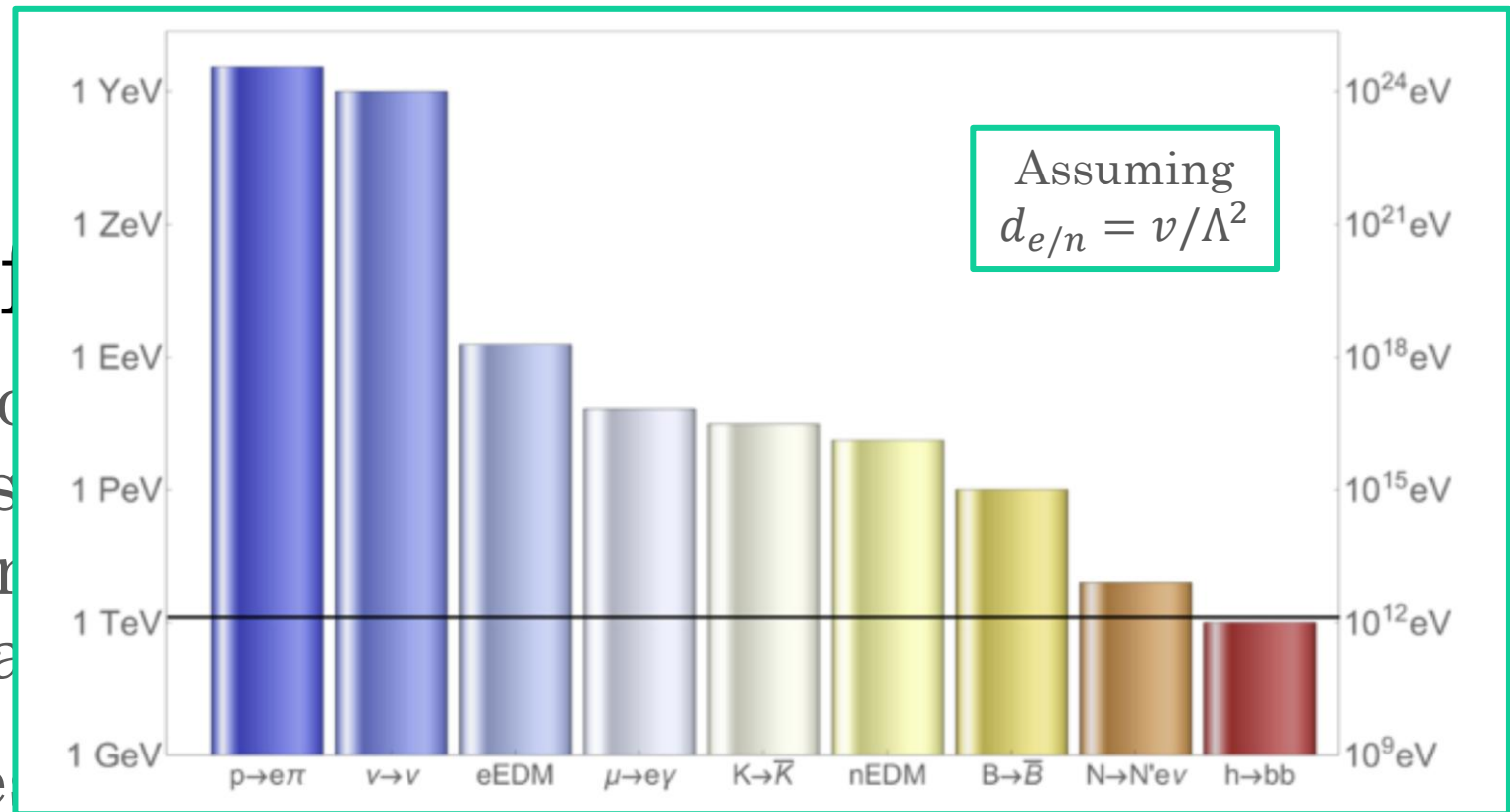


Part 1: hadronic CPV in paramagnetic EDMs

The state of

- Big problems: no neutrino masses, anti-matter asymmetry, dark matter, gravity

How can EDM re

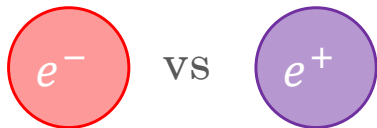


Plot by Adam Falkowski

1

Probe of CP violation

- Need BSM CPV to explain matter anti-matter asymmetry



- EDMs probe this (in case of electroweak baryogenesis)

2

“Model killer”

- Many BSM scenarios predict larger EDMs than SM
- Precise EDM bounds rule out large part of parameter space
- Probe high NP energies Λ



Dipole moments: magnetic & electric

$$L = -\frac{1}{2} \bar{\Psi} \sigma^{\mu\nu} (\mu + i\gamma^5 d) \Psi F_{\mu\nu}$$

Magnetic dipole moment (MDM)

Non-relativistically: $H = -\underbrace{\mu}_{\text{Magnetic dipole moment}} \cdot \underbrace{\vec{S}}_{\text{Spin}} \cdot \underbrace{\vec{B}}_{\text{Magnetic field}}$

Even under T: $-\mu (-\vec{S}) \cdot (-\vec{B}) = -\mu \vec{S} \cdot \vec{B}$

Electric dipole moment (EDM)

Non-relativistically: $H = -\underbrace{d}_{\text{Electric dipole moment}} \cdot \underbrace{\vec{S}}_{\text{Spin}} \cdot \underbrace{\vec{E}}_{\text{Electric field}}$

Odd under T: $-d (-\vec{S}) \cdot (\vec{E}) = +d \vec{S} \cdot \vec{E}$

Through CPT-conservation:
EDM violates CP, MDM doesn't

Also: MDM comes in at tree-level in SM, EDM only at higher loop level $\rightarrow$ MDM much larger than EDM in SM

Paramagnetic: total electron spin non-zero
 Diamagnetic: total electron spin = zero
 Systems: atoms, molecules

Tracing the CP violation

hope to one day

We measure EDMs of:

- Leptons
- Nuclei
- Neutron
- Paramagnetic systems
- Diamagnetic systems

Higher energy

Molecular, atomic, nuclear theory

Heavy-baryon chiral perturbation theory (π, p, n, γ, e)

C_P

C_T

C_{SP}

d_n

$\bar{g}_{\pi/\eta}$

d_p

- Quark EDM
- Quark chromo-EDM
- Gluon chromo-EDM
- 4-quark operators
- Lepton-quark operators
- Lepton EDM (d_e)

Elementary particle level: SM, SMEFT, other BSM

From these fundamental sources of CPV:

CKM phase

QCD $\bar{\theta}$ term

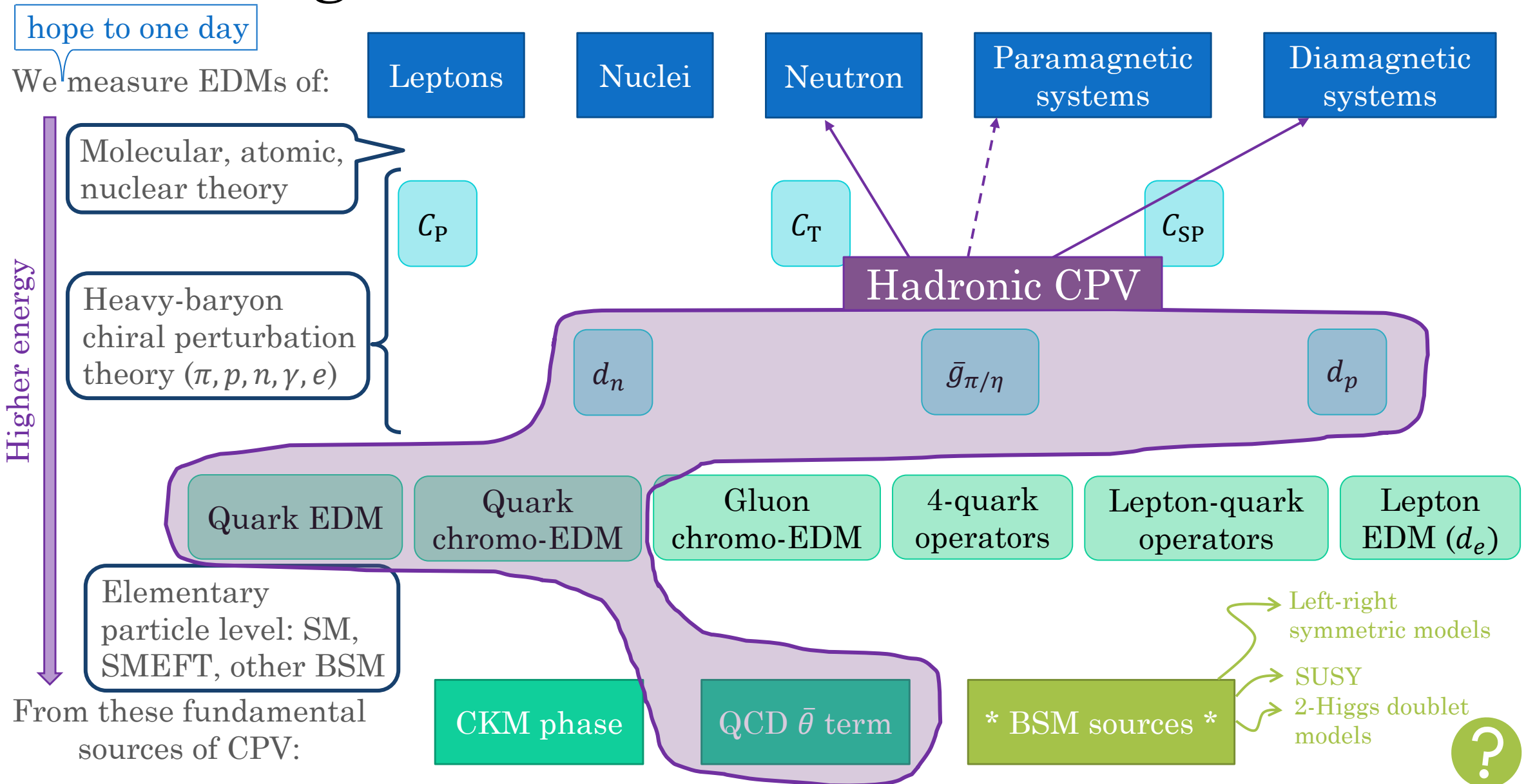
* BSM sources *

- Left-right symmetric models
- SUSY
- 2-Higgs doublet models

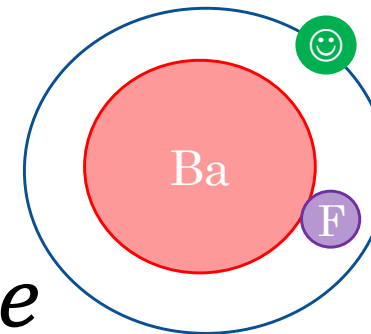


Paramagnetic: total electron spin non-zero
 Diamagnetic: total electron spin = zero
 Systems: atoms, molecules

Tracing the CP violation



Paramagnetic EDMs: C_{SP} and d_e



$$\mathcal{L}_\theta = \bar{\theta} m_* \bar{q} i \gamma^5 q$$

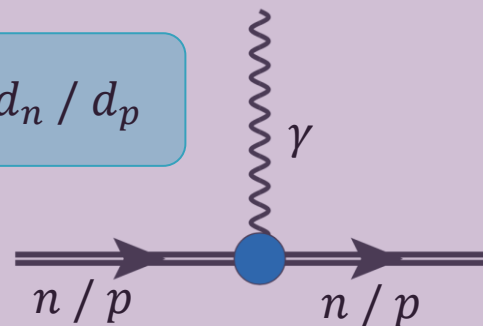
QCD $\bar{\theta}$ term

Quark EDM

Quark chromo-EDM

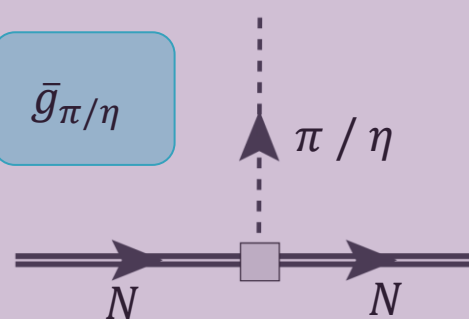
Hadronic CPV

d_n / d_p



$$\mathcal{L} = 2\bar{N}(d_0 + d_1\tau^3)v^\mu S^\nu N F_{\mu\nu}$$

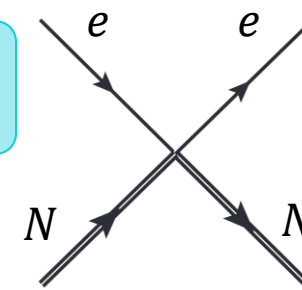
$\bar{g}_{\pi/\eta}$



$$\mathcal{L}_{\pi/\eta NN} = \bar{g}_0 \bar{N} \tau^a N \pi^a + \bar{g}_1 \bar{N} N \pi^0 + \bar{g}_{0\eta} \bar{N} N \eta$$

Semileptonic operator

C_{SP}

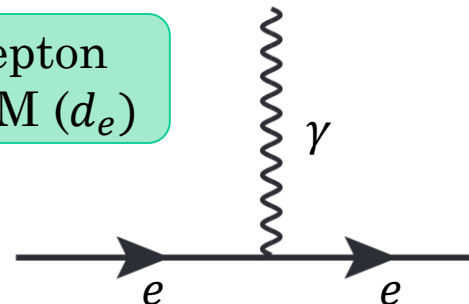


$$\mathcal{L} = C_{SP}^N \frac{G_F}{\sqrt{2}} \bar{e} i \gamma_5 e \bar{N} N$$

Negligible

Leptonic operator

Lepton EDM (d_e)



EDMs of paramagnetic systems

- Paramagnetic EDM = $a C_{SP} + b d_e$
- Usually, bounds on paramagnetic EDMs are interpreted as bounds on d_e
- But! Paramagnetic EDM bounds are now so stringent, that we can derive relevant bounds on $\bar{\theta}, d_p, d_n$ from them too

Precision development

Last ~20 years:

- nEDM bounds improved by factor ~4
- eEDM bounds (from paramagnetic systems) improved by factor ~400
 ↪ $\text{HfF}^+, \text{ThO}, \text{YbF}, \text{Tl} \dots$

Diamagnetic atoms, e.g. ^{199}Hg and ^{225}Ra , improve slowly

Currently:

- $|d_n| \lesssim 10^{-26} e \text{ cm}$

- $|\bar{\theta}|_{\text{nEDM}} \lesssim 10^{-10}$

PSI 2020, J. Liang et al. 2301.04331

- $|d_p|_{\text{Hg}} \lesssim 10^{-25} e \text{ cm}$

Graner et al. 2016, Dmitriev & Sen'kov 2003



TOGETHER

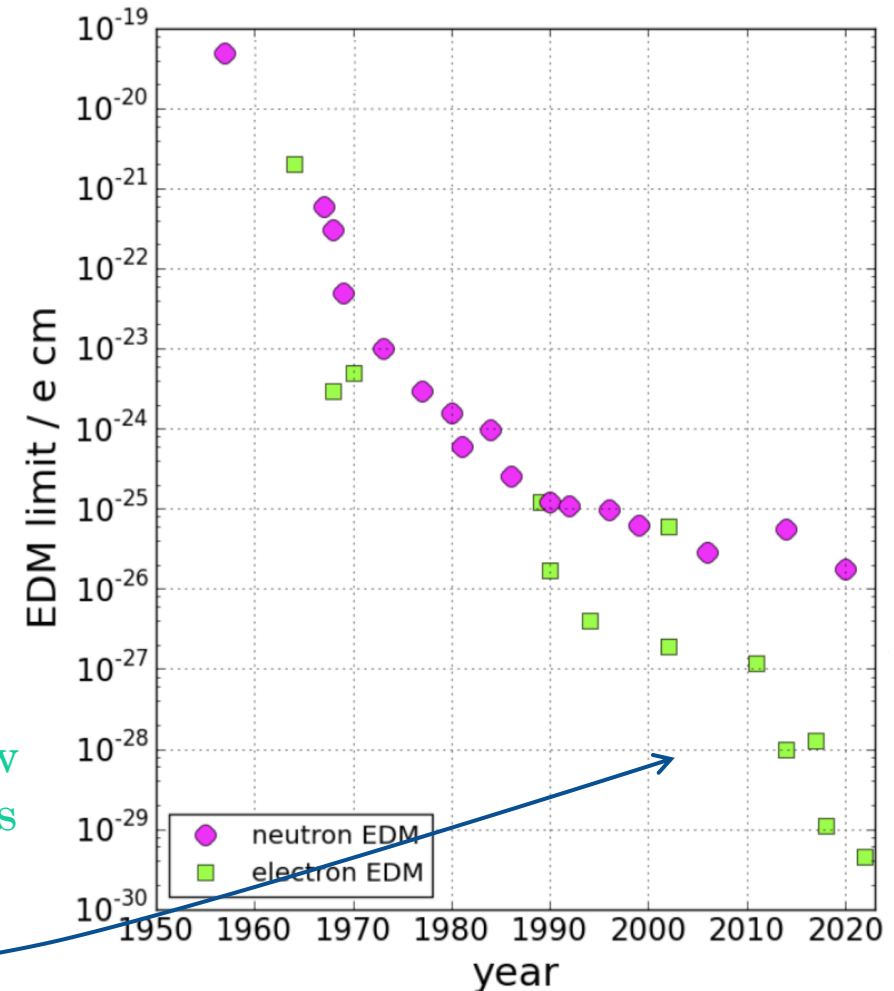
Paramagnetic bounds are now 2 - 3 orders less stringent. But:

1

2

Disentangling → next slide...

Plot by Kseniia Svirina,
ECT* slides 6/3/24



Disentangling

or

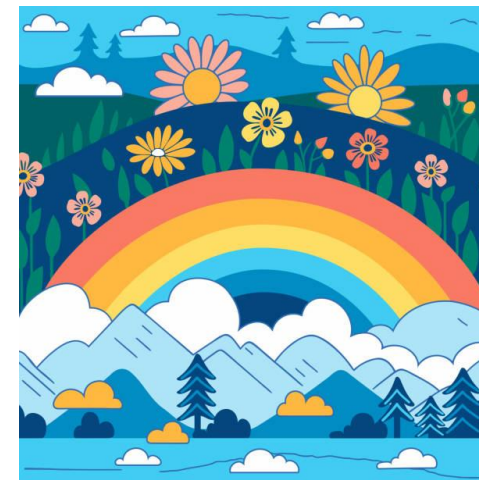
Suppose: non-zero nEDM(**and**)paramagnetic EDM measurements.

What can we say about the underlying CPV?

- We know $d_n = c_1 \bar{\theta}$ and $\omega = c_2 \bar{\theta}$ from theory
- Can check d_n/ω from theory against experiment
- Same for other CPV sources: $d_n = c_3 d_q$ and $\omega = c_4 d_q$, gives different d_n/ω

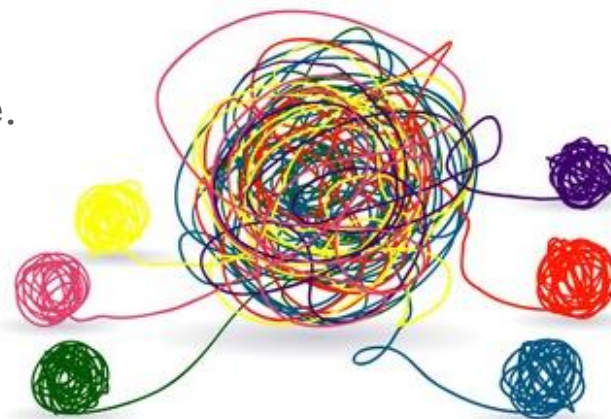
So: disentangle possible sources!

Also useful if precision is not (yet) the same.



$$\omega = a \textcircled{C_{\text{SP}}} + b d_e$$

Hadronic CPV:
 $\bar{\theta}, d_q \dots$



Recap slide

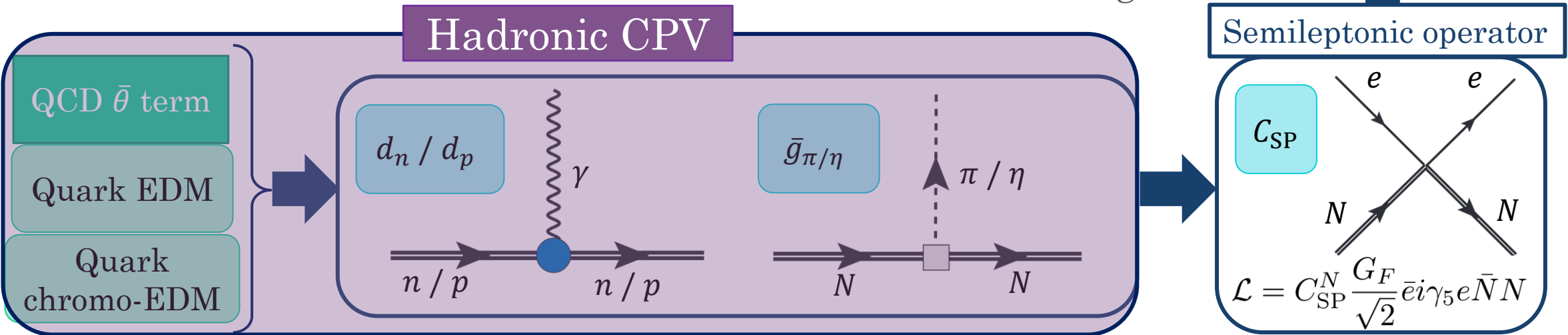
- EDMs are a powerful probe of CP-violation
- The EDM field spans over a wide range of energies, from fundamental CP-violating sources to experiments
- Measurable EDMs can be sourced by multiple CP-odd sources, through different ‘avenues’
- Hadronic CP-violating sources are e.g. d_p , d_n , $\bar{\theta}$
- Paramagnetic EDM searches have improved a lot in sensitivity over the past decades
- We want to investigate hadronic CP violation in paramagnetic EDMs. Why?
 - 1 Bound d_p , d_n , $\bar{\theta}$...
 - 2 Disentangle possible sources

as opposed to leptonic

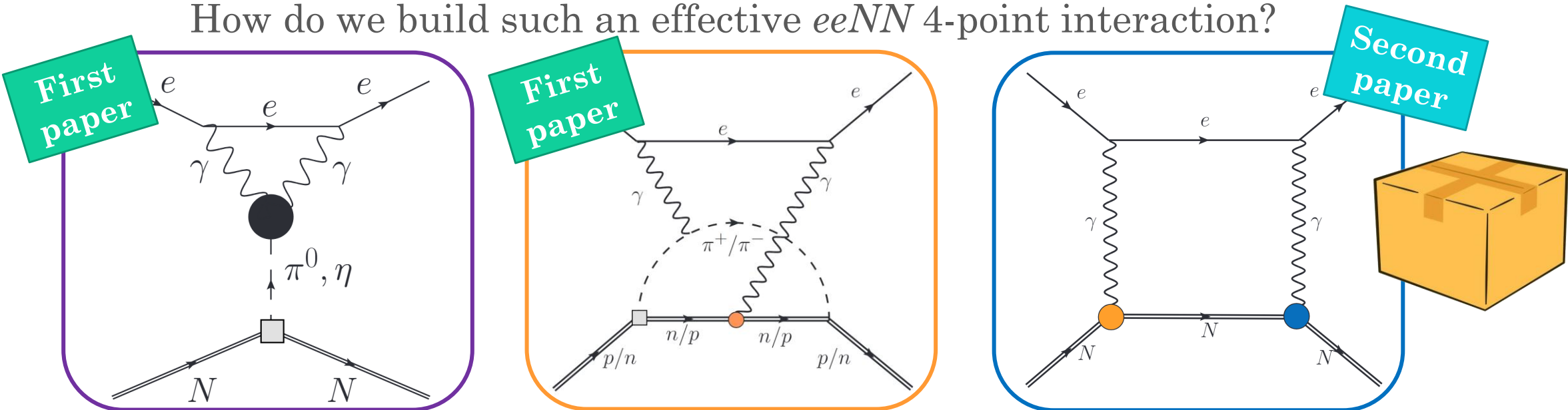
Part 2: electron-nucleus interactions

A closer look at C_{SP}

Recall the ‘chain’ of CP violation we want to investigate:



How do we build such an effective $eeNN$ 4-point interaction?

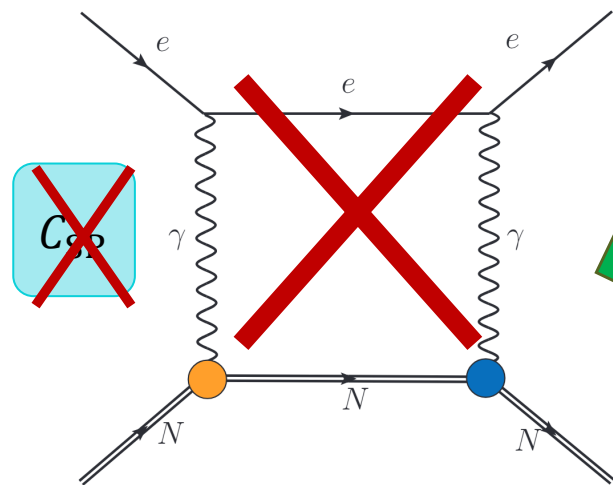
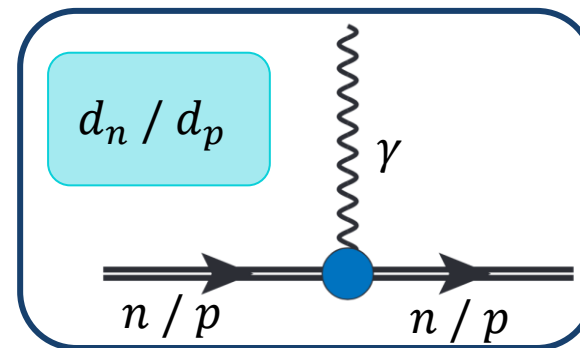


Simple?

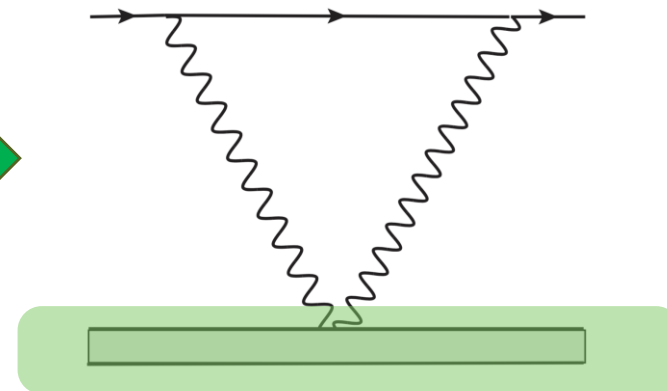
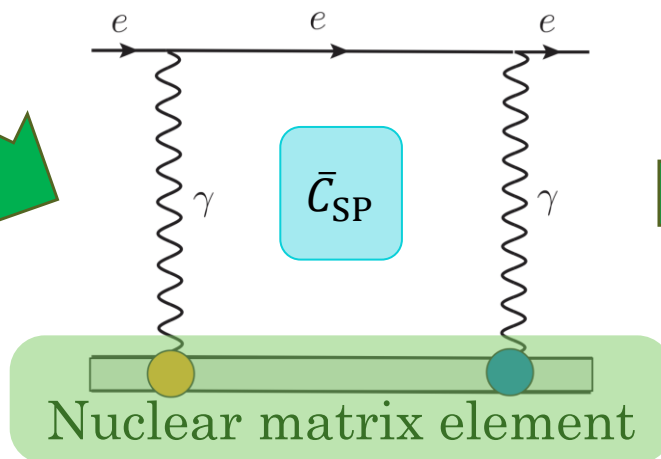


Tackling boxes

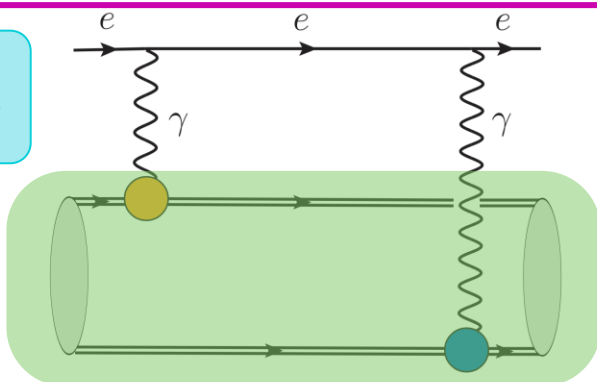
- **Ultrasoft** photon momenta $q_\gamma^0 \sim |\vec{q}_\gamma| \sim m_\pi^2/m_N$
- Probe **nuclear** structure



- Difficult integrals... use EFT: $E_n \gg m_e$
- Integrate out nuclear excited states

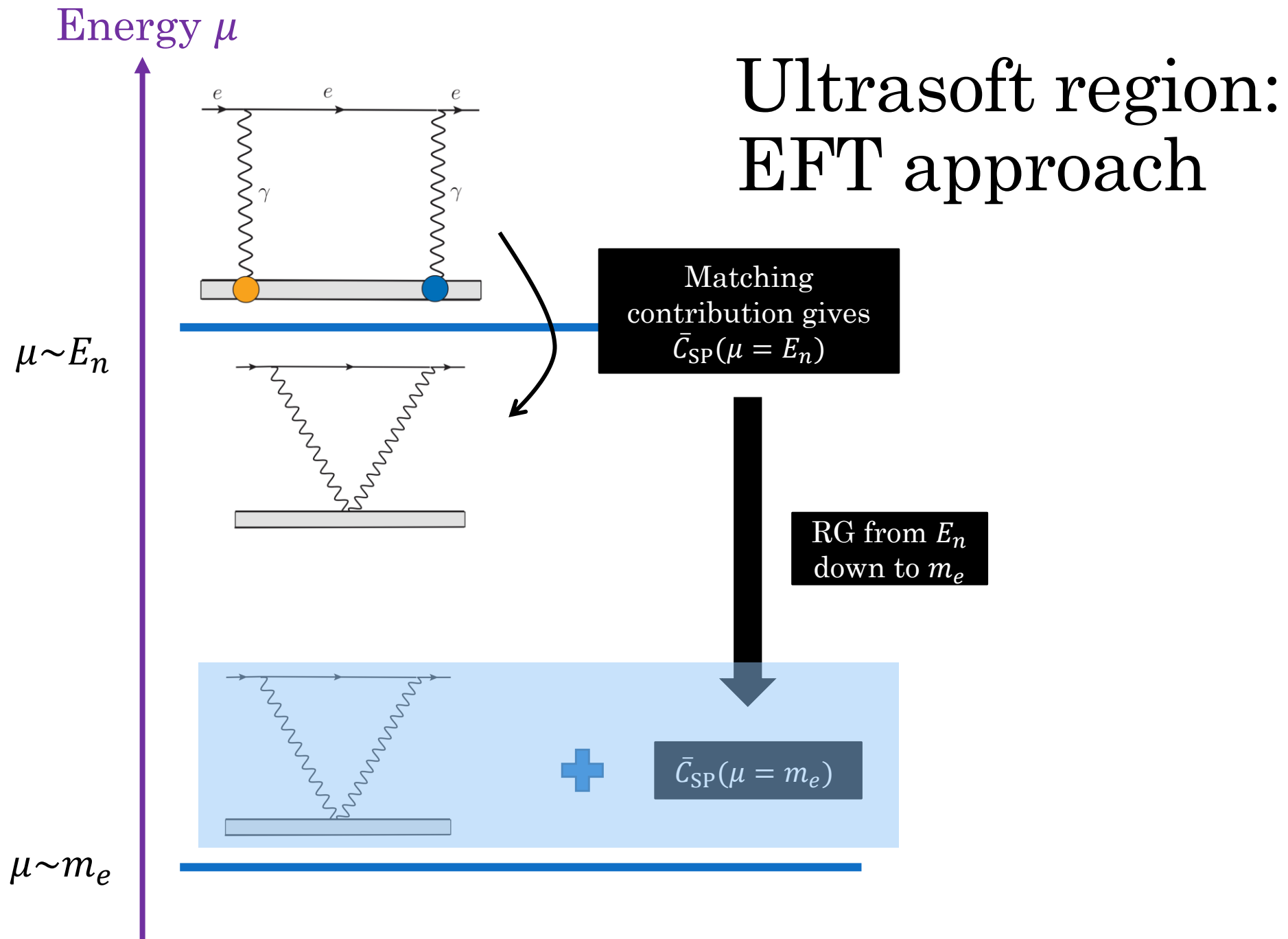


$\bar{C}_{SP}$



Another relevant momentum region:

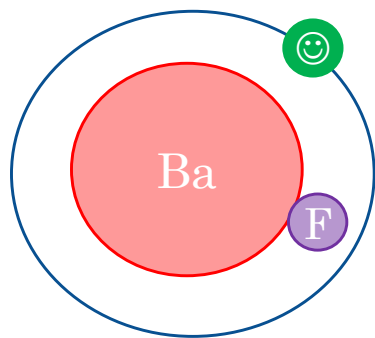
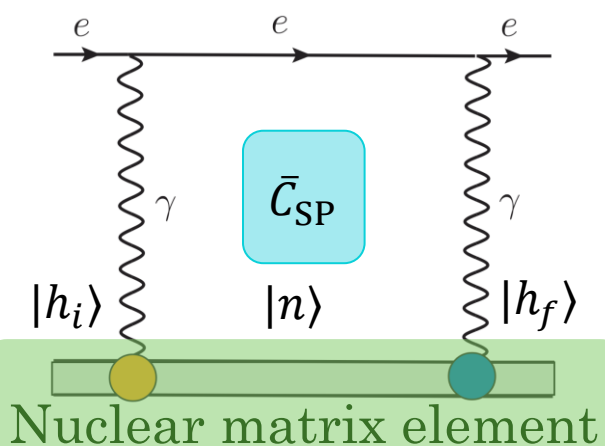
- **Potential** photon momenta $q_\gamma^0 \sim \vec{q}_\gamma^2/m_N$, $|\vec{q}_\gamma| \sim m_\pi$
- Expand in small scales like $q_\gamma^0/|\vec{q}_\gamma|$
- Only two-nucleon contributions



Results from the nuclear side



Ultrasoft region



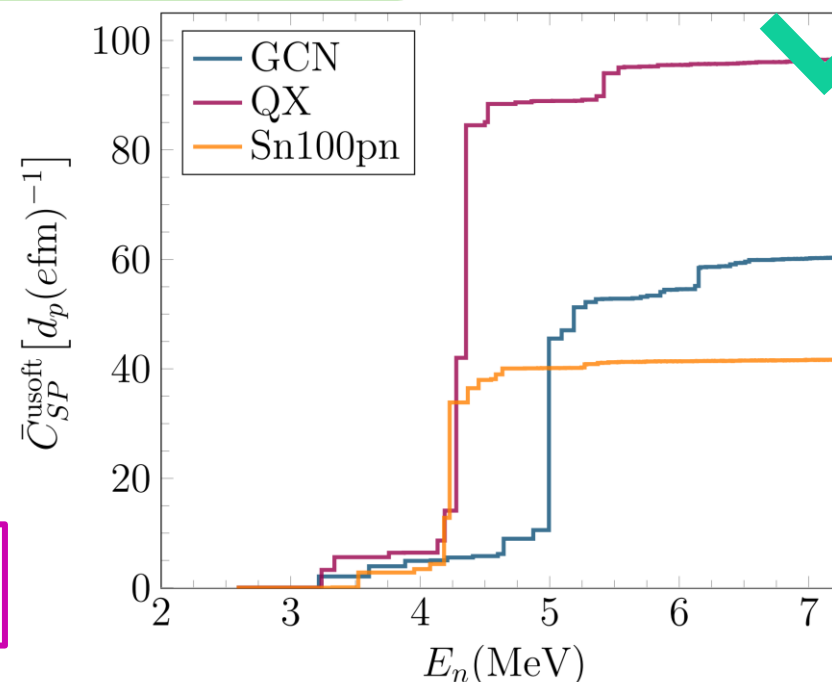
We obtain:

$$\bar{C}_{SP}^{\text{usoft}} = -\frac{\sqrt{2}\alpha^2 m_e}{3m_N G_F} M_{SP}^{\text{usoft}} \quad \text{EDM} \quad \text{MDM} \quad \Delta_n = E_n - E_i$$

$$M_{SP}^{\text{usoft}} = \sum_n \frac{\langle h_f | D^{(i)} \vec{\sigma} | n \rangle \cdot \langle n | \mu^{(i)} \vec{\sigma} | h_i \rangle}{\Delta_n} \left(1 + 3 \ln \frac{4\Delta_n^2}{m_e^2} \right)$$

- Calculate M_{SP}^{usoft} and Δ_n for Ba-138 in nuclear shell model
- Check approximation $\Delta_n \gg m_e$
- Ba-138 has magic neutron number $\rightarrow$ only d_p contributes

$$\bar{C}_{SP}^{\text{usoft}} = (67 \pm 28) d_p (e \text{ fm})^{-1}$$



Results from the nuclear side



We obtain:

$$\bar{C}_{SP}^{\text{pot}} = -\frac{\sqrt{2}}{G_F} \langle h_i | V | h_i \rangle$$

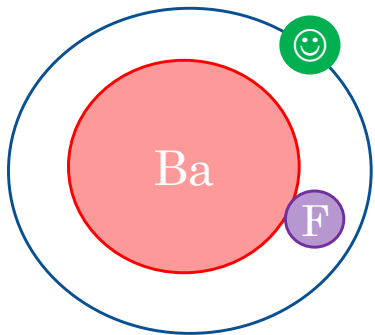
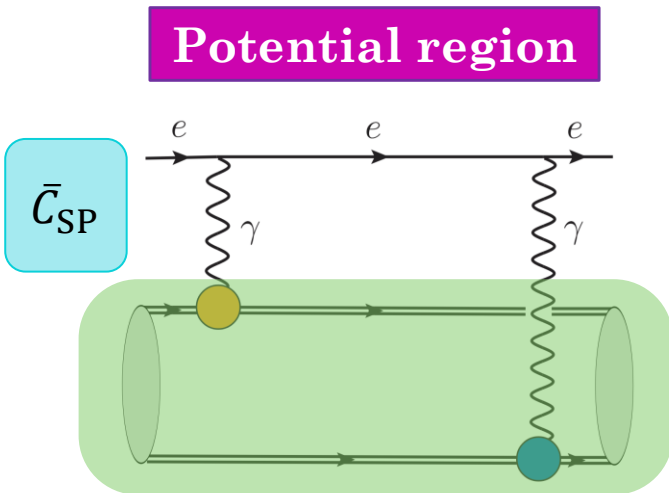
$$V(\vec{r}) = -\frac{e^4 m_e}{18\pi m_N} \sum_{i \neq j} \mu^{(i)} D^{(j)} |\vec{r}| \left[\sigma^{(i)} \cdot \sigma^{(j)} + \frac{1}{16} S^{(ij)}(r) \right]$$

~~Small~~

Again: calculate NME for Ba-138 in nuclear shell model

- Valence **and** core nucleons contribute
- Coherence: count p and n pairs. Consequences:
 1. $\text{NME} \sim c_1 Z d_p + c_2 N d_n$
 2. Potential > ultrasoft
 3. Potential less dependent on nuclear structure $\rightarrow$ small error

$$\bar{C}_{SP}^{\text{pot}} = [(-433 \pm 5) d_p + (387 \pm 0.4) d_n] (e \text{ fm})^{-1}$$



Scaling!

- Explore scaling of potential NME with Z and N

$$M_{SP}^{\text{pot}} = m_{SP}^{\text{pot},p} d_p + m_{SP}^{\text{pot},n} d_n$$

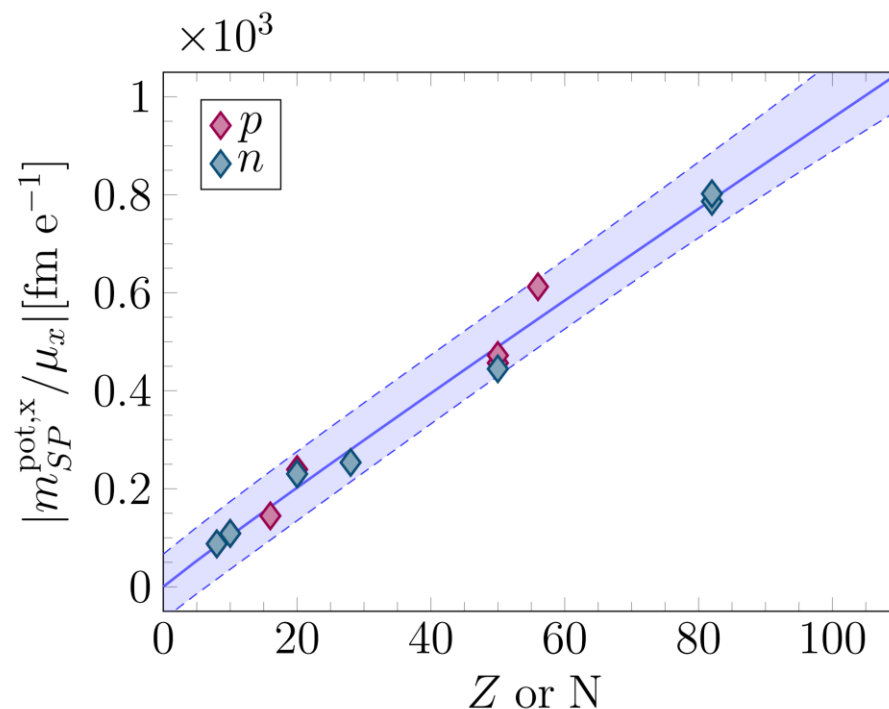
- Lighter nuclei confirm scaling (up to ~20%)
- What about heavier nuclei: Th, Hf?

Connect to experiments:

$$\omega = a C_{SP} + b d_e \rightarrow \frac{\omega}{b} = d_e + \frac{a}{b} C_{SP}$$

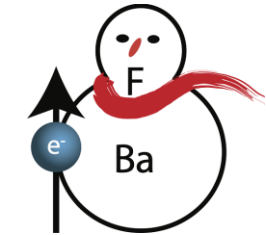
$$d_e^{\text{equiv}} \equiv r_{\text{mol}} \bar{C}_{SP} / A$$

$$d_e^{\text{equiv}} = \frac{\sqrt{2} e^4 m_e}{18 \pi G_F m_N} \frac{r_{\text{mol}}}{A} (-9.74) [Z \mu_p d_p + N \mu_n d_n] \frac{\text{fm}}{e}$$



Extrapolate from
BaF to HfF⁺

Results: bounds



- Projected: $|d_e|_{\text{BaF}} < 10^{-30} e \text{ cm}$ $\left\{ \begin{array}{l} |d_p|_{\text{BaF}} < 8.4 \cdot 10^{-24} e \text{ cm} \\ |d_n|_{\text{BaF}} < 8 \cdot 10^{-24} e \text{ cm} \end{array} \right.$

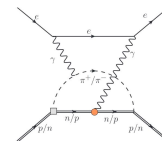
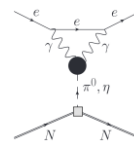
d_p and d_n
bounds

- With scaling function & $|d_e|_{\text{HfF}^+} < 4.1 \cdot 10^{-30} e \text{ cm}$ $\left\{ \begin{array}{l} |d_p|_{\text{HfF}^+} < 1.6 \cdot 10^{-23} e \text{ cm} \\ |d_n|_{\text{HfF}^+} < 1.6 \cdot 10^{-23} e \text{ cm} \end{array} \right.$

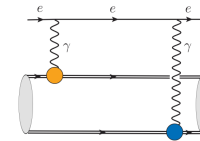
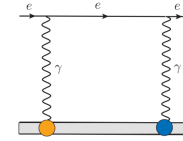
- Recall current limits $\left\{ \begin{array}{l} |d_p|_{\text{Hg}} \lesssim 10^{-25} e \text{ cm} \\ |d_n| \lesssim 10^{-26} e \text{ cm} \end{array} \right.$

$\bar{\theta}$ bound

- Unfortunate cancellation between

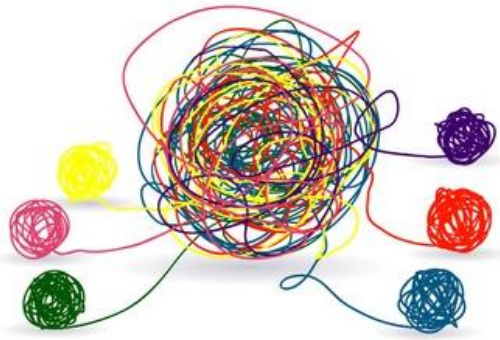


and



$\rightarrow \bar{\theta}$ limit weaker? More precision needed

Results: disentangle



Recall:

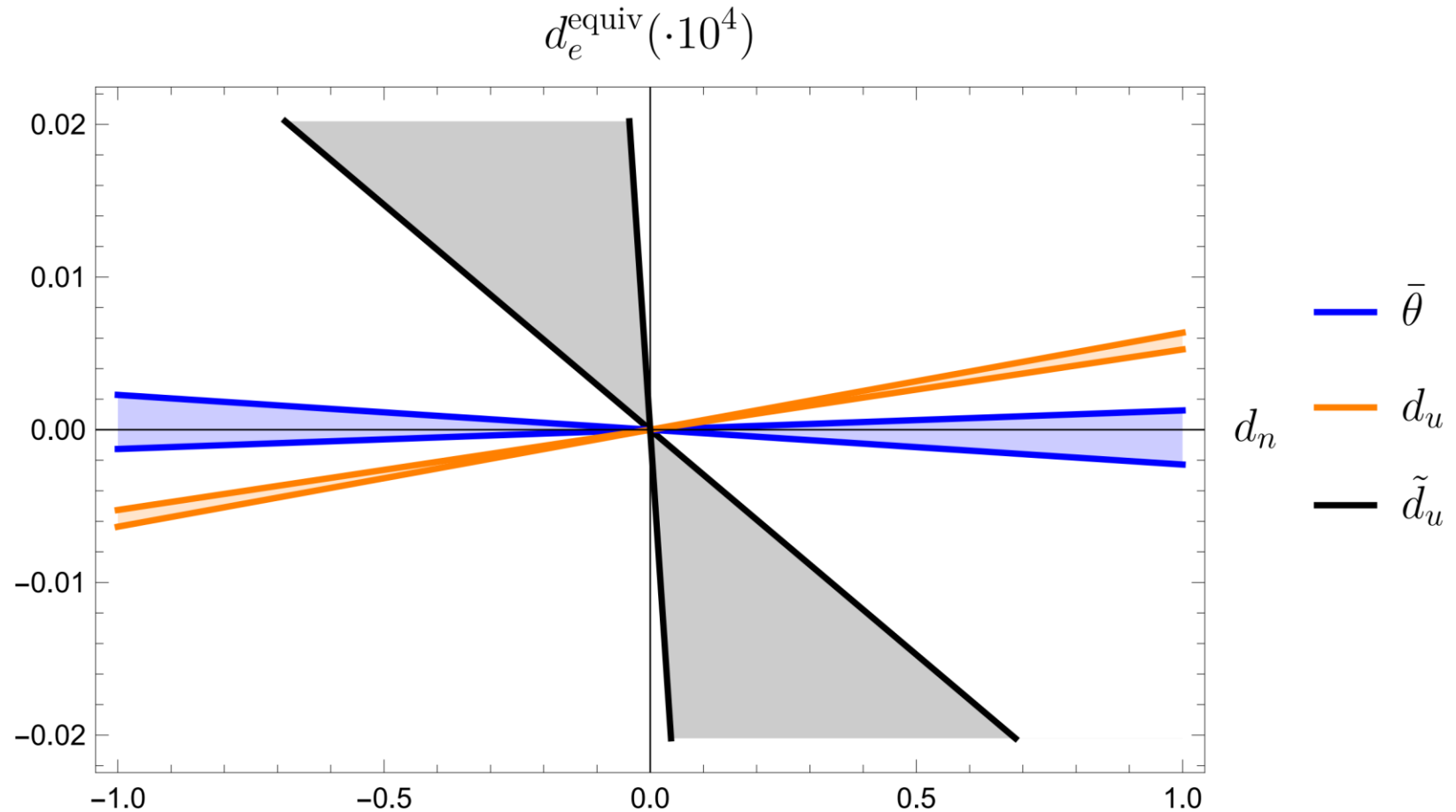
$$d_n = c_1 \bar{\theta} \text{ and } \omega = c_2 \bar{\theta}$$

$$d_n = c_3 d_q \text{ and } \omega = c_4 d_q$$



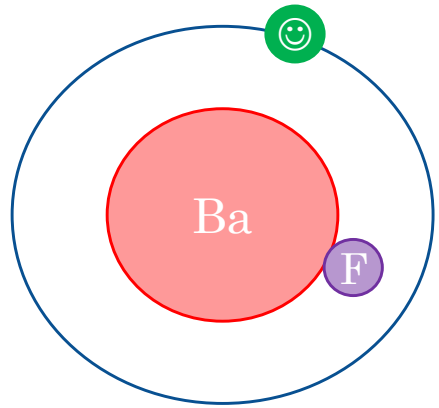
$$d_n/\omega$$

or actually d_n/d_e^{equiv}



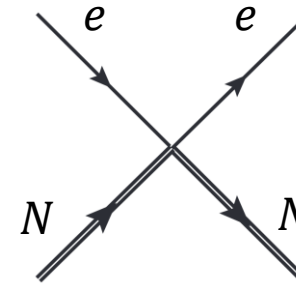
Summary

Part 1: hadronic CPV in paramagnetic EDMs

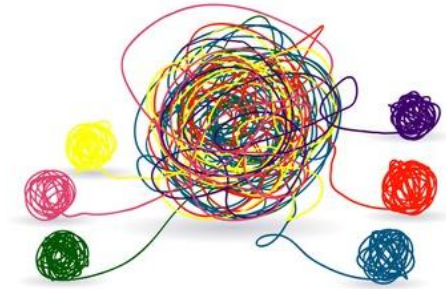


$$\omega = a C_{SP} + b d_e$$

Hadronic CPV



$$\mathcal{L} = C_{SP}^N \frac{G_F}{\sqrt{2}} \bar{e} i \gamma_5 e \bar{N} N$$

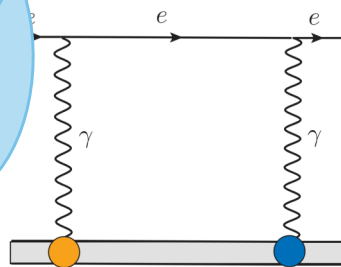


- CP-odd e-N interactions $\rightarrow$ paramagnetic molecule EDMs probe hadronic CPV
- Paramagnetic experimental precision advance $\rightarrow$ independent bounds crucial for disentangling possible CP-odd sources

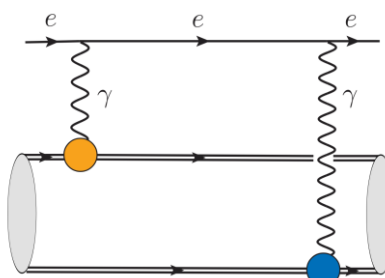
Part 2: electron-nucleus interactions

Outlook

- NME calculations in heavier systems (Th, Hf)
- Higher-order chiral corrections, e.g. to potential expression



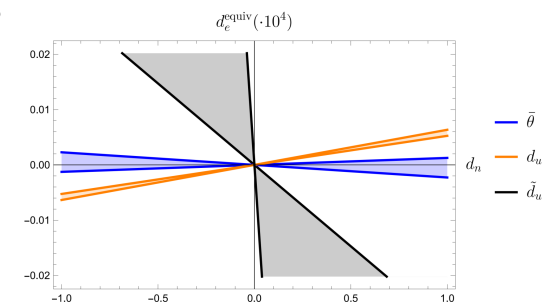
Ultrasoft: EFT without nuclear excited states



Potential: coherent scaling with Z & N $\rightarrow$ dominant

Results:

- Independent bounds on $d_p, d_n, \bar{\theta}$
- Disentangling of CP-odd sources



Thank you for listening!
Any questions?